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# Structural health monitoring with non-linear sensor measurements robust to unknown non-stationary input forcing

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## Abstract

Bayesian filtering based structural health monitoring algorithms typically assume stationary white Gaussian noise models to represent an unknown input forcing. However, typical structural damages occur mostly under the action of extreme loading conditions, like earthquake or high wind/waves, which are characteristically non-stationary and non-Gaussian. Clearly, this invalidates this basic assumption, causing these algorithms to perform poorly under non-stationary noise conditions. This paper extends an existing interacting filtering algorithm to efficiently estimate structural damages while being robust to unknown non-stationary non-Gaussian input forcing. Furthermore, this approach is generalized beyond linear measurements to encompass the case of non-linear measurements such as strains. The joint estimation of state and parameters is performed by combining Ensemble Kalman filtering, for non-linear system state estimation, and Particle filtering to estimate changes in the structural parameters. The robustness against input forcing is achieved through an output injection approach embedded in the state filter equation. Numerical simulations for two kinds of response measurements (acceleration and strain) are performed on a 3D frame structure under different damage location and severity scenarios. The sensitivity with respect to noise and the impact of different sensor combinations have also been investigated.

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## 1. Introduction

The objective of this paper is to monitor mechanical structures excited by some unknown time-varying input forces, by analyzing the measurements collected by means of some linear and non-linear sensors. To ascertain safety in structures, damages due to strong forces or extreme service conditions should be detected immediately after their occurrence. Traditionally, structural health monitoring (SHM) employs deterministic approaches for real-time structural damage detection with no consideration about uncertainties originating from unavoidable model inaccuracies, sensor noises and unknown external disturbances. This limits the utility of the methods for real-field applications. Bayesian filtering has been proved as an efficient alternative in this regard.

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30 For dynamics of mechanical structure, idealized as a Markov process and defined in state-space form  
31 with state variables  $\mathbf{x}_{1:k}$  observed for a time span of  $[1 : k]$  through a measurement sequence  $\mathbf{y}_{1:k}$ , Bayesian  
32 filtering employs two probabilistic models: i) for system state propagation  $\mathbf{p}(\mathbf{x}_k|\mathbf{x}_{k-1})$  formulated by the  
33 Chapman–Kolmogorov equation as  $\mathbf{p}(\mathbf{x}_k, \mathbf{x}_{k-1}) = \mathbf{p}(\mathbf{x}_{k-1}) \times \mathbf{p}(\mathbf{x}_k|\mathbf{x}_{k-1})$  and, ii) likelihood estimation model  
34 of states, i.e.,  $\mathbf{p}(\mathbf{y}_k|\mathbf{x}_k)$  for state estimate correction, to estimate recursively the states using the measurement  
35 sequence  $\mathbf{y}_{1:k}$ . The sequence  $\mathbf{y}_{1:k}$  can be linear, non-linear or mixed structural response through which  
36 the structural health, parameterized with a set of location-based structural health indices (**HI**s), can be  
37 interpreted. Eventually, SHM with Bayesian filtering is posed as a joint probability estimation of state  
38 and parameters to obtain  $\mathbf{p}(\mathbf{x}_k, \boldsymbol{\theta}_k|\mathbf{y}_{1:k})$ , in order to detect, localize and quantify the damages; where the  
39 additional states,  $\boldsymbol{\theta}_k$ , signify the **HI**s.

40 There exist successful applications of Bayesian filtering in SHM wherein **HI**s in  $\boldsymbol{\theta}_k$  are tracked to localize  
41 any deterioration in structural health [31, 36, 38, 42]. In the related literature,  $\boldsymbol{\theta}_k$ s are mostly augmented  
42 in the state definition as  $\mathbf{X}_k = [\mathbf{x}_k; \boldsymbol{\theta}_k]$  to estimate them alongside  $\mathbf{x}_k$  [5, 9–11]. Yet, owing to the induced  
43 non-linearity and/or the loss of observability, this approach is reported to cause divergence, leading to  
44 false or infeasible solutions, especially for time varying systems [11]. Recently, an interacting filtering  
45 strategy has emerged as a reliable alternative to tackle time varying systems with moderate state size.  
46 With this approach, a conditional posterior distribution estimation for the system states is followed by a  
47 marginal posterior distribution of the system parameters (also known as Rao-Blackwellisation) [4, 13, 35, 38].  
48 The advantage of the interacting approach, in terms of computational burden and stability, over the joint  
49 estimation approach has been discussed in [8].

50 Nevertheless, the likelihood estimation function,  $\mathbf{p}(\mathbf{y}_k|\boldsymbol{\theta}_k)$ , for parameter estimation problems, is typically  
51 a non-linear mapping of  $\mathbf{x}_k / \boldsymbol{\theta}_k$  (e.g., finite element (FE) models) for which only non-linear filter variants  
52 (e.g., Extended (EKF) [21], Unscented (UKF) [23, 30], Ensemble Kalman filter (EnKF), Particle filter (PF),  
53 etc.) are applicable. PF [18] has been recognized as a powerful approach in this endeavor [2, 7, 9] with some  
54 concerns regarding its computational expense [39]. To manage the computational expense, the Interacting  
55 Particle-Kalman filter (IPKF) [36, 38] was introduced, that efficiently handles the linear state estimation  
56 with linear measurements (e.g. acceleration, displacement, etc.) using a standard KF while PF handles the  
57 non-linear parameter estimation. However, because of the use of KF, the applicability of IPKF is limited to  
58 linear systems (linear state propagation and measurements) only.

59 To generalize the application of IPKF to non-linear systems, KF therefore should be replaced with non-  
60 linear filter variants like EKF/UKF/EnKF. As an alternative, EKF employs approximate local lineariza-  
61 tion through first order Taylor approximation. The associated Jacobian calculation, even being compute-  
62 intensive, does not usually hinder estimation for moderate sized systems. Nevertheless, some comparative  
63 studies [19, 20] identified that EKF’s approximate closure scheme may lack accuracy and hamper the de-  
64 tection promptness. With UKF, on the other hand, uncertainty propagation through a sparse set of “sigma

65 points” limits its performance for severely non-linear systems [20, 40]. EnKF employs a set of ensembles –  
66 realized from the entire domain of states – for uncertainty propagation while preserving the non-linearity in  
67 state transition. It also offers flexibility to enhance the accuracy through the employment of bigger ensemble  
68 pools. The selection of EnKF in this study takes basis on the comparative study of [20].

69 Filtering based SHM techniques typically idealize the unknown input force as a Stationary White Gaus-  
70 sian Noise (SWG N). Contrarily, the exogenous forces, that can potentially damage civil infrastructures (e.g.,  
71 seismic forcing or heavy wind/waves), are mostly unforeseen and do not satisfy this assumption. Clearly,  
72 to ensure estimation accuracy and subsequently the structural safety, the adopted SHM technique has to  
73 be robust against input forcing. [24] proposed an unbiased minimum-variance linear state estimation filter  
74 that does not require prior information about the unknown input. This was later improved for practical  
75 application by [22]. Joint state and input estimation has been employed by [16] for a system without a direct  
76 transmission term, and later updated for a system with a direct transmission term [17]. To avoid numerical  
77 instability especially for the systems that are redundantly instrumented, [26] proposed similar filters that  
78 jointly estimate states and inputs. However, it was mentioned that the unobservability in the system may  
79 lead to an estimation instability, for this augmentation strategy.

80 [3] proposed a dual PF for exogenous force estimation for time invariant systems. [1] combined the  
81 parameters and the inputs together in a very large state vector in order to estimate the seismic excitation  
82 acting on a linear time invariant (LTI) system. The estimation of the input alongside the system matrices  
83 and states have been done by [38]. An analogous approach focusing on the estimation of the input statistics  
84 has been presented in [37] for systems with varying noise level. [15] presents an auto-covariance least square  
85 based method for estimating noise covariances online for linear time varying (LTV) as well as non-linear  
86 systems.

87 [6] proposed an UKF based algorithm that employs a time varying auto-regressive model to jointly  
88 estimate the structural parameters and the unknown inputs. For time varying systems, [29] generalized  
89 their earlier proposal [28] of a smoothing algorithm for joint estimation of state, parameters and input.  
90 With a similar objective, [12] proposed a dual filtering approach in which the structural parameters, as  
91 augmented states, are jointly estimated with the response states conditioned on an estimate for the input  
92 force. Instead of an explicit reconstruction of the input forces, [34] estimated the input force model through  
93 the parameters of a Gaussian process within a Bayesian framework.

94 Recently, for LTV systems with known system matrices, instead of estimating the input time histories  
95 [12, 29, 38], or its statistics [34], [43] developed a robust and stable linear state estimator unaffected by  
96 unknown inputs. This unknown input rejection approach is used in Kalman filter environment with linear  
97 measurements. With an intent to generalize earlier works of IPKF [42] for non-linear systems with an added  
98 robustness against unknown external disturbances, PF has been coupled with EnKF. While the authors’  
99 previous works [38] approached this input robustness through an explicit estimation of inputs online, the

100 present article rejects the effect of input variability through output injection for a modified state transition  
101 equation following [43].

