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¹ Retrieval of Mars surface physical properties from ² OMEGA hyperspectral images using Regularized ³ Sliced Inverse Regression

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4 **Abstract.** Hyperspectral remote sensing, also known as imaging spec-
5 troscopy, is a promising space technology regularly selected by agencies with
6 regard to the exploration and observation of planets, to earth's geology or
7 to the monitoring of the environment. It allows to collect for each pixel of
8 a scene, the intensity of light energy reflected from planets as it varies across
9 different wavelengths. More than one hundred spectels in the visible and near
10 infra-red are typically recorded, making it possible to observe a continuous
11 spectrum for each image cell. Usually, in space exploration, the analysis of
12 these spectral signatures allows to retrieve the physical, chemical or miner-
13 alogical properties of surfaces and of atmospheres that may help to under-
14 stand the geological and climatological history of planets.

15 We propose in this paper a statistical method to evaluate the physical prop-
16 erties of surface materials on Mars from hyperspectral images collected by
17 the OMEGA instrument aboard the Mars express spacecraft. The approach
18 we develop is based on the estimation of the functional relationship F be-
19 tween some physical parameters and observed spectra. For this purpose, a
20 database of synthetic spectra is generated by a physical radiative transfer
21 model and used to estimate F . The high dimension of spectra is reduced by
22 using Gaussian regularized sliced inverse regression (GRSIR) to overcome
23 the curse of dimensionality and consequently the sensitivity of the inversion
24 to noise (ill-conditioned problems). Compared with a naive spectrum match-
25 ing approach such as the k -nearest neighbors algorithm, estimates are more
26 accurate and realistic.

1. Introduction

Visible and near infrared imaging spectroscopy is a key remote sensing technique to study and monitor planets. It allows the detection, mapping and characterization of minerals, as well as volatile species, that often constitute the first step toward the resolution of key climatic and geological issues [Murchie *et al.*, 2007; Bibring *et al.*, 2004a; Brown *et al.*, 2004; Carlson *et al.*, 1992]. These tasks are carried out by the spectral analysis of solar light reflected in different directions by the materials forming the top few millimeters or centimeters of the ground. Physical properties of the surface, such as chemical composition, granularity, texture, and physical state are some of the most important parameters that characterize the morphology of the hundred thousand spectra that typically constitute an image. Modeling the direct link between these parameters and observable spectra is usually called the forward problem in classical physics. It can be evaluated numerically by radiative transfer models, simulating the propagation of solar light through the atmosphere and surface to the sensor [Douté *et al.*, 2007; Shkuratov *et al.*, 1999; Douté and Schmitt, 1998; Hapke, 1993]. It allows us, for given values of the model parameters, to simulate the spectra that should be observed. Conversely, deducing the physical model parameters from the observed spectra is called an inverse problem. It generally cannot be solved analytically, and the use of optimization or statistical methods becomes necessary. From a general point of view, solving inverse problems requires that the fundamental physics are adequately understood, such that a function, G , may be specified between the spectra X and the parameters Y such that:

$$X = G(Y). \tag{1}$$

27 Once this relationship G is known, different methods can be used to deduce the parameters
28 Y from the observable X . They can be roughly divided in three categories (for further
29 details and comparisons, see [*Kimes et al.*, 2000; *Pragnère et al.*, 1999]):

30 • *Optimization algorithms* are the most traditional approaches. They consist of min-
31 imizing an objective function that expresses the similarity between an observed and a
32 simulated spectrum. For example, this function may be the mean square errors function,
33 where the errors are the differences between observed and simulated reflectances at each
34 wavelength. These methods involve numerical optimization techniques (Powell's method,
35 Simplex method, quasi-Newton method...) starting with an initial guess of the param-
36 eters and searching for the optimum parameters via an iterative process that minimizes
37 the objective function. In fact, estimations are unstable since inverse problems are often
38 ill-posed (solutions are not unique and a small change in the data can lead to enormous
39 differences in the estimations). A probabilistic formalism can then be used to regularize
40 inverse problems by introducing a prior distribution on model parameters [*Aster et al.*,
41 2005; *Tarantola*, 2005; *Mosegaard and Tarantola*, 2002]. These approaches are compu-
42 tationally expensive since they iteratively simulate new spectra. Therefore, they cannot
43 be applied in the case of the inversion of an image with several hundred thousand pix-
44 els. Moreover, they can sometimes fall into local minima if the objective function is not
45 convex.

46 • *Look-up table (LUT) / k-nearest neighbors approach (k-NN)*: In order to replace a
47 heavy runtime computation with a simpler look-up operation, a large database (look-up
48 table) is generated by radiative transfer for many parameter values and stored on disks.
49 To inverse a hyperspectral image, the spectrum at each pixel is then compared with the

50 look-up table spectra in order to find the best match (the nearest neighbor), according to
51 an objective function minimization. Parameters are then deduced from the look-up table
52 best match spectrum [*Carlson et al.*, 2005; *Philpot et al.*, 2004; *Combal et al.*, 2002; *Weiss*
53 *et al.*, 2000; *Douté et al.*, 2001]. In comparison with traditional optimization methods,
54 the speed gain is significant, since retrieving a value from memory is often faster than
55 undergoing an expensive computation. The main disadvantages of this approach are the
56 multiplicity of solutions and their instability (see section 3).

- *Training approaches* assume that there exists a functional relationship

$$Y = F(X) \tag{2}$$

57 between spectra and parameters. This relationship corresponds to the inverse of the
58 physical model G in the forward problem (1). The idea is to use a look-up table to
59 estimate the underlying mathematical relationship F . This relationship then allows us to
60 estimate the parameters of new spectra. Computation time can then be very competitive.
61 Neural networks [*Kamgar-Parsi and Gualtieri*, 1990] or support vector machines (SVM)
62 [*Scholkopf and Smola*, 2002] can be used to learn the function F . The advantage of a such
63 a training approach is that, once the relationship has been established, it can be used for
64 very large sets and for all new images with the same physical model.

65 Inverse problems occur in many branches of sciences such as geophysics, hydrology,
66 medical imaging, pedology, remote sensing, astronomy or planetology. The choice of
67 one methodology among the others is generally determined according to the degree of
68 knowledge of the studied physical phenomenon. It is also determined by the quality
69 and quantity of measurements. In geophysics or meteorology, probabilistic approaches
70 combined to optimization techniques are most of the time preferred for the following main

71 reasons: the inverse problem can be linearized, observations are numerous and physicists
72 do have an *a priori* about the parameters variations. In planetary sciences, massive inverse
73 problems have only been tackled recently since they intervene in the last step of a physical
74 study of the huge collections of hyperspectral images. Moreover, inverting hyperspectral
75 imaging on planets has to be designed with the following constraints:

76 • Dealing with extremely large datasets and various modeling require simple and fast
77 methodologies,

78 • Inverting high dimensional spectra comes up against the curse of dimensionality,

79 • Observed spectra always contain some noise.

80 In that context, LUT and training approaches are generally preferred since they are fast
81 and allow to deal with non linear inverse problems:

82 • The LUT/ k -NN methodology is currently used by physicists to study planetary bod-
83 ies [*Carlson et al.*, 2005; *Philpot et al.*, 2004; *Weiss et al.*, 2000; *Douté et al.*, 2001]. It
84 provides an instinctive, simple and fast way to retrieve the model parameters. A good
85 match between a simulated and an observed spectrum also reinforces the idea that the
86 chosen physical modeling is realistic. However, we will show in section 3 that LUT/ k -NN
87 approaches lead to very unstable estimates.

88 • Training approaches have been recently considered in the study of terrestrial vege-
89 tation. Among them, neural networks [*Combal et al.*, 2002; *Pragnère et al.*, 1999] seem
90 promising but the underlying learning process still remains time consuming. Since the
91 study of Mars is not focused on a particular type of soil, many physical models are ex-
92 plored, and a faster methodology is required. Support vector machines [*Durbha et al.*,

93 2007] could be used but their disadvantage is that they are difficult to interpret statisti-
94 cally and physically.

