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Real-time realistic illumination and shading of stratiform clouds

Antoine Bouthors, Fabrice Neyret, Sylvain Lefebvre
EVASION-GRAVIR/IMAG-INRIA

Abstract

Realistic rendering of clouds involves solving the complex interaction of light within the cloud and with its environment. Interactive methods achieve efficient cloud rendering by ignoring several lighting effects. However, these effects are visually important, and removing them strongly reduces realism.

We present a novel approach for capturing the important effects of multiple anisotropic Mie scattering within cloud layers (i.e., stratiform clouds), and the inter-reflections between the ground and the cloud base under sun and sky illumination. Our model maps well to graphics hardware, enabling the real-time rendering of animated cloud skies over landscapes.

1. Introduction

The visual aspect of a cloud depends on complex light interactions: individual cloud droplets scatter light according to a very irregular diagram, the Mie function [BH98], which is responsible for glories and fogbows around the antisolunar direction. Inside the cloud, light is anisotropically scattered multiple times before leaving. Despite this accumulated diffusion, anisotropy remains strong at many places such as along the silhouette (thus the “silver lining”) or in the direction of the “pseudo-specular” reflection. The ground-clouds interaction introduces further lighting effects. Since a cloud’s albedo is high, the cloud’s bottom reflects the ground. Inuits can find a channel in ice shelf or find the land from the sea by observing “water sky” and “ice blink” on the cloud bottom [Art].

The complete simulation of a cloud appearance including all these effects has never been done. Several interactive cloud simulations have been proposed, but they increase performance by ignoring most light interaction modes. Our model revisits the cloud illumination problem and accounts for new important lighting effects (see Fig. 2, right). It enables the real-time rendering of animated cloudy layers over a landscape. In this paper, we focus on stratiform clouds such as stratocumulus, i.e., not vertically developed like cumulus or cumulonimbus (see Fig. 1). Also, we do not consider clouds made of ice. Our goal is to render such clouds in real time from a viewpoint located either under the clouds (e.g., from the ground) or above them (e.g., from a plane). We are not considering flying trough or inside the clouds.

Our contributions are:

- An improved representation of the scattering model, based on the exact Mie function (whereas previous approaches use course approximations) and droplet size distributions.
- Reflectance and transmittance shaders to be applied to the cloud surface. These shaders account for local scattering and estimate the effects of multiple light scattering through the cloud.
- A GPU-friendly algorithm to efficiently compute inter-reflections between the clouds base and the ground (accounting for shadows, sky, and landscape texture).

Our approach consists in obtaining functions giving the intensity and direction distribution of light on each side of a homogeneous plane-parallel slab. To do so, we separately treat the different orders of scattering within this slab (see Fig. 2, left). We use analytical solutions for orders 0 and 1, and approximate solutions for orders 2 and ≥ 3. We propose a modified Mie model allowing us to treat only meaningful scattering events, i.e., events that cause a significant change in the photon’s path direction. We apply the resulting functions as a shader applied on the cloud surface, locally considered as a slab.

We account for the inter-reflection between the cloud base and the ground (considered as parallel planes) by introducing an efficient GPU-compliant algorithm performing radiosity iterations: density and color variations on the two planes are stored in textures, as well as their dynamic light-
2. Previous Work

An accurate way of simulating the multiple anisotropic Mie scattering of light in a cloud is to rely on photon tracing, i.e., Monte Carlo integration. However, this would require an unacceptably long time to converge. Various simplifications have been proposed in the literature:

- Low albedo / low density hypothesis: one single scattering is considered [Kv84, Bl82].
- Simpler scattering functions: Rayleigh [HL01, Kv84], Gaussian [PAT*04], Heneyy-Greenstein [Max94, NND96], which are easier and faster to compute and whose smooth variations present directional artifacts.
- Simplification into simple diffusion solvers: isotropic scattering [DKY*00], diffusion approximation [Sta95, JMLH01].
- Forward-dominant hypothesis: single-pass algorithms [HL01], very convenient for GPU-enhanced interactive rendering.

Only a few papers actually make use of Mie scattering to render atmospheric effects [REK*04, JG97]. Only [REK*04] use it for clouds, and only for local effects such as the glory (the light transport is assumed to be forward-dominant).