102 Rejection of unknown inputs is usually addressed by unknown input observers, with a rich literature on  
103 this topic (see [43] and references therein). While most unknown input observers are limited to linear systems,  
104 they do not consider parameter estimation in addition to state estimation. The methods for designing such  
105 unknown input observers cannot be applied to EnKF and PF. The method for unknown input rejection  
106 adopted in this paper is not limited to a particular unknown input observer design. It rejects unknown  
107 inputs by simply transforming the system model, so that state and/or parameter estimation algorithms  
108 can be applied as if the rejected unknown inputs no longer exist. This output injection has therefore  
109 been embedded within the EnKF environment to yield input-robust response state estimates for non-linear  
110 systems. In this process, the present article also overcomes the limitation of IPKF [38] of using linear  
111 measurements (e.g. acceleration) only, and extend its applicability to non-linear measurements (e.g. strain)  
112 as well. The novel contributions of the study can therefore be listed as: a novel noise robust Interacting  
113 Particle-Ensemble Kalman Filter (rIP-EnKF) algorithm in which, i) PF coupled with EnKF extends the  
114 reach of [38] to non-linear systems, ii) noise robust EnKF ensures rejection of unknown non-stationary  
115 excitation and finally, iii) linear and non-linear measurements can be dealt with, simultaneously.

116 The proposed algorithm is predictor model-based that involves a precise model (preferably a calibrated  
117 FE model) for state propagation. At least a few accelerometer measurement channels are assumed to be  
118 available for perfect functioning of the algorithm. Also, sensor noise statistics corresponding to accelerome-  
119 ters and strain gauges are assumed to be available. The system dynamics is assumed to remain linear even  
120 after damage.

121 It should be noted here, that stability of state estimation algorithms with unknown input rejection is  
122 studied in the literature of unknown input observers for LTI systems and for some affine parameter varying  
123 systems. To our knowledge, the only stability analysis of such algorithms for general LTV systems has been  
124 proposed in [43]. The generalization of this analysis to EnKF with output injection, applied in this paper to  
125 address non-linearities, is certainly an important and difficult task clearly outside of the scope of this paper.

126 In the following, based on the non-linear state-space modeling detailed in Section 2, output injection is  
127 demonstrated in Section 3, followed by the detailed proposal in Section 4. The proposed approach is tested  
128 on a numerical simulation in Section 5.

## 129 **2. Modeling and system dynamics**

130 The stiffness and damping matrices of LTV mechanical system, i.e.,  $\mathbf{K}(\boldsymbol{\theta}(t))_{n \times n}$  and  $\mathbf{C}(\boldsymbol{\theta}(t))_{n \times n}$ , are  
131 functions of time varying location based **HI**s represented by the time varying parameter vector  $\boldsymbol{\theta}(t)_{N_\theta \times 1}$ .  
132 The dynamics of such a system under seismic excitation, in state space domain, can be described by a time  
133 invariant mass matrix  $\mathbf{M}$ , a time varying stiffness matrix  $\mathbf{K}(\boldsymbol{\theta}(t))$  and damping matrix  $\mathbf{C}(\boldsymbol{\theta}(t))$  [42]. Defining

134  $\mathbf{B}_c = \begin{bmatrix} \mathbf{0}_{n \times m} \\ \mathbf{M}^{-1} \end{bmatrix}_{2n \times m}$ ,  $\mathbf{E}_c = \begin{bmatrix} \mathbf{0}_{n \times r} \\ \boldsymbol{\tau} \end{bmatrix}_{2n \times r}$  and  $\mathbf{F}(t) = \begin{bmatrix} \mathbf{0}_{n \times n} & \mathbf{I}_{n \times n} \\ -\mathbf{M}^{-1}\mathbf{K}(\boldsymbol{\theta}(t)) & -\mathbf{M}^{-1}\mathbf{C}(\boldsymbol{\theta}(t)) \end{bmatrix}_{2n \times 2n}$ , the  
 135 dynamic equation [1] is defined as follows,

$$\dot{\mathbf{x}}(t) = \mathbf{F}(t)\mathbf{x}(t) + \mathbf{B}_c\mathbf{u}(t) + \mathbf{E}_c\ddot{\mathbf{a}}^g(t) \quad (1)$$

136 where system state,  $\mathbf{x}(t) = [\mathbf{q}(t) \quad \dot{\mathbf{q}}(t)]_{2n \times 1}^T$ , with  $\mathbf{q}(t)$  and  $\dot{\mathbf{q}}(t)$  representing the displacement and velocity  
 137 responses.  $\mathbf{0}$  and  $\mathbf{I}$  are null and identity matrices of mentioned dimensions, respectively.  $\mathbf{u}(t)_{m \times 1}$  is encom-  
 138 passing both the process noise and the ambient force acting on the structure and will be collectively defined  
 139 as process noise from now on. It is assumed that,  $\mathbf{u}(t)_{m \times 1}$  has known statistics and can be modelled as an  
 140 SWGN of constant covariance  $\mathbf{Q}_{m \times m}$ , which takes into account both ambient forces and model uncertainty.  
 141  $\ddot{\mathbf{a}}^g(t)_{r \times 1}$  represents an unknown disturbance (e.g. seismic ground motion) which is an *unknown arbitrary*  
 142 function of  $t$ , without any assumed statistical property.  $r$  is the number of channels for the disturbance  
 143 input.

144 The measurement  $\mathbf{y}(t)$  can be a linear mapping (denoted here as linear measurement,  $\mathbf{y}^l(t)$ , e.g., relative  
 145 acceleration, displacement, etc.) or a non-linear mapping (denoted here as non-linear measurement,  $\mathbf{y}^{nl}(t)$ ,  
 146 e.g., dynamic strain, etc.) of the state variable  $\mathbf{x}(t)$ . The adopted nomenclature for this measurement is in  
 147 line with [25, 32, 33, 44]. A mix of  $\mathbf{y}^l(t)$  or  $\mathbf{y}^{nl}(t)$  is also possible. The present study adopts a combination  
 148 of acceleration, as  $\mathbf{y}^l(t)$ , and strain, as  $\mathbf{y}^{nl}(t)$ , as measurements. The equation for the relative accelerations  
 149  $\ddot{\mathbf{q}}_p(t)$ , at  $p$  accelerometers, can be written as,

$$\ddot{\mathbf{q}}_p(t) = \mathfrak{L}[\mathbf{H}(t)\mathbf{x}(t) + \mathbf{D}\mathbf{u}(t)] + \mathbf{L}\ddot{\mathbf{a}}^g(t) + \mathbf{w}^a(t) \quad (2)$$

150 LTV measurement model,  $\mathbf{H}(t)_{n \times 2n} = \begin{bmatrix} -\mathbf{M}^{-1}\mathbf{K}(\boldsymbol{\theta}(t)) & -\mathbf{M}^{-1}\mathbf{C}(\boldsymbol{\theta}(t)) \end{bmatrix}$  and LTI direct transmission ma-  
 151 trix,  $\mathbf{D}_{n \times m} = \mathbf{M}^{-1}$ , maps  $(2n \times 1)$  order states and  $(m \times 1)$  order inputs to the corresponding  $n$  order  
 152 acceleration response,  $\ddot{\mathbf{q}}(t)$ , at every *dof* (Degree of Freedom).  $\mathbf{L}_{p \times r}$  maps the direct impact of  $r$  distur-  
 153 bance channels to  $p$  output channels. The location matrix  $\mathfrak{L}_{p \times n}$  selects  $p$  measured *dofs* from  $n$ .  $\mathbf{w}^a(t)_{p \times 1}$ ,  
 154 representing the measurement noise in  $p$  accelerometers, is an SWGN process of covariance  $\mathbf{R}_{p \times p}^a$ .

155 In practice, for a base excited structure fitted with accelerometers patched onto its surface, the relative  
 156 acceleration,  $\ddot{\mathbf{q}}(t) / \ddot{\mathbf{q}}_p(t)$ , can never be measured due to lack of fixed reference. Clearly, for such cases, the  
 157 measured acceleration  $\mathbf{y}^l(t)$  is a summation of  $\ddot{\mathbf{q}}_p(t)$  with a contribution from the base acceleration  $\ddot{\mathbf{a}}^g(t)$ .  
 158 The measurement equation, with respect to measured acceleration,  $\mathbf{y}^l(t)$ , can be presented as,

$$\mathbf{y}^l(t) = \ddot{\mathbf{q}}_p(t) - \mathbf{L}\ddot{\mathbf{a}}^g(t) = \mathfrak{L}\{\mathbf{H}(t)\mathbf{x}(t) + \mathbf{D}\mathbf{u}(t)\} + \mathbf{w}^a(t) \quad (3)$$

159 Sampled at discrete time instants indexed by  $k = 0, 1, 2, \dots$ , Equations (1) and (3) then lead to the discrete

160 time state-space model [1],

$$\begin{aligned}\mathbf{x}_k &= \mathbf{F}_k \mathbf{x}_{k-1} + \mathbf{B}_k \mathbf{u}_k + \mathbf{E}_k \ddot{\mathbf{a}}_k^g; \\ \mathbf{y}_k^l &= \mathbf{H}_k \mathbf{x}_k + \mathbf{D}_k \mathbf{u}_k + \mathbf{w}_k^a\end{aligned}\quad (4)$$

161 where  $\mathbf{F}_k$ ,  $\mathbf{B}_k$ ,  $\mathbf{D}_k$ ,  $\mathbf{E}_k$ ,  $\mathbf{H}_k$ ,  $\mathbf{x}_k$ ,  $\mathbf{y}_k^l$  and  $\mathbf{w}_k^a$  are the discrete time counterparts of the corresponding continuous  
162 time entities described above, obtained through zero-order-hold technique. The location matrix  $\mathcal{L}$  has been  
163 incorporated into  $\mathbf{H}_k$ ,  $\mathbf{D}_k$  and  $\mathbf{w}_k^a$  for simplicity.