95 In this paper, we propose to develop a new training approach: the Gaussian regularized
96 sliced inverse regression (GRSIR). It is based on a dimension reduction technique: the
97 sliced inverse regression [Li, 1991]. To the best of our knowledge, this method has never
98 been used in the context of inverse problems in planetology. It has the advantage to be fast,
99 stable, statistically and physically interpretable and can also help to select an appropriate
100 look-up table for inversion. For the sake of validation, this approach is compared all along
101 with the naive but currently used k -NN approach. In section 2, we first introduce the
102 datasets we are working with. Then, the limits of the LUT/ k -NN approach are discussed
103 in section 3. The Gaussian regularized sliced inverse regression (GRSIR) [Bernard-Michel
104 *et al.*, 2007b] is introduced and detailed in section 4. The choice of an appropriate look-
105 up table is discussed in section 5. Two applications of GRSIR methodology are then
106 presented:

- 107 • A validation on simulations, in section 6, allows to quantify the relevance of GRSIR
108 according to two criteria,
- 109 • In section 7 the retrieval of some physical parameters for the south permanent polar
110 cap of Mars: proportions of CO₂ ice, dust, water ice and grain sizes of CO₂ ice and water
111 ice.

2. Data

112 The datasets used in the next sections can be divided in three categories: The hyper-
113 spectral images observed on Mars, a look-up table simulated by radiative transfer models

114 according to the physical modeling of these images and a test data in order to quantify
115 the estimates accuracy. In any case, these datasets consist of:

116 • Some spectra $X \in \mathbb{R}^d$, where d is the number of considered wavelengths and will be
117 called the dimension of the problem. In the presented application d is equal to 184 and
118 and consequently leads to the resolution of a high-dimensional problem.

119 • Some associated parameters $Y \in \mathbb{R}^p$, where p is the number of free parameters in
120 the physical model. In this paper, five parameters ($p = 5$) will be considered.

2.1. Hyperspectral Images from Mars

121 We concentrate on datasets collected by the spectro-imaging instrument: OMEGA
122 (observatoire pour la minéralogie, l'eau, la glace et l'activité). OMEGA is one of the
123 seven scientific instruments aboard the European spacecraft Mars express mission, sent
124 in orbit around Mars in 2003. It has been developed by IAS and LESIA (Observatoire de
125 Paris) with the support of CNES, and a participation of IFSI (Italy) and IKI (Russia).
126 This visible and infrared mapping spectrometer should scan most of Mars from orbits in
127 order to observe the gases and dust in the atmosphere and to look for signs of specific
128 materials such as silicates, carbonates, ices at the surface. It records the visible and
129 infrared light reflected from the planet in the wavelength range 0.5-5.2 microns and with
130 a ground resolution varying from 350 m to 10 km. Our study will refer to three OMEGA
131 hyperspectral images acquired during orbits 41, 61 and 103. They cover the high southern
132 latitudes of Mars. The spatial resolution is about 2km per pixel and we considered 184
133 wavelengths in the range 0.95-4.15. For each image, a preprocessing aiming at correcting
134 the atmospheric gaseous contribution in the spectra has been applied [Douté *et al.*, 2007].
135 After a first analysis, these OMEGA observations revealed [Bibring *et al.*, 2004b] that

136 the south polar region of Mars mainly contains water ice, carbon dioxide ice and dust. A
137 detailed qualitative mapping of H₂O and CO₂ ices during the local summer shows that
138 the permanent south polar region is dominated by superficial CO₂ ice on the bright cap
139 except at its edges where water ice appears in extended areas. Examining the coexistence
140 modes (geographical or granular) between H₂O, CO₂ ice and dust that best explain the
141 morphology of the spectra has then led to the implementation of a physical modeling
142 of individual spectra with a surface reflectance model [Douté *et al.*, 2007]. This model
143 allows the generation of synthetic spectra with the corresponding pairs of parameters that
144 constitute a look-table [Douté *et al.*, 2007]. We will focus on the terrain unit of strong
145 concentration of CO₂: the bright permanent south polar cap. This unit has been carefully
146 mapped by a classification method based on wavanglets [Schmidt *et al.*, 2007]. For each
147 image, the CO₂ ice areas contain about 10000 to 20000 spectra.

2.2. Look-up table

148 Because we do not have any *a priori* about the range of values in which each parameter
149 is varying, we chose to simulate a large look-up table, containing parameters values that
150 are probably not found in our three considered images. The look-up table has been
151 simulated for five parameters varying spatially: the grain sizes of water and CO₂ ices, the
152 proportions of water ice, CO₂ ice and dust. The sum of the proportions is constrained to
153 one. The other parameters intervening in the physical model, such as the incidence and
154 emergence angles or the grain size of dust have been fixed to a constant. Details about
155 the sampling strategy are given in Table 1. We will denote by:

- 156 • n the number of simulated spectra ($n = 31500$),
- 157 • $x_i \in \mathbb{R}^d, i \in \{1, \dots, n\}$ the spectra,

158 • $y_i \in \mathbb{R}^p, i \in \{1, \dots, n\}$ the p associated parameters.

2.3. Test data

159 For the validation, the use of a test dataset is required that we will denote by "Tdata".
160 The one we chose has been simulated by radiative transfer modeling for random values of
161 the proportions of dust and CO₂ ice and the grain sizes of CO₂ and water ices. They have
162 been simulated according to the same physical model than the one used for generating
163 the look-up table. Details about the ranges of variation can be found in Table 1. In
164 order to work in realistic conditions, a multiGaussian noise of dimension 184 has been
165 added to all the spectra of the test dataset. The noise has been simulated with a mean
166 fixed to zero for all wavelengths and with a covariance matrix determined experimentally
167 from the OMEGA image acquired during orbit 41. A small portion of the image, very
168 homogeneous in terms of composition and physical properties, is chosen so that we can
169 assume that much of the variability comes from the noise. The latter is then evaluated
170 by a statistics after applying a shift difference on the selected portion. Afterward, we will
171 denote:

172 • n_T the number of spectra from Tdata ($n_T = 3500$),

173 • $x_i^T \in \mathbb{R}^d, i \in \{1, \dots, n_T\}$ the spectra from Tdata,

174 • $y_i^T \in \mathbb{R}^p, i \in \{1, \dots, n_T\}$ the associated parameters.

3. Limits of the k -NN approach

Let $o \in \mathbb{R}^d$ be the spectral reflectance of one observed pixel on Mars.

The nearest neighbor algorithm (also known as the k -NN algorithm) consists of searching the k nearest spectra of o in the learning database minimizing the mean square errors

between the observed spectrum o and the simulated spectra. The k nearest spectra of the learning database are selected by sorting in ascending order the values of the objective function $\Phi(i)$ defined by:

$$\Phi(i) = \|o - x_i\|^2. \quad (3)$$

175 where $i \in \{1, \dots, n\}$.

176 If $k = 1$, the estimated parameters for the pixel are then the ones associated to the nearest
177 spectrum selected in the learning database. If k is greater than 1, then one can choose to
178 estimate parameters by the mode, or the mean of the k best matches, but also to keep all
179 the possible estimates. Most of the time, only the best match is retained.