The problem is that visual features of clouds (see Fig. 5) are due to different lighting modes directly or indirectly connected to the particular shape of the Mie scattering diagram (see Section 3.3): Glories and fogbows are essentially due to single scattering and strongly depend on the backward peaks of the phase function. “Pseudo-specular” reflection and bright silhouettes are caused by the forward anisotropy and are the effect of a few scattering events (2 to 10). The overall brightness of clouds comes from diffuse illumination (high orders of scattering cancels directionalities in the diagram). Thus, considering only diffusion, or only forward scattering, or smooth scattering functions, forbids realism. And simply adding a Mie single scattering on top of a forward [REK*04] or isotropic diffusion is not sufficient either.

The study of most probable paths [PAT*04] is an interesting way to focus on important features. This accounts for the low-order scatterings, but it misses the diffusion effect. Moreover, [PAT*04] relies on Gaussian phase function.

Cloud optics has been widely studied in other fields (meteorology, applied optics, atmospheric sciences etc.). In particular, the study of cloud droplet phase functions [Dei64] coupled with cloud droplet size distribution [Cla74, MVC00] are of interest to us.

3. Cloud optics

3.1. Multiple scattering in clouds

Multiple scattering has been extensively described in the literature (see for instance [JC98]). Let us recall the main principles of interest to us (the following notation will be used throughout the paper).

The material of a homogeneous cloud can be described by the effective radius of its droplets \( r_e \) and the density of droplets per cubic meter \( \rho_0 \). These define the extinction cross-section \( \sigma = 4\pi r_e^2 \), the extinction coefficient \( \kappa = \rho_0 \sigma \), and the extinction function (i.e., the cloud transparency) \( \tau(x) = e^{-\kappa x} \) which is the probability to traverse the cloud along a path of length \( x \) without hitting a droplet. Thus, the probability of hit within \( x \) meters is \( 1 - \tau(x) \), and its derivative \( s(x) = \kappa e^{-\kappa x} \) is the probability density of a hit exactly at the end of a free path of length \( x \). This also yields the mean free path \( \ell_0 = 1/\kappa \).

When hitting a particle, light can either be absorbed or scattered. The albedo of cloud droplets is very close to 1 for visible light, i.e., light is not absorbed in clouds: all the light that hits is distributed elsewhere. The probability that light hitting a droplet from direction \( \omega \) is scattered in direction \( \omega' \) is determined by the phase function \( P(\theta) \) with \( \theta = \omega, \omega' \).

Stratiform clouds usually hold between \( N_0 = 10^6 \) and \( 10^9 \) droplets per m\(^3\) having a radius \( r_e = 2 \, \mu m \) to 15 \( \mu m \). This yields a mean free path \( 1/\kappa \) of several meters (20m for \( N_0 = 3.10^6 \, m^{-3} \) and \( r_e = 7 \, \mu m \)). Since a cloud spans hundreds to thousands of meters in each direction, most rays are
scattered several times before exiting the volume (for the example above, the no-hit probability is $10^{-2}$ for 100 m and $10^{-10}$ for 500 m).

### 3.2. Droplet size distributions in clouds

The size of water droplets has a huge visual impact: for a given amount of water, the smaller the droplets the more opaque the cloud (half the size doubles the extinction coefficient). The Mie function also varies with respect to the droplets size. Low- and mid-level clouds are a mixture of water droplets of different sizes. Literature in applied optics and atmospheric sciences rely on droplet size distribution (DSD) models. The modified Gamma distribution [Lev58, Dei69, Cla74] is well suited for cloud droplets. It is defined by:

$$N(r) = \frac{N_0}{\Gamma(\gamma/r_n)} \left(\frac{r}{r_n}\right)^{\gamma-1} e^{-\frac{r}{r_n}}$$

![Figure 3: Example DSD.](image)

This function describes the density $N(r)$ of droplets of radius $r$, with $r_n$ the characteristic radius of the distribution, $\gamma$ representing its broadness, and $N_0$ the total density of droplets. $\Gamma$ is the Gamma function. A cloud distribution is thus only described by the three parameters $N_0$, $r_n$, and $\gamma$. In terms of optical properties, the effective radius corresponding to this distribution is $r_e = (\gamma + 2)/r_n$.