164 Following the time indexing scheme employed in [41], the inputs  $\mathbf{u}_k$  and  $\ddot{\mathbf{a}}_k^g$ , sampled at the  $(k-1)^{th}$   
165 time instant, take part in state transition from  $\mathbf{x}_{k-1}$  (past) to  $\mathbf{x}_k$  (current) which are observed at the current  
166 time instant as  $\mathbf{y}_k^l$ .  $\mathbf{F}_k$  and  $\mathbf{H}_k$ , defined at the current time instant, are functions of  $\boldsymbol{\theta}_k$  that remain constant  
167 over the time interval  $k-1$  to  $k$ . The process uncertainty is carried over to the  $k^{th}$  instant and eventually  
168  $\mathbf{y}_k^l$  is contaminated with  $\mathbf{w}_k^a$  at the  $k^{th}$  instant. The same time index formalism is also used in [38].

169 Strain measurements are additionally included in this study for system health estimation. The motivation  
170 behind this consideration is the relatively low cost of strain gauges over accelerometers. Discrete time  
171 strain response,  $\boldsymbol{\varepsilon}_k^{xx}$ , denoted as discrete non-linear measurement  $\mathbf{y}_k^{nl}$ , measured at  $s$  strain gauges, can be  
172 obtained by mapping the nodal displacements,  $\mathbf{q}_k$  (a subset of  $\mathbf{x}_k$ ) through the non-linear strain-displacement  
173 relationship  $\mathbf{y}_k^{nl} = f(\mathbf{x}_k) + \mathbf{w}_k^g$ ; with  $\mathbf{w}_k^g$  being the sensor noise in strain gauges, modelled as an SWGN  
174 of covariance  $\mathbf{R}_{s \times s}^g$ . A specific case of strain-displacement function for an Euler-Bernoulli beam has been  
175 demonstrated in Appendix A.

176 The strain response  $\boldsymbol{\varepsilon}_k^{xx}$  is then added into the state vector in order to estimate them alongside  $\mathbf{x}_k$   
177 resulting in an extended state vector,  $\mathbf{X}_k = \begin{bmatrix} \mathbf{x}_k & \boldsymbol{\varepsilon}_{k-1}^{xx} \end{bmatrix}^T$ . With a non-linear state mutation function  
178  $\tilde{F}_k(\mathbf{X}_{k-1}) = \begin{bmatrix} \mathbf{F}_k \mathbf{x}_{k-1} \\ f(\mathbf{x}_{k-1}) \end{bmatrix}$ , where  $f(\bullet)$  is the non-linear strain-displacement mapping mentioned earlier, the  
179 process and measurement equations can be redefined as,

$$\begin{aligned}\mathbf{X}_k &= \tilde{F}_k(\mathbf{X}_{k-1}) + \tilde{B}_k \mathbf{U}_k + \tilde{E}_k \mathcal{S}_k^g \\ \mathbf{Y}_k &= \tilde{H}_k \mathbf{X}_k + \tilde{D}_k \mathbf{U}_k + \mathbf{W}_k\end{aligned}\quad (5)$$

180 where the linear system matrices have been defined as  $\tilde{B}_k = \begin{bmatrix} \mathbf{B}_k \\ 0 \end{bmatrix}$ ,  $\tilde{E}_k = \begin{bmatrix} \mathbf{E}_k \\ 0 \end{bmatrix}$ ,  $\tilde{H}_k = \begin{bmatrix} \mathbf{H}_k & \mathbf{0}_{p \times s} \\ \mathbf{0}_{s \times 2n} & \mathbf{I}_{s \times s} \end{bmatrix}$ ,  
181  $\tilde{D}_k = \begin{bmatrix} \mathbf{D}_k \\ 0 \end{bmatrix}$ , and the random processes as  $\mathbf{Y}_k = \begin{bmatrix} \mathbf{y}_k^l \\ \mathbf{y}_k^{nl} \end{bmatrix}$ ,  $\mathbf{U}_k = \begin{bmatrix} \mathbf{u}_k \\ 0 \end{bmatrix}$ ,  $\mathcal{S}_k^g = \begin{bmatrix} \ddot{\mathbf{a}}_k^g \\ 0 \end{bmatrix}$ ,  $\mathbf{W}_k = \begin{bmatrix} \mathbf{w}_k^a \\ \mathbf{w}_k^g \end{bmatrix}$ .

182 **3. Unknown input rejection from system dynamics**

183 While the process noise  $\mathbf{U}_k$  can be well modeled as a random noise, typically with an assumed Gaussian  
 184 distribution, the unknown input  $\mathcal{S}_k^g$  (e.g. seismic excitation) is totally arbitrary, without any prior statistical  
 185 property, and not necessarily random. A major goal of this paper is to design an SHM approach robust to  
 186 the unknown input  $\mathcal{S}_k^g$  without assuming or estimating its statistical properties. This study thus neither  
 187 reconstructs the unknown input  $\mathcal{S}_k^g$  as in [12, 26, 27, 29, 38], nor estimates its statistics as in [15, 34].  
 188 Following the ideas of [43], the proposed approach ensures robustness against the unknown input  $\mathcal{S}_k^g$  by  
 189 rejecting its effect from the system dynamics by means of an output injection, as demonstrated in the  
 190 following.

191 Owing to the output equation in (5), the following holds true with an arbitrary bounded matrix  $\mathbf{G}_k \in$   
 192  $\mathbb{R}^{2n+s \times p+s}$ :

$$0 = \mathbf{G}_k \left( \mathbf{Y}_k - \tilde{H}_k \mathbf{X}_k - \tilde{D}_k \mathbf{U}_k - \mathbf{W}_k \right) \quad (6)$$

193 Setting,

$$\mathcal{L}_k = \mathbf{I}_{2n+s} - \mathbf{G}_k \tilde{H}_k, \quad (7)$$

194 the state equation in (5) is then rewritten as

$$\begin{aligned} \mathbf{X}_k &= \tilde{F}_k(\mathbf{X}_{k-1}) + \tilde{B}_k \mathbf{U}_k + \tilde{E}_k \mathcal{S}_k^g + \mathbf{G}_k \left( \mathbf{Y}_k - \tilde{H}_k \mathbf{X}_k - \tilde{D}_k \mathbf{U}_k - \mathbf{W}_k \right) \\ &= \mathcal{L}_k \tilde{F}_k(\mathbf{X}_{k-1}) + \mathcal{L}_k \tilde{E}_k \mathcal{S}_k^g + \mathcal{L}_k \tilde{B}_k \mathbf{U}_k + \mathbf{G}_k (\mathbf{Y}_k - \tilde{D}_k \mathbf{U}_k - \mathbf{W}_k) \end{aligned} \quad (8)$$

195 By choosing  $\mathbf{G}_k$  so that the matrix  $\mathcal{L}_k$  defined in (7) satisfies,

$$\mathcal{L}_k \tilde{E}_k = 0, \quad (9)$$

196 Equation (8) gets decoupled from  $\mathcal{S}_k^g$  and can be rewritten as,

$$\mathbf{X}_k = \mathcal{F}_k(\mathbf{X}_{k-1}) + \mathcal{B}_k \mathbf{U}_k + \mathbf{G}_k \mathbf{Y}_k + \mathbf{V}_k \quad (10)$$

197 with  $\mathcal{F}_k(X) = \mathcal{L}_k \tilde{F}_k(X)$ ,  $\mathcal{B}_k = \mathcal{L}_k \tilde{B}_k - \mathbf{G}_k \tilde{D}_k$  and  $\mathbf{V}_k = -\mathbf{G}_k \mathbf{W}_k$  modelled as an SWGN process of variance  
 198  $\mathbf{G}_k \begin{bmatrix} \mathbf{R}^a & \mathbf{0} \\ \mathbf{0} & \mathbf{R}^g \end{bmatrix} \mathbf{G}_k^T$ . The unknown input  $\mathcal{S}_k^g$  has disappeared from Equation (10), owing to an appropriate  
 199 injection of the known output  $\mathbf{Y}_k$  through  $\mathbf{G}_k$ . For Equation (9) to be valid,  $\mathbf{G}_k$  is chosen as:  $\mathbf{G}_k =$   
 200  $\tilde{E}_k (\tilde{H}_k \tilde{E}_k)^\dagger$ , where  $\dagger$  denotes Moore-Penrose Pseudo-inverse operation. It is assumed that the inverse of the  
 201 square matrix  $\tilde{E}_k^T \tilde{H}_k^T \tilde{H}_k \tilde{E}_k$  exists and is bounded, so that the Penrose Pseudo-inverse of  $\tilde{H}_k \tilde{E}_k$  is upper



202 bounded. In the particular case of time invariant matrix product  $\tilde{H}\tilde{E}$ , it is sufficient that  $\tilde{H}\tilde{E}$  is full column  
 203 rank. This assumption implies that  $p \geq r$ .