180

181 The problem of k -NN approach is that it usually leads to very unstable estimates and
182 it is then difficult to choose a judicious k . Let us consider the example of two spectra
183 S1 and S2 for which parameters are known and let us add a reasonable noise to these
184 spectra. Applying the k -NN methodology ($k = 1$) leads to estimations given in Table 2.
185 We can notice here that estimation errors are relatively small for spectrum S1 whereas
186 they are much greater for spectrum S2. One first thinks that a good idea to improve
187 results would be to consider more than 1 neighbor. In fact, figure 1 shows how unstable
188 the estimates are. It presents the relative estimation error as a function of k , where k is
189 the k th neighbor. First, we can see that considering the k th neighbor or the $(k + 1)$ th
190 neighbor generally leads to very different estimations. Moreover, for spectrum S1, errors
191 are increasing when k is increasing whereas for spectrum S2 errors are decreasing when k
192 is increasing. That means that for spectrum S1, very few neighbors would be necessary
193 to estimate the parameters properly whereas with spectrum S2, a lot of neighbors would

194 be required. In fact, for spectrum S2, more than one hundred neighbors are necessary
195 to have at least a realistic estimation of the grain size of CO₂ ice. In these conditions,
196 one can see that estimating parameters by a k -NN approach can lead to significant errors
197 especially when there is no *a priori* on the solutions.

4. Proposed approach

198 We propose here to estimate the functional relationship F between the spectra $X \in \mathbb{R}^d$
199 and each physical parameter $Y^j = F(X) \in \mathbb{R}$ with $j \in \{1, \dots, p\}$. Since each parameter
200 is studied individually, we will consider in the next sections that $p = 1$ and we will always
201 omit j for the sake of simplicity. Because of the curse of dimensionality, estimating a
202 d -variate function from the data is not realistic and dimension reduction techniques have
203 to be applied. They rely on the assumption that the predictor X can be replaced by its
204 projection onto a subspace of smaller dimension L , called the effective dimension reduction
205 space (EDR), without loss of information. Denoting by β_1, \dots, β_L the basis of this sub-
206 space, the functional relationship $Y = F(X)$ can be rewritten as $Y = f(\beta_1^t X, \dots, \beta_L^t X)$
207 where f is now a L -variate function. In most applications, L is smaller than 3 making the
208 estimation of f realistic. The estimation of the projection axes β_1, \dots, β_L is discussed in
209 the next paragraph.

4.1. Dimension reduction step: Gaussian regularized sliced inverse regression

Principal component analysis (PCA) is the most current technique that could be used
to choose the β_i 's [Bishop, 2006, chapter 12]. The basic idea in PCA is to define the
orthogonal projection of the spectra onto a lower dimensional linear space such that the
total variance of the projected spectra is maximized. It amounts to diagonalizing the

spectra covariance matrix

$$\Sigma = \frac{1}{n} \sum_{i=1}^n (x_i - \bar{x})(x_i - \bar{x})^t \text{ where } \bar{x} = \frac{1}{n} \sum_{i=1}^n x_i \quad (4)$$

210 or equivalently to calculating the eigenvectors of Σ . However, in the case of a regres-
211 sion problem, PCA is generally not satisfying since only the explanatory variables X are
212 considered and the dependent variable Y is not taken into account. Specific dimension
213 reduction techniques have been developed for regression problems, and among them Sliced
214 Inverse Regression (SIR) is the most popular [Li, 1991]. It consists first of sorting the
215 parameter values in increasing order and to divide them into a chosen number H of non
216 overlapping slices $S_h, h \in \{1, \dots, H\}$. For each slice S_h , a mean spectrum m_h is calcu-
217 lated by averaging the spectra for which the parameter values are varying in the slice (see
218 equation (5)). The basic idea in SIR is then to define the orthogonal projection of the
219 mean spectra onto a lower dimensional linear space such that the variance of the projected
220 mean spectra is maximized under the constraint that the total projected variance is 1. It
221 amounts to diagonalizing $\Sigma^{-1}\Gamma$ where Γ is the mean spectra covariance matrix given in
222 equation (6). This principle is illustrated on figure 2 showing that GRSIR consists of max-
223 imizing the between-slice variance of the reduced spectra or equivalently in minimizing
224 the within-slice variance.

Since inverse problems are generally ill-posed [Aster *et al.*, 2005; Tarantola, 2005], Σ is ill-conditioned making its inversion difficult. We thus propose to compute a Gaussian regularized version of Sliced Inverse Regression (GRSIR). Theoretical details can be found in [Bernard-Michel *et al.*, 2007a]. The concept of this method is to incorporate some prior information on the projections in order to dampen the effect of the noise in the input data, and to make the solution more regular or smooth. The ill-posed problem is then

replaced by a slightly perturbed well-posed problem that depends on a parameter δ , called the regularization parameter. Similarly, GRSIR consists of computing the L eigenvectors corresponding to the L largest eigenvalues of $(\Sigma^2 + \delta I_p)^{-1} \Sigma \Gamma$ where I_p is the $d \times d$ identity matrix. In practice, GRSIR then requires the four following computational steps:

Step 1: Sorting $y_i, i \in \{1, \dots, n\}$ in increasing order and dividing them into H non overlapping slices $S_h, h = 1, \dots, H$. If the Look-up table has been simulated for random values of Y , then the slices are chosen such that each slice contains the same number of observations. If the look-up table has been simulated for a fixed number of distinct parameters, slices are these distinct values.

Step 2: Computing the slice means. For each slice:

$$m_h = \frac{1}{n_h} \sum_{y_i \in S_h} x_i \quad (5)$$

where n_h denotes the number of observations in the slice S_h .

Step 3: Computing the "between slices" means covariance matrix:

$$\Gamma = \sum_{h=1}^H \frac{n_h}{n} (m_h - \bar{x})(m_h - \bar{x})^t \quad (6)$$

Step 4: Estimating the GRSIR axes $\beta_1(\delta), \dots, \beta_L(\delta)$ by computing the eigenvectors of $(\Sigma^2 + \delta I_p)^{-1} \Sigma \Gamma$. The first GRSIR axis $\beta_1(\delta)$ is determined by the eigenvector corresponding to the largest eigenvalue, the second GRSIR axis $\beta_2(\delta)$ is determined by the eigenvector corresponding to the second largest eigenvalue and so on.

In section 6.3, we will give a criterion to choose the dimension L . We will see that in practical cases, only the first GRSIR axis is sufficient to retrieve parameters. Let us notice that this first axis can be viewed as a weight function of the wavelengths, giving

233 some important weights to the wavelengths that carry much more information about the
 234 parameter. Afterward, we will denote by reduced spectra the projections of the spectra
 235 on this axis $\beta_1(\delta)$.

4.2. Estimation of the functional relationship

236 Once a relationship has been revealed between reduced spectra and the parameter (see
 237 for example figure 2), then the question comes to estimate this relation f . We propose to
 238 use a piecewise linear interpolation on the set of data points $(m_h^{proj}, m_h^{param}), h = 1, \dots, H$
 239 where H denotes the number of slices S_h , $m_h^{proj} = \langle x_i, \beta_1(\delta) \rangle$ denotes the average of the
 240 projection of the spectra $x_i, y_i \in S_h$ for slice S_h and $m_h^{param} = \frac{1}{n_h} \sum_{y_i \in S_h} y_i$ denotes the
 241 average parameter value for slice S_h .

242 For each new spectrum x with a projection $t = \langle \beta_1(\delta), x \rangle$, the estimated parameter
 243 value \hat{y} is then given by:

$$\hat{y} = \begin{cases} m_1^{param} & \text{if } t \in] - \infty, m_1^{proj}] \\ \sum_{h=1}^{H-1} \left[m_h^{param} + (t - m_h^{proj}) \left(\frac{m_{h+1}^{param} - m_h^{param}}{m_{h+1}^{proj} - m_h^{proj}} \right) \right] & \text{if } t \in] m_h^{proj}, m_{h+1}^{proj}] \\ m_H^{param} & \text{if } t \in] m_H^{proj}, +\infty [\end{cases}$$

244 An example of the application of GRSIR to the look-up table is given in Figure 2. It
 245 shows the relationship between reduced spectra and the grain size of CO₂ ice and its
 246 estimation by a linear interpolation.

5. Choice of the Look-up table

247 In the simulated look-up table, we can imagine that many spectra are not necessary
 248 to estimate properly the parameters of an observed image. For example, if the studied
 249 image is the test data Tdata, we should remove from the look-up table all the spectra
 250 with parameters values that are not included in the range of values of Tdata's parameters.

251 We would then obtain a new look-up table that leads to better estimations for both k -
252 NN and GRSIR [*Bernard-Michel et al.*, 2007a]. The difficulty is that in practical cases,
253 we do not have any *a priori* about the range of variations of the parameters. The best
254 approach in that case is then to take a huge look-up table and to reduce it after the first
255 estimations by GRSIR or k -NN. This approach suffers from the fact that if the database
256 is too important, then estimations are strongly deteriorated. In the case of k -NN, these
257 deteriorations are due to the similarity of some spectra with very different parameters
258 that lead to multiple solutions. In the case of GRSIR, these deteriorations are mainly
259 due to the spectra that are on the opposite too different from the observed one. In
260 order to visualize the adequacy of the chosen look-up table with the observed spectra,
261 we propose to use principal component analysis. We consider that we can retrieve the
262 parameters of the observed spectra from a look-up table only if their projections on the
263 m first PCA axes (deduced from the application of PCA to the look-up table) coincide
264 with the projection of the look-up table itself. For the sake of visibility and because in
265 our applications two components explain until 88% of the total variance, we chose $m = 2$.
266 In some cases, a greater value for m could be required. To select the most appropriate
267 spectra for inversion in the look-up table, our idea is to retain the spectra from the look-
268 up table whose projections on PCA axes are close to projections of the spectra from
269 the observed image. In the plane spanned by the m first PCA axes, we calculate the
270 distance of each projected spectrum from the look-up table with its nearest neighbor in
271 the observed image. The histogram of these distances allows to distinguish a mixture of
272 n_{class} Gaussian densities $\mathcal{N}(\mu_i, \Sigma_i)$ characterized by parameters such as the proportions
273 of the mixture $\pi_i, i \in \{1, \dots, n_{class}\}$, the expectation μ_i and the covariance matrix Σ_i

274 of each class $i, i \in \{1, \dots, nclass\}$. The number of classes is chosen by the user and
275 the parameters can be estimated by the Expectation-Maximization algorithm (EM) after
276 being initialized randomly or by the K-means algorithm (for further details about mixture
277 models and EM, see *Bishop* [see 2006, chapter 9]). Roughly summarizing, EM algorithm
278 is a simple iterative optimization process for computing the maximum likelihood estimates
279 of the parameters. One can then calculate and maximize the posterior probability that a
280 spectrum belongs to a class observing the distances. This leads to a classification of the
281 spectra into $nclass$ classes, generally two classes: the class of spectra from the look-up
282 table that are far from the observed spectra, and conversely the class of spectra that are
283 close. A statistics calculated separately for the spectra belonging to each class allows to
284 refine the estimation of π_i, μ_i , and $\Sigma_i, i \in \{1, \dots, nclass\}$. An example is given in Figure
285 3 with the image observed from orbit 41. The histogram allows to distinguish 3 classes
286 from which the third class corresponds quite well to the class we are looking for.

287 We can also notice that some of the observed spectra are outside the projected look-up
288 table. We can wonder if we should keep these spectra for the inversion step, because either
289 the simulated look-up table is too small, or the chosen physical model is not relevant for
290 them. We propose to select and remove these spectra using the previous methodology. In
291 the PCA plane, the histogram of the distance of each observed spectrum with its nearest
292 neighbor in the selected sub-look-up table allows to distinguish two classes: the invertible
293 spectra and the non invertible spectra (figure 3). The non invertible spectra correspond
294 to pixels at the boundary of the CO₂ bright area and we can then suppose that another
295 physical model should be used to inverse them. Other examples can be found in the work

296 of *Bernard-Michel et al.* [2007a]. It is also shown that, on simulations, estimations are
297 more accurate when applying such a selection.

6. Analyzing results on simulations

298 In this section, we propose to quantify the accuracy of GRSIR estimations by using
299 the test data presented in section 2. We also show how to choose the regularization
300 parameter and the EDR dimension. At last, results are presented comparing GRSIR and
301 k -NN methods.

6.1. Validation criteria

302 To assess the relevance of GRSIR methodology, we focus on two aspects of the valida-
303 tion: the accuracy of the estimates and the quality of the relationship between reduced
304 spectra and parameters. To this end, we define two validation criteria. The first one,
305 denoted Normalized Root Mean Square Errors (NRMSE), quantifies the importance of
306 estimation errors, that are the differences between the estimations \hat{y}_i^T and the real values
307 y_i^T . The second one, denoted SIR Criterion (SIRC) is the ratio between the "between-
308 slices" variance of GRSIR projections and the total variance. It quantifies the quality
309 of the relationship between projected spectra and parameters. A third validation crite-
310 rion that could have been introduced is the RMSE between the test spectrum x_i^T and
311 the spectrum that can be reconstructed by running the radiative transfer model with the
312 estimated parameters. However, this criterion would have faced the ill-posed problem
313 that two spectra can be very close in terms of RMSE even with very different parame-
314 ters values. The consequences are that a good reconstruction does not always mean that
315 parameters are well estimated.

The Normalized Root Mean Square Errors (NRMSE) criterion is defined as:

$$NRMSE = \sqrt{\frac{\sum_{i=1}^{n_T} (\hat{y}_i^T - y_i^T)^2}{\sum_{i=1}^{n_T} (y_i^T - \bar{y}^T)^2}} \text{ with } \bar{y}^T = \frac{1}{n_T} \sum_{i=1}^{n_T} y_i^T.$$

317 This measure is invariant with respect to position and scale and therefore it is possible to
 318 deduce if one parameter is better estimated than another. The NRMSE is close to zero
 319 when the predicted values are accurate and conversely becomes larger when predictions
 320 are poor.

321 **6.1.2. the SIRC**

The SIR criterion for an axis $\beta_\ell(\delta)$ is defined as the ratio between the "between-slices"
 variance $\beta_\ell^t(\delta)\Gamma\beta_\ell(\delta)$ of the projections of $x_i, i \in \{1, \dots, n\}$ on $\beta_\ell(\delta)$ and the total variance
 $\beta_\ell^t(\delta)\Sigma\beta_\ell(\delta)$ of these same projections:

$$SIRC = \frac{\beta_\ell^t(\delta)\Gamma\beta_\ell(\delta)}{\beta_\ell^t(\delta)\Sigma\beta_\ell(\delta)} \quad (7)$$

322 It indicates the quality of the functional relationship between projections of the spectra
 323 on the ℓ th GRSIR axis and the parameters. In the SIR criterion, the total variance can
 324 be interpreted as the sum of the between slices variance and the within slice variance. In
 325 this case, one can see the quality of the relationship is perfect when the within variance
 326 is close to zero, or equivalently when SIRC is close to 1. Finally, the closer SIRC is to 1,
 327 the better the relationship is.

6.2. Choice of the regularization parameter

328 In the case of Regularized Sliced Inverse Regression, estimations \hat{y}_i^T and consequently
 329 NRMSE criterion as well as the SIRC score depend on a regularization parameter δ .

330 When the regularization parameter increases, the functional relationship between pro-
331 jected spectra and parameters is getting worse and consequently estimation errors are
332 increasing. From another point of view, if the regularization parameter is too small, then
333 estimation errors are huge in presence of noise because we are dealing with an ill-posed
334 problem. So, there is a compromise to reach between deteriorating the functional rela-
335 tionship and improving estimations by regularizing. In fact, this compromise lies in the
336 choice of the regularization parameter. We propose here to choose this regularization
337 parameter, for each parameter individually, minimizing the NRMSE criterion calculated
338 between the parameter values from a test data and their estimations. The test data can
339 be the look-up table itself spoiled by a multiGaussian noise supposed representative of the
340 one affecting the OMEGA hyperspectral images. Let us notice that if there is no noise
341 in the data, or in other words, if the observed data exactly corresponds to spectra that
342 could be simulated by radiative transfer model, then no regularization is required and
343 minimizing the NRMSE criterion for GRSIR would yield a value δ close to zero. Figure
344 4 shows how the choice of the regularization parameter changes according to the noise:

- 345 • the stronger the noise is, the greater the NRMSE are,
- 346 • when the noise is increasing, the chosen regularization parameter, that corresponds
347 to the minimum of the NRMSE, is also increasing.

348 One can easily see that making a mistake on the estimation of the noise leads to a
349 mistake in the estimation of the regularization parameter and consequently can lead to
350 much more uncertain estimations.

6.3. Choice of the EDR dimension

351 In order to choose the EDR dimension, we propose to calculate the SIRC for each GRSIR
352 axis $\beta_\ell(\delta)$ with $\ell \in \{1, \dots, d\}$. As ℓ is increasing, the SIRC decreases very rapidly and is
353 almost zero after very few iterations of the index ℓ . In the next sections, we will fix the
354 dimension of the EDR to the index L from which the SIRC is significantly decreasing. For
355 example, If we apply GRSIR to the look-up table, the SIRC calculated for the proportion
356 of CO₂ ice on the first GRSIR axis is 0.985. This SIRC, very close to one, indicates that
357 the relationship of interest is well revealed by the first axis. Calculated with the second
358 axis, the SIRC drops to 0.1799 and then is close to zero for other axes. In this case,
359 we will consider that taking into account these axes in the regression step do not bring
360 sufficient information and is not necessary, and will fix the EDR dimension to one. In
361 fact, analyzing the SIRC for all parameters and for two different physical models reveals
362 that only one dimension is sufficient to retrieve the model parameters. We found only one
363 situation for which two GRSIR axes have to be kept: the thickness of the compact slab of
364 nearly pure CO₂ ice making the seasonal deposits of the south pole of Mars during spring
365 [*Douté et al.*, 2008]. For this parameter, the SIRC is 0.89 for the first axis, 0.81 for the
366 second axis and close to zero for the third one. Because the necessity to keep more than
367 one GRSIR axis seems rare, L is set to one in the next sections.

6.4. Analyzing the GRSIR axis

368 The spectral variability of the individuals x_i constituting the look-up table results from
369 a complex interplay of the different input parameters y_i of the model. In particular two
370 parameters can have a similar influence on the spectra for a certain range of values. As a
371 consequence it is crucial to understand how the GRSIR method can untangle the depen-
372 dencies in order to find independent one to one functional relationships between spectra

373 and parameters. We thus try to correlate each axis $\beta_1(\delta)$ - which is a vector of spectral
374 weights - and the variability specifically induced by the variation of the corresponding
375 physical parameter. The other parameter values are kept constant at mean values. The
376 weights are individually applied to very different level of reflectance depending on wave-
377 length when calculating the reduced spectrum $\beta_1^t x_i$. Hence we find it more illustrative
378 to represent the componentwise multiplication (wavelength by wavelength) of each axis
379 with a representative spectrum. An illustration is given in figure 5 for the grain size of
380 CO₂ ice. Additionally we overplot in figure 6 the absolute value (for the sake of clarity) of
381 the proportions of water ice and dust and of the grain sizes of water and CO₂ ice. From
382 our analysis we make the following observations: (i) The criteria applied when seeking an
383 axis (maximizing the between-slice variance while minimizing the within-slice variance)
384 implies that GRSIR does put weight on key spectral points that do not display necessary
385 the highest variability when we change the value of the corresponding parameter. (ii) In
386 the case of the physical model that we consider in this paper (intimate mixture of H₂O,
387 CO₂, and dust), the spectral range from 0.92 to 2.75 microns is the most determinant
388 for evaluating the parameters. (iii) The axis for H₂O abundance, dust abundance, and
389 grain size of CO₂ ice are quite similar: we have a strong weight put on wavelength 1.43
390 micron (the bottom of a very reliable and narrow CO₂ ice absorption band) and, to a
391 lesser extent, on the interval between 1.47 and 1.54 microns (H₂O ice feature), and on
392 wavelengths 1.87 and 2.29 microns (weak CO₂ ice bands but very distinct). Besides the
393 three axes have in common a moderate value at 2.00 microns, at the bottom of a strong
394 CO₂ ice band. The latter is not very diagnostic because it often saturates. Thus, it is
395 satisfying that GRSIR does not put too much emphasis on it. The three axes also shows

396 differences that allow differentiation between the functional relationships. A relatively
397 strong weight at 2.38 microns on the right wing of the 2.34 micron band is a specificity of
398 the proportion of water ice. We can note in figure 5 that the reflectance at this wavelength
399 is greatly variable when water ice abundance or CO₂ ice grain size varies. The axis for the
400 proportion of dust stresses very much the variation of the spectrum continuum around 1
401 micron and 1.77 micron that indeed strongly changes with dust concentration. No other
402 parameter affects so much the continuum. The axis for the proportion of dust minimizes
403 the influence of the small plateau at 2.62 microns contrary to the axis for the grain size
404 of CO₂ ice. This is the solution GRSIR finds to differentiate axes of the proportion of
405 dust and the grain size of CO₂ ice because the plateau level varies very much with both
406 dust concentration and CO₂ ice grain size. Furthermore the axis for grain sizes of CO₂
407 ice induces less sensitivity to the spectrum continuum than the axis for the proportion
408 of dust. In addition, it does not put a noticeable weight on the right wing of the 2.34
409 micron band contrary to the axis for the proportion of water ice. (iv) Axis 5 which is
410 linked to H₂O grain size is very different from the others since it emphasizes, on the one
411 hand, the continuum level around 1 micron, but not around 1.75 microns and, on the
412 other hand, it emphasizes the whole spectral range between 2.35 and 2.65 microns. The
413 principal conclusions we can draw from our study is that, at least for the model presented
414 in this paper, the GRSIR method succeeds in decorrelating the evaluation of the differ-
415 ent parameters. Indeed it has the ability to find a unique set of wavelengths where the
416 variability of the spectrum is the most pertinent (but not necessary the highest) for the
417 evaluation of a given parameter. One must note that these key wavelengths not only fall
418 on the bottom of specific absorption bands for both H₂O and CO₂ ices but also on specific

419 parts of the continuum as well as on specific band wings. The GRSIR method thus reveals
420 and uses all the pertinent information contained in the spectra that allows to estimate
421 the parameters. This information is more complete than if we were focusing only on band
422 depths as one could expect in the first place. Finally, knowing at the key wavelengths
423 the characteristics (signal to noise ratio, accuracy of the absolute radiometric calibration,
424 etc.) of the instrument that acquires the spectra (the imaging spectrometer OMEGA in
425 our case) is crucial to assess the quality of the solutions.