Measured parameters for all kinds of clouds can be found in the literature. For stratiform clouds, we refer to [MVC00]. In the examples of this paper, we used values of $N_0 = 300 \text{cm}^{-3}$, $\gamma = 2$ and $r_e = 7 \mu m$. The resulting DSD is shown on figure 3. We consider this distribution to be constant in space and time.

![Figure 4: RGB: efficient Mie phase function (log). Gray: pure Mie ($r = 20 \mu m$). Note the undesired high frequency ripples. Inset: Polar plot of the function (log).](image)

### 3.3. Scattering behavior

The phase function $P_\lambda(\theta)$ characterizing the scattering caused by a water droplet of size $r$ at wavelength $\lambda$ is obtained by the Mie theory [BH98]. It is quite expensive to compute, highly dependent on $r$ and $\lambda$, and its intensity oscillates strongly when $\theta$ varies. This is what makes the Mie function so impractical, and probably why it is avoided in CG cloud simulations.

The effective phase function $P(\theta)$ corresponding to a DSD can be obtained by weighting and summing the phase functions of the various droplet sizes [Dei64]. We used Philip Laven’s MiePlot software [Lav] to pre-compute the effective phase function for our droplet distributions (see Section 3.2).

Besides the more accurate physics—which shows especially on the glory [MSPS04]—considering the DSD has a very fortunate consequence: the high frequency ripples in $P_\lambda(\theta)$ vanish in $P(\theta)$ (and, indeed, they are not observed in real clouds). So, the effective phase function is a lot more practical than the pure one.

Each feature of this function directly corresponds to an important visual phenomenon [Min54, LL95, Cow] in the cloud:

- **Narrow forward:** there is a strong peak in the forward direction. It corresponds to $5^\circ$ concentrating 51% of the function. It is responsible for the silver lining on the silhouette when the sun is behind the cloud, since a few scattering events occur for these light paths (see fig. 5, right).
- **Wide forward:** This large lobe represents 48% of the scattering. This part produces a fairly isotropic distribution when several scattering events are accumulated. It also causes the “pseudo-specular” reflection in the antisolar azimuthal direction on the top of the cloud layer, since less scattering events are required for light paths in this direction to get out of the layer (see fig. 5, middle).
- **39° backward peak:** This peak is responsible for the fogbow (or cloudbow), a bow where the contrast of the cloud layer is enhanced. It is visually similar to the rainbow with poor or no colors. This effect is caused by single scattering event occurring close to the cloud surface and represents very little energy. It is more visible on mist and thin layers (see fig. 5, left).
- **Backward peak:** This $2^\circ$-wide peak is responsible for the glory, a colored circle around the antisolar point which is often seen from plane, and the antisolar spot, which is generally masked by the shadow of the observer. These features are caused by single scattering events occurring close to the cloud surface (see fig. 5, left).
4. Our model

As stated in the overview, we consider different categories of light paths. First, we separate the strong forward scattering events $F$ from other scattering events $S$. We then approximate light paths $\sum_i (F^i S) F^i$ as $F^i + \sum (S')$, where $S'$ corresponds to a modified Mie model: the forward peak is embedded into the extinction function (see Section 4.1).

Furthermore, we study the contributions of $S^1, S^2, S^3 \geq 3$ paths for a given outgoing direction of reflection or transmission. The study of these modes is done on plane-parallel layers of thickness $H$ with a homogeneous droplet distribution. Clouds are assumed to behave locally like slabs of such thickness. We are looking for the intensity $I$ outgoing in direction $\vec{V}$ when a slab is hit from the sun direction $\vec{T}$.

The 0-scatter corresponds to the transparency and strong forward scattering (Section 4.2). The single scattering contribution can be obtained analytically (Section 4.3). The double scattering is integrated thanks to an approximation (Section 4.4). Note that integrating the energy of these 3 modes with our modified Mie model is equivalent to averaging over integrating paths with up to 5 scatters with the standard Mie model. This is due to the fact that the strong forward scattering event are embedded in the extinction function. We treat the remainder (the $3^\geq$ mode) as diffuse, and we obtain its importance via a simple 1D model (Section 4.5).

