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To cite this version:

Judit Chamorro-Servent, Laura Bear, Josselin Duchateau, Mark Potse, Rémi Dubois, et al.. Do we need to enforce the homogeneous Neumann condition on the torso for solving the inverse electrocardiographic problem?. Computing in Cardiology, Sep 2016, Vancouver, Canada. 2016. hal-01378822

HAL Id: hal-01378822 <https://inria.hal.science/hal-01378822>

Submitted on 10 Oct 2016

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Do we need to enforce the homogeneous Neumann condition on the torso for solving the inverse electrocardiographic problem?

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References

Methods

Acknowledgements

Introduction Results

Conclusions

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- \triangleright The results of MFS without Neumann condition are in better agreement with the original data. Furthermore, removing the Neumann conditions reduces the forward and inverse computational cost, since the size of the linear system is divided by 2 and the normal vectors do not need to be computed. The ill-posedness of the problem is reduced and consequently the problem is less sensible to the regularization.
- The Lagrange multipliers of *method 2* were significantly higher on the electrodes on the front of the torso than on the ones on the back*.* This seems to indicate that partial or weighed Neumann conditions (depending on the position of the electrodes) can be explored further in order to refine the inverse problem and improve the resolution of the ECGI problem.
- \triangleright Note that the results presented here only apply to the MFS method. However, preliminary work seems indicate that homogeneous Neumann conditions at the torso surface have a positive impact on Boundary element method.

 \triangleright Finally, we used only 5 ventricular datasets obtained by a numerical model. The case of complex atrial activations, and clinical data must be considered for a complete evaluation of the techniques proposed.

This study received financial support from the French Government as part of the « Investissement d'avenir » program managed by the National Research Agency (ANR), Grant reference ANR-10-IAHU-04 from the Conseil Régional Aquitaine as part of the project « Assimilation de données en cancérologie et cardiologie ».

- \triangleright We can remark that:
- The contribution of the submatrix *B¹* and *B⁰* in the whole solution is related to their respectives norms.
- Due to the meshing of the torso, non-phisiological conditions are applied to the top, bottom of the torso and the arms.
- Accurate computation of the normals at torso surface are required.

 $Y_{E}(\lambda) = u_0 + \sum_{i} u_i J \parallel_{i} N \lambda = 1$ $E_i Y_j - 1$ $T \cup T$ E

, $\chi\in\Gamma_{_{E}}$,)

 $= a_0 + \sum_{i=1}^{\infty} a_i f(||x - y_i||), x \in \Gamma_E, y_i \in \Gamma_T \cup \Gamma_E$ (3)

- The $\varepsilon = |B_1|/||B_0||$ depends only on the geometry and location of sources (for which we used a fixed rule defined in [2]). We found $\varepsilon = 0.0013$, indicating predominance of B_{0} compared to B_{1} in the standard MFS.
- Choosing $\lambda = 1/(1+\varepsilon) = 0.9987$ in equation (4) leads to a balance between the contribution of the two submatrices.
- **Comparison between the reconstructions**

 $-di\nu(\phi)$,in Ω . Under: Dirichlet conditions: $\phi = \phi_\tau$ on Γ_τ Homogeneous Neumann conditions: $\partial_n \phi = 0$ on Γ_{τ} (1)

The aim of the inverse problem is to find ϕ = ϕ_{ε} on Γ_{ε}

 Robust calculations of the inverse electrocardiographic problem may require accurate specification of boundary conditions at the torso and cardiac surfaces [1].

- **2. Minimization with the Neumann conditions as constraint**
	- The Neumann conditions can be introduced as a homogeneous equality constraint. In this case, we look for $\boldsymbol{a} \in \mathbb{R}^{1+N_s}$ that minimizes

- Five **activation patterns** were simulated : *1 single site pacing* (LV lateral endocardial lasts 500ms), and *4 single spiral waves* with different conductivity coefficients (lasts 3000ms). Simulations were conducted using [4]:
- i. The propagation of electrical activity in a membrane-based **realistic-geometry** computer model, using the governing **monodomain reaction-diffusion equation**. The simulations provide both the theoretical, *in-silico* $\phi_{\scriptscriptstyle\! T}^{}$ and $\phi_{\scriptscriptstyle\! E}^{}$ every 1ms.
- ii. The torso model used had **heterogeneous conductivity**.
- Correlation coefficients (CC) and relative errors (RE) for activation times and electrograms were calculated for solutions of *method* 1 with $\lambda = 0$, $\lambda = 0.5$, $\lambda = \frac{1}{1+\epsilon}$ where $\varepsilon = ||B_1||/||B_0||$, and *method* 2. \mathbf{I} 0, $\lambda = 0.5$, $\lambda = \frac{1}{2}$ where $\varepsilon = ||B_1||/2$ 1 λ = 0, λ = 0.5, λ = $\frac{1}{\lambda}$ where ε = $\left\|B_{_{1}}\right\|/\left\|B_{_{0}}\right\|$ ${\cal E}$ $=$ 0, λ = 0.5, λ = $\frac{1}{\lambda}$ where ε = $\ddot{}$