6.5. Results

426 In this section, we finally compare the naive k -NN approach with the GRSIR method-
427 ology on the test data Tdata. We quantify the gain of one approach from the other
428 according to the SIRC and NRMSE criteria but we also compare estimations to real val-
429 ues with the help of graphic tools such as scatter plots or histograms. For both methods,
430 the look-up table defined in section 2.2 has been used except that a selection of the spectra
431 described in section 5 has been performed before applying the GRSIR methodology. It
432 has no interest to apply this selection to the k -NN methodology because in this selection,
433 most of the spectra leading to multiple solutions for parameters are kept. The results,
434 presented in Table 3, show that for most parameters, estimations are much more accurate
435 with GRSIR. Only one parameter, the grain size of water ice, is better estimated on av-
436 erage by k -NN. This may be due to the fact that too few values have been considered for
437 simulations in the look-up table (only 5). In all other cases, we can see an important gain
438 on the NRMSE which can be 3 times smaller with GRSIR than with k -NN. An illustra-
439 tion of this gain is highlighted in Figure 7 which shows the scatter plot of the estimated
440 proportions of CO₂ ice for both methods versus the real proportions. In this graphic, we

441 can notice that the range of variation of the estimated proportions of CO₂ ice is much
442 larger with k -NN than with GRSIR. In fact, we can easily see, plotting the associated
443 densities, that GRSIR reproduces better the density of the original proportions of CO₂ ice
444 simulated in Tdata than k -NN (see figure 8). It is interesting to point out that k -NN does
445 not give any information about the range of variation of the parameters of the test data
446 and that estimations often vary in the same ranges of variation that the simulated values
447 from the look-up table. On the other hand, GRSIR, as a first step, allows to evaluate in
448 which ranges the parameters of an observed image are varying. Then a more appropriate
449 look-up table can be built in the estimated ranges with a higher density of parameter val-
450 ues. Applying GRSIR with this new extracted look-up table leads to even more accurate
451 estimations given in Table 4. Finally, the SIRC, always close to 1, is very satisfying.

7. Retrieval of the physical parameters for the south polar cap of Mars

452 In this section, we present the first model inversions obtained by k -NN and GRSIR for
453 hyperspectral images acquired by OMEGA during orbits 61, 103 and 41 (see section 2).
454 The study of these images by k -NN and GRSIR leads to 30 different maps that we will not
455 present in this paper but that can be found in the work of *Bernard-Michel et al.* [2007a].
456 We draw up here the main conclusions.

457 The model inversion on the image observed from orbit 103 by GRSIR and k -NN shows
458 that GRSIR gives very smooth mappings for all sets of parameters. Moreover, GRSIR
459 estimations vary continuously while k -NN estimations assume a small number of values.
460 An example is given in Figure 9, where we see that the proportion of dust is nearly al-
461 ways estimated to a constant equal to 0.0003 with k -NN. With GRSIR, the map is more
462 detailed. In the same way, figure 10 shows the proportion of dust estimated by k -NN and

463 GRSIR on the portion of the polar cap observed during orbit 41. With k -NN, estimations
464 assume now 8 different values and seem to indicate that, at the very center of the bright
465 cap, very few dust is observed. The abundance of dust increases significantly when getting
466 closer from the boundaries. The estimated map with GRSIR is more detailed and leads
467 to slightly different conclusions. In particular, the area presenting a poor proportion of
468 dust is more extended than with k -NN. Nevertheless if, globally, maps are much smoother
469 and detailed with GRSIR, they never differ entirely from k -NN's.

470 Another interesting remark concerns the estimation of the grain size of CO₂ ice in images
471 61 and 103. In fact, images 61 and 103 represent approximately the same portion of the
472 polar cap and should consequently give close estimations for each studied parameter. But
473 we observed that estimations of the grain size of CO₂ ice are strongly different for both
474 images and more especially with k -NN methodology (see figure 11). Observations 61 and
475 103 were acquired approximately 12 Mars days apart during a period of the year when
476 the solar illumination was declining over the south pole. Consequently, neither the state
477 of the atmosphere (e.g. the dust content) nor the acquisition geometry by the sensor were
478 strictly equivalent but likely comparable. These differences translate to small changes
479 on the overall level of reflectance of the spectra and on the CO₂ ice bands intensity and
480 shape. This is very much comparable to the introduction of noise in a relative sense be-
481 tween observations 61 and 103. Since our remote sensing problem is partly ill-conditioned,
482 a non regularized inversion method such as the k -NN will be much more sensitive to this
483 noise than a regularized one such as the GRSIR. This is exactly what one can notice.
484 The discrepancy that remains between the densities of the estimated grain sizes for orbit
485 61 and 103 is mainly due to the conjugated effect of an increasingly grazing illumination

486 over a slightly dusty atmosphere that diminishes the spectral contrast of the observations.
487 Such an effect was not taken into account in the modeling.

488 Finally, for the sake of validation, we compared estimated maps with the Wavanglet
489 approach developed in [*Schmidt et al.*, 2007]. "Wavanglet" is a supervised automatic de-
490 tection method that identifies in hyperspectral images spectral features and thus produces
491 distribution maps of chemical compounds. It uses three steps: A. Selection of a library
492 composed of reference spectra (the signature of the compounds to be detected); B. Appli-
493 cation of a Daubechies wavelet transform to referenced spectra and determination of the
494 wavelet subspace that best separates all referenced spectra; C. In this selected subspace,
495 calculation of the spectral angle between each spectrum of an observation and a given ref-
496 erence spectrum. In particular this angle, that we will denote the wavanglet angle, allows
497 to quantify in a relative sense the spatial variations of the different compound abundances
498 at the surface. We present in figure 12 the cosine of the wavanglet angle that makes each
499 spectrum of the images 41 and 103 with a reference spectrum of martian dust. The closer
500 it is from one, the greater the dust proportion is. We have a similar map for the water
501 proportion. Globally, estimates of the latter quantities are more noisy with wavanglet
502 than with GRSIR and k-NN, especially for the dust proportion but they generally are in
503 agreement with some exceptions. For example, in the image observed during orbit 103,
504 both *k*-NN and wavanglet methodologies display an area with strong proportion of dust
505 in the lower right part of figure 12(B), that is absent with GRSIR. In order to see if
506 this area really contains more dust than other areas, we selected the spectrum (denoted
507 by A in figure 12) corresponding to the greatest proportion of dust. We compared it to
508 two other spectra from the cap. The first spectrum (B) has been chosen in a pure CO₂

509 area with very few dust. The second one (C) has been chosen in a area containing dust.
510 According to k -NN and wavglet results, this area should however contain less dust than
511 the area where we extracted spectrum A. The spectra A, B and C are presented in figure
512 13. We can see that as, expected, the spectrum B is really different from the spectrum A.
513 On the other hand, spectra A and C are really similar showing that the area of interest
514 does not contain as much dust as the wavglet methodology predicts. On the opposite,
515 estimations given by GRSIR are more coherent with this spectral analysis.

8. Conclusion

516 In this paper, we propose a regularized version of Sliced Inverse Regression in order
517 to retrieve the physical parameters that best explain the spectra observed on Mars by
518 the OMEGA imaging-spectrometer. To the best of our knowledge, this methodology
519 has never been used in the domain of remote sensing and more particularly in planetary
520 sciences. Results on simulations are promising showing that estimations are accurate
521 and most of the time better than the ones given by the k -nearest neighbors algorithm.
522 On a real data, maps are much smoother than with k -NN and seem to give a coherent
523 mapping if we compare the inversion of different hyperspectral images of the same portion
524 of surface of Mars. Moreover, GRSIR is a fast algorithm that calculates once and for all
525 the relationship between spectra and parameters for a determined physical model. Thus,
526 it is then really easy to inverse each new observed spectrum if it is coherent with the model.
527 The main limit of GRSIR approach is that we currently do not give any uncertainties of our
528 estimations when inverting a real image. We could calculate experimental uncertainties
529 based on simulations, but it supposes that the noise in the spectra has been well evaluated.
530 If not, uncertainties will probably be underestimated. Some improvements could also be

531 proposed to choose the regularization parameter and a more complete analysis of the
532 influence of the noise in the GRSIR methodology would be interesting. A comparison
533 with other training approaches such as SVM or neural networks would allow to better
534 validate results. Finally, the development of a multivariate regularized GRSIR under
535 constraint is conceivable in order to estimate proportions simultaneously.