The total reflectance and transmittance are used to obtain the amount of reflected/transmitted environment light (coming from the sky and from the ground). The computation of inter-reflections between the cloud base and the ground is treated in Section 5.

Notation: Let $\vec{V}, \vec{T}$ and $\vec{N}$ be the view direction, the light direction, and the normal to the slab (i.e., the up or down direction) at the considered illuminated location. $\theta_0, \theta_1, \theta_2$ are the angles between $\vec{V}$ and $\vec{N}$, $\vec{T}$ and $\vec{N}$, and $\vec{V}$ and $\vec{T}$. $\mu_0 = \cos \theta_0$, $\mu_1 = \cos \theta_1$, $H_1 = H/\mu_0$ and $H_2 = H/\mu_1$ are the thickness of the slab along the light and view direction, respectively. The subscripts 0, 1, 2, 3 and $ms$ refer to light paths undergoing no scattering, single scattering, double scattering, three scattering events or more, or multiple scattering (i.e., $0 + 1 + 2 + 3^\geq$), respectively. The superscripts $r$ and $t$ indicate that we are computing the light reflected by the slab or transmitted through the slab. Typically, we will consider $I^r$ when the viewer and the light are on the same side of the layer, and $I^t$ otherwise (see Figure 2).

In the following sections, we estimate the contributions $I^0_0, I^0_1, I^1_1, I^2_2, I^2_3, I^3_3$, leading to $I^r = I^0_0 + I^0_1 + I^1_1 + I^2_2 + I^2_3 + I^3_3$ and $I^t = I^0_0 + I^0_1 + I^1_1$. If $\vec{V} \cdot \vec{T} \geq 0$ then $I^r = I^t$, otherwise $I^r = I^t$.

4.1. Modified-Mie: forward scattering embedding

In order to treat $F$ events separately from $S$ events, we define a modified Mie model embedding the forward scattering. Let $P_F = \frac{1}{4\pi} \int_0^{\theta_{\max}} 2\pi \sin(\theta) P(\theta) d\theta$ be the weight of the forward peak for $\theta < \theta_T$. For typical cloud values, $P_F \approx 51\%$ for $\theta_T = 5^\circ$. In paths containing at least one $S$ event we neglect the direction change for $F$ events, i.e., scattering events with $\theta < \theta_T$. This is equivalent to decreasing the extinction cross section according to $P_F$ and canceling the peak of the phase function. The phase function and extinction parameters presented in 3.1 are thus modified as follows:

\[ P_S(\theta) = \frac{P(\theta)}{(1 - P_F)^2} \]  
\[ \sigma_S = (1 - P_F) \sigma_S^0, \quad \tau_S(x) = e^{\sigma_S x}, \quad s_S(x) = \kappa_S e^{-\kappa_S x} \]

For $P_F = 51\%$, half of the events are implicitly accounted for in the modified extinction, so we only have to treat explicitly the other half. This doubles the mean free path and halves the number of scattering events per light ray.

The consequence of this approximation is a $5^\circ$ error at each extinction event. But since multiple scattering tends to smooth out directionality, this error is negligible for high orders of scattering. Figure 6 shows a comparison between a Monte Carlo simulation done with the original and the modified Mie scheme. Note that we do not apply this approximation for paths containing no $S$ event, since the deviation would not be negligible. $F^r$ paths (called “0-scatter”) are instead incorporated into the transparency, see Section 4.2.

In the following section we use our modified-Mie model in place of the original one. That is, notation $P(\theta), \sigma, \kappa, \tau(x), s(x)$ represents $P_S, \sigma_S, \kappa_S, \tau_S, s_S$. $P_F$ stands for the strong forward peak function, with $P_F(\theta) = \frac{P(\theta)}{P_F}$ if $\theta < \theta_T$ otherwise 0.

4.2. Transparency and 0-scattering

Transparency is obtained using the extinction function. Using the original $\tau(x)$ function provides the direct transmittance

\[ I^t_0 = \tau(H_0) I_0 \theta - \tau \]

with $\delta(x)$ the Kronecker function. Using $\tau_S(x)$ adds the effect of strong forward scattering. We prefer to provide this...
term separately, taking into account the $5^\circ$ deviation of the forward peak:  
\[ I_0' = T_2(H_2) P_2(\theta_2) \delta(H_2) \]

Note that $P_2(\theta_2) \delta(H_2)$ is the strong forward peak convolved by the mean number of scattering, which equals $\delta H_i$. We have $I_0' = I_0'd + I_0'\phi$.