0 0 1 0 β_1 0 $\bigcup \mu$ 0 0 $B_0^T B_0 + \alpha I d \quad B_1^T \parallel a \parallel _ \parallel B_0^T \phi_T$ *B* α Id B_1^T α $B_0^T \phi$ ₁ μ $\begin{pmatrix} B_0^T B_0 + \alpha I d & B_1^T \\ B_0^T & 0 & I \end{pmatrix} = \begin{pmatrix} B_0^T \phi_T \\ 0 & 0 \end{pmatrix}$ $\begin{pmatrix} B_1 & 0 \end{pmatrix} \begin{pmatrix} \mu \end{pmatrix}^{\equiv} \begin{pmatrix} 0 \end{pmatrix}$ (8)

- We introduced a weighting coefficient in the cost function (*Method 1*), so as to change the balance between the contributions of the Neumann and Dirichlet conditions.
- We proposed to consider the Neumann conditions as a constraint (*Method 2*), so that it is resolved exactly.
- \triangleright These remarks drove us to the question: Do we need to enforce the homogeneous Neumann conditions in the MFS inverse problem?
- **REFERENCE** $\lambda = 0$ $\lambda = 0.5$ 180 $\lambda = 0.5$ $\lambda = 0$ 160 140 120 100 CC=0.99 CC=0.98 RE=0.02 RE=0.03 80 $\lambda = 0.9987$ **CONSTRAINED** 60 40 CC=0.33 CC=0.83 RE=0.51 RE=0,44

Respective contribution of B0 and B1

ACTIVATION MAPS CC MAPS

0.8

 0.7

0.6

Geometrical meshes used: The body is cut at the top and bottom of the torso and the arms.

Simplified view of geometrical boundaries and pseudo-boundaries took into account in the MFS formulation [2]

1

M

To evaluate the contribution of both boundary conditions:

j

 $=$

1. Minimization with weighed boundary conditions

 $\mathcal{L}(\pmb{\chi}) = \pmb{\alpha}_0 + \sum \pmb{\alpha}_j \pmb{f}\big(\big\|\pmb{\chi} - \pmb{\mathsf{y}}_j\big\|\big).$

 $\phi_{\varepsilon}(x) = a_0 + \sum a_i f(||x - y_i||)$, $x \in \Gamma_{\varepsilon}$, y

The solution of (4) was found by solving the following system using Tikhonov regularization

- $\big(1\!-\!\lambda\,\big)\pmb{B}_{_{\!0}}$ 1 0 $\left(B_{0}\right)$ α $\left(\lambda\phi_{\tau}\right)$ *a B* λ) B_{0}) ($\lambda \phi_{1}$ $\lambda\mathsf{l}$ $\begin{pmatrix} (1-\lambda)B_0 \\ \lambda B_1 \end{pmatrix}$ $\boldsymbol{a} = \begin{pmatrix} \lambda \phi_{\tau} \\ 0 \end{pmatrix}$
- Choosing:
	- λ = 0.5 corresponds to standard MFS.
	- $\lambda = 0$ corresponds to MFS without Neumann conditions.
	- $\lambda = \frac{1}{n}$, where $\|\varepsilon = \|B_1\|/\|B_0\|$ provides a balance between both submatrices. 1 1 $\mathcal{A}=\frac{\bot}{\bot}$, where $\mathcal{E}=\left\|\mathcal{B}_{1}\right\|/\left\|\mathcal{B}_{0}\right\|$ ${\cal E}$ $\ddot{}$

This is a saddle-point problem and we can introduce its Lagrangian

The solution is found by solving the block linear system

We solve this problem with SDPT3, a primal-dual path-following algorithm available on CVX package for convex optimization [3].

Comparisons

1

 $(a) = \frac{1}{2} \left(\left(1 - \lambda\right)^2 \|B_0 a - \phi_T\|_2^2 + \lambda \|B_1 a\|_2^2 + \alpha \|a\|_2^2 \right), \lambda \in [0,1]$ 2 2 2 2 0^{α} $\mathcal{V}_T ||_2$ $\mathcal{V}_2 ||_2$ $\mathcal{V}_3 ||_2$ $\mathcal{V}_4 ||_2$ 1 $(1 - \lambda)^2 \|B_0 a - \phi_\tau\|_2^2 + \lambda \|B_1 a\|_2^2 + \alpha \|a\|_2^2, \lambda \in [0,1]$ 2 $J_{\lambda}(a) = \frac{1}{2} \left(\left(1 - \lambda\right)^2 \|B_0 a - \phi_T\|_2^2 + \lambda \|B_1 a\|_2^2 + \alpha \|a\|_2^2 \right), \ \lambda \in [0,1]$ (4)

(5)

$$
J_0(a) = \frac{1}{2} (||B_0a - \phi_\tau||_2^2 + \alpha ||a||_2^2), \text{s.t. } B_1a = 0 \text{ in } \mathbb{R}^{N_\tau}
$$
 (6)

$$
L(a,\mu) = \frac{1}{2} \|B_0 a - \phi_\tau\|_2^2 + \frac{1}{2} \alpha \|a\|_2^2 + \mu^T B_1 a, \ \mu \in \mathbb{R}^{N_\tau}
$$
 (7)

Figure 1. Activation maps for electrograms, single site pacing.

0.5 $\lambda = 0.9987$ **CONSTRAINED** 0.4 0.3 0.2 0.1 *Figure 2. CC maps for electrograms, single site pacing.* The best activation time and CC maps were achieved by $\lambda = 0$ (MFS without Neumann conditions) and $\lambda = 0.5$ (Standard MFS).

COMPARISON OF ELECTROGRAMS

\triangleright To solve this question:

