



**HAL**  
open science

## On conditional McKean Lagrangian stochastic models

Mireille Bossy, Jean-Francois Jabir, Denis Talay

► **To cite this version:**

Mireille Bossy, Jean-Francois Jabir, Denis Talay. On conditional McKean Lagrangian stochastic models. [Research Report] RR-6761, 2008, pp.35. inria-00345524v3

**HAL Id: inria-00345524**

**<https://inria.hal.science/inria-00345524v3>**

Submitted on 1 Mar 2009 (v3), last revised 10 Jul 2009 (v4)

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



INSTITUT NATIONAL DE RECHERCHE EN INFORMATIQUE ET EN AUTOMATIQUE

*On conditional McKean Lagrangian stochastic  
models*

Mireille Bossy — Jean-François Jabir — Denis Talay

N° 6761

December 2008

Thème NUM



*Rapport  
de recherche*





# On conditional McKean Lagrangian stochastic models

Mireille Bossy , Jean-François Jabir , Denis Talay

Thème NUM — Systèmes numériques  
Projets Tosca

Rapport de recherche n° 6761 — December 2008 — 36 pages

**Abstract:** This work is devoted to the well-posedness of a particle position-velocity system which is nonlinear in the sense of McKean. As the dynamics of the velocity depends on the conditional expectation w.r.t. its position, we have to deal with a singular interaction kernel between particles. This study is motivated by a new class of SDEs-PDEs systems, the so called Lagrangian stochastic models which are commonly used in the simulation of turbulent flows.

After a short presentation of these systems, we prove existence and uniqueness results for a simplified model, and a propagation of chaos result for the corresponding interacting particle system.

**Key-words:** Lagrangian stochastic model, propagation of chaos, conditional McKean nonlinearity

# Sur les modèles lagrangiens stochastiques McKean conditionnels

**Résumé :** Dans ce travail, nous nous étudions le caractère bien-posé d'un système de particules en (position, vitesse) nonlinéaire au sens de McKean. Dans ce type de système, l'équation de la vitesse fait apparaître une espérance conditionnelle par rapport à sa position courante. Ainsi, la nonlinéarité au sens de McKean est singulière.

Cette étude est motivée par une classe de modèles, appelés modèles lagrangiens stochastiques, utilisés pour la simulation d'écoulements turbulents en mécanique des fluides.

Après une courte présentation de ces modèles, nous montrons un résultat d'existence et d'unicité pour un modèle lagrangien simple, ainsi qu'un résultat de propagation du chaos pour le système de particules associé.

**Mots-clés :** modèle stochastique lagrangien, propagation du chaos, non-linéarité conditionnelle au sens de McKean

## 1 Introduction

In this paper, we prove the well-posedness, and construct an interacting particle approximation for the simplified Lagrangian stochastic model describing the time evolution of the position and velocity of a fluid-particle: for finite horizon time  $T > 0$ , we consider a  $d$ -dimensional standard Brownian motion  $(W_t; t \in [0, T])$ , and a  $\mathbb{R}^{2d}$ -valued r.v.  $(X_0, \mathcal{U}_0)$  independent of  $W$ , and  $((X_t, \mathcal{U}_t); t \in [0, T])$  satisfying

$$\begin{cases} X_t = X_0 + \int_0^t \mathcal{U}_s ds, \\ \mathcal{U}_t = \mathcal{U}_0 + \int_0^t B[X_s, \mathcal{U}_s; \rho_s] ds + \int_0^t \sigma(s, X_s, \mathcal{U}_s) dW_s, \\ \rho_t \text{ is the density distribution of } (X_t, \mathcal{U}_t) \text{ for all } t \in (0, T]. \end{cases} \quad (1.1)$$

Here,  $B$  is the mapping from  $\mathbb{R}^d \times \mathbb{R}^d \times L^1(\mathbb{R}^{2d})$  to  $\mathbb{R}^d$  defined by

$$B[x, u; \gamma] = \begin{cases} \frac{\int_{\mathbb{R}^d} b(v, u) \gamma(x, v) dv}{\int_{\mathbb{R}^d} \gamma(x, v) dv} & \text{if } \int_{\mathbb{R}^d} \gamma(x, v) dv \neq 0, \\ 0 & \text{elsewhere,} \end{cases} \quad (1.2)$$

where  $b : \mathbb{R}^d \times \mathbb{R}^d \rightarrow \mathbb{R}^d$  is a bounded interaction kernel. Formally, the drift component of (1.1) involves the function

$$(x, u) \mapsto \mathbb{E} \left[ b(\mathcal{U}_t, u) \middle/ X_t = x \right]. \quad (1.3)$$

Such nonlinearity is typical of Lagrangian stochastic models. A Lagrangian stochastic model is a system of SDEs describing characteristics of a generic fluid-particle in a turbulent flow. These characteristics include the position  $X_t$  and the velocity  $\mathcal{U}_t$  of a particle.

The present work is strongly motivated by the Lagrangian stochastic approach for the simulation of turbulent flows and the related Probability Density Function (PDF) methods for turbulent flows. The Lagrangian stochastic models for turbulent flows have a dramatic complexity (see description, Section 2). To our knowledge, the mathematical study of these models is an open problem. Since the model aims to be an alternative approach to the Navier-Stokes equations for turbulent flows, it is not surprising that its study presents difficulties. Several recent works separately face some of the difficulties. For example,

Bossy & Fontbona [2] study the Poisson PDE (2.2) and its relation with the incompressibility of the mean field velocity; Bossy and Jabir [3] study (1.1) with boundary condition. In this paper, we initiate the theoretical justification of these Lagrangian models. We have started their theoretical study by showing that the simplified model (1.1) is well-posed. Part of our analysis consists in proving a propagation of chaos result for the solution of a particle system approaching the solution of (1.1) which is a first step to prove the convergence of stochastic particle methods. Although simple, the model (1.1) inherits two important features of the Lagrangian stochastic model. First, the infinitesimal generator of the process is non-elliptic due to the Langevin dynamic. Second, the drift coefficient of the velocity involve a conditional expectation w.r.t. the particle position. Consequently, Equation (1.3) is a degenerate and singular nonlinear equation in the sense of McKean.

We prove that (1.1) has a unique solution constructed as the limit of solutions of smoothed systems (see Theorem 3.2). The existence of the solutions of the smoothed systems is obtained by means of the propagation of chaos of interacting particle systems (see Section 3). The uniqueness results are proven in Section 4. In the next section, we briefly present the Lagrangian stochastic models in fluid dynamics, and list some references on the models and their numerical issues.

## 2 A brief description of Lagrangian stochastic models for turbulent flows

We start this section with a short reminder on the notion of statistical solutions to the Navier–Stokes equation for turbulent flows. For the sake of simplicity, we limit our presentation to monophasic flows. These statistical solutions are random fields, the velocity and the pressure, which are decomposed into their averaged and fluctuating parts. The so called Reynolds decomposition of the Eulerian velocity  $U$  is

$$U(t, x, \omega) = \langle U \rangle(t, x) + \mathbf{u}(t, x, \omega),$$

where  $\langle U \rangle$  is the (ensemble) averaged part, and  $\mathbf{u}$  is the fluctuating part. The Reynolds average  $\langle \cdot \rangle$  is a linear operator applied to the random fields, which

is assumed to commute with spatial and times derivatives. Formally applying the Reynolds average to the Navier–Stokes equation, one obtains the so called Reynolds Averaged Navier–Stokes (RANS) equations:

$$\begin{cases} \nabla_x \cdot \langle U \rangle = 0, \\ \partial_t \langle U^{(i)} \rangle + \langle U \rangle \cdot \nabla_x \langle U^{(i)} \rangle = -\frac{1}{\varrho} \nabla_x \langle \mathcal{P}^{(i)} \rangle + \nu \Delta_x \langle U^{(i)} \rangle - \partial_{x_j} \langle \mathbf{u}^{(i)} \mathbf{u}^{(j)} \rangle, \\ \langle U \rangle(0, x) = \langle U_0 \rangle(x). \end{cases} \quad (2.1)$$

The averaged pressure  $\langle \mathcal{P} \rangle$  solves the Poisson equation

$$-\frac{1}{\varrho} \Delta_x \langle \mathcal{P} \rangle = \partial_{x_i x_j}^2 \langle U^{(i)} \rangle \langle U^{(j)} \rangle + \partial_{x_i x_j}^2 \langle \mathbf{u}^{(i)} \mathbf{u}^{(j)} \rangle. \quad (2.2)$$

The Reynolds stress tensor stands for the covariance of velocity components:

$$\langle \mathbf{u}^{(i)} \mathbf{u}^{(j)} \rangle = \langle U^{(i)} U^{(j)} \rangle - \langle U^{(i)} \rangle \langle U^{(j)} \rangle.$$

These terms are not closed in Equation (2.1). This problem has led to the introduction of closure models based on Kolmogorov’s theory for turbulent flows and experimental observations. For example, the  $k - \mathcal{E}$  closures consist in a set of closed equations for the turbulent kinetic  $k$  and the dissipation rate  $\mathcal{E}$  defined as

$$\begin{aligned} k(t, x) &= \frac{1}{2} \langle \mathbf{u}^{(i)} \mathbf{u}^{(i)} \rangle(t, x), \\ \mathcal{E}(t, x) &= \nu \langle \partial_{x_j} \mathbf{u}^{(i)} \partial_{x_j} \mathbf{u}^{(i)} \rangle(t, x), \end{aligned}$$

(see, e.g., Mohammadi & Pironneau [10], Pope [13]).

Lagrangian stochastic models have been successfully proposed to provide an alternative approach to the numerical resolution of RANS equations combined with closure models to simulate complex flows for which PDE solvers are inefficient.

In a series of papers initiated in 1985, Stephen B. Pope has proposed Lagrangian stochastic models to describe the position and the instantaneous velocity  $(X_t, \mathcal{U}_t)$  of a fluid–particle. Depending on the flow, other Lagrangian characteristics of the turbulence are added to the model. For a fluid with



constant mass density  $\varrho$ , Lagrangian and Eulerian quantities are related as follows: for all suitable measurable function  $g : \mathbb{R}^d \rightarrow \mathbb{R}^d$ ,

$$\langle g(U) \rangle(t, x) = \mathbb{E}[g(\mathcal{U}_t) / X_t = x].$$

Assuming that  $(X, \mathcal{U})$  is a diffusion process, the coefficients of its generator are designed such that the Lagrangian laws are consistent with closed RANS equations and other relevant Physic laws in turbulence theory (see Pope [12], [13] for details). This methodology is known as *PDF method for turbulent flows* in the literature. The simplest model proposed by Pope is the so called simplified Langevin model (see [13]):

$$\begin{cases} X_t = X_0 + \int_0^t \mathcal{U}_s ds, \\ \mathcal{U}_t = \mathcal{U}_0 - \frac{1}{\varrho} \int_0^t \nabla_x \langle \mathcal{P} \rangle(s, X_s) ds + \nu \int_0^t \Delta_x \langle U \rangle(s, X_s) ds \\ \quad + C_1 \int_0^t \frac{\mathcal{E}(s, X_s)}{k(s, X_s)} (\langle U \rangle(s, X_s) - \mathcal{U}_s) ds + \int_0^t \sqrt{C_2 \mathcal{E}(s, X_s)} dW_s. \end{cases}$$

Here,  $C_1$  and  $C_2$  are positive constants, and  $(W_t; t \geq 0)$  is a standard  $\mathbb{R}^3$ -valued Brownian motion. The Poisson equation (2.2) provides the averaged pressure  $\langle \mathcal{P} \rangle(t, x)$ , and  $k$  and  $\mathcal{E}$  are assumed to be known.

A less elementary model was proposed by Dreeben & Pope [7] where

$$k(t, x) = \mathbb{E} \left( (\mathcal{U}_t^{(i)})^2 / X_t = x \right) - \left( \mathbb{E} \left( \mathcal{U}_t^{(i)} / X_t = x \right) \right)^2,$$

and  $\mathcal{E}$  is defined as  $\mathcal{E}(t, x) = \langle \omega \rangle(t, x)k(t, x)$  where  $\langle \omega \rangle(t, x) = \mathbb{E}(\omega_t / X_t = x)$  and  $(\omega_t; t \geq 0)$  is the solution of the following SDE:

$$\begin{aligned} \omega_t = \omega_0 + C_3 \int_0^t \langle \omega \rangle(s, X_s) (\langle \omega \rangle(s, X_s) - \omega_s) ds \\ - \int_0^t \omega_s S_\omega(s, X_s) \langle \omega \rangle(s, X_s) ds + \int_0^t \sqrt{C_4 \omega_s (\langle \omega \rangle(s, X_s))^2} d\widetilde{W}_s. \end{aligned}$$

Here,  $C_3, C_4$  are positive constants,  $\widetilde{W}$  is a one-dimensional standard Brownian motion independent of  $W$ , and

$$S_\omega(t, x) = C_{\omega 2} + C_{\omega 1} \frac{(\langle U^{(i)} U^{(j)} \rangle(t, x) - \langle U^{(i)} \rangle(t, x) \langle U^{(j)} \rangle(t, x)) \partial_{x_j} \langle U^{(i)} \rangle(t, x)}{\langle \omega \rangle(t, x) k(t, x)}.$$

A description of the numerical issues can be found in Pope & Xu [14]. A recent application to meteorology is developed by one of the authors: see, e.g., Bossy *et al.* [1].