4.3. Single scattering

The amount of light getting out of the slab in direction $\bar{V}$ after one scatter can be integrated analytically along the ray as in [Bli82] (see Fig. 7, left):  
\[ I'_1 = \int_0^H s \left( \frac{1}{\mu_t} \right) P(\theta_1) \tau(H_1 - 1 - \tau(H_1 + H_0)) \frac{dH}{\mu_t} \]

\[ I'_2 = \int_0^H s \left( \frac{1}{\mu_t} \right) \int_0^{\bar{V}} \left( \bar{P}(\theta, \psi) \tau(H_1 + H_0) \psi d\psi \right) \frac{dH}{\mu_t} \]

In order to approximate the two last integrals, we use a common approximation [REK*04]: the weight of all the paths scattered by $\psi$ is considered to be constant over $\psi$. The double integral becomes $\tau(H_1 + H_0) \int_0^{\bar{V}} \psi d\psi P(\theta, \psi) \tau(H_1 + H_0) \psi d\psi$, which yields simply the convolution $P_2 * P$ denoted $P_2'(\theta_1)$. Therefore, we can rewrite $I'_2$ as  
\[ I'_2 = \frac{\kappa^2 P_2'(\theta_1) \mu_t}{\mu_{t+H_0} + \mu_0} \left( 1 - \tau(H_1 + H_0) \right) \]

\[ I'_2 = \frac{\kappa^2 P_2'(\theta_1) \mu_t}{\mu_{t-H_0} - \mu_t} \left( \tau(H_0) - \tau(H_1) \right) \]

4.4. Double scattering

In double scattering, photons are first scattered with an angle $\psi$, then follow a path $L_{1,2}$ before being re-scattered with an angle $\theta_{1,2} - \psi$ toward the viewer along a path $L_{2,exit}$: (see Fig. 7, right)

\[ I'_3 = \int_0^H \int_0^{\bar{V}} \int_0^{\bar{V}} P(\psi) s(L_{1,2}) P(\theta_1 - \psi) \tau(L_{2,exit}) \frac{d\psi d\theta_1 dH}{\mu_t} \]

To account for the effects of ground-clouds interaction, we propose an efficient GPU-compliant radiosity algorithm between the ground and the clouds plane. The following sections describe our approach.

5.1. Light and reflectance data

We consider that the clouds bottom and the ground are two parallel planes at distance $h$. The input parameters are:

- Sun and sky colors (i.e., R,G,B intensities), and sun location. Sky color might be defined by a sky shader (view and sun dependent).
- The ground texture, and a coefficient $I_2$ estimating its original illumination in order to interpret the texture as a color reflectance field $\rho_C(u,v)$.
- The texture of cloud bottom diffuse reflectance $\rho_C(u,v)$. It is obtained from the cloud height field using the diffuse reflectance shader $R_{diff}$ described in the previous section.
5.3 Integration of light over the hemisphere

For each point \( p = (u_0, v_0) \) of the receiver (i.e., each pixel of the buffer), we have to sum the emitted radiosity of all the pixels \((u, v)\) of the other buffer \(dE\) weighted by the pixel form factor \(F(u_0, v_0, u, v)\). Of course, this 4-fold loop is too costly. Fortunately, integrating textures is a common problem for which numerous solutions are available, such as MIP-mapping.