204 Notice that when no linear measurements are available (i.e.,  $p = 0$ ), Equation (6) yields  $\mathbf{G}_k = 0$ , thus  
 205 input rejection can not be achieved. With linear measurements only, rejecting  $\mathcal{S}_k^g$  is possible following the  
 206 lines of [43]. With both linear/non-linear measurements, as in the current state transition function in (10),  
 207 the rejection depends on the available number of linear measurements only with a condition of  $p \geq r$  to  
 208 ensure perfect robustness. Yet, due to the presence of non-linear measurements, the approach has to be  
 209 modified since standard KF is no longer an option. Eventually, the part of Equation (10) (excluding the  
 210 strain states) responsible for input rejection can be isolated as,

$$\mathbf{x}_k = \bar{\mathbf{F}}_k \mathbf{x}_{k-1} + \bar{\mathbf{B}}_k \mathbf{u}_k + \mathbf{G}_k^x \mathbf{y}_k + \bar{\mathbf{v}}_k \quad (11)$$

211 where,  $\bar{\mathbf{F}}_k = (\mathbf{I}_{2n} - \mathbf{G}_k^x \mathbf{H}_k) \mathbf{F}_k$ ,  $\bar{\mathbf{B}}_k = (\mathbf{I}_{2n} - \mathbf{G}_k^x \mathbf{H}_k) \mathbf{B}_k - \mathbf{G}_k^x \mathbf{D}_k$  and  $\bar{\mathbf{v}}_k = -\mathbf{G}_k^x \mathbf{w}_k^a$ . The new process  
 212 noise  $\bar{\mathbf{v}}_k$  in Equation (11) can still be defined as SWGN with an altered co-variance  $\mathbf{G}_k^x \mathbf{R}_k^a \mathbf{G}_k^{xT}$ , where  
 213  $(\mathbf{I}_{2n} - \mathbf{G}_k^x \mathbf{H}_k) \mathbf{E}_k = 0$ . Since the strain response states are non-linear explicit functions of the state subset  
 214  $\mathbf{x}_k$ , the induced robustness in  $\mathbf{x}_k$  will also ensure robustness in the strain response states. Notice that any  
 215 linear or non-linear measurement can be used in place of accelerations or strains respectively. This paper  
 216 thus extends the assumptions of [5] where only linear measurements are considered.

#### 217 4. Robust Interacting Particle-Ensemble Kalman Filtering

218 To generalize the IPKF methods [37, 38] to non-linear systems, the present work replaces KF within  
 219 IPKF with EnKF to handle the non-linear state estimation problem as in Equation (5). In the modified  
 220 interacting particle-ensemble Kalman filter (IP-EnKF), a set of EnKFs runs within an envelop of PF.  
 221 Within PF, each particle represents a parameter instance that defines the structural matrices to be used  
 222 in the process equation of the EnKF. Thus, both filters interact to obtain the conditional estimates for the  
 223 response states and parameters simultaneously.

##### 224 4.1. Envelop Parameter Filter

225 Bayesian belief propagation requires an explicit analytical integration to be performed over the entire  
 226 state domain. This task is simple with Gaussian states evolving through a linear state transition. The  
 227 current parameter estimation problem is, however, non-linear for which an explicit analytical integration  
 228 over the entire parameter space is not possible. PF attempts a particle approximation of this integration  
 229 by representing and propagating the parametric uncertainty through a cloud of  $N_p$  independent particles  
 230  $\Xi_k = [\xi_k^1, \xi_k^2, \dots, \xi_k^{N_p}]$  listing all individual particles as  $\xi_{k N_\theta}^i$ , where  $N_\theta$  denotes the number of health  
 231 parameters (**HI**) that are to be estimated. Additionally, no presumption on the stochastic nature of the

parameter states is enforced. The time evolution of these particles is a random perturbation around their current position  $\boldsymbol{\xi}_{k-1}^j$  along with a forced shift towards the current particle mean estimate  $\bar{\boldsymbol{\xi}}_{k-1}$  [38],

$$\boldsymbol{\xi}_k^j = \alpha \boldsymbol{\xi}_{k-1}^j + \mathbb{N}(\delta \boldsymbol{\xi}_k; \boldsymbol{\sigma}_k^\xi) \quad (12)$$

where a Gaussian blurring is performed on  $\boldsymbol{\xi}_{k-1}^j$  with a shift  $\delta \boldsymbol{\xi}_k = (1 - \alpha) \bar{\boldsymbol{\xi}}_{k-1}$  and a spread of  $\boldsymbol{\sigma}_k^{\xi 1}$ .  $\alpha$  is a hyper-parameter that controls the turbulence in the estimation. The evolved particles are assigned a weight which gets updated on each iteration based on their likelihood, detailed later in this article.

#### 4.2. Nested robust state filter

Unlike the typical EnKF formalism, the EnKF in the present study does not model the unknown input disturbance,  $\mathcal{S}_k^g$ , as a case of SWGN. Instead, rejection of this unknown input through output injection is approached (as detailed in Section 3) to gain robustness against  $\mathcal{S}_k^g$ . To accommodate this modification, the state and measurement equations within EnKF have been redefined accordingly.

For any arbitrary  $j^{\text{th}}$  particle  $\boldsymbol{\xi}_k^j$  among the  $N_p$  parameter particles available at the  $k^{\text{th}}$  time step,  $N_e$  state ensemble ( $\{\mathbf{X}_{k|k}^{i,j}\}; i = 1, 2, \dots, N_e$ ) are propagated over time using EnKF. Using a set of simulated SWGN processes  $\mathbf{U}_k^{i,j}$ ,  $\mathbf{V}_k^{i,j}$  and  $\mathbf{W}_k^{i,j}$ , with variances as defined in Section 3, the  $i^{\text{th}}$  ensemble is propagated and subsequently observed as follows as per Equations (5) and (10),

$$\begin{aligned} \mathbf{X}_{k|k-1}^{i,j} &= \mathcal{F}_k^j(\mathbf{X}_{k-1|k-1}^{i,j}) + \mathcal{B}_k^j \mathbf{U}_k^{i,j} + \mathbf{G}_k^j \mathbf{Y}_k + \mathbf{V}_k^{i,j} \\ \mathbf{Y}_{k|k-1}^{i,j} &= \tilde{H}_k^j \mathbf{X}_{k|k-1}^{i,j} + \tilde{D}_k^j \mathbf{U}_k^{i,j} + \mathbf{W}_k^{i,j}; \end{aligned} \quad (13)$$

The corresponding innovation for each of the  $i^{\text{th}}$  ensemble can be obtained as a departure of predicted measurement  $\mathbf{Y}_{k|k-1}^{i,j}$  from the actual measurement  $\mathbf{Y}_k$  as  $\boldsymbol{\epsilon}_{k|k-1}^{i,j} = \mathbf{Y}_k - \mathbf{Y}_{k|k-1}^{i,j}$ , with an ensemble mean  $\bar{\boldsymbol{\epsilon}}_{k|k-1}^j = \frac{1}{N_e} \sum_{i=1}^{N_e} \boldsymbol{\epsilon}_{k|k-1}^{i,j}$ . The cross-covariance between the state and the measurement prediction  $\mathbf{C}_k^{j,XY}$  and the innovation covariance  $\mathbf{S}_k^j$  can be computed as suggested in [14],

$$\begin{aligned} \mathbf{C}_k^{j,XY} &= \frac{1}{N_e - 1} \sum_{i=1}^{N_e} \{\mathbf{X}_{k|k-1}^j - \bar{\mathbf{X}}_{k|k-1}^j\} \{\mathbf{Y}_{k|k-1}^j - \bar{\mathbf{Y}}_{k|k-1}^j\}^T \\ \mathbf{S}_k^j &= \mathbf{C}_k^{j,YY} = \frac{1}{N_e - 1} \sum_{i=1}^{N_e} \{\mathbf{Y}_{k|k-1}^j - \bar{\mathbf{Y}}_{k|k-1}^j\} \{\mathbf{Y}_{k|k-1}^j - \bar{\mathbf{Y}}_{k|k-1}^j\}^T \end{aligned} \quad (14)$$

where  $\bar{\mathbf{X}}_{k|k-1}^j$  and  $\bar{\mathbf{Y}}_{k|k-1}^j$  are the ensemble mean estimates for the propagated states and predicted measurements, which can be obtained as:  $\bar{\mathbf{X}}_{k|k-1}^j = \frac{1}{N_e} \sum_{i=1}^{N_e} \mathbf{X}_{k|k-1}^{i,j}$  and  $\bar{\mathbf{Y}}_{k|k-1}^j = \frac{1}{N_e} \sum_{i=1}^{N_e} \mathbf{Y}_{k|k-1}^{i,j}$ . The EnKF gain can further be obtained as  $\mathbf{K}_k^j = \mathbf{C}_k^{j,XY} (\mathbf{S}_k^j)^{-1}$ . Finally, the state ensembles are updated as,

$$\mathbf{X}_{k|k}^{i,j} = \mathbf{X}_{k|k-1}^{i,j} + \mathbf{K}_k^j \boldsymbol{\epsilon}_{k|k-1}^{i,j} \quad (15)$$

<sup>1</sup>  $A + B\mathbb{N}(\mu; \sigma)$  means  $A + Bz$  where  $z$  follows  $\mathbb{N}(\mu; \sigma)$

253 Thus with EnKF, a set of prior state ensembles, i.e.,  $\{\mathbf{X}_{k-1|k-1}^{i,j}\}$  gets updated as  $\{\mathbf{X}_{k|k}^{i,j}\}$  with the ensemble  
 254 mean  $\mathbf{X}_{k|k}^j$  without being affected by  $\mathcal{S}_k^g$ .