**Notation.** Let  $0 < T < +\infty$  be fixed.

- For all  $t \in (0, T]$ , we set  $Q_t = (0, t) \times \mathbb{R}^{2d}$ .
- For all  $q \geq 1$ ,  $\mathcal{C}([0, T]; \mathbb{R}^q)$  denotes the space of  $\mathbb{R}^q$ -valued continuous functions equipped with the compact uniform metric  $\|\cdot\|_\infty$ .
- Given a metric space  $E$ ,  $\mathcal{C}_b^k(E)$  denotes the set of real-valued bounded functions defined on  $E$  with continuous derivatives up to order  $k$ ;  $\mathcal{C}_c^k(E)$  denotes the set of real-valued functions with continuous derivatives up to order  $k$  and with compact support.
- Given a metric space  $E$ ,  $\mathcal{M}(E)$  denotes the set of probability measures defined on  $E$ , equipped with the weak convergence topology.
- In all the paper,  $C$  is a constant which does not depend on the various parameters  $\epsilon, N, \dots$ , but does vary from line to line.

### 3 Main result

In the study of (1.1), difficulties come from the dependency of the drift coefficient on the conditional expectation (1.3). Related situations have been studied by Sznitman [17], Oelschläger [11] and Dermoune [4]: Sznitman [17] has considered the one-dimensional nonlinear SDE

$$d\zeta_t = p_t(\zeta_t) dt + dW_t,$$

where  $p_t$  is the Lebesgue density of  $\zeta_t$ . Oelschläger [9] has considered the family of models

$$d\zeta_t = F(\zeta_t, p_t(\zeta_t)) dt + dW_t,$$

where  $F : \mathbb{R}^d \times \mathbb{R}$  is a bounded Lipschitz function, and

$$d\zeta_t = \nabla p_t(\zeta_t) dt + dW_t.$$

Dermoune [4] has studied the system

$$d\zeta_t = \mathbb{E}(v(\zeta_0)/\zeta_t) dt + dW_t,$$

where  $v : \mathbb{R}^d \rightarrow \mathbb{R}^d$  is a bounded continuous function. Our situation substantially differs from the above: our drift coefficient depends on conditional density rather than the density and the infinitesimal generator of  $(X_t, \mathcal{U})$  is not strongly elliptic.

In the sequel, we suppose the following hypotheses:

- (H) •  $b$  is a bounded continuous function and the law  $\mu_0$  of  $(X_0, \mathcal{U}_0)$  is such that

$$\int_{\mathbb{R}^{2d}} (|x| + |u|^2) \mu_0(dx du) < +\infty.$$

- The velocity diffusion coefficient  $\sigma$  is bounded and strongly elliptic: for  $a(t, x, u) := \sigma(t, x, u)\sigma^*(t, x, u)$ , there exists  $\lambda > 0$  such that, for all  $t \in (0, T]$ ,  $x, u, v \in \mathbb{R}^d$ ,

$$\frac{|v|^2}{\lambda} \leq \sum_{i,j=1}^d a^{(i,j)}(t, x, u) v_i v_j \leq \lambda |v|^2. \quad (3.1)$$

- For all  $1 \leq i, j \leq d$ ,  $a^{(i,j)}$  is Hölder continuous in the following sense: there exist  $\alpha \in (0, 1]$  and  $K$  depending only on  $T$  and  $d$  such that, for all  $(s, x, u), (t, y, v) \in (0, T] \times \mathbb{R}^d \times \mathbb{R}^d$ ,

$$|a^{(i,j)}(t, x, u) - a^{(i,j)}(s, y, v)| \leq C(|t-s|^{\frac{\alpha}{2}} + |x-y-v(t-s)|^\alpha + |u-v|^\alpha). \quad (3.2)$$

**Remark 3.1.** *The hypothesis (3.2) on  $a$  is classical in the literature on ultra-parabolic PDEs, see e.g. Theorem A.1 and Subsection A.2. In particular, it is satisfied when  $\sigma(t, x, u)$  is a function of  $(t, u)$  only.*

For fixed  $N \geq 1$  and  $\epsilon > 0$ , we consider the interacting particle system  $\{(X_t^{i,\epsilon,N}, \mathcal{U}_t^{i,\epsilon,N}; t \in [0, T]); 1 \leq i \leq N\}$  as defined by

$$\begin{cases} X_t^{i,\epsilon,N} &= X_0^i + \int_0^t \mathcal{U}_s^{i,\epsilon,N} ds, \\ \mathcal{U}_t^{i,\epsilon,N} &= \mathcal{U}_0^i + \int_0^t \frac{\sum_{j=1}^N b(\mathcal{U}_s^{j,\epsilon,N}, \mathcal{U}_s^{i,\epsilon,N}) \phi_\epsilon(X_s^{i,\epsilon,N} - X_s^{j,\epsilon,N})}{\sum_{j=1}^N (\phi_\epsilon(X_s^{i,\epsilon,N} - X_s^{j,\epsilon,N}) + \epsilon)} ds \\ &\quad + \int_0^t \sigma(s, X_s^{i,\epsilon,N}, \mathcal{U}_s^{i,\epsilon,N}) dW_s^i, \quad i = 1, \dots, N. \end{cases} \quad (3.3)$$

Here,  $\{(X_0^i, \mathcal{U}_0^i), (W_t^i; t \in [0, T]); i \geq 1\}$  are independent copies of  $((X_0, \mathcal{U}_0), (W_t; t \in [0, T]))$ , and  $\{\phi_\epsilon; \epsilon > 0\}$  denotes a family of mollifiers of the type  $\phi_\epsilon(x) = \frac{1}{\epsilon^d} \phi(\frac{x}{\epsilon})$ , where  $\phi \in \mathcal{C}_c^1(\mathbb{R}^d)$  is such that  $\phi \geq 0$  and  $\int_{\mathbb{R}^d} \phi(z) dz = 1$ . As the drift coefficient of the particle system (3.3) is uniformly bounded, the well-posedness of (3.3) follows from Proposition 4.4 (see Section 4.1) and Girsanov's theorem.

In Section 5, we prove that the particles propagate chaos. In particular, as  $N$  tends to infinity,  $(X^{1,\epsilon,N}, \mathcal{U}^{1,\epsilon,N})$  converges weakly to the solution of

$$\begin{cases} X_t^\epsilon = X_0 + \int_0^t \mathcal{U}_s^\epsilon ds, \\ \mathcal{U}_t^\epsilon = \mathcal{U}_0 + \int_0^t B_\epsilon [X_s^\epsilon, \mathcal{U}_s^\epsilon; \rho_s^\epsilon] ds + \int_0^t \sigma(s, X_s^\epsilon, \mathcal{U}_s^\epsilon) dW_s, \\ \rho_t^\epsilon \text{ is the density of } (X_t^\epsilon, \mathcal{U}_t^\epsilon) \text{ for all } t \in (0, T], \end{cases} \quad (3.4)$$

where the kernel  $B_\epsilon [x, u; \gamma]$  is defined as follows: for all nonnegative  $\gamma \in L^1(\mathbb{R}^{2d})$ ,  $(x, u) \in \mathbb{R}^{2d}$ ,

$$B_\epsilon [x, u; \gamma] = \frac{\int_{\mathbb{R}^d} b(v, u) \phi_\epsilon \star \gamma(x, v) dv}{\int_{\mathbb{R}^d} \phi_\epsilon \star \gamma(x, v) dv + \epsilon},$$

where

$$\phi_\epsilon \star \gamma(x, u) = \int_{\mathbb{R}^d} \phi_\epsilon(x - y) \gamma(y, u) dy.$$

Our main result is as follows.

**Theorem 3.2.** *Assume (H).*

- (i) *For all  $\epsilon > 0$ , the sequence  $\{(X^{1,\epsilon,N}, \mathcal{U}^{1,\epsilon,N}); N \geq 1\}$  converges weakly to a weak solution  $(X^\epsilon, \mathcal{U}^\epsilon)$  of (3.4). This solution is unique and, if  $\mathbb{P}^\epsilon$  denotes the law of  $(X^\epsilon, \mathcal{U}^\epsilon)$ , the interacting particle system is  $\mathbb{P}^\epsilon$ -chaotic; that is, for every integer  $k \geq 2$  and every finite family  $\{\psi_l; l = 1, \dots, k\}$  of  $\mathcal{C}_b(\mathcal{C}([0, T]; \mathbb{R}^{2d}))$ ,*

$$\langle \mathbb{P}^{\epsilon,N}, \psi_1 \otimes \dots \otimes \psi_k \otimes \dots \rangle \longrightarrow \prod_{l=1}^k \langle \mathbb{P}^\epsilon, \psi_l \rangle, \text{ when } N \longrightarrow +\infty.$$

(ii) When  $\epsilon$  tends to 0,  $(X^\epsilon, \mathcal{U}^\epsilon)$  converges weakly to the unique solution  $(X, \mathcal{U})$  of (1.1).

Weak solutions of (3.4) and (1.1) are defined by appropriate martingale problems in the next section (see Definition 4.1 and Definition 4.2).

The rest of the paper is organized as follows: in Section 4, we prove weak uniqueness results for Equations (3.4) and (1.1). In Section 5, we get existence: we show that, for all  $\epsilon > 0$ , the law of the particle system (3.3) converges weakly and we identify the limit as the weak solution of (3.4); we then get existence of a weak solution of (1.1) by letting  $\epsilon$  decreases to 0.

## 4 Uniqueness results

We introduce the notions of weak solutions to (1.1) and (3.4). Let  $((x_t, u_t); t \in (0, T])$  be the canonical processes in the sample space  $\mathcal{C}([0, T]; \mathbb{R}^{2d})$ . The martingale problem related to the smoothed particle system (3.4) is stated as follows.

**Definition 4.1.** *A probability measure  $\mathbf{P}^\epsilon$  on the canonical space  $\mathcal{C}([0, T]; \mathbb{R}^{2d})$  is a weak solution of (3.4), or equivalently, a solution to the martingale problem  $(MP_\epsilon)$  if:*

- (i)  $\mathbf{P}^\epsilon \circ (x_0, u_0)^{-1} = \mu_0$ .
- (ii) For all  $t \in (0, T]$ , the time marginal  $\mathbf{P}^\epsilon \circ (x_t, u_t)^{-1}$  has a density  $\rho_t^\epsilon$  w.r.t. Lebesgue measure on  $\mathbb{R}^{2d}$ .
- (iii) For all  $f \in \mathcal{C}_b^2(\mathbb{R}^{2d})$ , the process

$$f(x_t, u_t) - f(x_0, u_0) - \int_0^t \mathcal{A}_{\rho_s^\epsilon}^\epsilon f(s, x_s, u_s) ds$$

is a  $\mathbf{P}^\epsilon$ -martingale where, for all  $\gamma \in L^1(\mathbb{R}^{2d})$ ,  $\mathcal{A}_\gamma^\epsilon$  is the differential operator

$$\begin{aligned} \mathcal{A}_\gamma^\epsilon f(t, x, u) := & u \cdot \nabla_x f(x, u) + B_\epsilon[x, u; \gamma] \cdot \nabla_u f(x, u) \\ & + \frac{1}{2} \sum_{i,j=1}^d a^{(i,j)}(t, x, u) \partial_{u_i, u_j}^2 f(x, u). \end{aligned}$$

The martingale problem related to (1.1) is stated as follows.

**Definition 4.2.** *A probability measure  $\mathbf{P}$  on the canonical space  $\mathcal{C}([0, T]; \mathbb{R}^{2d})$  is a weak solution of (1.1), or equivalently, a solution to the martingale problem (MP) if*

- (i)  $\mathbf{P} \circ (x_0, u_0)^{-1} = \mu_0$ .
- (ii) For all  $t \in (0, T]$ , the time marginal  $\mathbf{P} \circ (x_t, u_t)^{-1}$  has a positive density  $\rho_t$  w.r.t. Lebesgue measure on  $\mathbb{R}^{2d}$ .
- (iii) For all  $f \in \mathcal{C}_b^2(\mathbb{R}^{2d})$ , the process

$$f(x_t, u_t) - f(x_0, u_0) - \int_0^t \mathcal{A}_{\rho_s} f(s, x_s, u_s) ds$$

is a  $\mathbf{P}$ -martingale, where, for each positive  $\gamma \in L^1(\mathbb{R}^{2d})$ ,  $\mathcal{A}_\gamma$  is the differential operator

$$\begin{aligned} \mathcal{A}_\gamma f(t, x, u) := & u \cdot \nabla_x f(x, u) + (B[x, u; \gamma] \cdot \nabla_u f(x, u)) \\ & + \frac{1}{2} \sum_{i,j=1}^d a^{(i,j)}(t, x, u) \partial_{u_i u_j}^2 f(x, u). \end{aligned}$$

We prove the following uniqueness result.

**Proposition 4.3.** *There is at most one weak solution to Equation (1.1) and one weak solution to Equation (3.4).*

For a weak solution  $\mathbf{P}$  of (1.1), and a weak solution  $\mathbf{P}^\epsilon$  of (3.4), we consider the densities  $\rho_t$  and  $\rho_t^\epsilon$  as in Definitions (4.2) and (4.1). We prove that  $\rho_t$  and  $\rho_t^\epsilon$  are the unique solutions of nonlinear mild equations (see Lemma 4.5) which implies Proposition 4.3. A preliminary step consists in studying the linear case ( $b = 0$ ).

## 4.1 Study of a Langevin system

For  $(y, v) \in \mathbb{R}^{2d}$ , consider the pair of processes  $(Y_t^{s,y,v}, V_t^{s,v}; t \geq s \geq 0)$  solution of the Langevin equation

$$\begin{cases} Y_t^{s,y,v} &= y + \int_s^t V_\theta^{s,y,v} d\theta, \\ V_t^{s,y,v} &= v + \int_s^t \sigma(\theta, Y_\theta^{s,y,v}, V_\theta^{s,y,v}) dW_\theta. \end{cases} \quad (4.1)$$

The following result is a slight extension of a theorem due to Di Francesco and Pascucci [5]. We postpone the statement of this theorem and the proof of Proposition 4.4 in Subsection A.2.

**Proposition 4.4.** *Assume (H)–(3.1) and (H)–(3.2). There exists a unique weak solution to (4.1). In addition, this solution admits a density  $\Gamma(s, y, v; t, x, u)$  w.r.t. Lebesgue measure such that:*

- (i) *For all  $(t, x, u) \in (0, T] \times \mathbb{R}^{2d}$ ,  $1 \leq i, j \leq d$ , the derivatives  $\partial_{v_i} \Gamma(s, y, v; t, x, u)$  exist and are continuous for all  $(s, y, v) \in \mathbb{R} \times \mathbb{R}^{2d}$  such that  $(s, y, v) \neq (t, x, u)$ .*
- (ii) *Let  $f : \mathbb{R}^{2d} \rightarrow \mathbb{R}$  be a bounded continuous function. Then the function  $G_{t,f}$  defined by*

$$G_{t,f}(s, y, v) = \int_{\mathbb{R}^{2d}} \Gamma(s, y, v; t, x, u) f(x, u) dx du,$$

*is the unique classical solution of the Cauchy problem*

$$\begin{cases} \partial_s G_{t,f} + \mathcal{L}G_{t,f} = 0 & \text{in } [0, t) \times \mathbb{R}^{2d}, \\ \lim_{s \rightarrow t^-} G_{t,f}(s, y, v) = f(y, v) & \text{in } \mathbb{R}^{2d}, \end{cases} \quad (4.2)$$

*where*

$$\mathcal{L}\psi(s, y, v) := v \cdot \nabla_x \psi(s, y, v) + \frac{1}{2} \sum_{i,j=1}^d a^{(i,j)}(s, y, v) \partial_{v_i, v_j}^2 \psi(s, y, v). \quad (4.3)$$

- (iii) *There exists a constant  $C > 0$  depending only on  $T$  and  $\lambda$  such that*

$$\sup_{(y,v) \in \mathbb{R}^{2d}} \int_{\mathbb{R}^d} |\nabla_v \Gamma(s, y, v; t, x, u)| dx du \leq \frac{C}{\sqrt{t-s}}, \quad \forall 0 \leq s < t \leq T. \quad (4.4)$$

Let us now identify mild equations satisfied by  $(\rho_t; t \in (0, T])$  and  $(\rho_t^\epsilon; t \in (0, T])$ .

## 4.2 Mild equations for the densities of $(X, \mathcal{U})$ and $(X^\epsilon, \mathcal{U}^\epsilon)$

Consider a weak solution  $(X, \mathcal{U})$  of (1.1). For all  $f \in \mathcal{C}_b(\mathbb{R}^{2d})$ , since  $G_{t,f}$  is a classical solution of (4.2), Itô's formula leads to

$$\begin{aligned} \mathbb{E}_{\mathbb{P}} [f(X_{t \wedge \tau_M}, \mathcal{U}_{t \wedge \tau_M})] &= \mathbb{E}_{\mathbb{P}} [G_{t,f}(0, X_0, \mathcal{U}_0)] \\ &+ \mathbb{E}_{\mathbb{P}} \left[ \int_0^{t \wedge \tau_M} (\nabla_v G_{t,f}(s, X_s, \mathcal{U}_s) \cdot B[X_s, \mathcal{U}_s; \rho_s]) ds \right], \end{aligned} \quad (4.5)$$

where  $\{\tau_M; M \geq 1\}$  is the sequence of stopping times

$$\tau_M = \inf \{t > 0; |X_t| + |\mathcal{U}_t| \geq M\}.$$

The boundedness of  $b$  and  $\sigma$  implies that  $\lim_{M \rightarrow +\infty} \tau_M = +\infty$ ,  $\mathbb{P}$ -a.s. By the dominated convergence theorem, the left-hand side of (4.5) converges to  $\mathbb{E}_{\mathbb{P}} [f(X_t, \mathcal{U}_t)]$  as  $M$  tends to infinity. For the right-hand side, Proposition 4.4 shows that, for  $s \neq t$

$$\nabla_v G_{t,f}(s, y, v) = \int_{\mathbb{R}^{2d}} \nabla_v \Gamma(s, y, v; t, x, u) f(x, u) dx du,$$

and,  $\mathbb{P}$ -a.s.,

$$\sup_{M \geq 1} \int_0^{t \wedge \tau_M} |\nabla_v G_{t,f}(s, X_s, \mathcal{U}_s) \cdot B[X_s, \mathcal{U}_s; \rho_s]| ds \leq \|b\|_{\infty} \|f\|_{\infty} \int_0^t \frac{C}{\sqrt{t-s}} ds.$$

Letting  $M$  tends to infinity, we get

$$\begin{aligned} \int_{\mathbb{R}^{2d}} f(x, u) \rho_t(x, u) dx du &= \int_{\mathbb{R}^{2d}} G_{t,f}(0, y, v) \mu_0(dy dv) \\ &+ \int_{Q_t} (\nabla_v G_{t,f}(s, y, v) \cdot \rho_s(y, v) B[y, v; \rho_s]) ds dy dv. \end{aligned} \quad (4.6)$$

We denote by  $(S_{t,s}^*; 0 \leq s < t \leq T)$  the adjoint of the transition operator of  $(Y_t^{s,y,v}, V_t^{s,y,v})$ , that is

$$S_{t,s}^*(f)(x, u) = \int_{\mathbb{R}^{2d}} \Gamma(s, y, v; t, x, u) f(y, v) dy dv.$$



In view of Proposition 4.4, for all  $0 \leq s < t \leq T$ ,  $S_{t,s}^*$  is a linear operator from  $\mathcal{M}(\mathbb{R}^{2d})$  to  $L^1(\mathbb{R}^{2d})$ . In particular,  $S_{t,0}^*(\mu_0) \in L^1(\mathbb{R}^{2d})$  and the first term in the right-hand side in (4.6) can be rewritten as

$$\int_{\mathbb{R}^{2d}} S_{t,0}^*(\mu_0)(x, u) f(x, u) dx du.$$

In addition, for all  $0 \leq s < t \leq T$ , we define the operator  $S'_{t,s} : L^1(\mathbb{R}^{2d}; \mathbb{R}^d) \rightarrow L^1(\mathbb{R}^{2d}; \mathbb{R})$  by

$$S'_{t,s}(h(\cdot))(x, u) = \int_{\mathbb{R}^{2d}} (\nabla_v \Gamma(s, y, v; t, x, u) \cdot h(y, v)) dy dv.$$

In particular, Proposition 4.4 shows that

$$\int_0^t \|S'_{t,s}(h(s, \cdot))\|_{L^1(\mathbb{R}^{2d})} ds \leq \int_0^t \frac{C}{\sqrt{t-s}} \|h(s, \cdot)\|_{L^1(\mathbb{R}^{2d})} ds \quad (4.7)$$

for all  $h \in L^\infty((0, T); L^1(\mathbb{R}^{2d}))$ . Thus the second term in the right-hand side of (4.6) writes

$$\int_{\mathbb{R}^{2d}} f(x, u) \int_0^t S'_{t,s}(\rho_s(\cdot)B[\cdot; \rho_s])(x, u) ds dx du.$$

Therefore the marginal distributions  $(\rho_t; t \in (0, T])$  of the solution of (1.1) satisfy the mild equation

$$\forall t \in (0, T], \rho_t = S_{t,0}^*(\mu_0) + \int_0^t S'_{t,s}(\rho_s(\cdot)B[\cdot; \rho_s]) ds, \quad \text{in } L^1(\mathbb{R}^{2d}). \quad (4.8)$$

The preceding calculations hold true when  $\rho$  and  $B[\cdot; \rho]$  are replaced by  $\rho^\epsilon$  and  $B_\epsilon[\cdot; \rho^\epsilon]$ . Therefore the marginal distributions  $(\rho_t^\epsilon; t \in (0, T])$  satisfy the mild equation

$$\forall t \in (0, T], \rho_t^\epsilon = S_{t,0}^*(\mu_0) + \int_0^t S'_{t,s}(\rho_s^\epsilon(\cdot)B_\epsilon[\cdot; \rho_s^\epsilon]) ds, \quad \text{in } L^1(\mathbb{R}^{2d}). \quad (4.9)$$

### 4.3 Uniqueness of the solutions to the mild equations (4.8) and (4.9)

**Lemma 4.5.** *There exists at most one positive solution  $(\rho_t)$  to Equation (4.8), and at most one solution  $(\rho_t^\varepsilon)$  to (4.9).*

*Proof.* Let  $(\rho_t^1)$  and  $(\rho_t^2)$  be two positive solutions of (4.8). We set

$$\bar{\rho}_t^i(x) := \int_{\mathbb{R}^d} \rho_t^i(x, u) du.$$

For a.e.  $(t, x)$  in  $(0, T] \times \mathbb{R}^d$ ,  $\bar{\rho}_t^i(x) > 0$  for  $i = 1, 2$  and so,  $B[x, u; \rho_t^i]$  in (1.2) is well-defined for a.e.  $(t, x, u) \in (0, T] \times \mathbb{R}^{2d}$ . In view of (4.8), we have

$$\|\rho_t^1 - \rho_t^2\|_{L^1(\mathbb{R}^{2d})} = \int_{\mathbb{R}^{2d}} \left| \int_0^t S'_{t,s} (\rho_s^1(\cdot) B[\cdot; \rho_s^1] - \rho_s^2(\cdot) B[\cdot; \rho_s^2]) (x, u) ds \right| dx du. \quad (4.10)$$

We aim to prove the following estimate which implies the uniqueness result by Gronwall's lemma: there exists  $C > 0$  such that, for all  $t \in (0, T]$ ,

$$\|\rho_t^1 - \rho_t^2\|_{L^1(\mathbb{R}^{2d})} \leq C \int_0^t \frac{\|\rho_s^1 - \rho_s^2\|_{L^1(\mathbb{R}^d)}}{\sqrt{t-s}} ds. \quad (4.11)$$

From (4.10), we have

$$\begin{aligned} \|\rho_t^1 - \rho_t^2\|_{L^1(\mathbb{R}^{2d})} &\leq \int_{Q_t} |S'_{t,s} ((\rho_s^1(\cdot) - \rho_s^2(\cdot)) B[\cdot; \rho_s^1]) (x, u)| ds dx du \\ &\quad + \int_{Q_t} |S'_{t,s} (\rho_s^2(\cdot) (B[\cdot; \rho_s^1] - B[\cdot; \rho_s^2])) (x, u)| ds dx du \\ &=: A_1 + A_2. \end{aligned} \quad (4.12)$$

