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INSTITUT NATIONAL DE RECHERCHE EN INFORMATIQUE ET EN AUTOMATIQUE

*Stability analysis of the Interior Penalty
Discontinuous Galerkin method for the wave
equation*

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Observation and Modeling for Environmental Sciences



*Report
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Stability analysis of the Interior Penalty Discontinuous Galerkin method for the wave equation

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Abstract: We show how to derive analytically the optimal value of the penalization parameter and we give the analytical expression of the CFL condition of the scheme. We restrict ourselves to the case of basis function from degree 1 to 5 but the study can be extended to higher polynomial degrees.

Key-words: Discontinuous Galerkin, penalization coefficient, CFL condition, wave equation

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Sur la condition CFL de la méthode de Galerkin Discontinue avec Pénalité Intérieure pour l'équation des ondes

Résumé : Nous montrons comment obtenir analytiquement la valeur optimale du paramètre de pénalisation et nous donnons l'expression analytique de la condition CFL du schéma. Nous nous restreignons au cas de fonctions de base de degrés 1 à 5 mais cette étude peut être étendue à des degrés polynomiaux plus élevés.

Mots-clés : Galerkin Discontinue, coefficient de pénalisation, condition CFL, équation des ondes

1 Introduction

Nowadays, in many domains as for instance medical imaging, seismic imaging, radar, earthquakes simulations, one requires the numerical solution to acoustic, elastodynamic or electromagnetic Wave Equations. In many cases, the equations have to be solved in very large domains with strong heterogeneities. Therefore, one has to use sophisticated space and time discretization methods to compute the most accurate solution for the smallest computational cost.

Finite Differences Methods, which are widely used because of their small computational cost and the simplicity of their implementation, are not adapted to deal with strong heterogeneities. Indeed, they rely on structured grids, which can not accurately approximate the shape of the various layers of the domain. Finite Elements Methods (FEM) are more adapted to this kind of problems since they allow for the use of unstructured grids. Among all existing FEM, the Spectral Element Method [13, 7, 12], is probably one of the most efficient to solve the Wave Equation since the resulting mass matrix is diagonal. Hence, it can be easily coupled to explicit time-schemes which requires the inversion of the mass matrix at each time step. However, SEM requires the use of quadrilateral (in 2D) or hexahedral (in 3D) meshes which can be difficult to generate for realistic applications. Let us however mention that SEM can be extended to handle triangular meshes [6], but it requires additional degrees of freedom and the implementation is more complex. Moreover, as far as we know, the extension to tetrahedral meshes has not been proposed yet.

Discontinuous Galerkin Methods are more and more popular for solving the Wave Equation since they lead to block-diagonal mass matrices without the help of quadrature formula. Moreover, they can be used with any type of meshes and even allow for the variation of the physical parameters inside the cells of the mesh (provided that the variation is polynomial). DGM are also naturally adapted to parallel computing since all volume integrals are computed locally and the communications between the cells are ensured by integrals over the faces of the elements. In [2], the authors provide a detailed review of the various Discontinuous Galerkin approximations of the Laplacian operator. They also show that the so-called Interior Penalty Discontinuous Galerkin Method (IPDGM), also known as Symmetric Interior Penalty (SIP) [4], is one of the most suitable since it is stable and adjoint consistent, which guarantees the optimal order of convergence of the scheme. This explains why this method has been successfully used to solve Helmholtz Equation [1, 5] and the wave equation [10, 11, 1]. Comparisons of the performances of IPDGM and SEM can be found in [8] and [3]. It is worth noting that the first paper concluded that the performances of SEM are better than IDPGM when IPDGM is applied to structured grids composed of squares, while the second paper show that the performances of IPDGM are better when it is applied to triangular meshes.

Nevertheless, in spite of all its interesting properties, IPDGM still suffers from two difficulties. The first one is the determination of the penalization parameter, which penalizes the discontinuities of the solution through the faces. The accurate determination of the optimal parameter is crucial, since a too small value leads to instabilities while a too large value could (strongly) hamper the CFL (Courant-Friedrichs-Lewy) condition, which gives the maximal time step that can be used to ensure the stability of the scheme. In [1], the authors conjectured a minimal value of the penalization parameters, depending on p ,

the polynomial degree of the basis functions and on the size of the elements. They proved their conjecture up to $p = 3$. The extension of this result to $p > 3$ and to unstructured meshes, is still to be done. The second difficulty is the determination of the CFL condition. It is well-known that this condition decreases when the penalization parameter increases, but no analytical formula has been proposed yet. The aim of this paper is a) to prove the conjecture of Ainsworth, Monk and Muniz up to $p = 5$; b) to propose a solution methodology to prove it for a given p ; and c) to provide an analytic formula linking the CFL condition to the penalization parameter. We restrict ourselves to the cases of structured meshes composed of segments (in 1D), squares (in 2D) or cubes (in 3D). In section 1, we recall the IPDG discretization of the Wave Equation. In section 2, we propose two theorems, the first one provides necessary stability conditions on the penalization parameter and the time step while the second one provides the necessary and sufficient stability conditions. The proof of these theorems in the one dimensional case is given in section 3. Section 4 is devoted to the proof in the three dimensional case and contains a discussion on the adaptation of the theorems to structured meshes composed of rectangles or parallelepipeds. We do not present the proof in the two dimensional case, but it can be adapted without any difficulties from the three dimensional case. Finally, we present numerical results in section 5 that illustrate the validity of our conditions.

2 Discretization of the acoustic wave equation

In this section, we present the so called Interior Penalty Discontinuous Galerkin method applied to the acoustic wave equation in homogeneous bounded media $\Omega \subset \mathbb{R}^d$, $d = 1, 2, 3$. For the sake of simplicity, we impose homogeneous Neumann boundary conditions on the boundary $\Gamma := \partial\Omega$ but this study can be extended to Dirichlet boundary conditions without major difficulties.

$$\left\{ \begin{array}{l} \text{Find } u : \Omega \times [0, T] \mapsto \mathbb{R} \text{ such that :} \\ \frac{1}{\mu} \frac{\partial^2 u}{\partial t^2} - \operatorname{div} \left(\frac{1}{\rho} \nabla u \right) = f \quad \text{in } \Omega \times]0, T], \\ u(x, 0) = u_0, \quad \frac{\partial u}{\partial t}(x, 0) = u_1 \quad \text{in } \Omega, \\ u = 0 \quad \text{on } \partial\Omega. \end{array} \right. \quad (1)$$

where u stands for the displacement, μ is the compressibility modulus, ρ is the density and f is the source term.

We introduce a triangulation \mathcal{T}_h of Ω and the following space of approximation with piecewise discontinuous polynomial functions :

$$V_h := \{v \in L^2(\Omega) : v|_K \in P^p(K), \forall K \in \mathcal{T}_h, p \geq 3\}.$$

The set of the mesh faces are denoted \mathcal{F}_h which is partitionned into two subsets \mathcal{F}_h^i and \mathcal{F}_h^b corresponding respectively to the interior faces and those located on the boundary. For $F \in \mathcal{F}_h^i$, we note arbitrarily K^+ and K^- the two elements sharing F and we define ν as the unit outward normal vector pointing from K^+ to K^- . Moreover, v^\pm represents the restriction of a function v to the element

K^\pm and we define the jump and the average of a piecewise smooth function $v \in V_h$ over $F \in \mathcal{F}_h^i$ such that:

$$\llbracket v \rrbracket = v^+ - v^-, \quad \{\!\!\{ v \}\!\!\} = \frac{v^+ + v^-}{2} \quad (2)$$

and for $F \in \mathcal{F}_h^b$ such that $\llbracket v \rrbracket = v$ and $\{\!\!\{ v \}\!\!\} = v$.
The IPDG discretization of (1) reads as

$$\begin{cases} \text{Find } u_h \in V_h \text{ such that, } \forall v_h \in V_h : \\ \sum_{K \in \mathcal{T}_h} \int_K \frac{1}{\mu} \partial_t^2 u_h v_h \, dx = -a_h(u_h, v_h) + \sum_{K \in \mathcal{T}_h} \int_K f v_h \, dx. \end{cases} \quad (3)$$

where a_h is a bilinear form defined by

$$a_h(u_h, v_h) = B_{\mathcal{T}_h}(u_h, v_h) - \mathcal{I}(u_h, v_h) - \mathcal{I}(v_h, u_h) + B_S(u_h, v_h),$$

with

$$B_{\mathcal{T}_h}(u_h, v) = \sum_{K \in \mathcal{T}_h} \int_K \frac{1}{\rho} \nabla u_h \nabla v, \quad \mathcal{I}(u_h, v) = \sum_{F \in \mathcal{F}_h^i} \int_F \llbracket v \rrbracket \left\{ \left\{ \frac{1}{\rho} \nabla u_h \cdot \boldsymbol{\nu} \right\} \right\},$$

$$B_S(u_h, v) = \sum_{F \in \mathcal{F}_h^i} \int_F \gamma \llbracket u_h \rrbracket \llbracket v \rrbracket.$$

The bilinear form B_S is devoted to enforce the coercivity of a_h and the penalization function γ is defined on each interior face F by

$$\gamma = \frac{\alpha}{\xi_F}$$

where α is a positive parameter. There are many definitions of the function ξ_F in the litterature. The most commonly used are:

- $\xi_F = h(F)$ where $h(F)$ denotes the diameter F . See for instance [2, 1, 11].
It is worth noting that that this definition does not make sense in 1D.
- $\xi_F = \min(h(K^+), h(K^-))$ where $h(K^\pm)$ is the diameter of K^\pm . See for instance [10].
- $\xi_F = \min(\rho_{K^+}, \rho_{K^-})$ where $\rho(K^\pm)$ is the diameter of the inscribed circle (or sphere) of K^\pm . See for instance [14].

Whatever the definition of ξ_F , the coercivity of a_h is ensured for $\alpha \geq \alpha_0$. Obviously, the optimal parameter α_0 depends on the choice of the basis functions of V_h , but also on ξ_F . It has been shown by Shabazi in [14] that the third definition was the most appropriate for triangular meshes.

At this point, we choose not to explicit the expression of ξ_F . This will be done in the next section.

We refer to [2, 1, 10] for more details on the properties of the bilinear form a_h .

Considering $\{\varphi_i\}_{i=1, \dots, m}$ the classical discontinuous Lagrange basis functions

of degree p of V_h , where m denotes the number of degrees of freedom of the problem, we obtain the following linear system:

$$\partial_t^2 U = M^{-1}KU + M^{-1}F \quad (4)$$

where

$$(M)_{i,j} = \sum_{K \in \mathcal{T}_h} \int_K \varphi_i \varphi_j, \quad (K)_{i,j} = a_h(\varphi_i, \varphi_j), \quad (F)_i = \sum_{K \in \mathcal{T}_h} \int_K f \varphi_i.$$

Now, we have to discretize in time. Using the well known Leap-Frog scheme, we obtain the following fully discretized scheme:

$$\frac{U^{n+1} - 2U^n + U^{n-1}}{\Delta t^2} = -M^{-1}KU^n + M^{-1}F^n. \quad (5)$$

3 Stability analysis

In this section, we first propose necessary conditions over γ and Δt ensuring the L^2 -stability of scheme (5). This theorem provides an explicit dependance of Δt with respect to γ and h . Next we propose a sufficient and necessary stability condition. In this second theorem, the dependance of Δt with respect to γ is no longer explicit. However the CFL condition can be numerically computed using the roots of a polynomial of degree $2p$. We assume here that the domain Ω is unbounded ($\Omega = \mathbb{R}^d$) and uniformly meshed by segments (if $d = 1$), squares (if $d = 2$) or cubes (if $d = 3$). The length of the edges of the elements is denoted by h .

The necessary stability conditions are given by the following theorem.