---

**Algorithm 1** rIP-EnKF algorithm

---

```

1: procedure RIP-ENKF( $\mathbf{y}_k, \mathbf{Q}, \mathbf{R}^a, \mathbf{R}^g$ ) ▷ Process and measurements noise covariances
2:   At  $k = 0$  initialize particles  $\{\xi_0^j\}$  and state estimates:  $\{\mathbf{X}_{0|0}^{i,j}\}$  and  $\{\mathbf{P}_{0|0}^{i,j}\}$  ▷ Initialization
3:   for <each  $k^{th}$  measurement  $\mathbf{y}_k$ > do
4:     procedure RIP-ENKF( $\{\xi_{k-1}^j\}, \{\mathbf{X}_{k-1|k-1}^{i,j}\}, \{\mathbf{P}_{k-1|k-1}^{i,j}\}$ )
5:       for <each particle  $\xi_k^j \in \{\xi_k^j\}$ > do
6:         evolve  $\{\xi_{k-1}^j\} \rightarrow \{\xi_k^j\}$  ▷ Particle evolution, as per Equation (12)
7:         Define system matrices  $\mathbf{F}_k(\theta_k), \mathbf{H}_k(\theta_k)$  with  $\theta_k = \xi_k^j$  and  $\mathbf{B}_k, \mathbf{D}_k, \mathbf{E}_k$ 
8:         procedure ROBUST ENSEMBLE KALMAN FILTER( $\xi_k^j$ ) ▷ For  $j^{th}$  particle
9:           Define  $\mathbf{G}_k, \bar{\mathbf{F}}_k, \bar{\mathbf{B}}_k$  and  $\bar{\mathbf{v}}_k$  and realize  $\mathbf{u}_k^i$  from  $\mathcal{N}(0, Q)$ 
10:          for <each ensemble  $\mathbf{X}_{k-1|k-1}^{i,j} \in \{\mathbf{X}_{k-1|k-1}^{i,j}\}$ > do
11:            Predict  $\mathbf{X}_{k|k-1}^{i,j}$  and  $\mathbf{Y}_{k|k-1}^{i,j}$  as per Equation (13)
12:          end for
13:          Calculate  $\mathbf{X}_{k|k-1}^j, \mathbf{Y}_{k|k-1}^j, \epsilon_{k|k-1}^{i,j}$ , and  $\epsilon_{k|k-1}^j$  as per Section 4.2
14:          Perform correction as per Equation (15)
15:        end procedure
16:      end for
17:      procedure PARTICLE RE-SAMPLING( $\{\xi_k^j\}$ )
18:        Calculate  $w(\xi_k^j)$  for each  $\xi_k^j \in \{\xi_k^j\}$  and re-sample ▷ see Equation (17)
19:        Calculate,  $\mathbf{X}_{k|k}$  and parameter estimate,  $\bar{\xi}_k$  ▷ see Equation (18)
20:      end procedure
21:    end procedure
22:  end for
23: end procedure

```

---

255 *4.3. Particle approximation*

256 Since  $\mathbf{X}_{k|k}^j$  is conditioned on particle  $\xi_k^j$ , the likelihood of  $\xi_k^j$ , i.e.,  $\mathcal{L}(\xi_k^j) = \mathbf{p}(\mathbf{Y}_k | \theta_k = \xi_k^j)$  can thus be  
 257 defined using the ensemble mean of innovation,  $\epsilon_{k|k-1}^j$ , and the error covariance,  $\mathbf{S}_k^j$ , as,

$$\mathcal{L}(\xi_k^j) = \mathbf{p}(\mathbf{Y}_k | \theta_k = \xi_k^j) = (2\pi)^{-(p+s)/2} |\mathbf{S}_k^j|^{-1/2} \exp^{-\frac{1}{2} \epsilon_{k|k-1}^{jT} \mathbf{S}_k^{j-1} \epsilon_{k|k-1}^j} \quad (16)$$

258 Using  $\mathcal{L}(\xi_k^j)$ , the normalized updated weight for  $\xi_k^j$  can be obtained as,

$$w(\xi_k^j) = \frac{w(\xi_{k-1}^j) \mathcal{L}(\xi_k^j)}{\sum_{l=1}^{N_p} w(\xi_{k-1}^l) \mathcal{L}(\xi_k^l)} \quad (17)$$

259 leading to the particle approximation for state and parameter estimates as,

$$\mathbf{X}_{k|k} = \sum_{j=1}^{N_p} w(\xi_k^j) \mathbf{X}_{k|k}^j \quad \text{and} \quad \bar{\xi}_k = \sum_{j=1}^{N_p} w(\xi_k^j) \xi_k^j. \quad (18)$$

260 The proposed algorithm is detailed in a pseudo-code in Algorithm 1.

261 **5. Numerical experiment**

262 The case study is built from a numerical FE model of a 3D two storey - one bay fixed base concrete  
 263 frame structure consisting of sixteen members and 72 *dofs* (48 free *dofs*)(cf. Figure 1). Each frame member  
 264 is modelled with twelve *dofs* Euler Bernoulli beam element as detailed in Figure A.10 in Appendix A. The  
 265 length and cross section for each member are assumed to be 3m and  $0.3m \times 0.3m$ , respectively. The beam  
 266 material is assumed to have a modulus of elasticity of 30 GPa, a modulus of rigidity of 10 GPa and a density  
 267 of  $2500 \text{ kg/m}^3$ . This frame is excited by the true recorded data of the bi-directional El-Centro earthquake  
 268 ground motion (May 18, 1940 in CA, USA, direction North-South and East-West) (Data source: <http://peer.berkeley.edu/research/motions/>). In addition, SWGN and/or non-stationary WGN (NSWGN)  
 269 have been applied on all *dofs* throughout the simulation time, detailed later in the manuscript.  
 270

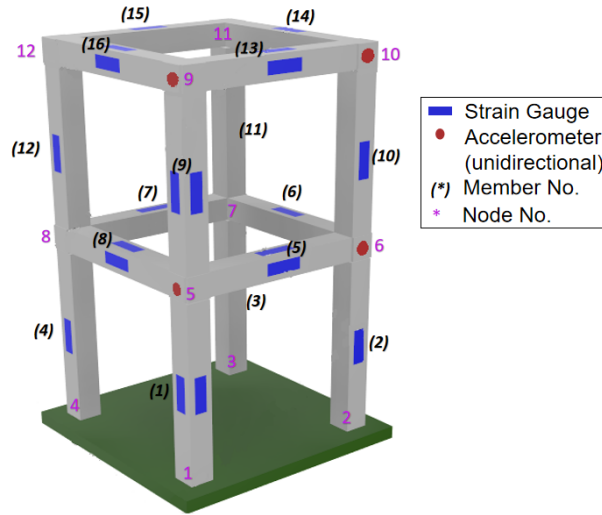


Figure 1: Schematic diagram of the sixteen member numerical frame

271 Strain measurements are collected from the strain gauges patched on to the top and the vertical external  
 272 surfaces for the beams and the vertical external surfaces of the columns at their midpoints. In addition,  
 273 horizontal accelerations are recorded at a set of nodes (cf. Figure 1). Various combinations of strain gauges  
 274 and accelerometers have been tried, with a maximum of 32 strain gauges (two for each of 16 members) and  
 275 8 accelerometers. Responses are sampled at a fixed sampling frequency of 50 Hz for 61.44 seconds to collect  
 276 the response time histories of length 3072. Noise contamination levels are defined in terms of signal-to-noise  
 277 ratio (SNR) in which the noise power of an SWGN has been scaled with respect to the structural response  
 278 against an SWGN vibration of variance 100 N. Various noise levels corresponding to different SNRs have  
 279 been experimented with.

280 A numerical reduction in elasticity is considered as damage for the simulations. To maintain consistency,

281 member 9 has been consistently assumed as damaged. A pool of 2000 particles are chosen for PF while  
 282 50 ensembles are chosen for EnKF. A better precision in estimation can be obtained with a bigger particle  
 283 pool or a bigger set of ensembles, which however comes at a higher computational cost. The particle pool  
 284 size and other tuning parameters relevant to the particle evolution have been standardized in [37, 38, 42].  
 285 For EnKF, 50 ensembles have been found to be sufficient with no significant loss in precision. A Gaussian  
 286 noise model,  $\mathcal{N}(1, 0.02)$ , has been chosen as the initial distribution for the parameter particles with  $\alpha$  chosen  
 287 as 0.98 (cf. Equation (12)). A lag has been introduced between the arrival of earthquake and the damage  
 288 occurrence to mimic the reality.