In view of (4.7) and the boundedness of  $b$ , we get

$$A_1 \leq \|b\|_\infty \int_0^t \frac{C}{\sqrt{t-s}} \|\rho_s^1 - \rho_s^2\|_{L^1(\mathbb{R}^{2d})} ds. \quad (4.13)$$

We now consider  $A_2$ . As

$$\frac{a_1}{b_1} - \frac{a_2}{b_2} = \frac{a_1 - a_2}{b_2} + \frac{a_1(b_2 - b_1)}{b_2 b_1}, \quad \forall a_1, a_2 \in \mathbb{R}, b_1, b_2 > 0, \quad (4.14)$$

we observe that

$$\begin{aligned} B[y, v; \rho_s^1] - B[y, v; \rho_s^2] &= \frac{1}{\bar{\rho}_s^2(y)} \int_{\mathbb{R}^d} b(v', v) (\rho_s^1(y, v') - \rho_s^2(y, v')) dv' \\ &\quad + \frac{\int_{\mathbb{R}^d} b(v', v) \rho_s^1(y, v') dv'}{\bar{\rho}_s^2(y) \bar{\rho}_s^1(y)} (\bar{\rho}_s^2(y) - \bar{\rho}_s^1(y)). \end{aligned}$$

Hence, for all  $0 \leq s < t$ ,

$$\begin{aligned} &\|\rho_s^2(\cdot) (B[\cdot; \rho_s^1] - B[\cdot; \rho_s^2])\|_{L^1(\mathbb{R}^{2d})} \\ &\leq 2\|b\|_\infty \int_{\mathbb{R}^{2d}} \frac{\rho_s^2(y, v)}{\bar{\rho}_s^2(y)} \int_{\mathbb{R}^d} |\rho_s^1(y, v') - \rho_s^2(y, v')| dv' dy dv \leq 2\|b\|_\infty \|\rho_s^1 - \rho_s^2\|_{L^1(\mathbb{R}^{2d})}. \end{aligned}$$

In view of (4.7) we thus have

$$A_2 \leq \|b\|_\infty \int_0^t \frac{C}{\sqrt{t-s}} \|\rho_s^1 - \rho_s^2\|_{L^1(\mathbb{R}^{2d})} ds. \quad (4.15)$$

It remains to gather (4.13) and (4.15) to get (4.11).

We now prove uniqueness for (4.9). First, let us observe that, by using (4.7) in Equation (4.9), one can find  $C > 0$  such that, for all solution  $(\rho_t^\epsilon)$  of (4.9),

$$\|\rho_t^\epsilon\|_{L^1(\mathbb{R}^{2d})} \leq C, \quad \forall t \in (0, T].$$

Next, consider two nonnegative solutions  $(\rho_t^{\epsilon,1})$  and  $(\rho_t^{\epsilon,2})$  of (4.9). As in (4.12), we have

$$\begin{aligned} \|\rho_t^{\epsilon,1} - \rho_t^{\epsilon,2}\|_{L^1(\mathbb{R}^{2d})} &\leq \int_{Q_t} |S'_{t,s} ((\rho_s^{\epsilon,1}(\cdot) - \rho_s^{\epsilon,2}(\cdot)) B_\epsilon[\cdot; \rho_s^{\epsilon,1}]) (x, u)| ds dx du \\ &\quad + \int_{Q_t} |S'_{t,s} (\rho_s^{\epsilon,2}(\cdot) (B_\epsilon[\cdot; \rho_s^{\epsilon,1}] - B_\epsilon[\cdot; \rho_s^{\epsilon,2}])) (x, u)| ds dx du \\ &=: A_1^\epsilon + A_2^\epsilon. \end{aligned}$$

Obviously,

$$A_1^\epsilon \leq \|b\|_\infty \int_0^t \frac{C}{\sqrt{t-s}} \|\rho_s^{\epsilon,1} - \rho_s^{\epsilon,2}\|_{L^1(\mathbb{R}^{2d})} ds.$$

In order to estimate  $A_2^\epsilon$ , we use again (4.14) and we observe that: for a.e.  $(s, y, v) \in Q_t$ ,

$$\begin{aligned} & \rho_s^{\epsilon,2}(y, v) (B_\epsilon[y, v; \rho_s^{\epsilon,1}] - B_\epsilon[y, v; \rho_s^{\epsilon,2}]) \\ &= \frac{\rho_s^{\epsilon,2}(y, v)}{\phi_\epsilon \star \bar{\rho}_s^{\epsilon,2}(y) + \epsilon} \int_{\mathbb{R}^d} b(v', v) (\phi_\epsilon \star \rho_s^{\epsilon,1}(y, v') - \phi_\epsilon \star \rho_s^{\epsilon,2}(y, v')) dv' \\ & \quad + \frac{\rho_s^{\epsilon,2}(y, v) \int_{\mathbb{R}^d} b(v', v) \phi_\epsilon \star \rho_s^{\epsilon,1}(y, v') dv'}{(\phi_\epsilon \star \bar{\rho}_s^{\epsilon,2}(y) + \epsilon) (\phi_\epsilon \star \bar{\rho}_s^{\epsilon,1}(y) + \epsilon)} (\phi_\epsilon \star \bar{\rho}_s^{\epsilon,2}(y) - \phi_\epsilon \star \bar{\rho}_s^{\epsilon,1}(y)) \\ & \leq \frac{\|b\|_\infty \rho_s^{\epsilon,2}(y, v)}{(\phi_\epsilon \star \bar{\rho}_s^{\epsilon,2}(y) + \epsilon)} \int_{\mathbb{R}^{2d}} \phi_\epsilon(y - y') |\rho_s^{\epsilon,1}(y', v') - \rho_s^{\epsilon,2}(y', v')| dy' dv'. \end{aligned}$$

In view of (4.7), it follows that

$$A_2^\epsilon \leq \|b\|_\infty \frac{\|\phi\|_\infty}{\epsilon^{d+1}} \int_{Q_t} \frac{C}{\sqrt{t-s}} \|\rho_s^{\epsilon,1} - \rho_s^{\epsilon,2}\|_{L^1(\mathbb{R}^{2d})} ds.$$

Hence

$$\|\rho_t^{\epsilon,1} - \rho_t^{\epsilon,2}\|_{L^1(\mathbb{R}^{2d})} \leq C \|b\|_\infty \left(1 + \frac{2\|\phi\|_\infty}{\epsilon^{d+1}}\right) \int_0^t \frac{\|\rho_s^{\epsilon,1} - \rho_s^{\epsilon,2}\|_{L^1(\mathbb{R}^{2d})}}{\sqrt{t-s}} ds.$$

We conclude on the uniqueness result for (4.9) by applying Gronwall's lemma.  $\square$

## 5 Existence results

In this section, we establish that Equations (3.4) and (1.1) admit a solution.

**Proposition 5.1.** *The martingale problem  $(MP_\epsilon)$  in Definition 4.1 has a unique solution  $\mathbb{P}^\epsilon$ . Furthermore, when  $\epsilon$  tends to 0,  $\mathbb{P}^\epsilon$  converges to a solution of the martingale problem  $(MP)$  introduced in Definition 4.2.*

The proof of the proposition 5.1 proceeds in two steps.

In the first step, we construct a weak solution to (3.4) by studying the interacting system (3.3) as the number of particles tends to infinity. As in Sznitman [18], we prove the relative compactness of the sequence of the empirical measures of the particles (see Lemma 5.3). We then show that the support of the limit probability measure is the set of solutions of the martingale problem  $(MP_\epsilon)$  (see Lemma 5.4). Using the uniqueness result in Proposition 4.3, we then get the propagation of chaos result.

The second step consists in exhibiting a solution to the martingale problem  $(MP)$  as the limit of the solution to martingale problem  $(MP_\epsilon)$  when  $\epsilon$  tends to 0.

## 5.1 A propagation of chaos result

Throughout this section we fix  $\epsilon > 0$ .

**Proposition 5.2.** *There exists a unique probability measure  $\mathbb{P}^\epsilon$  solution to the martingale problem  $(MP_\epsilon)$ . Moreover, the sequence of probability laws  $\{\mathbb{P}^{\epsilon,N}; N \geq 1\}$  of the processes  $\{(X^{i,\epsilon,N}, \mathcal{U}^{i,\epsilon,N}); 1 \leq i \leq N\}$  is  $\mathbb{P}^\epsilon$ -chaotic.*

Let  $\bar{\mu}^{\epsilon,N}$  be the empirical measure valued in  $\mathcal{M}(\mathcal{C}([0, T]; \mathbb{R}^{2d}))$  by

$$\bar{\mu}^{\epsilon,N} = \frac{1}{N} \sum_{i=1}^N \delta_{\{X^{i,\epsilon,N}, \mathcal{U}^{i,\epsilon,N}\}}.$$

Let  $\bar{\mathbb{P}}^{\epsilon,N}$  in  $\mathcal{M}(\mathcal{M}(\mathcal{C}([0, T]; \mathbb{R}^{2d})))$  be the probability law of  $\bar{\mu}^{\epsilon,N}$ .

**Lemma 5.3.** *The sequence  $\{\bar{\mathbb{P}}^{\epsilon,N}; N \geq 1\}$  is tight.*

*Proof.* We proceed as in Sznitman [18]: since the particle systems is exchangeable, the tightness of  $\bar{\mathbb{P}}^{\epsilon,N}$  is equivalent to the tightness of the probability laws of  $\{(X^{1,\epsilon,N}, \mathcal{U}^{1,\epsilon,N}); N \geq 1\}$ . Let  $((x_t^{(i)}, u_t^{(i)}); t \in (0, T], i = 1, \dots, N)$  be the canonical processes in the sample space  $\mathcal{C}([0, T]; \mathbb{R}^{2dN})$ . In view of the boundedness of  $b$  and  $\sigma$ ,

$$\mathbb{E}_{\mathbb{P}^{\epsilon,N}} [|u_t^{(1)} - u_s^{(1)}|^4] \leq C(t-s)^4 \quad \text{and} \quad \mathbb{E}_{\mathbb{P}^{\epsilon,N}} [|x_t^{(1)} - x_s^{(1)}|^2] \leq C(t-s)^2.$$

The result follows from the Kolmogorov criterion.  $\square$

We have shown that  $\{\bar{\mathbb{P}}^{\epsilon, N}; N \geq 1\}$  is relatively compact. We still denote by  $\{\bar{\mathbb{P}}^{\epsilon, N}; N \geq 1\}$  a weakly convergent subsequence. Let  $\bar{\mathbb{P}}^{\epsilon, \infty}$  be the limit of such a subsequence.

**Lemma 5.4.**  $\bar{\mathbb{P}}^{\epsilon, \infty}$  assigns full measure to the set of the solutions to the martingale problem  $(MP_\epsilon)$ .

*Proof.* We denote by  $m$  the canonical measure on  $\mathcal{M}(\mathcal{C}([0, T]; \mathbb{R}^{2d}))$ . Since  $\{(X_0^{i, \epsilon, N}, \mathcal{U}_0^{i, \epsilon, N}); 1 \leq i \leq N\}$  are  $\mu_0$ -i.i.d., it is easy to check that

$$\bar{\mathbb{P}}^{\epsilon, N} \circ (m \circ (x_0, u_0)^{-1})^{-1} = \delta_{\mu_0},$$

and the same holds for  $\bar{\mathbb{P}}^{\epsilon, \infty}$ , by the weak convergence of  $\bar{\mathbb{P}}^{\epsilon, N}$  to  $\bar{\mathbb{P}}^{\epsilon, \infty}$ . The property (i) of  $(MP_\epsilon)$  is then true for all  $m$ ,  $\mathbb{P}^{\epsilon, \infty}$ -a.e.