Theorem 3.1. *The scheme (5) is L^2 -stable only if, for $p \leq 5$,*

$$\gamma \geq \frac{p(p+1)}{2h}; \quad (6)$$

and, denoting $\alpha = \gamma/h$,

$$\sqrt{d} \frac{c\Delta t}{h} \leq \begin{cases} C_{1,p} & \text{if } \alpha \leq \alpha_{1,p} \\ C_{2,p}(\alpha) & \text{if } \alpha > \alpha_{1,p}. \end{cases} \quad (7)$$

where $\alpha_{1,p}$, $C_{1,p}$ and $C_{2,p}(\alpha)$ are defined with respect to the polynomial degree p such that:

p	$\alpha_{1,p}$	$C_{1,p}$	$C_{2,p}(\alpha)$
1	2	$\frac{\sqrt{3}}{3} \simeq 0.577$	$\frac{1}{\sqrt{3(\alpha-1)}}$
2	$\frac{27}{5} = 5.4$	$\frac{1}{\sqrt{15}} \simeq 0.258$	$\sqrt{\frac{2}{-15 + 6\alpha + (405 - 240\alpha + 36\alpha^2)^{1/2}}}$
3	$\frac{2\sqrt{1605} + 393}{49} \simeq 9.65$	$\sqrt{\frac{2}{45 + \sqrt{1605}}} \simeq 0.153$	$\sqrt{\frac{2}{-45 + 10\alpha + (4545 - 1320\alpha + 100\alpha^2)^{1/2}}}$
4	$\alpha_{1,4} \simeq 14.7$	$\sqrt{\frac{2}{3(35 + \sqrt{805})}} \simeq 0.103$	$\sqrt{\frac{1}{2\sqrt{5}g_{4,1}(\alpha)g_{4,2}(\alpha) + 5\alpha - 35}}$
5	$\alpha_{1,5} \simeq 20.8$	$\sqrt{\frac{1}{10\sqrt{133}\cos(g_5) + 70}} \simeq 0.074$	$\sqrt{\frac{1}{2\sqrt{7}g_{5,1}(\alpha)g_{5,2}(\alpha) + 7(\alpha - 10)}}$

where, for the case $p = 4$, we have

$$\begin{cases} g_{4,1}(\alpha) = (518 - 98\alpha + 5\alpha^2)^{\frac{1}{2}}, \\ g_{4,2}(\alpha) = \cos\left(\frac{1}{3}\arccos\left(\frac{1}{10}g_{4,3}(\alpha)\frac{\sqrt{5}}{g_{4,1}^3(\alpha)}\right)\right), \\ g_{4,3}(\alpha) = -47705 + 14574\alpha - 1470\alpha^2 + 50\alpha^3 \end{cases}$$

and for the polynomials of degree 5,

$$\begin{cases} g_5 = \frac{1}{3}\arccos\left(\frac{10447}{126350}\sqrt{133}\right) \\ g_{5,1}(\alpha) = (1555 - 200\alpha + 7\alpha^2)^{\frac{1}{2}}, \\ g_{5,2}(\alpha) = \cos\left(\frac{1}{3}\arccos\left(\frac{1}{14}g_{5,3}(\alpha)\frac{\sqrt{7}}{g_{5,1}^3(\alpha)}\right)\right), \\ g_{5,3}(\alpha) = -299825 + 61440\alpha - 4200\alpha^2 + 98\alpha^3. \end{cases}$$

Remark 3.2. • As it is was noted in [11], the stability condition on Δt

behaves as $C/\sqrt{\alpha}$ for large α . More precisely, $C = \sqrt{\frac{2}{(p+1)(p+2)}}$.

- This stability condition is constant for $\frac{p(p+1)}{2} \leq \alpha \leq \alpha_p^1$. This shows that it is not necessary to choose α too close from α_0 to improve the CFL condition.

However, the above condition is only necessary. To obtain the actual CFL, we need the following necessary and sufficient condition.

Theorem 3.3. *The scheme (5) is L^2 -stable if and only if, for $p \leq 5$, (6) and (7) are satisfied and*

$$\sqrt{d} \frac{c\Delta t}{h} \leq C_{3,p}(\alpha) \quad (8)$$

where $C_{3,p}(\alpha) = \min \left\{ \lambda \in \mathbb{R} : Q_{p,\alpha}(\lambda) = 0 \text{ and } |\tilde{Q}_{p,\alpha}(\lambda)| \leq 1 \right\}$ where $Q_{p,\alpha}(\lambda)$ is a polynomial of degree $2p$ and $\tilde{Q}_{p,\alpha}(\lambda)$ is a rational function. We give in appendix A the expressions of these two functions for $1 \leq p \leq 5$.

Remark 3.4.

- *This theorem does not provide an explicit CFL condition. However it can be computed numerically by the following algorithm:*
 1. *Compute all the roots of $Q_{p,\alpha}$,*
 2. *Select the real roots such that $|\tilde{Q}_{p,\alpha}(\lambda)| \leq 1$,*
 3. *Choose the minimum.*
- *The numerical study of condition (8) that we present at section 6 shows that the set V_α is actually empty except when α belongs to a small segment around α_1^p . This means that theorem 3.1 provides a sufficient and necessary stability condition when α is not in this segment. Moreover the remarks 3.2 are still valid.*
- *In [1], they authors proved (6) for $p = 0, \dots, 3$ and conjectured this relation for any p . Theorem (3.3) extends its validity until $p = 5$.*
- *The condition (6) does not depend on the dimension d . This would not have been the case if we had expressed γ as a function of the circumsphere (or circumsphere) diameter which is \sqrt{dh} . Since h is the diameter of the inscribed circle or sphere, we conjecture that the third definition of ξ_F is the most appropriate. We will strengthen when we discuss the extension of this theorem to meshes composed of rectangles or parallelepipeds.*

We were unfortunately unable to establish this theorem for any p and we have restricted ourselves to $p \leq 5$. The proofs in the one dimensional case are given in section 4 while its extension to $d = 3$ is the subject of section 5. The proof for $d = 2$ can be easily deduced from the case $d = 3$.

4 Study in the 1-Dimensional case

This section contains the proofs of theorems 3.1 and 3.3 in the one dimensional case. It consists of three steps. The first step is a Fourier analysis presented in section 4.1; the second step is devoted to the proof of conditions (6) and is presented in section 4.2; the last step concerns the proof of (7) and (8) in section 4.3. The proofs are detailed for $p = 3$ and the main steps of the proofs in the cases $1 \leq p \leq 5$ are given in appendix A.

Here, we assume that the domain is $\Omega = \mathbb{R}$ and is meshed by segments of length h . We consider a velocity $c^2 = \mu/\rho = 1$ but we can easily extend the proof

to other velocities by setting $\Delta t' = \Delta t/c$. We consider the scheme (5) without source term that is to say

$$M \frac{U^{n+1} - 2U^n + U^{n-1}}{\Delta t^2} + KU^n = 0. \quad (9)$$

Considering the equation on one element J of the mesh, we have $\forall J \in \mathcal{T}_h$

$$M_{1,p} \frac{U_J^{n+1} - 2U_J^n + U_J^{n-1}}{\Delta t^2} + (K_{1,p}^W)^T U_{J-1}^n + K_{1,p} U_J^n + K_{1,p}^W U_{J+1}^n = 0 \quad (10)$$

where U_J corresponds to the vector of unknowns U restricted to the element J and $M_{1,p}$, $K_{1,p}$ and $K_{1,p}^W$ are respectively the mass and stiffness matrices in dimension 1 considering polynomials of degree p and more precisely,

$$\begin{aligned} M_{1,p}(i, j) &= h \int_{[0,1]} \hat{\varphi}_i(\hat{x}) \hat{\varphi}_j(\hat{x}) d\hat{x} \\ K_{1,p}(i, j) &= \frac{1}{h} \int_{[0,1]} \frac{\partial \hat{\varphi}_i}{\partial \hat{x}}(\hat{x}) \frac{\partial \hat{\varphi}_j}{\partial \hat{x}}(\hat{x}) d\hat{x} + \frac{1}{2h} \hat{\varphi}_i(1) \frac{\partial \hat{\varphi}_j}{\partial \hat{x}}(1) + \frac{1}{2h} \hat{\varphi}_j(1) \frac{\partial \hat{\varphi}_i}{\partial \hat{x}}(1) \\ &\quad + \gamma \hat{\varphi}_i(1) \hat{\varphi}_j(1) - \frac{1}{2h} \hat{\varphi}_i(0) \frac{\partial \hat{\varphi}_j}{\partial \hat{x}}(0) - \frac{1}{2h} \hat{\varphi}_j(0) \frac{\partial \hat{\varphi}_i}{\partial \hat{x}}(0) \\ &\quad + \gamma \hat{\varphi}_i(0) \hat{\varphi}_j(0) \\ K_{1,p}^W(i, j) &= -\frac{1}{2h} \hat{\varphi}_i(1) \frac{\partial \hat{\varphi}_j}{\partial \hat{x}}(0) + \frac{1}{2h} \hat{\varphi}_j(0) \frac{\partial \hat{\varphi}_i}{\partial \hat{x}}(1) - \gamma \hat{\varphi}_i(1) \hat{\varphi}_j(0) \end{aligned} \quad (11)$$

where $\{\hat{\varphi}_i\}_{i=1, \dots, p+1}$ are the classical discontinuous Lagrange basis functions on the reference element $[0, 1]$.

4.1 Fourier Analysis of the IPDG scheme in 1D

In order to study the stability of the IPDG scheme, we have to introduce the discrete Fourier transform

$$\begin{aligned} \mathcal{F}_h : L_h^2 &\rightarrow L^2(K_h) \\ U &\rightarrow \tilde{U} = \mathcal{F}_h(U)(k) = \frac{h}{2\pi} \sum_{J \in \mathbb{Z}} U_J e^{-ikJh} \end{aligned}$$

with $K_h = [-\frac{\pi}{h}, \frac{\pi}{h}]$ and $L_h^2 = \left\{ U = (U_J)_{J \in \mathbb{Z}}, \sum_{J \in \mathbb{Z}} \|U_J\|^2 < +\infty \right\}$.

Now, applying this discrete Fourier transform to (10), we obtain, $\forall \beta \in [-\pi, \pi]$

$$M_{1,p} \frac{\tilde{U}_J^{n+1}(\beta) - 2\tilde{U}_J^n(\beta) + \tilde{U}_J^{n-1}(\beta)}{\Delta t^2} + K_\beta \tilde{U}_J^n(\beta) = 0 \quad (12)$$

where $\beta = hk$ and $K_\beta = (K_{1,p}^W)^T e^{-i\beta} + K_{1,p} + K_{1,p}^W e^{i\beta}$.

The L^2 -stability of (12), for all β , is equivalent to the L^2 -stability of (9), thanks to the Parseval equalities.

Since $M_{1,p}$ is positive definite and K_β is hermitian, all the eigenvalues of $N_\beta = M^{-1}K_\beta$ are real. Moreover, a classical stability analysis shows that (9) is stable if and only if

$$0 \leq \lambda \leq \frac{4}{\Delta t^2}$$

for $\lambda \in \Lambda(\beta)$ and for all $\beta \in [-\pi, \pi]$, where $\Lambda(\beta)$ denotes the set of the eigenvalues of N_β . Then, a necessary and sufficient condition of the stability of (7) is

$$\lambda_{\min} \geq 0 \quad \text{and} \quad \Delta t \leq \frac{2}{\sqrt{\lambda_{\max}}}$$

with $\lambda_{\min} = \min_{\beta \in [-\pi, \pi]} [\min(\Lambda(\beta))]$ and $\lambda_{\max} = \max_{\beta \in [-\pi, \pi]} [\max(\Lambda(\beta))]$.

In section 4.2, we show that the condition $\lambda_{\min} \geq 0$ is equivalent to (6) and in section 4.3, we show that the condition $\lambda_{\max} \leq \frac{4}{\Delta t^2}$ implies (7) and is equivalent to (7) and (8).

4.2 Study of the condition $\lambda_{\min} \geq 0$

In the following, we consider the change of variable $\alpha = h\gamma$ to simplify the presentation.

To show the equivalence between (5) and $\lambda_{\min} \geq 0$, we have to consider the characteristic polynomial of N_β :

$$q_\alpha(\lambda, \beta) = (-1)^{p+1} \lambda^{p+1} + \sum_{i=0}^p c_i(\alpha, \beta) \lambda^i.$$

The coefficients $c_i(\alpha, \beta)$ can be computed by a symbolic calculus software such as Maple. We present them in appendix B for $1 \leq p \leq 5$.

In order to study the sign of the eigenvalues of N_β , we will use the following lemma.