289 The member health is quantified using **HI**s.  $m$  **HI**s corresponding to  $m$  individual members are estimated  
 290 as individual elements of the  $m \times 1$  order parameter vector  $\boldsymbol{\theta}_k$ . With a damage induced in the structure in  
 291 terms of reduction in the initial elasticity  $\mathbb{E}_k = \{E_k^1 \ E_k^2 \cdots E_k^m\}$  of its members, the effect of **HI**s as  $\boldsymbol{\theta}_k$  on  
 292 member elasticity can be defined as,

$$\mathbb{E}_k^d(\boldsymbol{\theta}_k) = \langle \mathbb{E}_k \cdot \boldsymbol{\theta}_k \rangle \quad (19)$$

293 where  $\mathbb{E}_k^d(\boldsymbol{\theta}_k)$  denotes the reduced elasticity of the potentially damaged members and  $\langle \cdot \rangle$  denotes the  
 294 element-by-element multiplication operator. Since the structural stiffness  $\mathbf{K}(\boldsymbol{\theta}_k)$  is a linear function of the  
 295 member elasticity  $\mathbb{E}_k^d(\boldsymbol{\theta}_k)$ , **HI**s monitor the member health in terms of the ratio of the final to the initial  
 296 member stiffness within a range of 0 to 1, where 1 and 0 signify 0% and 100% damage levels. However,  
 297 instead of a reduction in elasticity, any other definition for damage can be applied. The impact of **HI**s to  
 298 the corresponding damaged stiffness is however required to be mapped. For instance, **HI** = 0.25 is roughly  
 299 equivalent to a 40% loss in beam depth for a conventional rectangular beam, which is practical for a real  
 300 world scenario.

### 301 5.1. Scenario description

302 A numerical investigation has been performed for undamaged, and therefore linear time invariant (LTI),  
 303 systems with constant system matrices. Further, linear time varying (LTV) systems are also investigated  
 304 for which the system matrices are varying because of the induced damage. These LTI and LTV systems are  
 305 experimented in combination with three forcing types (SWG N, NSWGN and earthquake (EQ)) with both  
 306 robust rIP-EnKF and non-robust IP-EnKF algorithms. This leads to 8 different scenarios (C1-C8) that have  
 307 been tested to validate the proposed method's relative efficiency in estimating damage over the non-robust  
 308 approach: i) LTI under SWGN and EQ (C1/2-LTI-EQ); ii) LTV under SWGN (C3/4-LTV-SWGN); iii) LTV  
 309 under NSWGN (C5/6-LTV-NSWGN); and finally, iv) LTV under SWGN and EQ (C7/8-LTV-EQ), with  
 310 odd and even numbered cases estimated with non-robust (NR) and robust (R) approaches, respectively. The  
 311 assumed SWGN of variance 100 N mimics the ambient excitation of known statistics. NSWGN is modelled  
 312 as two consecutive SWGNs (first one between 0-5 secs, second one between 5-62 secs) of variance 100 N and

313 1000 N. The bi-directional El centro earthquake excitation is adopted as EQ. The corresponding seismic  
314 excitation is presented in Figure 2. SWGNs are applied from the initiation of the simulation while NSWGN  
315 and EQs are introduced at the fifth second followed by damages (if any) at the eighth second. For all the  
316 cases, the responses are contaminated with SWGN of different SNR levels. With rIP-EnKF and IP-EnKF  
317 algorithms, denoted as R and NR, the case names in the figure are defined using the following formalism  
318 <case number>-<system>-<forcing>-<damage type>-<member number>-<estimation algorithm>.

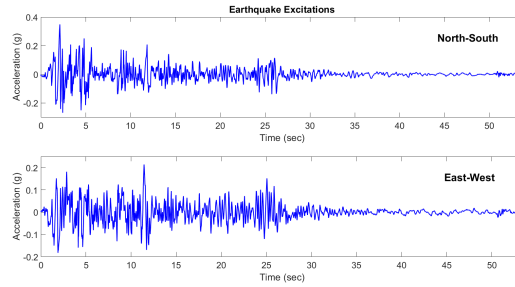


Figure 2: El centro seismic excitation in North-South (NS) and East-West (EW) direction

319 Further additional case studies (C9-C12) are performed to investigate the proposed method’s sensitivity  
320 to measurement noises through a numerical Noise Sensitivity Test (NST). The following eight case studies  
321 (C13-C20) investigate the effect of two sensor combinations (SC1 and SC2) under four different noise con-  
322 tamination levels (1/2/5/10% SNR). In this regard, additional results of ten numerical experiments (C-SC1  
323 – C-SC10) corresponding to ten other sensor combinations (SC1-SC10) are presented to explicitly demon-  
324 strate the effect of sensor densities under different particle and ensemble pool sizes. Case studies C21-C22  
325 demonstrate the capability of the proposed approach in detecting multiple damages while C-23 is included  
326 to illustrate the stability of the algorithm under a prolonged usage. A tabular description of each of the  
327 above mentioned scenarios has been given in Table 1 for better understanding. Each scenario has further  
328 been detailed in the subsequent sections.

### 329 5.2. Robust vs Non-robust IP-EnKF approach

330 The relative advantage of rIP-EnKF over the non-robust IP-EnKF is demonstrated through a compara-  
331 tive study presented in Figure 3. For the sake of brevity,  $\mathbf{HI}$  estimations are presented for two members only,  
332 one damaged (member 9) and one undamaged (member 1). Cases C1/2-LTI-EQ (Figure 3a) present that  
333 both IP-EnKF and rIP-ENKF perform equally good with the later being a little more stable. Further, in  
334 cases C3/4-LTV-SWGN (Figure 3b), better efficiency with rIP-EnKF becomes more evident. Till this phase,  
335 it is safe to conclude that neither of these two algorithms suffers from variations in the system matrices alone.  
336 Non-robust IP-EnKF starts performing poorly once the temporal variation is introduced in the input forcing  
337 in the case studies C5-LTV-NSWGN (cf. Figure 3c) and C7-LTV-EQ (cf. Figure 3d). For C5-LTV-NSWGN,  
338 the solution is not even converging (cf. Figure 3c), while for C7-LTV-EQ, the convergence is unstable and

Table 1: Details of numerical case studies

Objective	Case name	Algorithm	System	Forces	$d_i$	SNR	$N_s : N_a$
IP-EnKF vs rIPEnKF	C1-LTI-EQ	IP-EnKF	LTI	EQ+SWG	NA	1	32:8
	C2-LTI-EQ	rIP-EnKF	LTI	EQ+SWG	NA	1	32:8
	C3-LTV-SWGN	IP-EnKF	LTV	SWG	9	1	32:8
	C4-LTV-SWGN	rIP-EnKF	LTV	SWG	9	1	32:8
	C5-LTV-NSWGN	IP-EnKF	LTV	SWG+NSWGN	9	1	32:8
	C6-LTV-NSWGN	rIP-EnKF	LTV	SWG+NSWGN	9	1	32:8
	C7-LTV-EQ	IP-EnKF	LTV	EQ+SWG	9	1	32:8
	C8-LTV-EQ	rIP-EnKF	LTV	EQ+SWG	9	1	32:8
Noise sensitivity test	C9-NST	rIP-EnKF	LTV	EQ+SWG	9	1	32:8
	C10-NST	rIP-EnKF	LTV	EQ+SWG	9	2	32:8
	C11-NST	rIP-EnKF	LTV	EQ+SWG	9	5	32:8
	C12-NST	rIP-EnKF	LTV	EQ+SWG	9	10	32:8
Sensor combination test	C13	rIP-EnKF	LTV	EQ+SWG	9	1	32:4
	C14	rIP-EnKF	LTV	EQ+SWG	9	2	32:4
	C15	rIP-EnKF	LTV	EQ+SWG	9	5	32:4
	C16	rIP-EnKF	LTV	EQ+SWG	9	10	32:4
	C17	rIP-EnKF	LTV	EQ+SWG	9	1	32:8
	C18	rIP-EnKF	LTV	EQ+SWG	9	2	32:8
	C19	rIP-EnKF	LTV	EQ+SWG	9	5	32:8
	C20	rIP-EnKF	LTV	EQ+SWG	9	10	32:8
Double damage	C21	rIP-EnKF	LTV	EQ+SWG	1 & 5	1	32:8
	C22	rIP-EnKF	LTV	EQ+SWG	3 & 7	1	32:8
Stability check	C23	rIP-EnKF	LTV	EQ+SWG	9	1	32:8

$N_s$  and  $N_a$  denote the number of strain gauges and accelerometers, respectively.  $d_i$  and SNR denote damaged element/s and signal-to-noise ratio for the given case.