Now we prove that,  $\bar{\mathbb{P}}^{\epsilon, \infty}$ -a.e.,  $m$  satisfies the properties (ii) and (iii) of  $(MP_\epsilon)$ . We define  $\alpha : [0, T] \times \mathbb{R}^d \times \mathbb{R}^d \times \mathcal{M}(\mathcal{C}([0, T]; \mathbb{R}^{2d})) \mapsto \mathbb{R}^d$  by

$$\alpha(t, x, u, m) := \frac{\mathbb{E}_m [b(u_t, u) \phi_\epsilon(x - x_t)]}{\mathbb{E}_m [\phi_\epsilon(x - x_t)] + \epsilon},$$

For any  $f \in \mathcal{C}_c^2(\mathbb{R}^{2d})$ , for any  $0 \leq t_1 \leq \dots \leq t_n \leq s < t \leq T$ , and any finite family of functions  $\{\psi_j; 1 \leq j \leq n\}$  in  $\mathcal{C}_b(\mathbb{R}^{2d})$ , we set

$$F^\epsilon(m) := \left| \mathbb{E}_m \left[ \prod_{j=1}^n \psi_j(x_{t_j}, u_{t_j}) \left( f(x_t, u_t) - f(x_s, u_s) - \int_s^t (u_\theta \cdot \nabla_x f(x_\theta, u_\theta)) d\theta - \int_s^t (\alpha(\theta, x_\theta, u_\theta; m_\theta) \cdot \nabla_u f(x_\theta, u_\theta)) + \frac{1}{2} \sum_{i,j=1}^d a^{(i,j)}(\theta, x_\theta, u_\theta) \partial_{u_i, u_j}^2 f(x_\theta, u_\theta) d\theta \right) \right] \right|.$$

We aim to prove that  $F^\epsilon = 0$ ,  $\bar{\mathbb{P}}^{\epsilon, \infty}$ -a.s. Let us first check that, for  $\bar{\mathbb{P}}^{\epsilon, \infty}$ -a.e.  $m$ ,

$$\int_{\mathcal{C}([0, T]; \mathbb{R}^d)} \max_{t \in [0, T]} |u_t(\omega)| dm(\omega) < +\infty,$$

and therefore  $F^\epsilon < +\infty$ ,  $\bar{\mathbb{P}}^{\epsilon, \infty}$ -a.s. Indeed, by setting  $G(m) = \mathbb{E}_m [\max_{t \in [0, T]} |u_t|]$ , we observe that, from the boundedness of  $b$  and  $\sigma$ ,

$$\sup_{N \geq 1} \mathbb{E}_{\bar{\mathbb{P}}^{\epsilon, N}} [G(m)] = \sup_{N \geq 1} \mathbb{E}_{\mathbb{P}^{\epsilon, N}} \left[ G\left(\frac{1}{N} \sum_{i=1}^N \delta_{\{x^{(i)}, u^{(i)}\}}\right) \right] \leq C.$$

The upper bound for  $F^\epsilon$  follows by weak convergence.

It is clear that  $\alpha$  is a weakly continuous function of  $m$ , and therefore  $\mathbb{E}_{\bar{\mathbb{P}}^{\epsilon, N}}[F^\epsilon(m)]$  tends to  $\mathbb{E}_{\mathbb{P}^{\epsilon, \infty}}[F^\epsilon(m)]$ . Moreover, in view of (3.3),  $\mathbb{E}_{\bar{\mathbb{P}}^{\epsilon, N}}[F^\epsilon(m)] \leq C/N$ , and then  $F^\epsilon = 0$ ,  $\bar{\mathbb{P}}^{\epsilon, \infty}$ -a.s. Consequently,

$$\begin{aligned} f(x_t, u_t) - f(x_0, u_0) - \int_0^t (u_\theta \cdot \nabla_x f(x_\theta, u_\theta)) d\theta \\ - \int_s^t (\alpha(\theta, x_\theta, u_\theta, m) \cdot \nabla_u f(x_\theta, u_\theta)) + \frac{1}{2} \sum_{i,j=1}^d a^{(i,j)}(\theta, x_\theta, u_\theta) \partial_{u_i, u_j}^2 f(x_\theta, u_\theta) d\theta \end{aligned}$$

is a  $m$ -martingale. As  $\alpha$  is uniformly bounded, by Cameron–Martin’s formula, we easily check that  $m \circ (x_\theta, u_\theta)^{-1}$  has a density  $\rho_\theta^\epsilon$ , so that

$$f(x_t, u_t) - f(x_0, u_0) - \int_0^t \mathcal{A}_{\rho_\theta^\epsilon} f(\theta, x_\theta, u_\theta) d\theta$$

is a  $m$ -martingale. This ends the proof of (ii) and (iii).  $\square$

Proposition 4.3 ensure that  $\bar{\mathbb{P}}^{\epsilon, \infty}$  is reduced to a Dirac mass, and we denote  $\mathbb{P}^\epsilon$  the point such that  $\bar{\mathbb{P}}^{\epsilon, \infty} = \delta_{\{\mathbb{P}^\epsilon\}}$ . This implies that  $\mathbb{P}^\epsilon$  is the unique solution to  $(MP_\epsilon)$ . Moreover, this implies the  $\mathbb{P}^\epsilon$ -chaoticity of the particle system (see Sznitman [18, Prop. 2.2]).

## 5.2 Convergence of the smoothed system

In this subsection, we prove that the probability measure  $\mathbb{P}^\epsilon$ , solution of the martingale problem  $(MP_\epsilon)$ , converges to the solution to the martingale problem  $(MP)$ . We start with studying the probability measure  $\tilde{\mathbb{P}}^\epsilon$  defined on  $\mathcal{C}([0, T]; \mathbb{R}^{3d})$  by

$$\tilde{\mathbb{P}}^\epsilon = \mathbb{P}^\epsilon \circ \left( (x_t, u_t, u_t - u_0 - \int_0^t B_\epsilon[x_s, u_s; \rho_s^\epsilon] ds); t \in [0, T] \right)^{-1}.$$

As in the proof of Lemma 5.3, using (3.4), the Kolmogorov criterion implies that the sequence  $\{\tilde{\mathbb{P}}^\epsilon; \epsilon > 0\}$  is relatively compact in  $\mathcal{C}([0, T]; \mathbb{R}^{3d})$ . Let  $\tilde{\mathbb{P}}^\epsilon$  be a converging subsequence and denote its limit by  $\tilde{\mathbb{P}}$ . Let us characterize

the support of  $\tilde{\mathbb{P}}$ . To this aim, we introduce the subset  $\mathcal{H}_{\|b\|_\infty}$  of  $\mathcal{C}([0, T]; \mathbb{R}^{3d})$  defined by

$$\mathcal{H}_{\|b\|_\infty} = \left\{ \begin{array}{l} (x, u, D) \text{ in } \mathcal{C}([0, T]; \mathbb{R}^{3d}), \text{ s.t. } x(t) = x(0) + \int_0^t u(s) ds, \text{ and} \\ u(t) - u(0) - D(t) = \int_0^t \beta(s) ds, \text{ for a measurable function } \beta : [0, T] \rightarrow \mathbb{R}^d \\ \text{s.t. } \sup_{t \in [0, T]} |\beta(t)| \leq \|b\|_\infty. \end{array} \right\}$$

We now prove that

**Lemma 5.5.**  $\tilde{\mathbb{P}}$  has full measure on  $\mathcal{H}_{\|b\|_\infty}$ .

*Proof.* According to the Portemanteau theorem, the weak convergence of  $\tilde{\mathbb{P}}^\epsilon$  to  $\tilde{\mathbb{P}}$  yields that, for all closed subset  $F$  of  $\mathcal{C}([0, T]; \mathbb{R}^{3d})$ ,

$$\limsup_{\epsilon \rightarrow 0} \tilde{\mathbb{P}}^\epsilon(F) \leq \tilde{\mathbb{P}}(F).$$

Since  $\tilde{\mathbb{P}}^\epsilon(\mathcal{H}_{\|b\|_\infty}) = 1$  for all  $\epsilon > 0$ , it suffices to show that  $\mathcal{H}_{\|b\|_\infty}$  is a closed subset of  $\mathcal{C}([0, T]; \mathbb{R}^{3d})$ . Let  $\{(x_n, u_n, D_n); n \in \mathbb{N}\}$  be a sequence of  $\mathcal{H}_{\|b\|_\infty}$  converging to  $(x, u, D)$  in  $\mathcal{C}([0, T]; \mathbb{R}^{3d})$ . Set  $A_n(t) = u_n(t) - u_n(0) - D_n(t)$  and  $A(t) = u(t) - u(0) - D(t)$ . By uniform convergence, it holds that

$$\begin{aligned} x(t) &= x(0) + \int_0^t u(s) ds, \quad \forall t \in [0, T], \\ \text{and } \lim_{n \rightarrow +\infty} \max_{t \in [0, T]} |A_n(t) - A(t)| &= 0. \end{aligned}$$

To prove that  $(x, u, D)$  belongs to  $\mathcal{H}_{\|b\|_\infty}$ , it remains to show that  $A$  is a.e. differentiable with a time derivative uniformly bounded by  $\|b\|_\infty$ . By the Riesz representation theorem, it is enough to prove that

$$\left| \int_0^T A(t) f'(t) dt \right| \leq \|b\|_\infty \int_0^T |f(t)| dt, \quad \forall f \in \mathcal{C}_c^1([0, T]; \mathbb{R}^d).$$

As  $A_n(t) = \int_0^t \beta_n(s) ds$  for some measurable function  $\beta_n$  satisfying  $\sup_{t \in [0, T]} |\beta_n(t)| \leq \|b\|_\infty$ , an integration by parts allows us to write

$$\left| \int_0^T A(t) f'(t) dt \right| \leq \left| \lim_{n \rightarrow +\infty} \int_0^T \beta_n(t) f(t) dt \right| \leq \|b\|_\infty \int_0^T |f(t)| dt,$$

which ends the proof. □



Consider the marginal distribution  $\mathbb{P}$  of  $\tilde{\mathbb{P}}$  on  $\mathcal{C}([0, T]; \mathbb{R}^{2d})$ , defined by

$$\mathbb{P} = \tilde{\mathbb{P}} \circ ((x_t, u_t); t \in [0, T])^{-1}.$$

We have the following result:

**Proposition 5.6.**  $\mathbb{P}$  solves the martingale problem (MP) in Definition 4.2.

*Proof.* The point (i) of (MP) is obvious. To prove (ii), consider  $((x_t, u_t, D_t); t \in [0, T])$  the canonical processes of  $\mathcal{C}([0, T]; \mathbb{R}^{3d})$ . In view of Lemma 5.5, we know that,  $\tilde{\mathbb{P}}$ -a.s., for all  $t \in [0, T]$ ,

$$\begin{aligned} x_t &= x_0 + \int_0^t u_s ds, \\ u_t &= u_0 + \int_0^t \beta_s ds + D_t, \end{aligned}$$

with  $\sup_{t \in [0, T]} |\beta_t| \leq \|b\|_\infty$ . Since  $\sigma$  is bounded continuous, the weak convergence of  $\tilde{\mathbb{P}}^c$  to  $\tilde{\mathbb{P}}$  yields that, for all function  $f$  in  $\mathcal{C}_b^2(\mathbb{R}^d)$ ,

$$f(D_t) - f(D_0) - \frac{1}{2} \sum_{i,j=1}^d \int_0^t a^{(i,j)}(s, x_s, u_s) \partial_{u_i, u_j}^2 f(D_s) ds$$

is a  $\tilde{\mathbb{P}}$ -martingale. We deduce that  $D_t = \int_0^t \sigma(\theta, x_\theta, u_\theta) dw_\theta$  for some  $d$ -dimensional Wiener process  $(w_t; t \in [0, T])$ .

In view of (3.1), Girsanov's theorem allows one to define a new probability  $\mathbb{Q}$  absolutely continuous to  $\tilde{\mathbb{P}}$  on  $\mathcal{C}([0, T]; \mathbb{R}^{3d})$  such that  $\mathbb{Q} \circ (x_t, u_t)^{-1}$  is the law of the Langevin system

$$(y_t, v_t) = \left( x_0 + \int_0^t v_s ds, u_0 + \int_0^t \sigma(s, y_s, v_s) dw_s \right).$$

In view of Proposition 4.4, for all  $t \in [0, T]$  the law of  $(y_t, v_t)$  is absolutely continuous w.r.t. to Lebesgue measure. Thus the measure  $\mathbb{P} \circ (x_t, u_t)^{-1}$  has also a density  $\rho_t$  which satisfies

$$\gamma_t(x, u) = \rho_t(x, u) \mathbb{E}_{\mathbb{P}} \left( Z_t / (x_t, u_t) = (x, u) \right), \text{ for a.e. } (x, u) \in \mathbb{R}^{2d},$$

where  $Z_t$  is the restriction to  $\mathcal{C}([0, t]; \mathbb{R}^{3d})$  of the density  $\frac{d\mathbb{Q}}{d\mathbb{P}}$ , and

$$\gamma_t(x, u) := \int_{\mathbb{R}^{2d}} \Gamma(0, y, v; t, x, u) \mu_0(y, v) dy dv,$$

where  $\Gamma(0, y, v; t, x, u)$  is defined in Proposition 4.4. We now recall the following estimate (see Di Francesco & Polidoro [6]): there exists  $\eta > 0$ , depending only on  $\lambda$  and  $d$ , such that

$$\Gamma(t, x, u; s, y, v) \geq \Gamma_\eta(t, x, u; s, y, v),$$

where  $\Gamma_\eta$  is defined in (A.1). Hence, the function  $\rho_t(x, u)$  is strictly positive. We thus have proven (ii).