In the spirit of [SP89], we consider areas of similar contribution on which we assume the weight is constant. We improve their method by choosing 3 square rings \(S_1, S_2, S_3\) of exterior size \(h, 2h, 3h\) centered on \(p\). The form factors for these areas are \(F(p, S_1) = 0.5, F(p, S_2) = 0.3, F(p, S_3) = 0.18\). In fact we will use weights \(f_1 = 0.5, f_2 = 0.3, f_3 = 0.2\) to avoid loosing energy. Let \(A_1, A_2, A_3\) be squares of size \(h, 2h, 3h\). We have \(S_1 = A_1, S_2 = A_2 - A_1, S_3 = A_3 - A_2\). The average \(dE(A_1), dE(A_2), dE(A_3)\) of the texture \(dE\) over the squares is obtained in one texture fetch \(dE(u_0, v_0)\) using the appropriate level \(i\) in the MIP-map pyramid. Note that using buffers with pixels of size \(h\) is sufficient. Then, the levels 0.1 and 3 will be used. The total integral is thus \(E = f_1dE(A_1) + f_2(\frac{2}{3}dE(A_2) - \frac{1}{3}dE(A_1)) + f_3(\frac{4}{3}dE(A_3) - \frac{2}{3}dE(A_2))\), i.e., \(E = v_1dE_0(u_0, v_0) + v_2dE_1(u_0, v_0) + v_3dE_3(u_0, v_0)\) with \(v_1 = 0.4, v_2 = 0.387, v_3 = .213\).

6. Implementation and Results

Figure 9 validates our analytical reflectance model on two different scenes. Note how the combination of our 3 modes closely approximates the Monte-Carlo simulation. The error mostly comes from assuming that the \(3^+\) scattering is isotropic and from the double scattering approximation. Figure 8 validates our 1D reflectance/transmittance model (the fitting of our parameters is near perfect). We validated our radiosity algorithm against an analytical case (square cloud).

For the cloud field, we rely on 512 x 512 animated textures precomputed on the CPU using advected textures [Ney03]. We display the cloud surface as a height field, \(R\) and \(dR\) are updated simultaneously using multiple render targets. MIP-mapped pyramids are automatically generated by the GPU as well.

Note that all these textures are low-resolution, and only a few iterations are necessary: at the \(i^{th}\) iteration the amount of energy exchanged is roughly \(p^i\), where \(p\) is the average reflectance of the ground (assuming the cloud albedo is 1). Typical country grounds have reflectance ranging from 0.1 to 0.2, so 2 iterations are often enough. Even with a snow landscape requiring many inter-reflection passes, the radiosity computation remains fast enough for interactive cloud animation (we used 8 for Figure 11.E).

The text contains details on the implementation and results of a method for simulating light interactions with stratiform clouds. It includes the integration of light over the hemisphere, with a focus on how to efficiently compute the radiosity by integrating across texture rings. The method uses a MIP-mapped pyramid to handle textures of different resolutions, and it validates the results using both analytical and numerical simulations.
(i.e., a geometric patch) that is build on the fly on the CPU from the texture. The clouds height vary from 0 to 500m and they span several km. When the observer is on the ground and see the clouds from below, we display the cloud bottom using a textured quad. When viewed from below, the shader used at each pixel of this mesh renders

$$I' C_{sun} + T_{ns} C_{sky} + R_{ns} C_{ground}$$

When viewed from above, the shader renders

$$I' C_{sun} + R_{ns} C_{sky} + T_{ns} C_{ground}$$

$I'$ sums $f_i'$ (section 4.2), $f_1'$ (section 4.3), $f_2'$ (section 4.4) and $f_3'$ (section 4.5). The latter depends on $T_{R}$, which uses $T_{ns}$, $T_0$, $T_1$ and $T_2$ (appendix A). The same dependencies apply for the calculation of $I'$. The radiosity algorithm makes use of $R_{ns}$, $T_0$, $T_1$, $T_2$, $T_3$, and $T_{ns}$ (appendix A). The parameters sent to the shadors in order to compute these values are:

- $\nabla$, $\mathcal{T}$, and $\mathcal{N}$. Since single and double scattering are local events, we use the local norm of the surface as $\mathcal{N}$ for $I_1$ and $I_2$. Because multiple scattering is a global process, we use $\mathcal{N} = \mathcal{Z}$ for $I_3$.

- The texture $h(u,v)$ used to estimate $H$, plus scaling factors in the 3 directions.

- Sun, Sky colors $C_{sun}$, $C_{sky}$ and ground illumination texture $C_{ground}$.

- Phase functions $P_3(\theta)$, $P_F(\theta)$ and $P^2(\theta)$ as textures.