Lemma 4.1. *Let P be a polynomial of degree n with n real roots such that $P(X) = \sum_{i=0}^n c_i X^i$. All the roots of P are non negative if and only if*

$$(-1)^i c_i \geq 0.$$

Proof. The proof of this lemma is given in appendix C. □

Hence, we have to find a condition over α such that, $\forall i \in \{0, \dots, p+1\}$, $\forall \beta \in [-\pi, \pi]$,

$$(-1)^i c_i(\alpha, \beta) \geq 0.$$

Here we only detail the computations in the case $p = 3$ and we give the expression of the characteristic polynomials for $p \neq 3$ in the appendix C. For $p = 3$, we have

$$\begin{cases} c_3(\alpha, \beta) = \frac{8}{h^2} ((15 - \alpha) \cos(\beta) - 4\alpha) \\ c_2(\alpha, \beta) = \frac{240}{h^4} (\cos^2(\beta) - (23 + \alpha) \cos(\beta) + (18\alpha - 65)) \\ c_1(\alpha, \beta) = \frac{2880}{h^6} (4 \cos^2(\beta) + (65 - 3\alpha) \cos(\beta) + (141 - 32\alpha)) \\ c_0(\alpha, \beta) = \frac{100800}{h^8} (3 \cos^2(\beta) + 2(3 - \alpha) \cos(\beta) + (2\alpha - 9)). \end{cases}$$

- Let us first study the condition over c_3 . We have, $\forall \beta \in [-\pi, \pi]$,

$$-c_3(\alpha, \beta) \geq 0 \Leftrightarrow (\alpha - 15) \cos(\beta) + 4\alpha \geq 0.$$

It is clear that this condition is satisfied for all β if and only if

$$\begin{cases} (\alpha - 15) + 4\alpha \geq 0, \\ (15 - \alpha) + 4\alpha \geq 0 \end{cases} \quad (13)$$

which implies that

$$\begin{cases} \alpha \geq 3, \\ \alpha \geq -5. \end{cases} \quad (14)$$

Consequently, $-c_3(\alpha, \beta) \geq 0$, whatever the choice of β , if and only if

$$\alpha \geq 3 \quad (15)$$

- Let us now consider the condition over c_2 .

Setting $X = \cos(\beta)$, this condition is equivalent to,

$$f_\alpha(X) := X^2 - (23 + \alpha)X + (18\alpha - 65) \geq 0, \forall X \in [-1, 1]. \quad (16)$$

This second order equation admits two roots:

$$\begin{cases} X_1 = \frac{1}{2} \left(23 + \alpha + \left((\alpha - 13)^2 + 620 \right)^{1/2} \right), \\ X_2 = \frac{1}{2} \left(23 + \alpha - \left((\alpha - 13)^2 + 620 \right)^{1/2} \right). \end{cases}$$

We know that $f_\alpha(X)$ is a second-order polynomial on the variable X and its head coefficient is positive. Thus, to have f_α non negative $\forall X \in [-1, 1]$, we need one of the following conditions:

1. the two roots are in $] -\infty; -1]$ i.e. $X_1 \leq -1$,
2. the two roots are in $]1; +\infty[$ i.e. $X_2 \geq 1$,
3. $X_1 = X_2$.

Since $(\alpha - 13)^2 + 620 > 0$, $X_1 > X_2$ and 3. is impossible.

The case $X_1 \leq -1$ is also impossible since $X_1 \geq 0$ when $\alpha \geq 0$ so we just have to consider the case $X_2 \geq 1$, which leads to the inequality

$$23 + \alpha - \left((\alpha - 13)^2 + 620 \right)^{1/2} \geq 2,$$

which is equivalent to

$$\alpha \geq \frac{87}{17}. \quad (17)$$

Finally, $c_2(\alpha, \beta) \geq 0$, whatever the choice of β , if and only if (17) holds.

- Now, let us study the sign of $c_1(\alpha, \beta)$.
Making the change of variable $X = \cos(\beta)$, the condition $-c_1(\alpha, \beta) \geq 0, \forall \beta \in [-\pi, \pi]$ is equivalent to

$$f_\alpha(X) := -4X^2 + (3\alpha - 65)X + (32\alpha - 141) \geq 0, \forall X \in [-1, 1].$$

The polynomial f_α admits the two following roots:

$$\begin{cases} X_1 = \frac{-1}{8} \left(65 - 3\alpha + \left(\left(3\alpha + \frac{61}{3} \right)^2 + \frac{14000}{9} \right)^{1/2} \right), \\ X_2 = \frac{-1}{8} \left(65 - 3\alpha - \left(\left(3\alpha + \frac{61}{3} \right)^2 + \frac{14000}{9} \right)^{1/2} \right). \end{cases}$$

Since, the head coefficient of the polynomial f_α is negative, we need $X_1 \leq -1$ and $X_2 \geq 1$.

The condition $X_1 \leq -1$ implies that

$$65 - 3\alpha + \left(\left(3\alpha + \frac{61}{3} \right)^2 + \frac{14000}{9} \right)^{1/2} \geq 8$$

which leads to

$$\alpha \geq \frac{80}{29}. \quad (18)$$

In the same way, the condition $X_2 \geq 1$ is equivalent to

$$\alpha \geq 6. \quad (19)$$

Consequently, $-c_1(\alpha, \beta) \geq 0, \forall \beta \in [-\pi, \pi]$ if and only if $\alpha \geq 6$.

- Finally, let us look at the positivity of $c_0(\alpha, \beta), \forall \beta \in [-\pi, \pi]$.
Once again, using the change of variable $X = \cos(\beta)$, we have

$$f_\alpha(X) := 3X^2 + 2(3 - \alpha)X + 2\alpha - 9 \geq 0.$$

This polynomial function f_α admits the two following roots:

$$\begin{cases} X_1 = \frac{1}{3}(2\alpha - 9), \\ X_2 = 1. \end{cases}$$

In the same way than previously, we need $X_1 \geq 1$ which leads to the condition

$$\alpha \geq 6. \quad (20)$$

In conclusion, taking into account the conditions (15), (17), (18), (19) and (20), we have

$$\lambda_{\min} \geq 0 \Leftrightarrow \alpha \geq 6. \quad (21)$$

We used the same technique to derive a condition over α for all polynomial degrees p from $p = 1$ to $p = 5$. Since the calculations are very similar, we do not detail them here and we just present them in Tab. 1.

p	c_0	c_1	c_2	c_3	c_4	c_5
1	$\alpha \geq 1$	$\alpha \geq 1$				
2	$\alpha \geq 3$	$\alpha \geq \frac{30}{11}$	$\alpha \geq 2$			
3	$\alpha \geq 6$	$\alpha \geq 6$	$\alpha \geq \frac{87}{17}$	$\alpha \geq 3$		
4	$\alpha \geq 10$	$\alpha \geq \frac{543}{55}$	$\alpha \geq \frac{325}{34}$	$\alpha \geq \frac{581}{73}$	$\alpha \geq 4$	
5	$\alpha \geq 15$	$\alpha \geq 15$	$\alpha \geq \frac{336}{23}$	$\alpha \geq \frac{1185}{86}$	$\alpha \geq \frac{124}{11}$	$\alpha \geq 5$

Table 1: Conditions over α for each coefficient c_i and each polynomial degree p

From these results, we can easily deduce the smallest penalization parameters ensuring the stability of the scheme (see Tab. 2). It is clear that, for $1 \leq p \leq 5$, the stability is guaranteed if and only if

$$\alpha \geq \frac{p(p+1)}{2h}.$$

p	1	2	3	4	5
	$\alpha \geq 1$	$\alpha \geq 3$	$\alpha \geq 6$	$\alpha \geq 10$	$\alpha \geq 15$

Table 2: Stability condition over α for each polynomial degree p

4.3 The CFL condition

Here, we propose to prove that the condition

$$\Delta t \leq \frac{2}{\sqrt{\lambda_{\max}}}$$

implies (7) and is equivalent to (7) and (8). To compute λ_{\max} , we use the implicit function theorem. Indeed, each eigenvalue $\lambda \in \Lambda(\beta)$ is an implicit function $\lambda(\beta)$ defined by $p_\alpha(\lambda(\beta), \beta) = 0$. Since $p_\alpha(\lambda, \beta)$ is a continuous and periodic function of β , so is $\lambda(\beta)$ and we only need to find all the values $\beta_0 \in [-\pi; \pi]$ such that $\lambda'(\beta_0) = 0$. Using the implicit functions theorem, we have

$$\frac{\partial p_\alpha}{\partial \lambda} \lambda'(\beta) + \frac{\partial p_\alpha}{\partial \beta} = 0.$$

Hence β_0 satisfies $\frac{\partial p_\alpha}{\partial \beta}(\lambda(\beta_0), \beta_0) = 0$ and the two steps of the proofs are, for a given α :

1. Find all the β_0 satisfying $\frac{\partial p_\alpha}{\partial \beta}(\lambda(\beta_0), \beta_0)$
2. Find the maximum value among all the $\lambda(\beta_0)$.

In the following, we only detail the case $p = 3$, since the proofs are similar for all $p = 1, \dots, 5$.

We have

$$\begin{aligned} \frac{\partial p_\alpha}{\partial \beta}(\alpha, \beta) &= \sin(\beta) \left[\frac{8}{h^2}(\alpha - 15)\lambda^3 + \frac{240}{h^4}(23 + \alpha - 2\cos(\beta))\lambda^2 \right. \\ &\quad \left. + \frac{2880}{h^6}(3\alpha - 65 - 8\cos(\beta))\lambda + \frac{100800}{h^8}(-6\cos(\beta) + 2(\alpha - 3)) \right] \\ &:= \sin(\beta) \tilde{p}_\alpha(\lambda, \beta). \end{aligned}$$

Consequently,

$$\frac{\partial p_\alpha}{\partial \beta}(\beta_0) = 0 \Leftrightarrow \begin{cases} \sin(\beta_0) = 0 \\ \text{or} \\ \tilde{p}_\alpha(\lambda, \beta_0) = 0. \end{cases} \quad (22)$$

First, we consider only the condition $\sin(\beta_0) = 0$ and we show that it is equivalent to (7). Then, to obtain the necessary and sufficient condition (8), we take into account the condition $\tilde{q}_\alpha(\lambda, \beta) = 0$.

- If $\beta_0 = 0$, we obtain the following eigenvalues:

$$0; \frac{60}{h^2}; \frac{90 + 20\alpha + 2g_1(\alpha)}{h^2}; \frac{90 + 20\alpha - 2g_1(\alpha)}{h^2}$$

where $g_1(\alpha) = (4545 - 1320\alpha + 100\alpha^2)^{\frac{1}{2}}$.

It is clear that, for $\alpha \geq 0$, the two greatest eigenvalues are $\lambda_1 = \frac{60}{h^2}$ and $\lambda_2 = \frac{1}{h^2}(90 + 20\alpha + 2g_1(\alpha))$. Studying the sign of the quantity

$$h^2(\lambda_1 - \lambda_2) = -150 + 20\alpha + 2g_1(\alpha)$$

we can easily obtain

$$\begin{cases} \lambda_{\max,0} = \lambda_2 & \text{if } \alpha \geq 6 \\ \lambda_{\max,0} = \lambda_1 & \text{if } \alpha < 6. \end{cases} \quad (23)$$

We have proved in section 4.2 that the condition $\alpha \geq 6$ is a necessary stability condition. Therefore, we only have to consider $\lambda_{\max,0} = \lambda_2$.

- If $\beta_0 = \pi$, we obtain the following eigenvalues:

$$\frac{2}{h^2}(45 + \sqrt{1605}); \frac{2}{h^2}(45 - \sqrt{1605}); \frac{2}{h^2}(-15 + 6\alpha + g_2(\alpha)); \frac{2}{h^2}(-15 + 6\alpha - g_2(\alpha))$$

where $g_2(\alpha) = (405 - 240\alpha + 36\alpha^2)^{\frac{1}{2}}$.

The two greatest eigenvalues are $\lambda_3 = \frac{2}{h^2}(45 + \sqrt{1605})$ and $\lambda_4 = \frac{2}{h^2}(-15 + 6\alpha + g_2(\alpha))$.