We now prove (iii). Observe that there exists  $C > 0$  such that

$$\sup_{\epsilon > 0} \mathbb{E}_{\mathbb{P}^\epsilon} \left[ \max_{t \in [0, T]} |u_t| \right] \leq C.$$

Therefore it suffices to prove that, for all  $f \in \mathcal{C}_c^2(\mathbb{R}^{2d})$  and all process  $(\Psi_s)$  of the form

$$\Psi_s = \Psi_s(x, u) := \prod_{j=1}^n \psi_j(x_{t_j}, u_{t_j}),$$

where the  $\psi_j$ 's are bounded continuous functions, one has

$$\mathbb{E}_{\mathbb{P}} \left[ \Psi_s \left( f(x_t, u_t) - f(x_s, u_s) - \int_s^t \mathcal{A}_{\rho_\theta} f(x_\theta, u_\theta) d\theta \right) \right] = 0, \quad (5.1)$$

Since  $\tilde{\mathbb{P}}^\epsilon$  converges weakly to  $\tilde{\mathbb{P}}$ , we have

$$\begin{aligned} \lim_{\epsilon \rightarrow 0^+} \mathbb{E}_{\mathbb{P}^\epsilon} \left[ \Psi_s \left( f(x_t, u_t) - f(x_s, u_s) - \int_s^t \mathcal{L} f(x_\theta, u_\theta) d\theta \right) \right] \\ = \mathbb{E}_{\mathbb{P}} \left[ \Psi_s \left( f(x_t, u_t) - f(x_s, u_s) - \int_s^t \mathcal{L} f(x_\theta, u_\theta) d\theta \right) \right], \end{aligned}$$

for  $\mathcal{L}$  defined in (4.3). To obtain (5.1), it thus remains to show

$$\begin{aligned} \lim_{\epsilon \rightarrow 0^+} \mathbb{E}_{\mathbb{P}^\epsilon} \left[ \Psi_s \int_s^t (B_\epsilon [x_\theta, u_\theta; \rho_\theta^\epsilon] \cdot \nabla_u f(x_\theta, u_\theta)) d\theta \right] \\ = \mathbb{E}_{\mathbb{P}} \left[ \Psi_s \int_s^t (B [x_\theta, u_\theta; \rho_\theta] \cdot \nabla_u f(x_\theta, u_\theta)) d\theta \right]. \end{aligned} \quad (5.2)$$

If  $\mathbb{P}^\epsilon$  were the law of a strongly elliptic diffusion process and the coefficient  $B_\epsilon$  would not depend on  $\rho^\epsilon$ , (5.2) would result from Stroock & Varadhan's results on limits of martingale problems: see Lemma 11.3.2 and Lemma 9.1.15 in [16]. In our situation, we prove that (5.2) holds true by adapting Stroock & Varadhan's techniques and by taking advantage of the mild equation (4.9). By adding and subtracting to (5.2) the terms

$$\begin{aligned} \mathbb{E}_{\mathbb{P}^\epsilon} \left[ \Psi_s \int_s^t (B_\xi [x_\theta, u_\theta; \rho_\theta^\epsilon] \cdot \nabla_u f(x_\theta, u_\theta)) d\theta \right], \\ \text{and } \mathbb{E}_{\mathbb{P}} \left[ \Psi_s \int_s^t (B_\xi [x_\theta, u_\theta; \rho_\theta] \cdot \nabla_u f(x_\theta, u_\theta)) d\theta \right], \end{aligned}$$

we obtain that

$$\begin{aligned} & \left| \mathbb{E}_{\mathbb{P}^\epsilon} \left[ \Psi_s \int_s^t (B_\epsilon [x_\theta, u_\theta; \rho_\theta^\epsilon] \cdot \nabla_u f(x_\theta, u_\theta)) d\theta \right] \right. \\ & \left. - \mathbb{E}_{\mathbb{P}} \left[ \Psi_s \int_s^t (B [x_\theta, u_\theta; \rho_\theta] \cdot \nabla_u f(x_\theta, u_\theta)) d\theta \right] \right| \\ & \leq \left| \mathbb{E}_{\mathbb{P}^\epsilon} \left[ \Psi_s \int_s^t (B_\xi [x_\theta, u_\theta; \rho_\theta^\epsilon] \cdot \nabla_u f(x_\theta, u_\theta)) d\theta \right] \right. \\ & \quad \left. - \mathbb{E}_{\mathbb{P}} \left[ \Psi_s \int_s^t (B_\xi [x_\theta, u_\theta; \rho_\theta] \cdot \nabla_u f(x_\theta, u_\theta)) d\theta \right] \right| \\ & \quad + \left| \mathbb{E}_{\mathbb{P}^\epsilon} \left[ \Psi_s \int_s^t (B_\epsilon [x_\theta, u_\theta; \rho_\theta^\epsilon] - B_\xi [x_\theta, u_\theta; \rho_\theta^\epsilon] \cdot \nabla_u f(x_\theta, u_\theta)) d\theta \right] \right| \\ & \quad + \left| \mathbb{E}_{\mathbb{P}} \left[ \Psi_s \int_s^t (B_\xi [x_\theta, u_\theta; \rho_\theta] - B [x_\theta, u_\theta; \rho_\theta] \cdot \nabla_u f(x_\theta, u_\theta)) d\theta \right] \right| \\ & =: I_{\epsilon, \xi} + J_{\epsilon, \xi} + K_\xi. \end{aligned}$$

We aim to show that

$$\forall \xi > 0, \quad \lim_{\epsilon \rightarrow 0^+} I_{\epsilon, \xi} = 0, \quad (5.3)$$

$$\lim_{\xi \rightarrow 0^+} K_{\xi} = 0, \quad (5.4)$$

$$\forall \xi > 0, \quad \limsup_{\epsilon \rightarrow 0^+} J_{\epsilon, \xi} = 0. \quad (5.5)$$

A crucial step to get these results is to prove the following proposition.

**Proposition 5.7.** *For all  $0 < t \leq T$ ,  $\rho_t^\epsilon$  converges to  $\rho_t$  in  $L^1(\mathbb{R}^{2d})$  when  $\epsilon \rightarrow 0^+$ .*

In view of Lemma A.2 in the Appendix, Proposition 5.7 results from the following lemma:

**Lemma 5.8.** *For all  $t \in (0, T]$ ,  $h, \delta \in \mathbb{R}^d$ , it holds that*

$$\lim_{|h|, |\delta| \rightarrow 0} \limsup_{\epsilon \rightarrow 0^+} \int_{\mathbb{R}^{2d}} |\rho_t^\epsilon(x+h, u+\delta) - \rho_t^\epsilon(x, u)| \, dx \, du = 0.$$

*Proof.* As  $\mathbb{P}^\epsilon$  is the unique solution to the martingale problem  $(MP_\epsilon)$ , its time marginals satisfy the mild equation (4.9). Thus, for all  $t \in (0, T]$ ,

$$\begin{aligned} & \limsup_{\epsilon \rightarrow 0^+} \int_{\mathbb{R}^{2d}} |\rho_t^\epsilon(x+h, u+\delta) - \rho_t^\epsilon(x, u)| \, dx \, du \\ & \leq \int_{\mathbb{R}^{2d}} |S_{t,0}^*(\mu_0)(x+h, u+\delta) - S_{t,0}^*(\mu_0)(x, u)| \, dx \, du \\ & \quad + \limsup_{\epsilon \rightarrow 0^+} \int_{Q_t} |S'_{t,s}(\rho_s^\epsilon(\cdot)B_\epsilon[\cdot; \rho_s^\epsilon])(x+h, u+\delta) - S'_{t,s}(\rho_s^\epsilon(\cdot)B_\epsilon[\cdot; \rho_s^\epsilon])(x, u)| \, ds \, dx \, du. \end{aligned}$$

Since  $S_{t,0}^*(\mu_0)$  belongs to  $L^1(\mathbb{R}^{2d})$ , Lemma A.3 implies that

$$\lim_{|h|, |\delta| \rightarrow 0} \int_{\mathbb{R}^{2d}} |S_{t,0}^*(\mu_0)(x+h, u+\delta) - S_{t,0}^*(\mu_0)(x, u)| \, dx \, du = 0.$$

In addition,

$$\begin{aligned} & \int_{Q_t} \left| S'_{t,s}(\rho_s^\epsilon(\cdot)B_\epsilon[\cdot; \rho_s^\epsilon])(x+h, u+\delta) - S'_{t,s}(\rho_s^\epsilon(\cdot)B_\epsilon[\cdot; \rho_s^\epsilon])(x, u) \right| ds dx du \\ & \leq \|b\|_\infty \int_{Q_t} \left( \int_{\mathbb{R}^{2d}} |\nabla_v \Gamma(s, y, v; t, x+h, u+\delta) - \nabla_v \Gamma(s, y, v; t, x, u)| dx du \right) \\ & \quad \times \rho_s^\epsilon(y, v) ds dy dv. \end{aligned}$$

Set

$$L_{h,\delta}(t, s, y, v) := \int_{\mathbb{R}^{2d}} |\nabla_v \Gamma(s, y, v; t, x+h, u+\delta) - \nabla_v \Gamma(s, y, v; t, x, u)| dx du.$$

As  $\tilde{\mathbb{P}}^\epsilon$  converges weakly to  $\tilde{\mathbb{P}}$ ,  $\rho_t^\epsilon$  converges weakly to  $\rho_t$  for all  $t \in [0, T]$  and

$$\lim_{\epsilon \rightarrow 0^+} \int_{Q_t} L_{h,\delta}(t, s, y, v) \rho_s^\epsilon(y, v) ds dy dv = \int_{Q_t} L_{h,\delta}(t, s, y, v) \rho_s(y, v) ds dy dv.$$

In addition, in view of (4.4), one has

$$\sup_{(y,v) \in \mathbb{R}^{2d}} L_{h,\delta}(t, s, y, v) \leq \frac{C}{\sqrt{t-s}}$$

for all  $t > s$ , from which

$$\int_{\mathbb{R}^{2d}} L_{h,\delta}(t, s, y, v) \rho_s^\epsilon(y, v) dy dv \leq \frac{C}{\sqrt{t-s}}.$$

It then remains to apply Lebesgue's dominated convergence theorem.  $\square$

Now we are in a position to prove (5.3), (5.4), and (5.5).

**Proof of (5.3).** For all fixed  $\xi > 0$ ,  $0 \leq s \leq t \leq T$ , the bounded functions  $\{f^{\epsilon', \xi}; \epsilon' > 0\}$  defined on  $\mathcal{C}([0, T]; \mathbb{R}^{2d})$  by

$$f^{\epsilon', \xi}(x., u.) := \Psi_s \int_s^t \left( B_\xi[x_\theta, u_\theta; \rho_\theta^{\epsilon'}] \cdot \nabla_u f(x_\theta, u_\theta) \right) d\theta$$

are equicontinuous (this latter property results from the definition of  $B_\xi$  and the proposition 5.7). Moreover, setting

$$f^\xi(x., u.) := \Psi_s \int_s^t (B_\xi[x_\theta, u_\theta; \rho_\theta] \cdot \nabla_u f(x_\theta, u_\theta)) d\theta,$$

we have

$$\begin{aligned} I_{\epsilon, \xi} &= |\mathbb{E}_{\mathbb{P}^\epsilon} f^{\epsilon, \xi} - \mathbb{E}_{\mathbb{P}} f^\xi| \\ &\leq \sup_{\epsilon' > 0} |\mathbb{E}_{\mathbb{P}^\epsilon} f^{\epsilon', \xi} - \mathbb{E}_{\mathbb{P}} f^{\epsilon', \xi}| + |\mathbb{E}_{\mathbb{P}} f^{\epsilon, \xi} - \mathbb{E}_{\mathbb{P}} f^\xi|. \end{aligned}$$

In view of Lemma A.4, the first term in the right-hand side tends to 0 with  $\epsilon$ . The second term tends also to 0 in view of Proposition 5.7. We thus have proven (5.3).

