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Asymptotic Modelling for 3D Eddy Current Problems with a Conductive Thin Layer

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Abstract

In this work we derive and analyze an equivalent model for 3D Eddy Current problems with a conductive thin layer of small thickness ϵ . In our model, the conductive sheet is replaced by its mid-surface and their shielding behavior is satisfied by an equivalent transmission conditions on this interface. The transmission conditions are derived asymptotically for vanishing sheet thickness ϵ .

Key words: Asymptotic Expansions, Eddy-Current Problems, Thin Conducting Layers, Transmission Conditions

1 Introduction

We denote by $\Omega = \Omega_-^\epsilon \cup \overline{\Omega_0^\epsilon} \cup \Omega_+^\epsilon \subset \mathbb{R}^3$ the domain of study, where Ω_-^ϵ corresponds to a non-conductive linear material, Ω_+^ϵ the exterior of the structure domain, and Ω_0^ϵ a conductive thin layer of constant thickness ϵ (see figure 1). The discretisation of the conducting sheet by FEM needs a very fine mesh due to the rapid decay of the field under high conductivity. For this, we approximate a new model defined in ϵ -independent domains. Let Σ be a smooth surface, we denote by $[v]_{\Sigma}$ and $\{v\}_{\Sigma}$ the jump and mean of v respectively across Σ

$$[v]_{\Sigma} = v_{|_{\Sigma^+}} - v_{|_{\Sigma^-}}, \quad \{v\}_{\Sigma} = \frac{1}{2}(v_{|_{\Sigma^+}} + v_{|_{\Sigma^-}}), \quad \text{for} \quad v \in (C^{\infty}(\Omega_{\pm}))^3.$$

We consider the eddy current problem as follows

$$\begin{cases} \operatorname{curl} H^{\epsilon} & = & \sigma^{\epsilon} E^{\epsilon} + J_{0} & \text{in} & \Omega \\ \operatorname{curl} E^{\epsilon} & = & i w \mu_{0} H^{\epsilon} & \text{in} & \Omega \\ \operatorname{div} (H^{\epsilon}) & = & 0 & \text{in} & \mathbb{R}^{3} \\ [E \times n] = [H \times n] & = & 0 & \text{on} & \Gamma^{\epsilon}_{\pm} \end{cases}$$

where
$$\sigma^{\epsilon} = \begin{cases} 0 & \text{in } \Omega_{\pm}^{\epsilon} \\ \sigma_0 = \epsilon^{-2} \bar{\sigma} & \text{in } \Omega_0^{\epsilon} \end{cases}$$

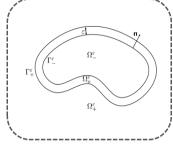


Figure 1: A cross Section of the domain Ω

Let u be a vector field on Γ , then we denote by $\gamma_D u = n \times (u \times n)$, and $\gamma_N u = \text{curl} u \times n$, the Dirichlet and Neumann data, respectively.

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2 Multiscale Expansion

Assuming that Γ is a smooth surface, then E^{ϵ} and H^{ϵ} can be expanded with an asymptotic expansion in power series of the small parameter ϵ . [1]

$$E^{\epsilon}(x) \approx E_0(x) + \epsilon E_1(x) + \epsilon^2 E_2(x) + \dots + O(\epsilon^k) \text{ in } \Omega^{\epsilon}_{\pm}$$

 $H^{\epsilon}(x) \approx \mathcal{H}_0(y_{\alpha}, \frac{h}{\epsilon}) + \epsilon \mathcal{H}_1(y_{\alpha}, \frac{h}{\epsilon}) + \dots + O(\epsilon^k) \text{ in } \Omega^{\epsilon}_0$

Here, $x \in \mathbb{R}^3$ are the cartesian coordinated, and (y_α,h) is the local normal coordinate system, $h \in (-\frac{\epsilon}{2},\frac{\epsilon}{2})$ is the normal coordinate to Γ . The term \mathcal{H}_j is a profile defined on $\Gamma \times (-\frac{1}{2},\frac{1}{2})$. The derivation is based on the expansion of the differential operators inside the thin layer Ω_0^ϵ , and the Taylor expansion of $E_j|_{\Gamma_+^\epsilon}$ around the mid-surface Γ .