The study of the sign of $\lambda_3 - \lambda_4$ implies

$$\begin{cases} \lambda_{\max, \pi} = \lambda_3 & \text{if } \alpha \leq 10 \\ \lambda_{\max, \pi} = \lambda_1 & \text{if } \alpha \geq 10. \end{cases} \quad (24)$$

Now we have to compare $\lambda_{\max, 0}$ and $\lambda_{\max, \pi}$. We can easily verify that

$$\begin{cases} \lambda_{\max, 0} \geq \lambda_{\max, \pi} & \text{if } \alpha \geq \frac{2\sqrt{1605} + 393}{49}, \\ \lambda_{\max, \pi} > \lambda_{\max, 0} & \text{if } \frac{2\sqrt{1605} + 393}{49} \leq \alpha. \end{cases} \quad (25)$$

Consequently, considering (23) and (25)

$$\frac{\Delta t}{h} \leq \begin{cases} C_{1,3} & \text{if } 6 \leq \alpha \leq \frac{2\sqrt{1605} + 393}{49}, \\ C_{2,3}(\alpha) & \text{if } \frac{2\sqrt{1605} + 393}{49} < \alpha \end{cases} \quad (26)$$

which corresponds to the necessary condition (6).

- Let us now find β_0 such that $\tilde{p}(\lambda, \beta_0) = 0$.

We can easily obtain that

$$\cos(\beta_0) = \frac{(\alpha - 15)h^6\lambda^3 + 30(23 + \alpha)h^4\lambda^2 + 360(3\alpha - 65)h^2\lambda + 25200(\alpha - 3)}{60(h^4\lambda^2 + 48h^2\lambda + 1260)} := \tilde{Q}_\alpha(\lambda).$$

Using this expression of $\cos(\beta_0)$ in the characteristic polynomial (37) we obtain that $\lambda = \lambda(\beta_0)$ is solution to

$$q_\alpha(\lambda, \beta_0) = -\frac{1}{15h^8(h^4\lambda^2 + 48h^2\lambda + 1260)} \sum_{i=0}^6 \lambda^i h^{2i} \tilde{c}_i(\alpha) = 0$$

or, equivalently,

$$\tilde{Q}_{p,\alpha}(\lambda) = \sum_{i=0}^6 \lambda^i h^{2i} \tilde{c}_i(\alpha) = 0$$

with

$$\begin{cases} c_0(\alpha) = 635040000(\alpha^2 - 12\alpha + 36), \\ c_1(\alpha) = 3628800(15\alpha^2 + 70\alpha - 96), \\ c_2(\alpha) = 86400(31\alpha^2 - 447\alpha + 5316), \\ c_3(\alpha) = 14400(8\alpha^2 - 135\alpha - 1728), \\ c_4(\alpha) = 180(17\alpha^2 - 442\alpha + 7740), \\ c_5(\alpha) = 60(\alpha^2 + 16\alpha - 357), \\ c_6(\alpha) = \alpha^2 - 30\alpha + 210. \end{cases}$$

Finally, the eigenvalues λ must satisfy $Q_\alpha(\lambda) = 0$ and $|\tilde{Q}_\alpha(\lambda)| \leq 0$ and $\lambda_{\max} = \min(C_{1,3}, C_{2,3}(\alpha), C_{3,3}(\alpha))$.

5 The d -dimensional case

In this section, we propose to adapt the technique proposed in [9] to extend the analysis from the 1D case to the d D case. Here, we only detail the three-dimensional case, since the technique is exactly the same for the two dimensional case.

First of all, we consider an infinite homogeneous 3D domain Ω and we mesh it uniformly by cubes of edge h .

Now, we need to introduce the same notations as in [9].

- $\Omega = \bigcup_{K_J \in \mathcal{T}_h} K_J$ where $K_J = \prod_{k=1}^3 S_{J_k} = \prod_{k=1}^3 [J_k h, (J_k + 1) h]$ and $J = (J_k)_{k=1, \dots, 3}$.
- On $\hat{K} [0, 1]^d$, we define the Lagrange basis functions $(\hat{\varphi}_1)_{1 \in \{1, \dots, p+1\}^3}$ by

$$\hat{\varphi}_1(\mathbf{x}) = \prod_{k=1}^3 \hat{\varphi}_{l_k}(x_k)$$

where $\hat{\varphi}_{l_k}$ are the 1D Lagrange basis functions.

- Since the mesh is uniform, the basis functions are defined thanks to the functions $(\hat{\varphi}_1)_{1 \in \{1, \dots, p+1\}^3}$ by

$$\varphi_{\mathbf{m}}^J(\mathbf{x}) = \hat{\varphi}_{\mathbf{m}} \left(\frac{x - Jh}{h} \right) \mathbb{1}_{K_J}(\mathbf{x})$$

where $\mathbb{1}_{K_J}$ is the indicator function of K_J . These functions can be written as a product of d 1D basis functions:

$$\varphi_{\mathbf{m}}^J(\mathbf{x}) = \prod_{k=1}^3 \hat{\varphi}_{m_k} \left(\frac{x - J_k h}{h} \right) \mathbb{1}_{S_{J_k}}(\mathbf{x}_k)$$

- The different faces of the reference element \hat{K} are denoted by an exponent C corresponding to the orientation of the face: North (N), South (S), East (E), West (W), in Front of (F) and in the Back (B) (cf. Fig 1).

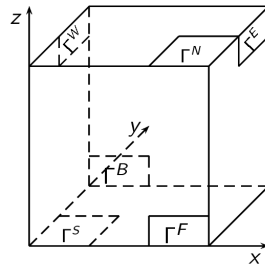


Figure 1: Notations of the faces in 3D

Since the mesh is uniform, we can rewrite the problem on an element $\mathbf{I} = \{I_1, I_2, I_3\}$:

$$\begin{aligned} M_{3,p} \delta^n U_{I_1, I_2, I_3} = & K_{3,p} U_{I_1, I_2, I_3} + K_{3,p}^E U_{I_1+1, I_2, I_3} + K_{3,p}^W U_{I_1-1, I_2, I_3} \\ & + K_{3,p}^F U_{I_1, I_2+1, I_3} + K_{3,p}^B U_{I_1, I_2-1, I_3} + K_{3,p}^N U_{I_1, I_2, I_3+1} \\ & + K_{3,p}^S U_{I_1, I_2, I_3-1} \end{aligned} \quad (27)$$

where

- $\delta^n U_{I_1, I_2, I_3} = \frac{U_{I_1, I_2, I_3}^{n+1} - 2U_{I_1, I_2, I_3}^n + U_{I_1, I_2, I_3}^{n-1}}{\Delta t^2}$

- $M_{3,p}$ is a block of the mass matrix M ,

$$M_{3,p}(\mathbf{i}, \mathbf{j}) = h^3 \int_{\hat{K}} \hat{\varphi}_{\mathbf{i}} \hat{\varphi}_{\mathbf{j}}, \quad \mathbf{i}, \mathbf{j} \in \{1, \dots, p+1\}^3$$

- $K_{3,p}$ is a diagonal block of the matrix K

$$\begin{aligned} K_{3,p}(\mathbf{i}, \mathbf{j}) = & h \int_{\hat{K}} \nabla \hat{\varphi}_{\mathbf{i}} \cdot \nabla \hat{\varphi}_{\mathbf{j}} - \sum_{C \in \{N, S, E, W, B, F\}} \frac{h}{2} \int_{\Gamma^C} (\hat{\varphi}_{\mathbf{i}} \nabla \hat{\varphi}_{\mathbf{j}} + \hat{\varphi}_{\mathbf{j}} \nabla \hat{\varphi}_{\mathbf{i}}) \nu_C \\ & + \sum_{C \in \{N, S, E, W, B, F\}} h^2 \int_{\Gamma^C} \alpha \hat{\varphi}_{\mathbf{i}} \hat{\varphi}_{\mathbf{j}}, \quad \mathbf{i}, \mathbf{j} \in \{1, \dots, p+1\}^3 \end{aligned}$$

where ν_C is the outward unit normal vector to the face Γ^C .

- $K_{3,p}^C$ is a block of the matrix K corresponding to the interactions between an element I and its neighbours on the face Γ^C :

$$\begin{aligned} K_{3,p}^E(\mathbf{i}, \mathbf{j}) = & \int_{[0,1]^2} \frac{1}{2} (\hat{\varphi}_{\mathbf{i}}(1, x_2, x_3) \nabla \hat{\varphi}_{\mathbf{j}}(0, x_2, x_3) + \hat{\varphi}_{\mathbf{j}}(0, x_2, x_3) \nabla \hat{\varphi}_{\mathbf{i}}(1, x_2, x_3)) \nu_E \\ & - \hat{\varphi}_{\mathbf{i}}(1, x_2, x_3) \hat{\varphi}_{\mathbf{j}}(0, x_2, x_3) \end{aligned}$$

$$\begin{aligned} K_{3,p}^F(\mathbf{i}, \mathbf{j}) = & \int_{[0,1]^2} \frac{1}{2} (\hat{\varphi}_{\mathbf{i}}(x_1, 1, x_3) \nabla \hat{\varphi}_{\mathbf{j}}(x_1, 0, x_3) + \hat{\varphi}_{\mathbf{j}}(x_1, 0, x_3) \nabla \hat{\varphi}_{\mathbf{i}}(x_1, 1, x_3)) \nu_F \\ & - \hat{\varphi}_{\mathbf{i}}(x_1, 1, x_3) \hat{\varphi}_{\mathbf{j}}(x_1, 0, x_3) \end{aligned}$$

$$\begin{aligned} K_{3,p}^N(\mathbf{i}, \mathbf{j}) = & \int_{[0,1]^2} \frac{1}{2} (\hat{\varphi}_{\mathbf{i}}(x_1, x_2, 1) \nabla \hat{\varphi}_{\mathbf{j}}(x_1, x_2, 0) + \hat{\varphi}_{\mathbf{j}}(x_1, x_2, 0) \nabla \hat{\varphi}_{\mathbf{i}}(x_1, x_2, 1)) \nu_N \\ & - \hat{\varphi}_{\mathbf{i}}(x_1, x_2, 1) \hat{\varphi}_{\mathbf{j}}(x_1, x_2, 0) \end{aligned}$$

$$K_{3,p}^W(i_1, i_2, i_3, j_1, j_2, j_3) = K_{3,p}^E(j_1, i_2, i_3, i_1, j_2, j_3)$$

$$K_{3,p}^B(i_1, i_2, i_3, j_1, j_2, j_3) = K_{3,p}^F(i_1, j_2, i_3, j_1, i_2, j_3)$$

$$K_{3,p}^S(i_1, i_2, i_3, j_1, j_2, j_3) = K_{3,p}^N(i_1, i_2, j_3, j_1, j_2, i_3)$$

Then, multiplying the equation (27) by the inverse of the mass matrix $M_{3,p}$, we obtain

$$\begin{aligned} \delta^n U_{I_1, I_2, I_3} = & N_{3,p} U_{I_1, I_2, I_3} + N_{3,p}^E U_{I_1+1, I_2, I_3} + N_{3,p}^W U_{I_1-1, I_2, I_3} \\ & + N_{3,p}^F U_{I_1, I_2+1, I_3} + N_{3,p}^B U_{I_1, I_2-1, I_3} + N_{3,p}^N U_{I_1, I_2, I_3+1} \\ & + N_{3,p}^S U_{I_1, I_2, I_3-1} \end{aligned} \quad (28)$$

where $N_{3,p} = M_{3,p}^{-1} K_{3,p}$ and $N_{3,p}^C = M_{3,p}^{-1} K_{3,p}^C$.

Now, we are interested in rewriting the matrices $N_{3,p}$ and $N_{3,p}^C$ with respect to the matrices we have obtained for the one dimensional case.

5.1 From the 3 dimensional case to the one dimensional case

The coefficients of $M_{3,p}$, $K_{3,p}$, $K_{3,p}^C$, $N_{3,p}$ and $N_{3,p}^C$ can be deduced from the coefficients of $M_{1,p}$, $K_{1,p}$, $K_{1,p}^W$, $N_{1,p}$ and $N_{1,p}^W$ thanks to the following theorem.