**Proof of (5.4).** We recall the notation

$$\bar{\rho}_\theta(x) := \int_{\mathbb{R}^d} \rho_\theta(x, v) dv.$$

Observe that

$$\begin{aligned} |B_\xi[x, u, \rho_\theta] - B[x, u; \rho_\theta]| &\leq \left| \frac{\int_{\mathbb{R}^d} b(v, u) \phi_\xi \star \rho_\theta(x, v) dv}{\phi_\xi \star \bar{\rho}_\theta(x) + \xi} - \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta(x, v) dv}{\bar{\rho}_\theta(x) + \xi} \right| \\ &\quad + \left| \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta(x, v) dv}{\bar{\rho}_\theta(x) + \xi} - \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta(x, v) dv}{\bar{\rho}_\theta(x)} \right|. \end{aligned}$$

Then, in view of (4.14), we get

$$\begin{aligned} &\left| \frac{\int_{\mathbb{R}^{2d}} b(v, u) \phi_\xi \star \rho_\theta(x, v) dv}{\phi_\xi \star \bar{\rho}_\theta(x) + \xi} - \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta(x, v) dv}{\bar{\rho}_\theta(x) + \xi} \right| \\ &\leq \left| \frac{\int_{\mathbb{R}^d} b(v, u) (\phi_\xi \star \rho_\theta(x, v) - \rho_\theta(x, v)) dv}{\bar{\rho}_\theta(x) + \xi} \right| \\ &\quad + \frac{\left| \int_{\mathbb{R}^d} b(v, u) \phi_\xi \star \rho_\theta(x, v) dv \right| \left| \int_{\mathbb{R}^d} (\phi_\xi \star \rho_\theta(x, v) - \rho_\theta(x, v)) dv \right|}{(\phi_\xi \star \bar{\rho}_\theta(x) + \xi) (\bar{\rho}_\theta(x) + \xi)} \\ &\leq \frac{2\|b\|_\infty}{\bar{\rho}_\theta(x) + \xi} \int_{\mathbb{R}^d} |\phi_\xi \star \rho_\theta(x, v) - \rho_\theta(x, v)| dv. \end{aligned} \tag{5.6}$$

Thus

$$\begin{aligned}
& \lim_{\xi \rightarrow 0^+} \int_{Q_T} |B_\xi[x, u, \rho_\theta] - B[x, u; \rho_\theta]| \rho_\theta(x, u) \, d\theta \, dx \, du \\
& \leq 2 \|b\|_\infty \lim_{\xi \rightarrow 0^+} \int_0^T \|\phi_\xi \star \rho_\theta - \rho_\theta\|_{L^1(\mathbb{R}^{2d})} \, d\theta \\
& \quad + \lim_{\xi \rightarrow 0^+} \int_{Q_T} \left| \frac{1}{\bar{\rho}_\theta(x) + \xi} - \frac{1}{\bar{\rho}_\theta(x)} \right| \int_{\mathbb{R}^d} b(v, u) \rho_\theta(x, v) \, dv \rho_\theta(x, u) \, d\theta \, dx \, du.
\end{aligned} \tag{5.7}$$

As

$$\|\phi_\xi \star \rho_\theta - \rho_\theta\|_{L^1(\mathbb{R}^{2d})} = \int_{\mathbb{R}^d} \phi(y) \int_{\mathbb{R}^{2d}} |\rho_\theta(x - \xi y, v) - \rho_\theta(x, v)| \, dx \, du \, dy,$$

Lemma A.3 implies that

$$\lim_{\xi \rightarrow 0^+} \int_0^T \|\phi_\xi \star \rho_\theta - \rho_\theta\|_{L^1(\mathbb{R}^{2d})} \, d\theta = 0. \tag{5.8}$$

The second limit in the right-hand side of (5.7) is null in view of the Lebesgue dominated convergence theorem and the positivity of  $(\rho_t; t \in (0, T])$ . That ends the proof of (5.4).

**Proof of (5.5).** Observe that

$$\begin{aligned}
J_{\epsilon, \xi} & \leq \|\Psi_s\|_\infty \|\nabla_u f\|_\infty \left( \int_{Q_t} |B_\xi[x, u; \rho_\theta^\epsilon] - B[x, u; \rho_\theta^\epsilon]| \rho_\theta^\epsilon(x, u) \, d\theta \, dx \, du \right. \\
& \quad \left. + \int_{Q_t} |B[x, u; \rho_\theta^\epsilon] - B_\epsilon[x, u; \rho_\theta^\epsilon]| \rho_\theta^\epsilon(x, u) \, d\theta \, dx \, du \right) \\
& =: \|\Psi_s\|_\infty \|\nabla_u f\|_\infty (J_{\epsilon, \xi}^1 + J_\epsilon^2).
\end{aligned}$$

In order to show that  $J_\epsilon^1$  tends to 0, observe that

$$\begin{aligned}
 & \int_{Q_t} |B_\xi[x, u, \rho_\theta^\epsilon] - B[x, u; \rho_\theta^\epsilon]| \rho_\theta^\epsilon(x, u) d\theta dx du \\
 & \leq \int_{Q_t} \left| \frac{\int_{\mathbb{R}^d} b(v, u) \phi_\xi \star \rho_\theta^\epsilon(x, v) dv}{\phi_\xi \star \bar{\rho}_\theta^\epsilon(x) + \xi} - \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta^\epsilon(x, v) dv}{\bar{\rho}_\theta^\epsilon(x) + \xi} \right| \rho_\theta^\epsilon(x, u) d\theta dx du \\
 & \quad + \int_{Q_t} \left| \frac{\int_{\mathbb{R}^{2d}} b(v, u) \rho_\theta^\epsilon(x, v) dv}{\bar{\rho}_\theta^\epsilon(x) + \xi} - \frac{\int_{\mathbb{R}^{2d}} b(v, u) \rho_\theta^\epsilon(x, v) dv}{\bar{\rho}_\theta^\epsilon(x)} \right| \rho_\theta^\epsilon(x, u) d\theta dx du.
 \end{aligned} \tag{5.9}$$

We now check that each term in the right-hand side of (5.9) tends to 0. Using (4.14) again, we get

$$\begin{aligned}
 & \left| \frac{\int_{\mathbb{R}^{2d}} b(v, u) \phi_\xi \star \rho_\theta^\epsilon(x, v) dv}{\phi_\xi \star \bar{\rho}_\theta^\epsilon(x) + \xi} - \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta^\epsilon(x, v) dv}{\bar{\rho}_\theta^\epsilon(x) + \xi} \right| \\
 & \leq \frac{2\|b\|_\infty}{\bar{\rho}_\theta^\epsilon(x) + \xi} \int_{\mathbb{R}^d} |\phi_\xi \star \rho_\theta^\epsilon(x, v) - \rho_\theta^\epsilon(x, v)| dv.
 \end{aligned}$$

Therefore the first term in the right-hand side of (5.9) is bounded from above by

$$C \int_{\mathbb{R}^d} |\phi_\xi \star \rho_\theta^\epsilon(x, v) - \rho_\theta^\epsilon(x, v)| dv.$$

Notice that

$$\begin{aligned}
 & \int_0^T \|\phi_\xi \star \rho_\theta^\epsilon - \rho_\theta^\epsilon\|_{L^1(\mathbb{R}^{2d})} d\theta \\
 & \leq \int_0^T \|\phi_\xi \star \rho_\theta^\epsilon - \phi_\xi \star \rho_\theta\|_{L^1(\mathbb{R}^{2d})} d\theta + \int_0^T \|\rho_\theta^\epsilon - \rho_\theta\|_{L^1(\mathbb{R}^{2d})} d\theta + \int_0^T \|\phi_\xi \star \rho_\theta - \rho_\theta\|_{L^1(\mathbb{R}^{2d})} d\theta \\
 & \leq 2 \int_0^T \|\rho_\theta^\epsilon - \rho_\theta\|_{L^1(\mathbb{R}^{2d})} d\theta + \int_0^T \|\phi_\xi \star \rho_\theta - \rho_\theta\|_{L^1(\mathbb{R}^{2d})} d\theta.
 \end{aligned}$$

Use (5.8) and apply Proposition 5.7. It results that the first term in the right-hand side of (5.9) tends to 0 with  $\epsilon$  and  $\xi$ .

To show that the second term in this right hand-side also tends to 0, we start with setting

$$A_{\xi, \epsilon}(\theta, x, u) := \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta^\epsilon(x, v) dv}{\bar{\rho}_\theta^\epsilon(x) + \xi} \rho_\theta^\epsilon(x, u) - \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta(x, v) dv}{\bar{\rho}_\theta(x) + \xi} \rho_\theta(x, u),$$



and inserting

$$\frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta^\epsilon(x, v) dv}{\bar{\rho}_\theta^\epsilon(x) + \xi} \rho_\theta(x, u).$$

We then apply (4.14) to

$$\frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta^\epsilon(x, v) dv}{\bar{\rho}_\theta^\epsilon(x) + \xi} - \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta(x, v) dv}{\bar{\rho}_\theta(x) + \xi}$$

and get

$$\begin{aligned} |A_{\xi, \epsilon}(\theta, x, u)| &\leq C \frac{\int_{\mathbb{R}^d} \rho_\theta^\epsilon(x, v) dv}{\bar{\rho}_\theta^\epsilon(x) + \xi} |\rho_\theta^\epsilon(x, u) - \rho_\theta(x, u)| + C \frac{\rho(x, u)}{\bar{\rho}_\theta^\epsilon(x) + \xi} \int_{\mathbb{R}^d} |\rho_\theta^\epsilon(x, v) - \rho_\theta(x, v)| dv \\ &\quad + C \frac{\rho(x, u)}{\bar{\rho}_\theta^\epsilon(x) + \xi} \frac{\int_{\mathbb{R}^d} \rho_\theta^\epsilon(x, v) dv}{\bar{\rho}_\theta^\epsilon(x) + \xi} |\bar{\rho}_\theta^\epsilon(x, u) - \bar{\rho}_\theta(x, u)|. \end{aligned}$$

We deduce that

$$\int_{Q_t} |A_{\xi, \epsilon}(\theta, x, u)| d\theta dx du \leq \int_0^T \|\rho_\theta^\epsilon - \rho_\theta\|_{L^1(\mathbb{R}^{2d})} d\theta$$

which tends to 0 in view of Proposition 5.7.

To end the proof that the second term in the right hand-side of (5.9) tends to 0, it now remains to observe that, owing to the positivity of  $(\rho_t; t \in [0, T])$ , we have

$$\begin{aligned} &\lim_{\xi \rightarrow 0^+} \limsup_{\epsilon \rightarrow 0^+} \int_{Q_t} \left| \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta^\epsilon(x, v) dv}{\bar{\rho}_\theta^\epsilon(x) + \xi} - \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta(x, v) dv}{\bar{\rho}_\theta(x)} \right| \rho_\theta^\epsilon(x, u) d\theta dx du \\ &\leq \lim_{\xi \rightarrow 0^+} \int_{Q_t} \left| \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta(x, v) dv}{\bar{\rho}_\theta(x) + \xi} - \frac{\int_{\mathbb{R}^d} b(v, u) \rho_\theta(x, v) dv}{\bar{\rho}_\theta(x)} \right| \rho_\theta(x, u) d\theta dx du = 0, \end{aligned}$$

by using Lebesgue's dominated convergence theorem. We thus have proven

$$\lim_{\xi \rightarrow 0^+} \limsup_{\epsilon \rightarrow 0^+} J_{\epsilon, \xi}^1 = 0. \quad (5.10)$$

Similarly, we have

$$\begin{aligned} & \lim_{\epsilon \rightarrow 0^+} \int_{Q_t} |B[x, u, \rho_\theta^\epsilon] - B_\epsilon[x, u; \rho_\theta^\epsilon]| \rho_\theta^\epsilon(x, u) d\theta dx du \\ & \leq 2\|b\|_\infty \lim_{\epsilon \rightarrow 0^+} \int_0^T \|\phi_\epsilon \star \rho_\theta^\epsilon - \rho_\theta^\epsilon\|_{L^1(\mathbb{R}^{2d})} d\theta \\ & \quad + \lim_{\epsilon \rightarrow 0^+} \int_{Q_t} \left| \frac{1}{\bar{\rho}_\theta^\epsilon(x) + \epsilon} - \frac{1}{\bar{\rho}_\theta^\epsilon(x)} \right| \left( \int_{\mathbb{R}^d} |b(v, u)| \rho_\theta^\epsilon(x, v) dv \right) \rho_\theta^\epsilon(x, u) d\theta dx du = 0, \end{aligned}$$

so that  $\lim_{\epsilon \rightarrow 0^+} J_\epsilon^2 = 0$ . In view of (5.10), the proof is completed.  $\square$

## 6 Conclusion and perspectives

In this paper we have studied a Lagrangian stochastic model and shown its well-posedness. We have proved that the unique weak solution is an hypoelliptic diffusion process whose dynamics depends on the conditional distribution of the velocity component knowing the position component. To our knowledge, this is the first theoretical result on the Lagrangian stochastic models modelling turbulent fluid particles. Bossy and Jabir [3] consider models with specular boundary conditions. See also Jabir [8]. A lot remains to be done to study the complex models developed by e.g. Pope [13].