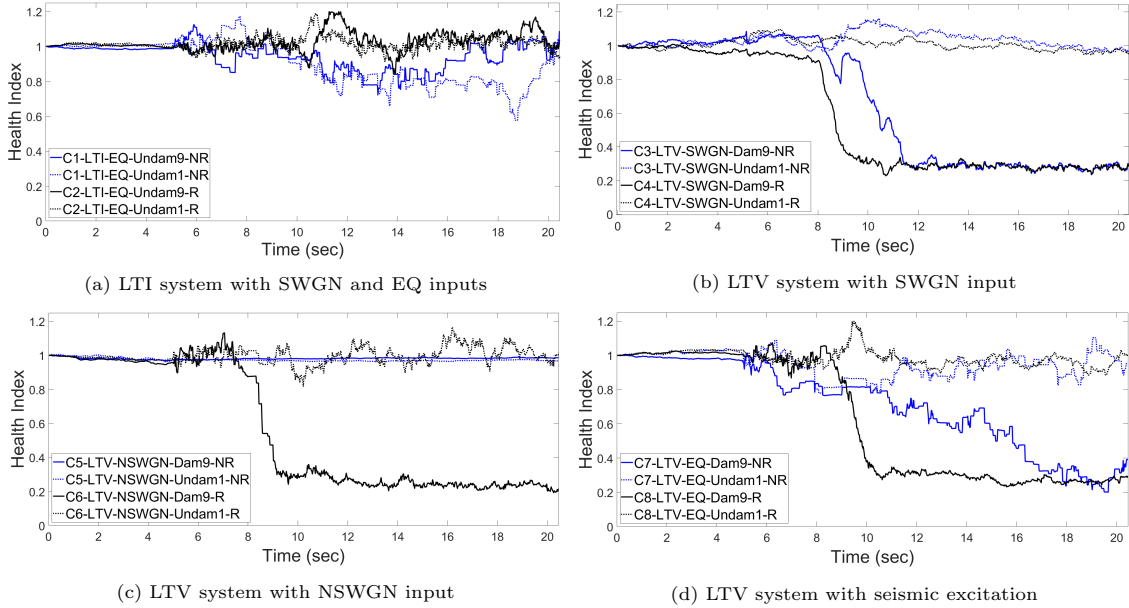


Figure 3: Comparative study between IP-EnKF vs rIP-EnKF under different system (LTI/LTV) and input (Stationary/non-stationary/seismic) conditions

339 not prompt. It takes more samples than rIP-EnKF to approach the actual value of the parameter. The  
 340 evolution is also not stable after convergence (cf. Figure 3d). Thus, the non-robust algorithm fails to handle  
 341 this variation at the onset of change in the input statistics. Putting the same dataset through rIP-EnKF for  
 342 C6-LTV-NSWGN (Figure 3c) and C8-LTV-EQ (Figure 3d), prompt and precise estimations are achieved.  
 343 The estimation history for C8-LTV-EQ is further demonstrated in terms of particle dispersion and particle

344 histogram in Figure 4.

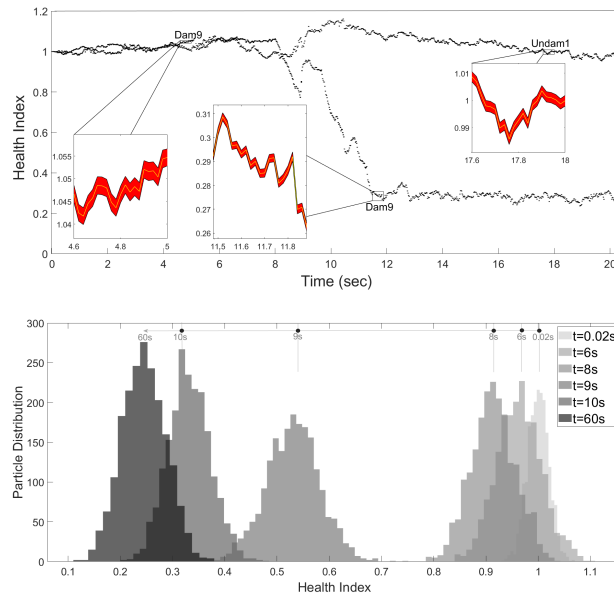


Figure 4: Temporal evolution of particles for damaged and undamaged members: Variation of standard deviation (top) and particle histogram (bottom)

345 *5.3. Robustness against measurement noise*

346 The proposed algorithm is tested for its robustness against measurement noise through four case studies  
 347 corresponding to four increasing SNR levels (1/2/5/10% SNR) while keeping the other parameters (damage  
 348 location, extent and forcing) similar to that taken for the case study C8. The earthquake signal is introduced  
 349 at the fifteenth second of simulation while damage is initiated at the twentieth second. Figure 5 presents  
 350 the results of the noise sensitivity test for four case studies (C9-NST – C12-NST) corresponding to the four  
 351 SNR levels adopted. As expected, with the increasing noise levels, the estimation promptness and precision  
 352 gradually degrades (cf. Figure 5). Yet, the algorithm is observed to be efficient till SNR noise level of  
 353 5%, beyond which (10% SNR), some undamaged elements are falsely identified as slightly damaged. This  
 354 illustrates the limiting noise contamination level for the algorithm to work precisely.

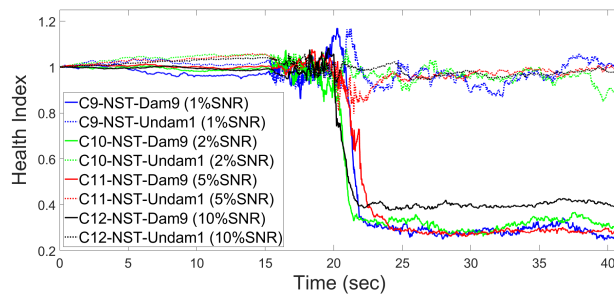


Figure 5: Measurement noise sensitivity of rIP-EnKF algorithm



355 5.4. Effect of measurement density and measurement type combinations

356 From a theoretical point of view, it can be perceived that robustness against input forcing depends on  
 357 the number of available linear measurement channels (acceleration in this case) while the detection precision  
 358 is governed by the overall instrumentation density. Positioning of sensors with respect to damage locations,  
 359 power in the recorded signal with respect to noise, size of particle and/or ensemble pools also play major  
 360 roles in defining detection certainty. Thus, judging the detection ability of the proposed algorithm merely  
 361 by sensor density may not be proper or practical. Yet a few test cases (C-SC1–C-SC10) are simulated for  
 362 different sensor combinations and particle and ensemble pool sizes. The results are presented in Table 2.

363 It is evident from Table 2, the proposed method efficiently estimates the damage location and severity  
 364 even with reduced sensor densities. It can also be observed that a lack of sensors can be complimented  
 365 with proper positioning of the sensors relative to the damage location and populating the ensemble and/or  
 366 particle pools. Additionally, **HI** estimation for two case studies (Cases C13-SC11 – C20-SC12) corresponding  
 367 to two sensor combinations (SC11: 32 strain gauges and 4 accelerometers and SC12: 32 strain gauges and 8  
 368 accelerometers) (other details are same as taken for C8) are presented in Figure 6a and 6b, respectively. These  
 369 case studies investigate the performance of rIP-EnKF under varying noise contamination levels (1/2/5/10%  
 370 SNR) under reduced instrumentation. Evidently, it can be observed that with reduced instrumentation, the  
 371 algorithm’s weakness to measurement noise increases.

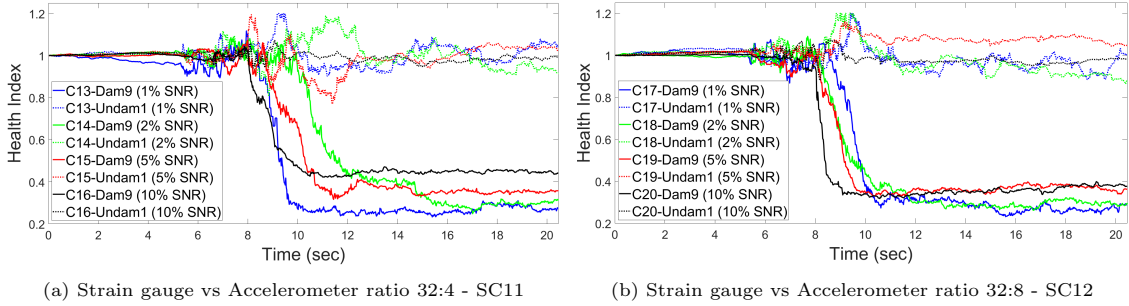


Figure 6: Performance of the algorithm under varying ratio of strain gauge and accelerometers

Table 2: Efficiency of varying ratio of strain gauge and accelerometers

No.	$N_s$	$N_a$	$N_e$	$N_p$	$D^{11}$	Acc
C-SC1	16	4	50	2000	✓	99.4%
C-SC2	16	2	50	2000	✓	99.4%
C-SC3	16	1	50	2000	✓	98.8%
C-SC4	8	4	50	2000	✓	73.3%
C-SC5	8	2	75	2000	✓	94.4%

No.	$N_s$	$N_a$	$N_e$	$N_p$	$D^{11}$	Acc
C-SC6	2	1	200	2000	✓	94.4%
C-SC7	4	1	100	2000	✓	88.9%
C-SC8	4	2	100	2000	✓	94.4%
C-SC9	4	1	75	2000	✓	88.9%
C-SC10	4	2	100	3000	✓	96.1%

$N_s, N_a, N_e, N_p$  denote the number of strain gauges, accelerometers, ensembles and particles, respectively.  $D^{11}$  and Acc denotes true detection and corresponding accuracies.