Theorem 5.1. For all $\mathbf{m} = (m_k)_{k=1, \dots, 3} \in \{1, \dots, p+1\}^3$ and $\mathbf{n} = (n_k)_{k=1, \dots, 3} \in \{1, \dots, p+1\}^3$, we have

$$\begin{aligned} 1. \quad M_{3,p}(\mathbf{m}, \mathbf{n}) &= \prod_{i=1}^3 M_{1,p}(m_i, n_i), \\ 2. \quad K_{3,p}(\mathbf{m}, \mathbf{n}) &= \sum_{i=1}^3 \left(K_{1,p}(m_i, n_i) \prod_{k=1, k \neq i}^3 M_{1,p}(m_k, n_k) \right) \\ 3. \quad K_{3,p}^C(\mathbf{m}, \mathbf{n}) &= K_{1,p}^W(m_{p_C}, n_{p_C}) \prod_{k=1, k \neq p_C}^3 M_{1,p}(m_k, n_k) \\ 4. \quad N_{3,p}(\mathbf{m}, \mathbf{n}) &= \sum_{p=1}^3 N_{1,p}(m_p, n_p) \prod_{k=1, k \neq p}^3 \delta_{m_k, n_k} \\ 5. \quad N_{3,p}^C(\mathbf{m}, \mathbf{n}) &= N_{1,p}^W(m_{p_C}, n_{p_C}) \prod_{k=1, k \neq p_C}^3 \delta_{m_k, n_k} \end{aligned} \quad (29)$$

$$\text{where } p_C = \begin{cases} 1 & \text{if } C \in \{E, W\}, \\ 2 & \text{if } C \in \{N, S\}, \\ 3 & \text{if } C \in \{B, F\}, \end{cases} \quad \text{and } N_{1,p} = M_{1,p}^{-1} K_{1,p} \text{ and } N_{1,p}^C = M_{1,p}^{-1} K_{1,p}^C.$$

Proof.

- *Proof of 1.*

Considering the notations and the results of the section 5, we have

$$\begin{aligned}
M_{3,p}(\mathbf{m}, \mathbf{n}) &= h^3 \int_{\hat{K}} \hat{\varphi}_{\mathbf{m}} \hat{\varphi}_{\mathbf{n}} d\mathbf{x} \\
&= h^3 \int_{\hat{K}} \prod_{i=1}^3 \hat{\varphi}_{m_i}(x_i) \prod_{i=1}^3 \hat{\varphi}_{n_i}(x_i) \prod_{i=1}^3 dx_i \\
&= \prod_{i=1}^3 h \int_{[0,1]} \hat{\varphi}_{m_i}(x_i) \hat{\varphi}_{n_i}(x_i) dx_i \\
&= \prod_{i=1}^3 M_{1,p}(m_i, n_i)
\end{aligned}$$

- *Proof of 2.*

We first have the following lemma for the volum term.

Lemma 5.2. *For all $\mathbf{m} = (m_k)_{k=1,\dots,3} \in \{1, \dots, p+1\}^3$ and $\mathbf{n} = (n_k)_{k=1,\dots,3} \in \{1, \dots, p+1\}^3$, we have*

$$\int_{\hat{K}} \nabla \hat{\varphi}_{\mathbf{m}} \cdot \nabla \hat{\varphi}_{\mathbf{n}} d\mathbf{x} = \sum_{i=1}^3 \left(\int_{[0,1]} \frac{\partial \hat{\varphi}_{m_i}}{\partial x_i}(x_i) \frac{\partial \hat{\varphi}_{n_i}}{\partial x_i}(x_i) dx_i \prod_{k=1, k \neq i}^3 M_{1,p}(m_k, n_k) \right).$$

Proof. Let us look at the term $h^3 \int_{\hat{K}} \nabla \hat{\varphi}_{\mathbf{m}} \cdot \nabla \hat{\varphi}_{\mathbf{n}} d\mathbf{x}$.

We know that, $\forall \mathbf{m} \in \{1, \dots, p+1\}^d$

$$\hat{\varphi}_{\mathbf{m}}(\mathbf{x}) = \prod_{k=1}^3 \hat{\varphi}_{m_k}(x_k)$$

which implies that

$$\frac{\partial \hat{\varphi}_{\mathbf{m}}}{\partial x_k}(\mathbf{x}) = \frac{\partial \hat{\varphi}_{m_k}}{\partial x_k}(x_k) \prod_{i=1, i \neq k}^3 \hat{\varphi}_{m_i}(x_i).$$

Then, using the same reasoning as previously,

$$\begin{aligned}
h^3 \int_{\hat{K}} \nabla \hat{\varphi}_{\mathbf{m}} \cdot \nabla \hat{\varphi}_{\mathbf{n}} d\mathbf{x} &= \sum_{i=1}^3 h \int_{[0,1]} \frac{\partial \hat{\varphi}_{m_i}}{\partial x_i}(x_i) \frac{\partial \hat{\varphi}_{n_i}}{\partial x_i}(x_i) dx_i \prod_{k=1, k \neq i}^3 h \int_{[0,1]} \hat{\varphi}_{m_k}(x_k) \hat{\varphi}_{n_k}(x_k) dx_k \\
&= \sum_{i=1}^3 \left(h \int_{[0,1]} \frac{\partial \hat{\varphi}_{m_i}}{\partial x_i}(x_i) \frac{\partial \hat{\varphi}_{n_i}}{\partial x_i}(x_i) dx_i \prod_{k=1, k \neq i}^3 M_{1,p}(m_k, n_k) \right)
\end{aligned}$$

which ends the proof. \square

Now, we have to deal with the surface terms.

Let us first remark that

$$\nabla \hat{\varphi}_{\mathbf{m}} \cdot \nu|_{\Gamma^C} = \frac{\partial \hat{\varphi}_{m_{p_C}}}{\partial x_{p_C}}(x_{p_C}) \nu_{1,C} \prod_{k=1, k \neq p_C}^3 \hat{\varphi}_{m_k}(x_k) \quad (30)$$

where $\nu_{1,C}$ is the outward unit normal vector in the one dimensional case defined by

$$\nu_{1,C} = \begin{cases} 1 & \text{if } C \in \{E, N, F\}, \\ -1 & \text{if } C \in \{W, S, B\} \end{cases}$$

and x_{p_C} is defined by

$$x_{p_C} = \begin{cases} 1 & \text{if } C \in \{E, N, F\}, \\ 0 & \text{if } C \in \{W, S, B\}. \end{cases}$$

Then, we can propose the following lemma.

Lemma 5.3. *For all 2-D surface, we have $\forall \mathbf{m}, \mathbf{n} \in \{1, \dots, p+1\}^d$*

$$\begin{aligned} h^2 \int_{\Gamma^C} \hat{\varphi}_{\mathbf{m}} (\nabla \hat{\varphi}_{\mathbf{n}} \cdot \nu) d\sigma &= \hat{\varphi}_{m_{p_C}}(x_{p_C}) \frac{\partial \hat{\varphi}_{n_{p_C}}}{\partial x_{p_C}}(x_{p_C}) \nu_{1,C} \prod_{k=1, k \neq p_C}^3 M_{1,p}(m_k, n_k) \\ h^2 \int_{\Gamma^C} \alpha \hat{\varphi}_{\mathbf{m}} \hat{\varphi}_{\mathbf{n}} d\sigma &= \alpha \hat{\varphi}_{m_{p_C}}(x_{p_C}) \hat{\varphi}_{n_{p_C}}(x_{p_C}) \prod_{k=1, k \neq p_C}^3 M_{1,p}(m_k, n_k) \end{aligned}$$

Proof. First of all, using (30), we have

$$\begin{aligned} h^2 \int_{\Gamma^C} \hat{\varphi}_{\mathbf{m}} (\nabla \hat{\varphi}_{\mathbf{n}} \cdot \nu) d\sigma &= h^2 \int_{\Gamma^C} \prod_{k=1}^3 \hat{\varphi}_{m_k}(x_k) \frac{\partial \hat{\varphi}_{n_{p_C}}}{\partial x_{p_C}}(x_{p_C}) \nu_{1,C} \prod_{k=1, k \neq p_C}^3 \hat{\varphi}_{n_k}(x_k) \prod_{k=1, k \neq p_C}^3 dx_k \\ &= \hat{\varphi}_{m_{p_C}}(x_{p_C}) \frac{\partial \hat{\varphi}_{n_{p_C}}}{\partial x_{p_C}}(x_{p_C}) \nu_{1,C} \prod_{k=1, k \neq p_C}^3 h^2 \int_{[0,1]} \hat{\varphi}_{m_k}(x_k) \hat{\varphi}_{n_k}(x_k) dx_k \end{aligned}$$

which can be rewritten as

$$\int_{\Gamma^C} \hat{\varphi}_{\mathbf{m}} (\nabla \hat{\varphi}_{\mathbf{n}} \cdot \nu) d\sigma = \hat{\varphi}_{m_{p_C}}(x_{p_C}) \frac{\partial \hat{\varphi}_{n_{p_C}}}{\partial x_{p_C}}(x_{p_C}) \nu_{1,C} \prod_{k=1, k \neq p_C}^3 M_{1,p}(m_k, n_k).$$

In the same way, for the penalization term, we have:

$$\begin{aligned} h^2 \int_{\Gamma^C} \alpha \hat{\varphi}_{\mathbf{m}} \hat{\varphi}_{\mathbf{n}} d\sigma &= \hat{\varphi}_{m_{p_C}}(x_{p_C}) \hat{\varphi}_{n_{p_C}}(x_{p_C}) h^2 \int_{\Gamma^C} \alpha \prod_{k=1, k \neq p_C}^3 \hat{\varphi}_{m_k}(x_k) \prod_{k=1, k \neq p_C}^3 \hat{\varphi}_{n_k}(x_k) \prod_{k=1, k \neq p_C}^3 d\sigma_k \\ &= \alpha \hat{\varphi}_{m_{p_C}}(x_{p_C}) \hat{\varphi}_{n_{p_C}}(x_{p_C}) \prod_{k=1, k \neq p_C}^3 h^2 \int_{[0,1]} \hat{\varphi}_{m_k}(x_k) \hat{\varphi}_{n_k}(x_k) dx_k \end{aligned}$$

which clearly implies that

$$\int_{\Gamma^C} \alpha \hat{\varphi}_{\mathbf{m}} \hat{\varphi}_{\mathbf{n}} d\sigma = \alpha \hat{\varphi}_{m_{p_C}}(x_{p_C}) \hat{\varphi}_{n_{p_C}}(x_{p_C}) \prod_{k=1, k \neq p_C}^3 M_{1,p}(m_k, n_k)$$

which ends the proof. \square

Finally, using the two lemmas, we obtain

$$K_{3,p}(\mathbf{m}, \mathbf{n}) = \sum_{i=1}^3 \left(K_{1,p}(m_i, n_i) \prod_{k=1, k \neq i}^3 M_{1,p}(m_k, n_k) \right). \quad (31)$$

- *Proof of 3.*

To rewrite the terms $K_{3,p}^C(\mathbf{m}, \mathbf{n})$ for all \mathbf{m}, \mathbf{n} , we use a similar reasoning as for $K_{3,p}$.

- *Proof of 4.*

To prove 4. and 5., we need the following lemma.

Lemma 5.4. *Let $\mathbf{m} = (m_p)_{p=1,\dots,3}$ and $\mathbf{n} = (n_p)_{p=1,\dots,3}$. We have*

$$M_{3,p}^{-1}(\mathbf{m}, \mathbf{n}) = \prod_{k=1}^3 M_{1,p}^{-1}(m_k, n_k).$$

Proof. Let A be the matrix defined $\forall \mathbf{m}, \mathbf{n} \in \{1, \dots, p+1\}^d$ by

$$A(\mathbf{m}, \mathbf{n}) = \prod_{k=1}^3 M_{1,p}^{-1}(m_k, n_k).$$

We have

$$\begin{aligned} (AM)(\mathbf{m}, \mathbf{n}) &= \sum_{l_1, \dots, l_3=1}^3 \prod_{k=1}^3 M_{1,p}^{-1}(m_k, l_k) \prod_{k=1}^3 M_{1,p}(l_k, n_k) \\ &= \sum_{l_1, l_2=1}^3 \prod_{k=1}^2 M_{1,p}^{-1}(m_k, l_k) \prod_{k=1}^2 M_{1,p}(l_k, n_k) \sum_{l_3=1}^3 M_{1,p}^{-1}(m_3, l_3) M_{1,p}(l_3, n_3) \\ &= \prod_{k=1}^3 \left(\sum_{l_k=1}^3 M_{1,p}^{-1}(m_k, l_k) M_{1,p}(l_k, n_k) \right) \end{aligned}$$

But,

$$\sum_{l_k=1}^3 M_{1,p}^{-1}(m_k, l_k) M_{1,p}(l_k, n_k) = (M_{1,p}^{-1} M_{1,p})(m_k, n_k) = \delta_{m_k, n_k}$$

so that

$$(AM)(\mathbf{m}, \mathbf{n}) = \prod_{k=1}^3 \delta_{m_k, n_k} = I(\mathbf{m}, \mathbf{n})$$

where I is the identity matrix, which ends the proof. \square

Let us now consider the matrix $N_{3,p} = M_{3,p}^{-1} K_{3,p}$.