We also emphasize another possible extension of our result. We conjecture the following PDE analysis result: the estimate (4.4) holds true under classical Hölder conditions rather than (3.2), possibly by using Maxwellian approximations rather than using the parametrix method.

## A Appendix

### A.1 Di Francesco and Pascucci's estimates on fundamental solutions of ultraparabolic PDEs

Before stating the estimate on fundamental solutions of ultraparabolic PDEs which are used in this paper, we need to introduce some new notation.

In [5] Di Francesco and Pascucci consider ultraparabolic PDEs of the type

$$-\partial_s \psi + \sum_{i,j} \bar{a}^{(i,j)} \partial_{v_i v_j}^2 \psi + (y, v) \cdot B \nabla_{(y,v)} \psi = 0,$$

where  $\bar{a}(s, y, v)$  and  $B$  are  $2d \times 2d$  matrices, and  $B$  has constant entries. The statement of their results for general matrices  $B$  require to introduce a pseudo-metric depending on  $B$  and some notational effort. We thus limit ourselves to our context where

$$B = \begin{pmatrix} 0 & 0 \\ \text{Id}_{\mathbb{R}^d} & 0 \end{pmatrix}.$$

In this context, Di Francesco and Pascucci's assumption on the coefficient  $\bar{a}$  writes as follows:  $\bar{a}$  is a bounded function and there exist  $\alpha \in (0, 1]$  and  $C > 0$  such that, for all  $(s, x, u), (t, y, v) \in (0, T] \times \mathbb{R}^d \times \mathbb{R}^d$ , the inequality (3.2) holds true.

For all  $\eta > 0$ , for  $s < t$ , let  $\Gamma_\eta(s, y, v; t, x, u)$  be the transition density function of

$$\left( y + \int_s^t v_\theta d\theta, v + \eta(W_t - W_s) \right)$$

that is,

$$\begin{aligned} & \Gamma_\eta(s, y, v; t, x, u) \\ &= \frac{\sqrt{3^d}}{(\sqrt{\pi}\eta(t-s))^{2d}} \\ & \times \exp \left\{ -\frac{6|x-y-(t-s)v|^2}{\eta(t-s)^3} + \frac{6(x-y-(t-s)v) \cdot (u-v)}{\eta(t-s)^2} - \frac{2|u-v|^2}{\eta(t-s)} \right\}. \end{aligned} \tag{A.1}$$

We are in a position to state the following theorem:

**Theorem A.1** (Di Francesco & Pascucci [5]). *Suppose that  $\bar{a}$  satisfies (3.1) and (3.2). There exists a fundamental solution  $\bar{\Gamma}(s, y, v; t, x, u)$  to the operator*

$$\bar{\mathcal{L}} = -\partial_s + \sum_{i,j} \bar{a}^{(i,j)}(s, y, v) \partial_{v_i v_j} + v \cdot \nabla_x$$

which satisfies

- (i) for all  $(t, x, u) \in (T_0, T] \times \mathbb{R}^{2d}$ ,  $1 \leq i, j \leq d$ , the derivatives  $\partial_{v_i} \bar{\Gamma}(s, y, v; t, x, u)$  exist and are continuous for all  $(s, y, v) \in (T_0, T] \times \mathbb{R}^{2d}$  such that  $(s, y, v) \neq (t, x, u)$ .
- (ii) Let  $f : \mathbb{R}^{2d} \rightarrow \mathbb{R}$  be a bounded continuous function. Then the function  $J_{T_0, T}$  defined by

$$J_{T_0, T}(t, x, u) = \int_{\mathbb{R}^{2d}} \bar{\Gamma}(T_0, y, v; t, x, u) f(y, v) dy dv, \quad t \in (T_0, T]$$

is the unique solution of the Cauchy problem

$$\begin{cases} \bar{\mathcal{L}} J_{T_0, T} = 0, \\ J_{T_0, T}(T_0, x, v) = f(x, v) \text{ in } \mathbb{R}^{2d}. \end{cases}$$

- (iii) There exists  $\eta > \lambda$  and a constant  $C > 0$  such that, for  $s < t$

$$|\nabla_v \bar{\Gamma}(s, y, v; t, x, u)| \leq \frac{C}{\sqrt{t-s}} \Gamma_\eta(s, y, v; t, x, u).$$

## A.2 Proof of Proposition 4.4

To get the existence of a weak solution, we adapt the proof of Theorem 6.3.2 in Stroock & Varadhan [16]: consider a sequence  $\{\sigma^n; n \in \mathbb{N}\}$  of  $\mathbb{R}^d \times \mathbb{R}^d$ -valued Lipschitz functions on  $[0, T] \times \mathbb{R}^d$  such that  $\lim_{n \rightarrow +\infty} \sigma^n = \sigma$  uniformly. For all  $n \in \mathbb{N}$ , one has existence of a strong solution  $(Y_t^{n, s, y, v}, V_t^{n, s, v}; s \leq t \leq T)$  to Equation (4.1) when one substitutes  $\sigma$  to  $\sigma^n$ . Then it is easy to check that  $(Y_t^{n, s, y, v}, V_t^{n, s, v}; s \leq t \leq T)$  converges in distribution to a weak solution of (4.1).

The uniqueness of the weak solution and the properties (i) to (iii) result from Theorem A.1: observe that  $G_{t, f}(s, y, v) = J_{0, t}(t-s, y, v)$ , where we have set  $\bar{a}(s, y, v) := a(t-s, y, v)$  in the definition of the operator  $\bar{\mathcal{L}}$ . Consequently, the density  $\Gamma(s, y, v; t, x, u)$  writes

$$\Gamma(s, y, v; t, x, u) = \bar{\Gamma}(0, x, u; t-s, y, v)$$

where  $\bar{\Gamma}(0, y, v; \theta, x, u)$  is the fundamental solution to  $\bar{\mathcal{L}}$ .

### A.3 Technical lemmas

For the reader's convenience we state three technical results which played a key role in our proofs.

The first lemma can be found in Stroock & Varadhan [16, Lemma 11.4.1].

**Lemma A.2.** *Let  $\{f_n; n \geq 1\}$  be a sequence of non-negative measurable functions such that  $\int_{\mathbb{R}^q} f_n(z) dz = 1$  and, for all  $h \in \mathbb{R}^q$ ,*

$$\limsup_{|h| \rightarrow 0} \sup_{n \geq 1} \int_{\mathbb{R}^q} |f_n(z+h) - f_n(z)| dz = 0.$$

*Suppose that there exists a density function  $f$  such that, for all function  $\psi \in \mathcal{C}_c(\mathbb{R}^q)$ ,*

$$\lim_{n \rightarrow +\infty} \int_{\mathbb{R}^q} f_n(z) \psi(z) dz = \int_{\mathbb{R}^q} f(z) \psi(z) dz.$$

*Then  $\{f_n\}$  converges to  $f$  in  $L^1(\mathbb{R}^q)$ .*

The next lemma can be found in Rana [15].

**Lemma A.3.** *Let  $1 \leq p < +\infty$ . For all  $f \in L^p(\mathbb{R}^q)$  and  $h \in \mathbb{R}^q$  we have*

$$\lim_{|h| \rightarrow 0} \int_{\mathbb{R}^q} |f(z+h) - f(z)|^p dz = 0. \quad (\text{A.2})$$

The last lemma can be found in Stroock & Varadhan [16, Cor.1.1.5].

**Lemma A.4.** *Let  $S$  be a Polish space and let  $\{F_{\epsilon'}, \epsilon' > 0\}$  be a uniformly bounded set of functions which are equicontinuous at each point of  $S$ . For all  $\{\mu_\epsilon; \epsilon > 0\}$  and  $\mu$  in  $\mathcal{M}(S)$  such that  $\lim_{\epsilon \rightarrow 0} \mu_\epsilon = \mu$  one has*

$$\limsup_{\epsilon \rightarrow 0} \sup_{\epsilon' > 0} \left| \int_S F_{\epsilon'} d\mu_\epsilon - \int_S F_{\epsilon'} d\mu \right| = 0.$$

## References

- [1] F. Bernardin, M. Bossy, C. Chauvin, P. Drobinski, A. Rousseau, and T. Salameh. Stochastic Downscaling Methods : Application to Wind Refinement. *Stochastic Environmental Research and Risk Assessment*, 2008. To appear.

- 
- [2] M. Bossy and J. Fontbona. Work in progress, 2008.
  - [3] M. Bossy and J.-F. Jabir. Confined Langevin systems for the mean no-permeability conditions, 2008. Submitted.
  - [4] A. Dermoune. Propagation and conditional propagation of chaos for pressureless gas equations. *Probab. Theory Related Fields*, 126(4):459–476, 2003.
  - [5] M. Di Francesco and A. Pascucci. On a class of degenerate parabolic equations of Kolmogorov type. *AMRX Appl. Math. Res. Express*, (3):77–116, 2005.
  - [6] M. Di Francesco and S. Polidoro. Schauder estimates, Harnack inequality and Gaussian lower bound for Kolmogorov-type operators in non-divergence form. *Adv. Differential Equations*, 11(11):1261–1320, 2006.
  - [7] T. D. Dreeben and S. B. Pope. Wall-function treatment in pdf methods for turbulent flows. *Physics of Fluids*, pages 2692–2703, 1997.
  - [8] J.F. Jabir. Modèles Stochastiques Lagrangiens de Type McKean–Vlasov Conditionnel et Leur Confinement, PhD thesis, University of Nice Sophia Antipolis, 2008.
  - [9] S. Méléard and S. Roelly-Coppoletta. A propagation of chaos result for a system of particles with moderate interaction. *Stochastic Process. Appl.*, 26(2):317–332, 1987.
  - [10] B. Mohammadi and O. Pironneau. *Analysis of the  $k-\varepsilon$  turbulence model*. Masson editions, 1995.
  - [11] K. Oelschläger. A law of large numbers for moderately interacting diffusion processes. *Z. Wahrsch. Verw. Gebiete*, 69(2):279–322, 1985.
  - [12] S. B. Pope. On the relationship between stochastic Lagrangian models of turbulence and second-moment closures. *Physics of Fluids*, pages 973–985, 1993.
  - [13] S. B. Pope. *Turbulent flows*. Cambridge Univ. Press, 2003.

- 
- [14] S. B. Pope and J. Xu. Assessment of numerical accuracy of pdf/Monte-Carlo methods for turbulent reacting flows. *Journal of computational physics*, pages 192–230, 1999.
  - [15] I. K. Rana. *An introduction to measure and integration*. American Mathematical Society, second edition, 2002.
  - [16] D. W. Stroock and S. R. S. Varadhan. *Multidimensional diffusion processes*. Springer-Verlag, 1979.
  - [17] A.-S. Sznitman. A propagation of chaos result for Burgers' equation. *Probab. Theory Relat. Fields*, pages 581–613, 1986.
  - [18] A.-S. Sznitman. Topics in propagation of chaos. In *École d'Été de Probabilités de Saint-Flour XIX—1989*, Lecture Notes in Math., pages 165–251. Springer, Berlin, 1991.



---

Unité de recherche INRIA Sophia Antipolis  
2004, route des Lucioles - BP 93 - 06902 Sophia Antipolis Cedex (France)

Unité de recherche INRIA Futurs : Parc Club Orsay Université - ZAC des Vignes  
4, rue Jacques Monod - 91893 ORSAY Cedex (France)

Unité de recherche INRIA Lorraine : LORIA, Technopôle de Nancy-Brabois - Campus scientifique  
615, rue du Jardin Botanique - BP 101 - 54602 Villers-lès-Nancy Cedex (France)

Unité de recherche INRIA Rennes : IRISA, Campus universitaire de Beaulieu - 35042 Rennes Cedex (France)

Unité de recherche INRIA Rhône-Alpes : 655, avenue de l'Europe - 38334 Montbonnot Saint-Ismier (France)

Unité de recherche INRIA Rocquencourt : Domaine de Voluceau - Rocquencourt - BP 105 - 78153 Le Chesnay Cedex (France)

---

Éditeur

INRIA - Domaine de Voluceau - Rocquencourt, BP 105 - 78153 Le Chesnay Cedex (France)

<http://www.inria.fr>

ISSN 0249-6399