3 Equivalent Model of Order 2

We introduce a problem satisfied by an approximation E_{ϵ}^k of the expression $E_0(x) + \epsilon E_1(x) + \epsilon^2 E_2(x) + ... + \epsilon^k E_k(x)$ up to a residual term $O(\epsilon^{k+1})$.

The second order approximate solution E_{ϵ}^{1} , solves

$$\begin{array}{llll} \operatorname{curl}(\operatorname{curl} E_{\epsilon}^{1}) & = & i\omega\mu J_{0} & & \operatorname{in} & \Omega_{\pm} \\ \left(\begin{array}{ccc} [\gamma_{D}E_{\epsilon}^{1}]_{\Gamma} \\ \{\gamma_{D}E_{\epsilon}^{1}\}_{\Gamma} \end{array} \right) & = & \epsilon \left(\begin{array}{ccc} C_{1} & 0 \\ 0 & C_{2} \end{array} \right) \left(\begin{array}{ccc} \{\gamma_{N}E_{\epsilon}^{1}\}_{\Gamma} \\ [\gamma_{N}E_{\epsilon}^{1}]_{\Gamma} \end{array} \right) & \operatorname{on} & \Gamma \end{array} \right|_{\Gamma} \left(\begin{array}{ccc} C_{1} & C_{1} & C_{2} \\ C_{1} & C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} \{\gamma_{N}E_{\epsilon}^{1}\}_{\Gamma} \\ C_{1} & C_{2} \end{array} \right) & = & C_{1} \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} \{\gamma_{N}E_{\epsilon}^{1}\}_{\Gamma} \\ C_{2} & C_{2} \end{array} \right) & = & C_{1} \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} \{\gamma_{N}E_{\epsilon}^{1}\}_{\Gamma} \\ C_{2} & C_{2} \end{array} \right) & = & C_{1} \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} \{\gamma_{N}E_{\epsilon}^{1}\}_{\Gamma} \\ C_{2} & C_{2} \end{array} \right) & = & C_{2} \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} \{\gamma_{N}E_{\epsilon}^{1}\}_{\Gamma} \\ C_{2} & C_{2} \end{array} \right) & = & C_{2} \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} \{\gamma_{N}E_{\epsilon}^{1}\}_{\Gamma} \\ C_{2} & C_{2} \end{array} \right) & = & C_{2} \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} \{\gamma_{N}E_{\epsilon}^{1}\}_{\Gamma} \\ C_{2} & C_{2} \end{array} \right) & = & C_{2} \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} \{\gamma_{N}E_{\epsilon}^{1}\}_{\Gamma} \\ C_{2} & C_{2} \end{array} \right) & = & C_{2} \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) & = & C_{2} \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} & C_{2} \end{array} \right) \left(\begin{array}{ccc} C_{1} & C_{2} \\ C_{2} &$$

where

$$C_1 = -1 + \frac{2\tanh(\frac{\gamma}{2})}{\gamma} , \quad C_2 = -\frac{1}{4} + \frac{\coth(\frac{\gamma}{2})}{2\gamma}$$
$$\gamma = exp(\frac{3i\pi}{4})\sqrt{\omega\mu_0\bar{\sigma}}.$$

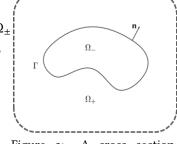


Figure 2: A cross section of the ϵ -independent subdomains

4 Numerical Results

Numerical experiments are performed to assess the accuracy of our model. The results are in particular compared to the model given in [2]. Complementary simulations will be conducted to study the robustness with respect to the sheet conductivity and the convergence of the modelling error.

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