First, we rewrite $K_{3,p}$ as

$$K_{3,p} = \sum_{q=1}^3 T_q,$$

with $T_q(\mathbf{m}, \mathbf{n}) = K_{1,p}(m_q, n_q) \prod_{k=1, k \neq q}^3 M_{1,p}(m_k, n_k)$.

Then, $M_{3,p}^{-1} K_{3,p} = \sum_{q=1}^3 M_{3,p}^{-1} T_q$ and, using lemma 5.4

$$\begin{aligned} (M_{3,p}^{-1} T_1)(\mathbf{m}, \mathbf{n}) &= \sum_{l_1, \dots, l_3=1}^{p+1} \left(\prod_{k=1}^3 M_{1,p}^{-1}(m_k, l_k) K_{1,p}(l_1, n_1) \prod_{k=2}^3 M_{1,p}(l_k, n_k) \right) \\ &= \sum_{l_1, \dots, l_3=1}^{p+1} M_{1,p}^{-1}(m_1, l_1) K_{1,p}(l_1, n_1) M_{1,p}^{-1}(m_2, l_2) M_{1,p}(l_2, n_2) M_{1,p}^{-1}(m_3, l_3) M_{1,p}(l_3, n_3) \\ &= \sum_{l_1=1}^{p+1} M_{1,p}^{-1}(m_1, l_1) K_{1,p}(l_1, n_1) \prod_{k=2}^3 \left(\sum_{l_k=1}^{p+1} M_{1,p}^{-1}(m_k, l_k) M_{1,p}(l_k, n_k) \right). \end{aligned}$$

So that

$$(M_{3,p}^{-1} T_1)(\mathbf{m}, \mathbf{n}) = N_{1,p}(m_1, n_1) \delta_{m_2, n_2} \delta_{m_3, n_3}.$$

Performing the same calculations for T_2 and T_3 , we obtain

$$N_{3,p}(\mathbf{m}, \mathbf{n}) = \sum_{p=1}^3 N_{1,p}(m_p, n_p) \prod_{k=1, k \neq p}^3 \delta_{m_k, n_k}.$$

- *Proof of 5.*

We apply the technique used to prove 4. to show that

$$N_{3,p}^C(\mathbf{m}, \mathbf{n}) = N_{1,p}^W(m_{p_C}, n_{p_C}) \prod_{k=1, k \neq p_C}^3 \delta_{m_k, n_k}$$

□

5.2 Consequences on the stability analysis

Let us now apply a Fourier transform in the three directions to (28) to obtain, for $\beta = [-\pi, \pi]^3$,

$$\delta^n \tilde{U}_{\beta_1, \beta_2, \beta_3} = N_\beta \tilde{U}_{\beta_1, \beta_2, \beta_3} \quad (32)$$

where the matrix N_β is defined by

$$\begin{aligned} N_\beta(\mathbf{m}, \mathbf{n}) &= \sum_{j=1}^3 N_{1,p}(m_j, n_j) \prod_{q=1, q \neq j}^3 \delta_{m_q, n_q} \\ &+ e^{i\beta_j} N_{1,p}^W(m_j, n_j) \prod_{q=1, q \neq j}^3 \delta_{m_q, n_q} \\ &+ e^{-i\beta_j} N_{1,p}^W(n_j, m_j) \prod_{q=1, q \neq j}^3 \delta_{m_q, n_q} \end{aligned}$$

which can be rewritten as:

$$\begin{aligned} N_{\beta}(\mathbf{m}, \mathbf{n}) &= \sum_{j=1}^3 \left((N_{1,p}(m_j, n_j) + e^{i\beta_j} N_{1,p}^W(m_j, n_j) + e^{-i\beta_j} N_{1,p}^W(n_j, m_j)) \prod_{q=1, q \neq j}^3 \delta_{m_q, n_q} \right) \\ &= \sum_{j=1}^3 N_{\beta_j}(m_j, n_j) \prod_{q=1, q \neq j}^3 \delta_{m_q, n_q}. \end{aligned}$$

Using the stability analysis as in section 4.1, the stability of the scheme is ensured if and only if

$$\lambda_{\min,3} \geq 0 \quad \text{and} \quad \lambda_{\max,3} \leq \frac{4}{\Delta t^2}$$

where $\lambda_{\min,3} = \min_{\beta \in [-\pi, \pi]^3} (\min \Lambda(N_{\beta}))$, $\lambda_{\max,3} = \max_{\beta \in [-\pi, \pi]^3} (\max \Lambda(N_{\beta}))$ and $\Lambda(N_{\beta})$ is the set of eigenvalues of N_{β} .

To compute these values, we use the following lemma.

Lemma 5.5. *Let $(\lambda_i^{\beta})_{i=1, \dots, p+1}$ denote the eigenvalues of N_{β} and $(v_i^{\beta})_{i=1, \dots, p+1}$ the associated eigenvectors. Then, the eigenvalues of N_{β} are given by*

$$\lambda_{\mathbf{i}}^{\beta} = \sum_{k=1}^3 \lambda_{i_k}^{\beta_k}$$

where $\lambda_{i_k}^{\beta_k}$ are the eigenvalues of N_{β_k} , and the associated eigenvectors by

$$v_{\mathbf{i}}^{\beta}(\mathbf{m}) = \prod_{k=1}^3 v_{i_k}^{\beta_k}(m_k). \quad (33)$$

Proof. Let $v_{\mathbf{i}}^{\beta}$ defined by (33), then

$$\begin{aligned} (N_{\beta} v_{\mathbf{i}}^{\beta})(\mathbf{m}) &= \sum_{n_1, \dots, n_3=1}^{p+1} \left(\sum_{q=1}^3 N_{\beta_q}(m_q, n_q) \prod_{k=1, k \neq q}^3 \delta_{m_k, n_k} \prod_{k=1}^3 v_{i_k}^{\beta_k}(n_k) \right) \\ &= \sum_{q=1}^3 \left(\left(\sum_{n_q=1}^{p+1} N_{\beta_q}(m_q, n_q) v_{i_q}^{\beta_q}(n_q) \right) \prod_{k=1, k \neq q}^3 v_{i_k}^{\beta_k}(m_k) \right) \\ &= \sum_{q=1}^3 \left(\lambda_{i_q}^{\beta_q} v_{i_q}^{\beta_q}(m_q) \prod_{k=1, k \neq q}^3 v_{i_k}^{\beta_k}(m_k) \right) \\ &= \left(\sum_{q=1}^3 \lambda_{i_q}^{\beta_q} \right) v_{\mathbf{i}}^{\beta}(\mathbf{m}) \end{aligned}$$

□

It is then clear that

$$\lambda_{\min,3} = 3\lambda_{\min} \quad \text{and} \quad \lambda_{\max,3} = 3\lambda_{\max}.$$

Hence, it is clear that the scheme (33) is stable, if and only if

$$\lambda_{\min} \geq 0 \quad \text{and} \quad \frac{c\Delta t}{h} \leq \frac{1}{\sqrt{3}} \sqrt{\frac{2}{\lambda_{\max}}}$$

which is equivalent to the stability conditions (6), (7) and (8).

5.3 Extension to rectangular or parallelepiped mesh

For the sake of simplicity, we restricted our theorem to the case of squared or cubic mesh. However, one can extend the proof to the case of rectangular or parallelepipeds meshes to show that a necessary stability condition is in 2D:

$$\alpha \geq \frac{p(p+1)}{2 \min(h_x, h_y)}$$

and in 3D:

$$\alpha \geq \frac{p(p+1)}{2 \min(h_x, h_y, h_z)}.$$

Here, h_x , h_y and h_z denote respectively the length of the edges of the elements in the x , y and z direction. The minimal value of h_x , h_y and h_z is actually the diameter of the inscribed sphere of each element. This remark confirms that the third definition of ξ_F using the diameter of the inscribed sphere or circle in 2D is the most appropriate.

The proof could also be extended to obtain a CFL condition, but its expression is complicated and does not add much insight.

6 Numerical results

In this section, we first represent the behaviour of the CFL condition with respect to α (section 6.1). This illustrates the fact that theorem 3.1 is actually necessary and sufficient for most of the values of α . Then, we compare our analytical CFL condition in infinite domain to the CFL condition computed numerically on finite meshes in order to illustrate the validity of theorem 3.3 (section 6.2).

6.1 Behaviour of the CFL condition with respect to α

In Fig. 2, 4, 6, 8 and 10 we plot the functions $C_{1,p}$ (blue line with diamonds), $C_{2,p}(\alpha)$ (red line with circles) and $C_{3,p}(\alpha)$ (black line) respectively for $p = 1, 2, 3, 4$ and 5.

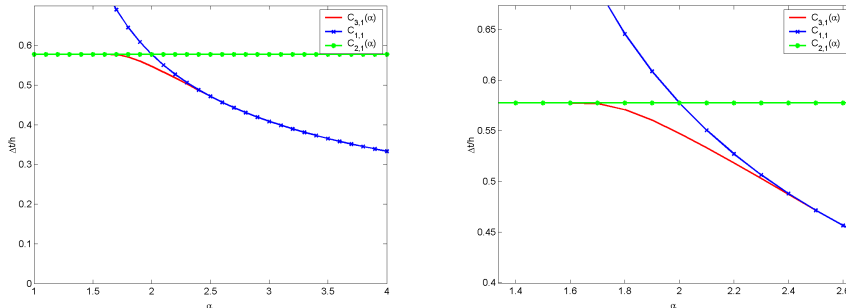
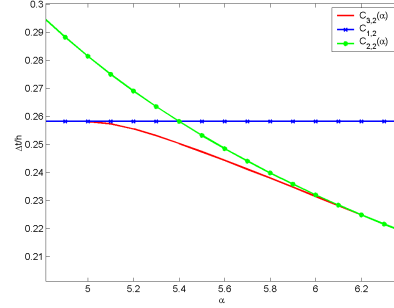
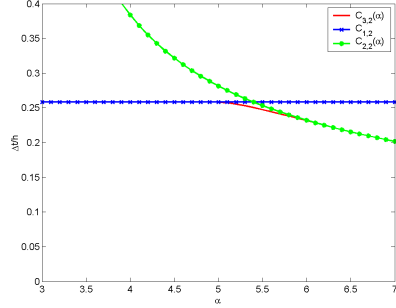
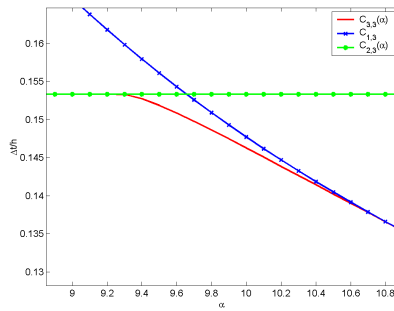
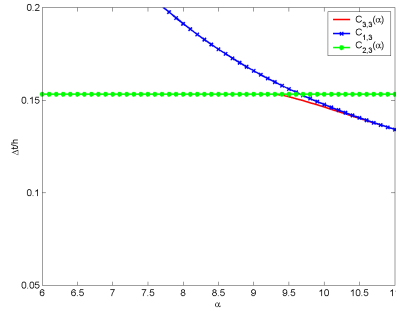
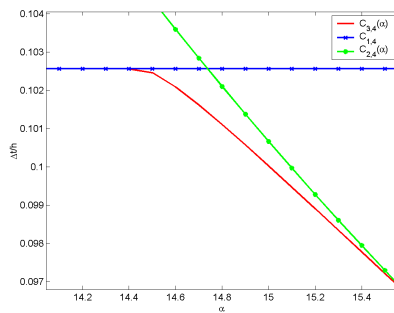
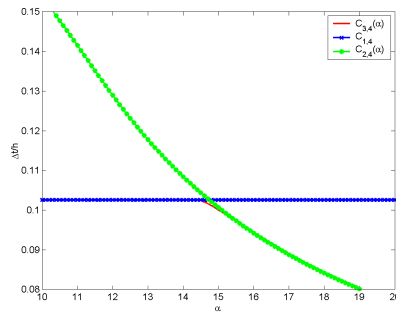


Figure 2: The 3 conditions for $p = 1$ Figure 3: Zoom on the 3 cond. for $p = 1$

Figure 4: The 3 conditions for $p = 2$ Figure 5: Zoom on the 3 cond. for $p = 2$ Figure 6: The 3 conditions for $p = 3$ Figure 7: Zoom on the 3 cond. for $p = 3$ Figure 8: The 3 conditions for $p = 4$ Figure 9: Zoom on the 3 cond. for $p = 4$

6.2 Comparison with numerical experiments

In this section, we compare the results we have obtained previously with numerical experiments. We consider the simulation of wave propagation in an homogeneous 1D domain $\Omega = [0, 10]$ with a velocity $c = (\mu/\rho)^{1/2} = 1 \text{ ms}^{-1}$. We impose also Dirichlet boundary conditions at the both ends of the domain and the length of the space step is $h = 0.1$.