- The whole set of $\theta_i$-dependent parameters needed for $T_i$ and $R_i$, passed as a 2D texture (we used the ones listed in appendix B). At each frame we construct and send the cloud scene to the GPU, we run on GPU the radiosity simulation described in section 5, we render the background (sky or ground), then the cloud layer using our color+transparency shaders.

The implementation is done on a PC with an NVIDIA Quadro FX 1400. At resolution 1024 x 768, an image in which 50% of the pixels are covered by clouds (intermediates between the two images of Fig. 11.C) runs at 18 FPS. One major bottleneck during clouds animation is the update of the 512x512 height field geometry. When clouds are not animated, the frame rate increases to 40 FPS.

Our results are illustrated in Figure 11 and in the companion video. They show the various expected features including “pseudo-specular” reflection (Fig. 11.C), “silver lining” (Fig. 11.A, 11.D), fogbow (fig. 10, left) and glory (fig. 10, right) , and the radiosity effects including water sky and ice blink(Fig. 11.A,right, 11.E). See figure 5 for comparison with real photographs.

Figure 9: Comparison of our backward multiple scattering approximated model against Monte-Carlo photon tracing. Left: $H = 100 m$, $\theta_i = 70^\circ$. Right: $H = 5 m$, $\theta_i = 80^\circ$. Areas: Contribution of each order of scattering. Blue line: Monte-Carlo result.

Figure 10: Close-up views on reproduced cloud features. Left: fogbow. Right: glory (6-fold zoom of the left image).

7. Conclusion and future work

We presented a method to render cloud layers in real-time providing all expected optical features, as well as inter-reflections with the ground, which was not feasible with existing methods. However, our model is restricted to stratiform clouds. We would like to carry on this approach to all kinds of animated cloudy skies (e.g., cumulonimbus storms or congestus fields), and from any point of view. This method should be easily transposable to ice-content stratiform clouds, where only the phase function differs.

We assumed that the rays always pass through the cloud bottom. This neglect interesting behaviors at the cloud sides. We plan to explore better estimation for these configurations.

Relying on geometry for clouds is very efficient as compared to the massive overwrite caused by methods with sliced volumes. Still, it lacks the blurry effect on the silhouette. We aim at adapting Gardner’s technique [Gar85] for that.

The accuracy of our model could be further improved: the pseudo-specular part is slightly underestimated (see fig. 9) and yields overdarkened parts on top of the cloud. This is because the diffuse mode $I_3$ does not fully match the real behavior well in thin areas, where it keeps some anisotropy.

Radiative transfer can be improved as well. We did not consider the inter-reflections between different parts of the cloud (which usually brightens creases), and between different clouds. Aerial perspective should be taken into account for distant interactions. Also, the interaction with a non-flat ground has yet to be treated.

Still, we were able to achieve real-time rendering on detailed scenes while reproducing the main features of clouds aspect. Such complex lighting effects where not possible with previous methods, and we believe our approach opens promising avenues of future work toward realistic real-time cloud rendering.

8. Acknowledgments

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A. **Top:** Final rendering of clouds top, $I_1^*$ contribution, sky contribution, $I_0 + I_2$ contribution, ground contribution. **Bottom:** Final rendering of clouds bottom, $I_1^*$, sky, $I_0 + I_1 + I_2$ and ground contributions. (see Figure 2)

B. View from under the clouds without (top) and with radiosity (bottom).

C. **Left:** Clouds in full day-light. **Right:** Clouds at sun-set. Notice the strong contribution of the scatter terms (pseudo-specular).

D. **Left:** Sun through the clouds, seen from the ground. **Right:** Clouds in full day light, seen from the ground.

E. **Left and right:** Two instances of the water sky / ice blink effect. Clouds located just above the water appear darker than above ice or sand.

**Figure 11:** Results obtained at 18 to 40 fps with $512 \times 512$ clouds animated heightfields and 3 to 10 pairs of radiosity passes.

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References

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[Dei04] DEIRMENDJIAN D.: Scattering and polarization properties of water clouds and hazes in the visible and infrared. Applied Optics, 34, 3 (Feb. 1964), 4–7. 2


Appendix A: 1D models

Figure 12: Left: 1D multiple scattering model, where X is the incoming intensity. Right: 1D single scattering model, 1D double scattering model.