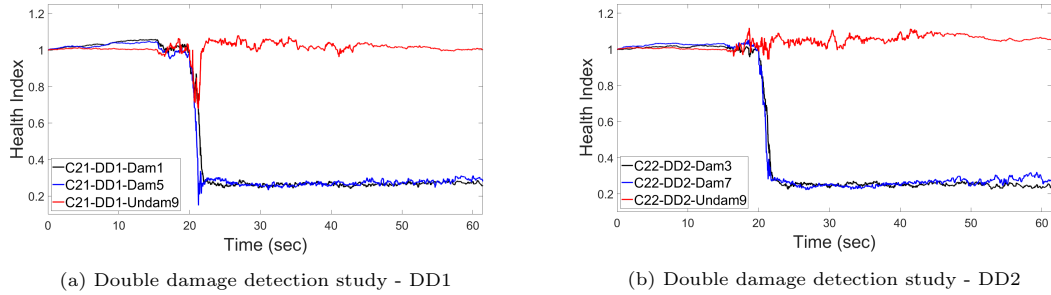


Figure 7: Performance of the algorithm to detect multiple damage location and stability

### 372 5.5. Detection of multiple damage scenario

373 Multiple damage scenarios are also tested with the proposed algorithm through two case studies (C21-  
 374 C22). Figure 7 presents two such cases with damage at two locations – case study C21-DD1: with damages  
 375 located at 1&5 (Figure 7a) and case study C22-DD2: with damages located at 3&7 (Figure 7b). A separate  
 376 case to check the stability of the algorithm under prolonged usage is also undertaken in C23-SC. The forcing  
 377 used in C9-NST – C12-NST is also used for the above mentioned cases. The algorithm performs promptly and  
 378 precisely in detecting multiple damages (cf. Figure 7). Case study, C23-SC, tests the proposed algorithm’s  
 379 stability for a period of 388 seconds against possible error accumulation or divergence. The algorithm is  
 380 observed to be stable for the mentioned duration showing no trend or significant turbulence (cf. Figure 8).

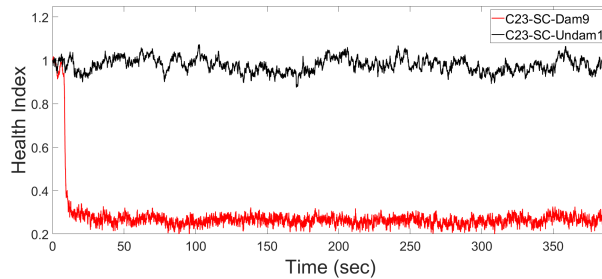


Figure 8: Stability check of the algorithm

381 The limiting value of damage that can be tracked with the proposed algorithm is investigated next. The  
 382 results are presented in Figure 9 corresponding to four different damage levels (75%, 30%, 20% and 10%).  
 383 It has been observed that rIP-EnKF is consistent in detecting the damage occurrence even for 10% damage  
 384 level ( $\mathbf{HI} = 0.9$ ). The precise estimation of the damage extent should however be limited to 20% damage  
 385 levels corresponding to  $\mathbf{HI}=0.80$ .

386 Overall, the proposed algorithm took about 5448 seconds of *cpu* time to process a 1024 long time series,  
 387 with 2000 particles and 50 ensembles for any case study discussed in this article. The employed computation  
 388 system is equipped with Intel(R) Xeon(R) Silver 4210 CPU @ 2.2GHz 2.19GHz (2 processors and 20 physical  
 389 cores with multi threading capability) with 64 GB RAM.

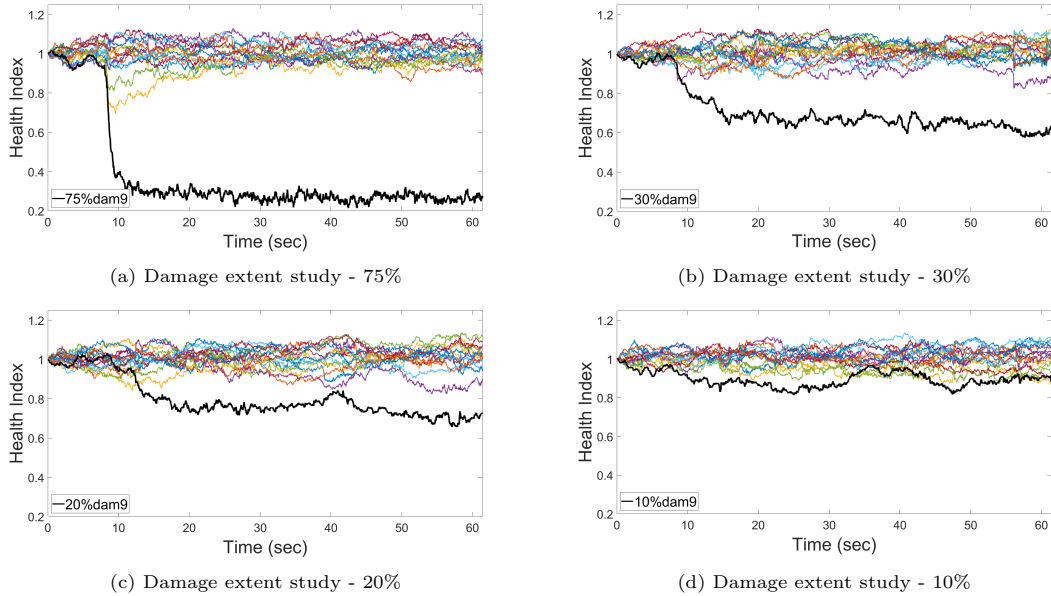


Figure 9: Performance of the algorithm to detect various damage levels

## 6. Conclusions

This paper has presented a Bayesian filtering-based structural health monitoring strategy robust to unknown and arbitrary input forcing using a mix of linear (acceleration) and non-linear (strain) sensor measurements. The considered framework leads to a non-linear estimation problem beyond the capability of classical Kalman filtering. The proposed algorithm is thus based on an interacting filtering strategy coupling Ensemble Kalman filters (EnKF), to track the evolution of the system states, with a Particle filter (PF), to track the changes in the system parameters due to damage. The input robustness is achieved through an output injection technique embedded within the EnKF formalism. The efficacy, robustness, stability and sensitivity of the proposed approach is validated numerically.

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## 494 Appendix A. Strain-displacement function for Euler-bernoulli beam

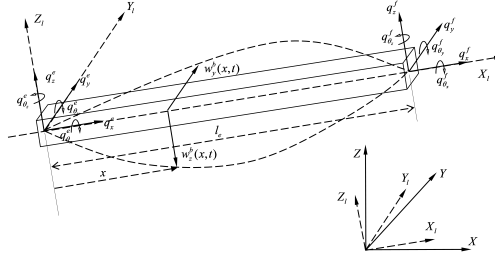


Figure A.10: Assumed Euler-Bernoulli Beam element with local and global degrees of freedom

495 A Strain-displacement mapping using FE model is performed using 3D Euler-Bernoulli beam elements  
496 schematically drawn in Figure A.10. The transverse deflection,  $w(x, t)$  at a distance  $x$  along beam orientation  
497 (i.e.,  $X_l$ ), with its components,  $w_y(x, t)$  and  $w_z(x, t)$  along  $Y_l$  and  $Z_l$ , can be defined as the interpolation of  
498 the nodal displacements as,

$$w(x, t) = \begin{bmatrix} w^b(x, t) \\ w^a(x, t) \end{bmatrix} = \begin{bmatrix} w_y^b(x, t) \\ w_z^b(x, t) \end{bmatrix} + \begin{bmatrix} w_y^a(x, t) \\ w_z^a(x, t) \end{bmatrix} = \boldsymbol{\psi}(x) \mathbf{q}^l(t) \quad (\text{A.1})$$

499  $\mathbf{q}^l(t) = \left[ q_x^e \quad q_y^e \quad q_z^e \quad q_{\theta_x}^e \quad q_{\theta_y}^e \quad q_{\theta_z}^e \quad q_x^f \quad q_y^f \quad q_z^f \quad q_{\theta_x}^f \quad q_{\theta_y}^f \quad q_{\theta_z}^f \right]^T$  can be retrieved from the nodal dis-  
500 placements,  $\mathbf{q}(t)$ , defined in Global Coordinate System (GCS) through coordinate transformation as  $\mathbf{q}^l(t) =$   
501  $\mathbf{T} \mathbf{q}(t)$  where  $\mathbf{T}$  is the coordinate transformation matrix.  $\boldsymbol{\psi}(x)$  is the interpolation function constituted with  
502 the associated shape functions. Similar to displacement, the slope  $\phi^b(x, t)$  and curvature  $\kappa^b(x, t)$  at  $x$  at time  
503  $t$  due to bending can be obtained as,  $\phi^b(x, t) = \boldsymbol{\psi}^b(x)' \mathbf{q}^l(t)$  and  $\kappa^b(x, t) = (1 + \phi^b(x, t)^2)^{\frac{3}{2}} \{ \boldsymbol{\psi}^b(x)'' \mathbf{q}^l(t) \}^{-1}$ .  
504 The curvature, measured through longitudinal strains,  $\boldsymbol{\varepsilon}^{xx}(x, t)$ , is eventually a function of  $\mathbf{q}^l(t)$  or  $\mathbf{q}(t)$ .  
505  $\mathbf{q}(t)$  being a subset of the state variable  $\mathbf{x}(t)$  yields the non-linear mapping  $\boldsymbol{\varepsilon}^{xx}(x, t) = f(\mathbf{x}(t))$  to describe  
506 strain-displacement relationship.