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Table 1. Sampling strategy for the simulation of the look-up table. The range of variation is given for each varying parameter as well as the number of distinct values simulated.

Parameters	Look-up table		Test data
	range	# distinct values	range
Proportion of water	[0.0001 0.0029]	15	[0.0006 0.002]
Proportion of CO ₂	[0.9942 0.9998]	29	[0.996 0.9988]
Proportion of dust	[0.0001 0.0029]	15	[0.0006 0.002]
Grain size of water	[50 450]	5	[100 400]
Grain size of CO ₂	[30000 165000]	28	[40000 105000]

Table 2. k -NN estimates for spectra S1 and S2 (described in section 3)

Parameters	Spectrum S1		Spectrum S2	
	Real values	k -NN	Real values	k -NN
Proportion of water	0.0019	0.0021	0.0013	0.0029
Proportion of CO ₂	0.9969	0.9966	0.9969	0.9942
Proportion of dust	0.0012	0.0013	0.0017	0.0029
Grain size of water	156	150	109	150
Grain size of CO ₂	65721	60000	78993	45000

Table 3. Validation criteria calculated on Tdata with GRSIR and k -NN methodology.

Parameters	k -NN	GRSIR	
		NRMSE	SIRC
Proportion of water	0.86	0.40	0.90
Proportion of CO ₂	0.88	0.30	0.98
Proportion of dust	0.44	0.17	0.99
Grain size of water	0.43	0.54	0.84
Grain size of CO ₂	0.53	0.22	0.95

Table 4. Validation criteria calculated on Tdata with GRSIR. Here the spectra that constitute the look-up table have been selected after a first inversion by GRSIR.

Parameters	GRSIR	
	NRMSE	SIRC
Proportion of water	0.27	0.92
Proportion of CO ₂	0.22	0.99
Proportion of dust	0.13	0.99
Grain size of water	0.37	0.92
Grain size of CO ₂	0.19	0.98

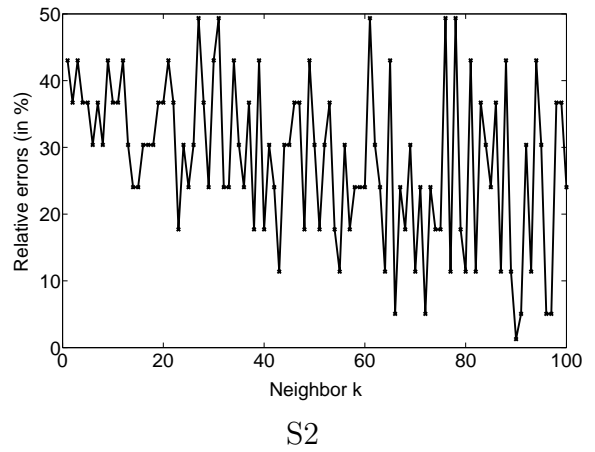
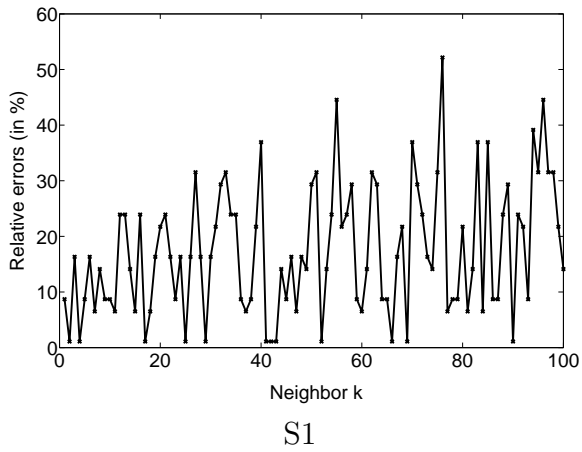


Figure 1. Relative errors on the grain size of CO₂ ice for the 100 nearest neighbors of spectra S1 and S2 in the lookup table. Horizontally: k th neighbor. Vertically: Relative errors on the grain size of CO₂ ice.

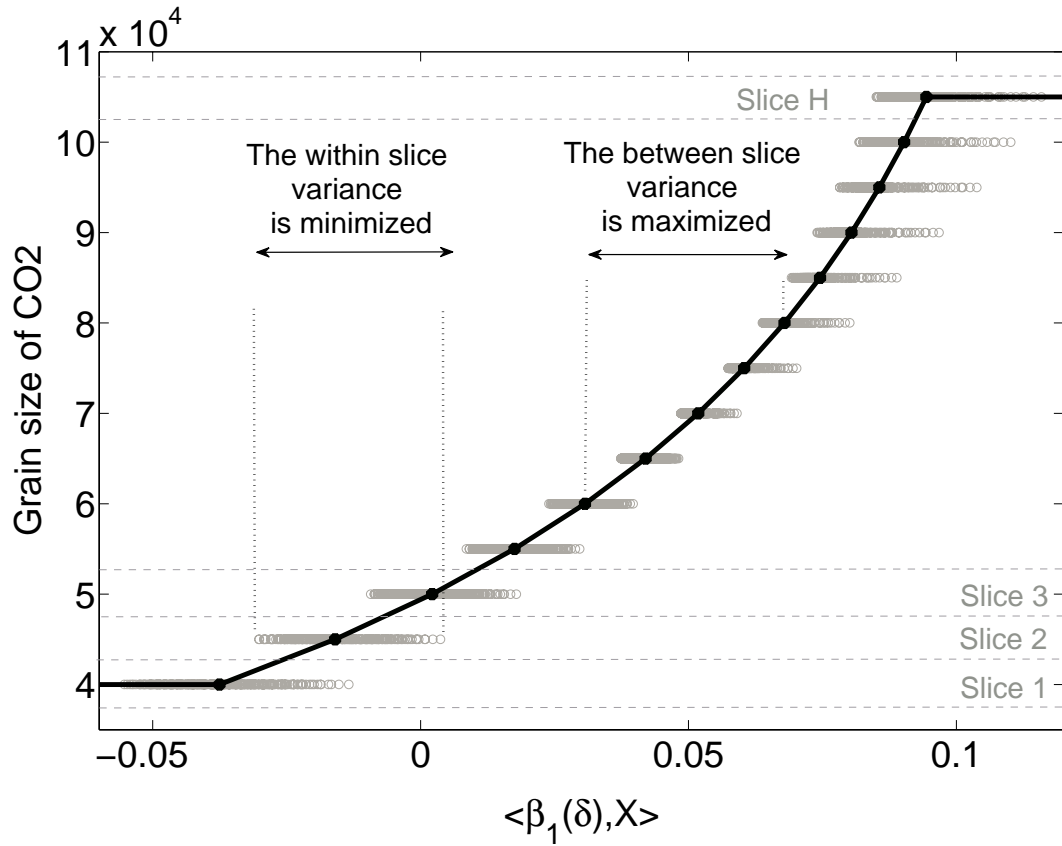


Figure 2. Functional relationship between reduced spectra on the first GRSIR axis and the grain size of CO₂ ice. This graphics illustrates the GRSIR methodology showing that the relationship is the best when the within slice variance is minimized or equivalently when the between slice variance is maximized. Horizontally: reduced spectra from the learning database on the first GRSIR axis. Vertically: Grain size of CO₂ ice.