We computed numerically the greatest eigenvalue of the matrix $M^{-1}K$ and we

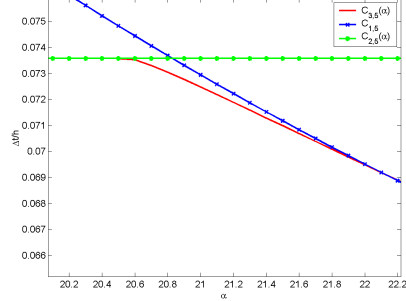
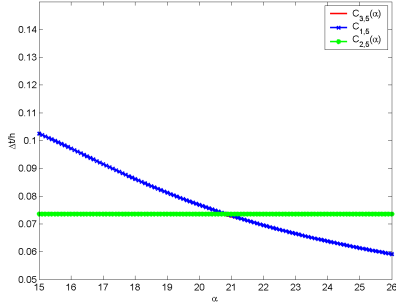


Figure 10: The 3 conditions for $p = 5$ $p = 5$ Figure 11: Zoom on the 3 cond. for

deduced the CFL condition of the scheme. In Fig. 12, 13, 14, 15 and 16 we compare the analytical CFL (red line) obtained by theorem 3.3 to the numerical CFL (triangles), respectively for $p = 1, 2, 3, 4$ and 5. All figures show a very good agreement between the analytical and the numerical CFL.

The function $C_{3,p}(\alpha)$ only modifies the CFL condition in a small segment around $\alpha_{1,p}$. The behaviour is confirmed in Fig. 3, 5, 7, 9 and 11 which represent a zoom around $\alpha_{1,p}$. Theses numerical results confirm the fact that theorem 3.1 provides actually a necessary and sufficient condition except for a small range of α . Moreover, the CFL condition remains constant for α from $\frac{p(p+1)}{2}$ to a close value of α_p^1 , which means that it is not necessary to choose $\alpha = \frac{p(p+1)}{2}$ to optimize the time step.

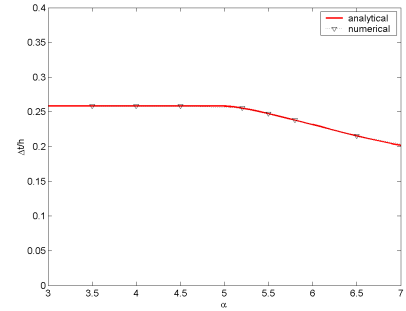
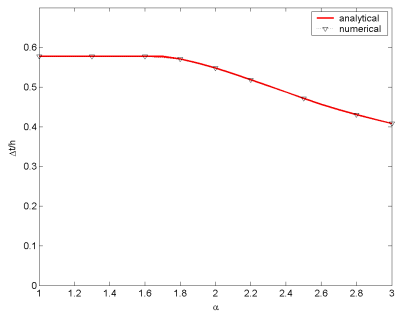
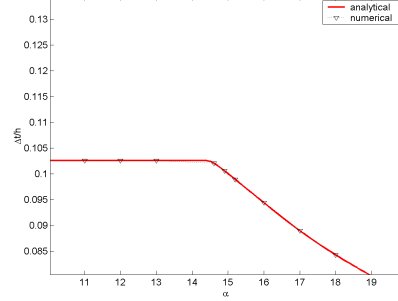
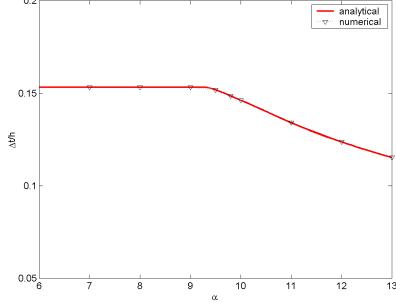
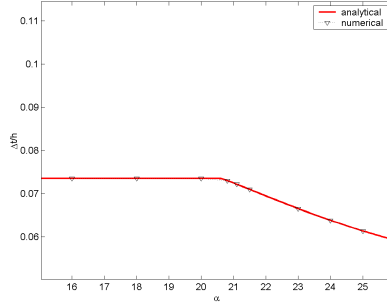


Figure 12: Numerical comparison in P^1 Figure 13: Numerical comparison in P^2

7 Conclusion

In this paper, we have proposed sufficient and necessary conditions of L^2 -stability of an IPDG method using regular meshes. These conditions confirm the conjecture of Ainsworth, Monk and Muniz up to $p = 5$. Moreover, we have observed that the CFL condition is constant with respect to α on a segment

Figure 14: Numerical comparison in P^3 Figure 15: Numerical comparison in P^4 Figure 16: Numerical comparison in P^5

$\left[\frac{p(p+1)}{2}, \tilde{\alpha} \right]$ and is decreasing as $\alpha^{-1/2}$ for $\alpha > \tilde{\alpha}$. This means that it is not necessary to choose α to close to $\frac{p(p+1)}{2}$ to improve the CFL condition. Finally, we have observed that a good choice for ξ_F should be the diameter of the inscribed circle (or sphere). This should be confirmed by an analysis on triangular meshes, which will be the topic of a future work.

A Definition of $Q_{p,\alpha}$ and $\tilde{Q}_{p,\alpha}$

We present here the expressions of the polynomial $Q_{p,\alpha}$ and the rational function $\tilde{Q}_{p,\alpha}$ for $1 \leq p \leq 5$.

- For the polynomials of degree 1, $\tilde{Q}_{p,\alpha}$ is defined by

$$\tilde{Q}_{p,\alpha}(\lambda) = \frac{h^2 \lambda}{2} \left(\frac{\alpha}{3} - 1 \right) + \alpha.$$

We have,

$$Q_{p,\alpha}(\lambda) = \sum_{i=0}^2 \lambda^i h^{2i} \tilde{c}_i(\alpha)$$

with

$$\begin{cases} \tilde{c}_0(\alpha) = 36(\alpha^2 - 2\alpha + 1), \\ \tilde{c}_1(\alpha) = 12(\alpha^2 - \alpha), \\ \tilde{c}_2(\alpha) = \alpha^2 - 6\alpha + 6, \end{cases}$$

- In the case $p = 2$, the definition of $\tilde{Q}_{p,\alpha}$ is

$$\tilde{Q}_{p,\alpha}(\lambda) = -\frac{(\alpha - 1)h^4\lambda^2 + 4(15 + 4\alpha)h^2\lambda + 240(\alpha - 3)}{24(h^2\lambda + 20)}.$$

The polynomial $Q_{p,\alpha}$ is such that

$$Q_{p,\alpha}(\lambda) = \sum_{i=0}^4 \lambda^i h^{2i} \tilde{c}_i(\alpha)$$

with

$$\begin{cases} \tilde{c}_0(\alpha) = 57600(\alpha^2 - 2\alpha + 1), \\ \tilde{c}_1(\alpha) = 1920(4\alpha^2 - 43\alpha + 39), \\ \tilde{c}_2(\alpha) = 16(46\alpha^2 - 342\alpha + 1521), \\ \tilde{c}_3(\alpha) = 8(4\alpha^2 + \alpha - 140), \\ \tilde{c}_4(\alpha) = \alpha^2 - 16\alpha + 56. \end{cases}$$

- For $p = 3$, we have

$$Q_{p,\alpha}(\lambda) = \sum_{i=0}^6 \lambda^i h^{2i} \tilde{c}_i(\alpha)$$

with

$$\begin{cases} \tilde{c}_0(\alpha) = 635040000(\alpha^2 - 12\alpha + 36), \\ \tilde{c}_1(\alpha) = 3628800(15\alpha^2 + 70\alpha - 96), \\ \tilde{c}_2(\alpha) = 86400(31\alpha^2 - 447\alpha + 5316), \\ \tilde{c}_3(\alpha) = 14400(8\alpha^2 - 135\alpha - 1728), \\ \tilde{c}_4(\alpha) = 180(17\alpha^2 - 442\alpha + 7740), \\ \tilde{c}_5(\alpha) = 60(\alpha^2 + 16\alpha - 357), \\ \tilde{c}_6(\alpha) = \alpha^2 - 30\alpha + 210. \end{cases}$$

and $\tilde{Q}_{p,\alpha}$ is defined by

$$\tilde{Q}_{p,\alpha}(\lambda) = \frac{(\alpha - 15)h^6\lambda^3 + 30(23 + \alpha)h^4\lambda^2 + 360(3\alpha - 65)h^2\lambda + 25200(\alpha - 3)}{60(h^4\lambda^2 + 48h^2\lambda + 1260)}.$$

- For the polynomials of degree 4, $\tilde{Q}_{p,\alpha}$ is defined by

$$\tilde{Q}_{p,\alpha}(\lambda) = -\frac{\tilde{A}_{p,\alpha}(\lambda)}{120(169344 + h^6\lambda^3 + 84h^4\lambda^2 + 5040h^2\lambda)}$$

with

$$\begin{aligned}\tilde{A}_{p,\alpha}(\lambda) = & \lambda^4 h^8 (\alpha - 24) + 12\lambda^3 h^6 (4\alpha + 287) + 1008\lambda^2 h^4 (3\alpha - 305) \\ & + 20160\lambda h^2 (8\alpha - 165) + 5080320 (\alpha - 1)\end{aligned}$$

And, we have

$$Q_{p,\alpha}(\lambda) = \sum_{i=0}^8 \lambda^i h^{2i} \tilde{c}_i(\alpha).$$

with

$$\left\{ \begin{array}{l} \tilde{c}_0(\alpha) = 25809651302400 (\alpha - 6)^2, \\ \tilde{c}_1(\alpha) = 204838502400 (8\alpha^2 - 357\alpha + 1854), \\ \tilde{c}_2(\alpha) = 81285120 (698\alpha^2 + 3882\alpha + 292185), \\ \tilde{c}_3(\alpha) = 203212800 (72\alpha^2 - 13791\alpha - 328), \\ \tilde{c}_4(\alpha) = 48384 (719\alpha^2 - 12750\alpha + 2419275), \\ \tilde{c}_5(\alpha) = 8064 (76\alpha^2 - 972\alpha - 286209), \\ \tilde{c}_6(\alpha) = 144 (58\alpha^2 - 5282\alpha + 201609), \\ \tilde{c}_7(\alpha) = 24 (4\alpha^2 + 241\alpha - 6972), \\ \tilde{c}_8(\alpha) = \alpha^2 - 48\alpha + 55. \end{array} \right.$$