Multiple scattering Let us assume we know the transmittance $T_{\text{trans}}(H)$ of a slab of height $H$ through multiple scattering. The reflectance is $R_{\text{trans}}(H) = 1 - T_{\text{trans}}(H)$. A 2H slab can be seen as two slabs of height $H$ stacked together. If $x$ is the amount of light traversing the interface between the two layers in the transmissive direction, we have $x = T_{\text{trans}}(H) + xR_{\text{trans}}(H)$ whose solution is $x = \frac{1}{T_{\text{trans}}(H)}$ (see Fig. 12, left). The amount of light transmitted by the 2H layer is what is left from $x$ after traversing the second layer. Thus $T_{\text{trans}}(2H) = T_{\text{trans}}(H)$. By induction, we get $T_{\text{trans}}(H) = \frac{\beta}{\beta - H}$ with $\beta = T_{\text{trans}}(1m)$. We find $\beta$ by fitting against a Monte Carlo simulation (see appendix B).

But isotropic approximation is valid for a large thickness $H$ only. Below a certain critical thickness (several hundred meters in our case), $T_{\text{trans}}(H)$ does not fit the real behavior which becomes strongly dependent on $\theta$. We found experimentally that this error can be well corrected by an exponential factor. We use $T_{\text{trans}}(H) = (b + (1 - b)e^{-cH}) \frac{\beta}{\beta - H}$ where $b$ and $c$ are correction parameters depending on $\theta$ that we fit against a Monte Carlo simulation (see appendix B and Fig. 8).

Single and double scattering For the single and double scattering approximation, we consider a 1D column holding particles with an extinction function $\tau_e(x) = e^{-kx}$ and...
\[ s_c(x) = \kappa_c e^{-\kappa_c x} \] (see Fig. 12, right). In the 1D case, light can only be scattered in 2 directions: reflective (coefficient \( r \)) or transmissive (coefficient \( t \)). With this model, the 1D single scattering solution is

\begin{align*}
T_1(H) &= \int_0^H s_c(l)\alpha_c \tau_c(H-l)dl = tHs_c(H) \\
R_1(H) &= \int_0^H s_c(l)\alpha_c \tau_c(l)dl = \frac{1}{2}(1 - \tau_c(2H))
\end{align*}

Transmittance from double scattering can happen either by double-backward or double-forward scattering.

\begin{align*}
T_2(H) &= T_{2a}(H) + T_{2b}(H) \\
T_{2a}(H) &= \int_0^H s_c(l)\tau_c(H-k)dk \int_0^H s_c(k)\tau_c(H-k)dl \\
&= \frac{1}{2} (\kappa_c tH)^2 \tau_c(H) \\
T_{2b}(H) &= \int_0^H s_c(k)r \int_0^H s_c(k-l)\tau_c(H-l)dl dk \\
&= \frac{1}{4} r^2 (1 + (2H\kappa_c - 1)\tau_c(2H)\tau_c(3H))
\end{align*}

Reflectance from double scattering can happen when being scattered backward then forward, or vice-versa.

\begin{align*}
R_2(H) &= 2 \int_0^H s_c(l)\tau_c(H-l)dl \int_0^H s_c(l-k)\tau_c(k)dk \\
&= \frac{1}{2} r t (1 - (2H\kappa_c + 1)\tau_c(2H))
\end{align*}

Here again, the \( \theta_l \)-dependent parameters \( \kappa_c, t \) and \( r \) are fit against Monte Carlo simulation (see appendix B and Fig. 8).

As for \( b \) and \( c \), these parameters account for the bias resulting from modeling a 3D phenomenon with a 1D model. That is why they vary with \( \theta_l \), and why \( r \neq 1 - t \).

Transparency is directly given by the extinction function:

\[ T_0(H) = \tau(H) \]

### Appendix B: Used values

Here is an example set of parameters used in appendix A for the DSD used in the examples, i.e., \( N_0 = 300 \text{cm}^{-3} \), \( \gamma = 2 \) and \( r_e = 7 \mu \text{m} \). We obtained them by running several photon tracing simulations at relevant values of \( H \) and \( \theta_l \) and fitting the parameters by optimization. The fitting is very good, as displayed on figure 8.

\[ \beta = 0.9961 \] (constant).

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