- For $p = 5$, $\tilde{Q}_{p,\alpha}$ is such that

$$\tilde{Q}_{p,\alpha}(\lambda) = \frac{\tilde{A}_{p,\alpha}(\lambda)}{210(39916800 + \lambda^4 h^8 + 128\lambda^3 h^6 + 12960\lambda^2 h^4 + 967680\lambda h^2)}$$

with

$$\begin{aligned}\tilde{A}_{p,\alpha}(\lambda) = & \lambda^5 h^{10} (\alpha - 35) + 70\lambda^4 h^8 (\alpha + 168) + 6720\lambda^3 h^6 (\alpha - 351) \\ & + 302400\lambda^2 h^4 (2\alpha + 385) + 8467200\lambda h^2 (5\alpha - 303) + 1676505600 (\alpha - 1)\end{aligned}$$

and we have

$$Q_{p,\alpha}(\lambda) = \sum_{i=0}^{10} \lambda^i h^{2i} \tilde{c}_i(\alpha)$$

with

$$\left\{ \begin{array}{l} \tilde{c}_0(\alpha) = 2810671026831360000(\alpha - 15)^2, \\ \tilde{c}_1(\alpha) = 28390616432640000(5\alpha + 168)(\alpha - 15), \\ \tilde{c}_2(\alpha) = 10241925120000(373\alpha^2 - 35067\alpha + 855423), \\ \tilde{c}_3(\alpha) = 3072577536000(24\alpha^2 - 895\alpha - 159240), \\ \tilde{c}_4(\alpha) = 1016064000(1151\alpha^2 - 11360\alpha + 24379995), \\ \tilde{c}_5(\alpha) = 67737600(257\alpha^2 - 3570\alpha - 9096540), \\ \tilde{c}_6(\alpha) = 2822400(76\alpha^2 - 2417\alpha + 2988895), \\ \tilde{c}_7(\alpha) = 67200(32\alpha^2 + 2383\alpha - 974820), \\ \tilde{c}_8(\alpha) = 140(131\alpha^2 - 37062\alpha + 2285235), \\ \tilde{c}_9(\alpha) = 140(\alpha^2 + 151\alpha - 5912), \\ \tilde{c}_{10}(\alpha) = \alpha^2 - 70\alpha + 1190. \end{array} \right.$$

B Necessary condition of stability

- In the case of discontinuous basis functions of degree 1, we can easily obtain the following characteristic polynomial associated to the matrix A_β

$$p_\alpha(\lambda, \beta) = \lambda^2 + c_1(\alpha, \beta)\lambda + c_0(\alpha, \beta) \quad (34)$$

with

$$\left\{ \begin{array}{l} c_1(\alpha, \beta) = \frac{4}{h^2}((3 - \alpha)\cos(\beta) - 2\alpha) \\ c_0(\alpha, \beta) = \frac{12}{h^4}(\cos^2(\beta) - 2\alpha\cos(\beta) + 2\alpha - 1). \end{array} \right.$$

- In the case of discontinuous basis functions of degree 2, the characteristic polynomial associated to the matrix A_β is

$$p_\alpha(\lambda, \beta) = -\lambda^3 + c_2(\alpha, \beta)\lambda^2 + c_1(\alpha, \beta)\lambda + c_0(\alpha, \beta) \quad (35)$$

with

$$\left\{ \begin{array}{l} c_2(\alpha, \beta) = -\frac{6}{h^2}((\alpha - 8)\cos(\beta) - 3\alpha) \\ c_1(\alpha, \beta) = \frac{12}{h^4}(-6\cos^2(\beta) - 2(15 + 4\alpha)\cos(\beta) + 4(24 - 13\alpha)) \\ c_0(\alpha, \beta) = -\frac{1440}{h^6}(\cos^2(\beta) + (\alpha - 3)\cos(\beta) + 2 - \alpha). \end{array} \right.$$

- In the case of discontinuous basis functions of degree 4, the characteristic polynomial associated to the matrix A_β is

$$p_\alpha(\lambda, \beta) = -\lambda^5 + c_4(\alpha, \beta)\lambda^4 + c_3(\alpha, \beta)\lambda^3 + c_2(\alpha, \beta)\lambda^2 + c_1(\alpha, \beta)\lambda + c_0(\alpha, \beta) \quad (36)$$

with

$$\begin{cases} c_4(\alpha, \beta) = -\frac{10}{h^2} ((\alpha - 24) \cos(\beta) - 5\alpha) \\ c_3(\alpha, \beta) = -\frac{120}{h^4} (5 \cos^2(\beta) + (4\alpha + 287) \cos(\beta) + 10(15\alpha - 88)\alpha) \\ c_2(\alpha, \beta) = -\frac{10080}{h^6} (5 \cos^2(\beta) + (3\alpha - 305) \cos(\beta) + 990 - 133\alpha) \\ c_1(\alpha, \beta) = -\frac{201600}{h^8} (15 \cos^2(\beta) + (8\alpha + 165) \cos(\beta) + 2(59\alpha - 468)) \\ c_0(\alpha, \beta) = -\frac{50803200}{h^{10}} (2 \cos^2(\beta) + (\alpha - 10) \cos(\beta) + 8 - \alpha). \end{cases}$$

- In the case of discontinuous basis functions of degree 5, the characteristic polynomial associated to the matrix A_β is

$$p_\alpha(\lambda, \beta) = \lambda^6 + c_5(\alpha, \beta) \lambda^5 + c_4(\alpha, \beta) \lambda^4 + c_3(\alpha, \beta) \lambda^3 + c_2(\alpha, \beta) \lambda^2 + c_1(\alpha, \beta) \lambda + c_0(\alpha, \beta) \quad (37)$$

with

$$\begin{cases} c_5(\alpha, \beta) = -\frac{12}{h^2} ((\alpha - 35) \cos(\beta) + 6\alpha) \\ c_4(\alpha, \beta) = -\frac{420}{h^4} (-3 \cos^2(\beta) + (2\alpha + 336) \cos(\beta) + 1155 - 134\alpha) \\ c_3(\alpha, \beta) = -\frac{40320}{h^6} (-4 \cos^2(\beta) + (2\alpha - 702) \cos(\beta) + 256\alpha - 2849\alpha) \\ c_2(\alpha, \beta) = -\frac{1814400}{h^8} (-9 \cos^2(\beta) + (4\alpha + 770) \cos(\beta) + 9343 - 326\alpha) \\ c_1(\alpha, \beta) = -\frac{101606400}{h^{10}} (-12 \cos^2(\beta) + (5\alpha - 303) \cos(\beta) + 94\alpha - 1170) \\ c_0(\alpha, \beta) = -\frac{10059033600}{h^{12}} (-5 \cos^2(\beta) + (2\alpha - 20) \cos(\beta) + 25 - 2\alpha). \end{cases}$$

C Proof of lemma 4.1

Herein, we just present the case where n is even but we can easily extend this proof to the odd case without any difficulty. In this case, we clearly have :

$$\begin{cases} \sigma_0^n = \prod_{i=1}^n \lambda_i \\ \sigma_p^n = \sum_{i_1=1}^n \left[\sum_{i_2 > i_1}^n \cdots \sum_{i_{p+1} > i_p}^n \prod_{\substack{j=1 \\ j \neq i_p, \forall p \in \{1, \dots, p+1\}}}^n \lambda_k \right] \text{ if } 1 \leq p \leq n-2 \\ \sigma_{n-1}^n = \sum_{i=1}^n \lambda_i \end{cases}$$

where $\lambda_i, i = 1, \dots, n$ denote the roots of the polynomial P .
 We can remark that $c_i = (-1)^i \sigma_i^n$ and then we have to show that

$$\lambda_i \geq 0 \iff \sigma_i^n \geq 0, \forall i$$

- It is clear that if all the roots λ_i are positive then for all i

$$\sigma_i^n \geq 0.$$

- Now, we assume that for all $i \in \{0, \dots, n-1\}, \sigma_i^n \geq 0$.

We assume that there exists a root strictly negative, we denote this root λ_n .

Then, we have

$$\sigma_0^{n-1} \leq 0.$$

Moreover, we can easily remark that

$$\sigma_p^n = \sigma_{p-1}^{n-1} + \lambda_n \sigma_p^{n-1}.$$

Let us prove by induction that $\sigma_p^{n-1} \leq 0$ for $0 \leq p \leq n-2$.

We have already seen that the property is checked for $p=0$.

Let us suppose that $\sigma_{p-1}^{n-1} \leq 0$. Then, as $\lambda_n < 0$ and $\sigma_p^n \geq 0$ by hypothesis, we have

$$\sigma_p^{n-1} \leq 0.$$

Consequently, in particular, we have

$$\sigma_{n-2}^{n-1} \leq 0,$$

but $\sigma_{n-2}^{n-1} = \sum_{i=1}^n \lambda_i$ which implies that

$$\sigma_{n-1}^n < 0$$

which is impossible by hypothesis then all the roots λ_i are positive.

Contents

1	Introduction	3
2	Discretization of the acoustic wave equation	4
3	Stability analysis	6
4	Study in the 1-Dimensional case	8
4.1	Fourier Analysis of the IPDG scheme in 1D	9
4.2	Study of the condition $\lambda_{\min} \geq 0$	10
4.3	The CFL condition	13
5	The d-dimensional case	16
5.1	From the 3 dimensional case to the one dimensional case	18
5.2	Consequences on the stability analysis	22
5.3	Extension to rectangular or parallelepiped mesh	24

6	Numerical results	24
6.1	Behaviour of the CFL condition with respect to α	24
6.2	Comparison with numerical experiments	25
7	Conclusion	26
A	Definition of $Q_{p,\alpha}$ and $\tilde{Q}_{p,\alpha}$	27
B	Necessary condition of stability	30
C	Proof of lemma 4.1	31

References

- [1] M. Ainsworth, P. Monk, and W. Muniz. Dispersive and dissipative properties of discontinuous galerkin finite element methods for the second-order wave equation. *Journal of Scientific Computing*, 27, 2006.
- [2] D. N. Arnold, F. Brezzi, B. Cockburn, and L.D. Marini. Unified analysis of discontinuous galerkin methods for elliptic problems. *SIAM J. Numer. Anal.*, 39:1749–1779, 2002.
- [3] C. Baldassari. Modélisation et simulation numérique pour la migration terrestre par équation d’ondes. *PhD Thesis*, 2009.
- [4] F. Bassi and S. Rebay. A high-order accurate discontinuous finite element method for the numerical solution of the compressible navier-stokes equations. *J. Comput. Phys.*, 131:267–279, 1997.
- [5] G. Benitez Alvarez, A.F. Dourado Loula, E.G. Dutra do Carmo, and A. Alves Rochinha. A discontinuous finite element formulation for helmholtz equation. *Comput. Methods. Appl. Mech. Engrg.*, 195:4018–4035, 2006.
- [6] S. Cohen, P. Joly, J.E. Roberts, and N. Tordjman. Higher-order triangular finite elements with mass-lumping for the wave equation. *SIAM J. Numer. Anal.*, Vol. 44, 6:2408–2431, 2006.
- [7] S. Cohen, P. Joly, and N. Tordjman. Higher-order finite elements with mass-lumping for the 1d wave equation. *Finite Elements in Analysis and Design*, Vol. 16, Issues 3-4:329–336, 1994.
- [8] J. D. De Basabe and M. K. Sen. Stability of the high-order finite elements for acoustic or elastic wave propagation with high-order time stepping. *Geophys. J. Int.*, 181:577–590, 2010.
- [9] S. Fauqueux. Eléments finis mixtes spectraux et couches absorbantes parfaitement adaptées pour la propagation d’ondes élastiques en régime transitoire. *PhD Thesis*, 2003.
- [10] M. J. Grote, A. Schneebeli, and D. Schötzau. Discontinuous galerkin finite element method for the wave equation. *SIAM J. on Numerical Analysis*, 44:2408–2431, 2006.

-
- [11] M. J. Grote and D. Schötzau. Convergence analysis of a fully discrete discontinuous galerkin method for the wave equation. *Preprint No. 2008-04*, 2008.
 - [12] D. Komatitsch and J. Tromp. Introduction to the spectral element method for three-dimensional seismic wave propagation. *Geophys. J. Int.*, 139:806–822, 1999.
 - [13] G. Seriani and E. Priolo. Spectral element method for acoustic wave simulation in heterogeneous media. *Finite Elements in Analysis and Design*, Vol. 16 Issues 3-4:337–348, 1994.
 - [14] K. Shahbazi. An explicit expression for the penalty parameter of the interior penalty method. *J. of Computational Physics*, 205:401–407, 2005.



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