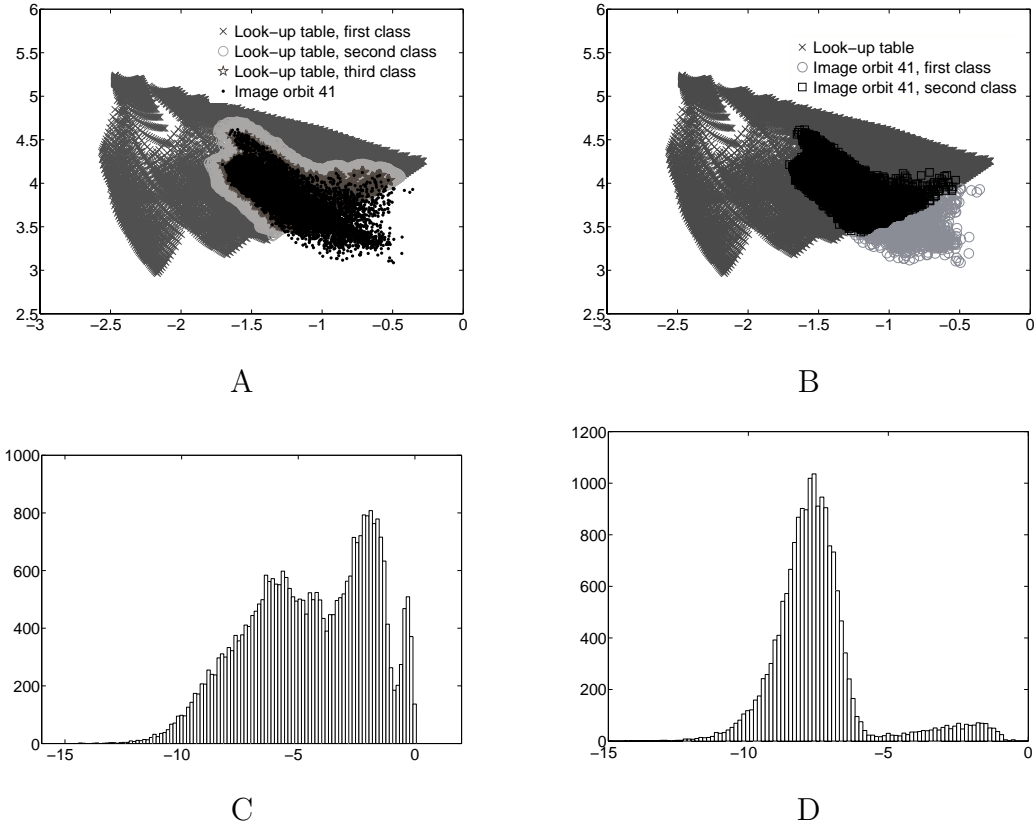


Figure 3. Selection of the “useful” spectra in the look-up table and of the “invertible” spectra in the observed data. A: Selection of a look-up table. Projections of the observed spectra from orbit 41 and the look-up table on the 2 first PCA axes (PCA applied to the look-up table). We distinguish 3 classes. The third class is the retained look-up table. B: Selection of invertible spectra in the image from orbit 41. The second class is retained for inversion. C: histogram of the distances between each spectrum of the look-up table and its nearest neighbor in the image from orbit 41. D: histogram of the distances between each spectrum of the image from orbit 41 and its nearest neighbor in the selected sub-look-up table.

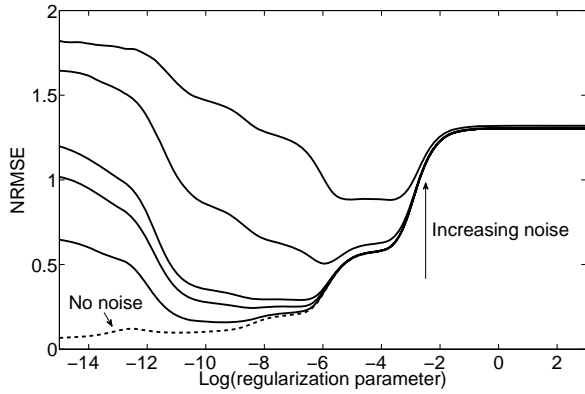


Figure 4. Influence of the noise in the choice of the regularization parameter. Horizontally: logarithm of the regularization parameter. Vertically: NRMSE for the grain size of CO₂ ice. A multiGaussian noise has been introduced in the test data with a diagonal covariance $C = \nu I_p$ proportional to the identity matrix I_p . The experience has been repeated five times with increasing values of ν . One test has been also realized without any noise in the data.

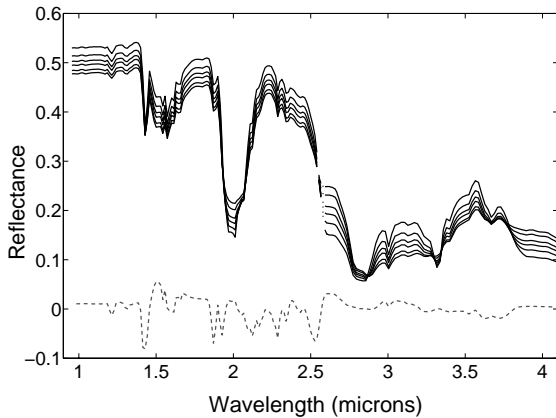


Figure 5. Plain line: series of spectra extracted from the lookup table for CO₂ ice grain size varying between 30 and 165 millimeters, the other parameters being kept constant at mean values. Dotted-dash line: Componentwise multiplication between GRSIR axis for the grain size of CO₂ ice and a representative spectrum (see section 6.4)

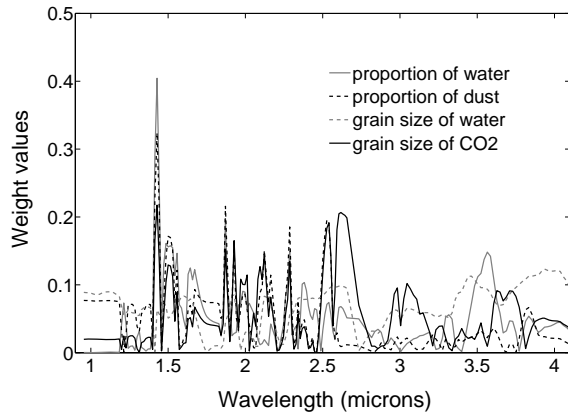


Figure 6. Horizontally: wavelengths. Vertically: GRSIR weights for the proportions of water ice and dust and the grain sizes of water and CO₂ ices.

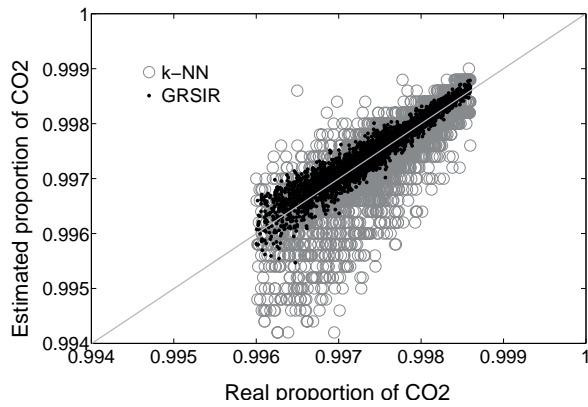


Figure 7. Scatter plot of the proportion of CO₂ ice from the test data versus estimated values. Horizontally: proportion of CO₂ ice. Vertically: estimated proportions of CO₂ ice by GRSIR and *k*-NN.

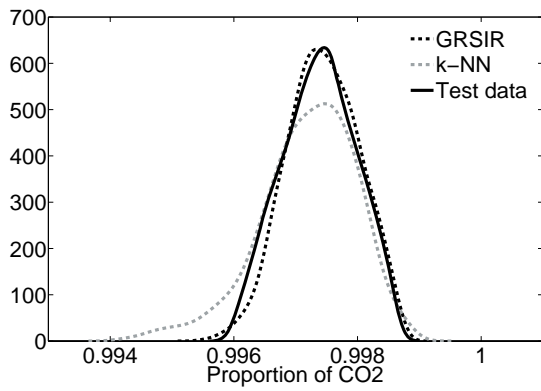


Figure 8. Density of the estimated proportions of CO₂ ice from the test data by k -NN and GRSIR compared to the density of real proportions.

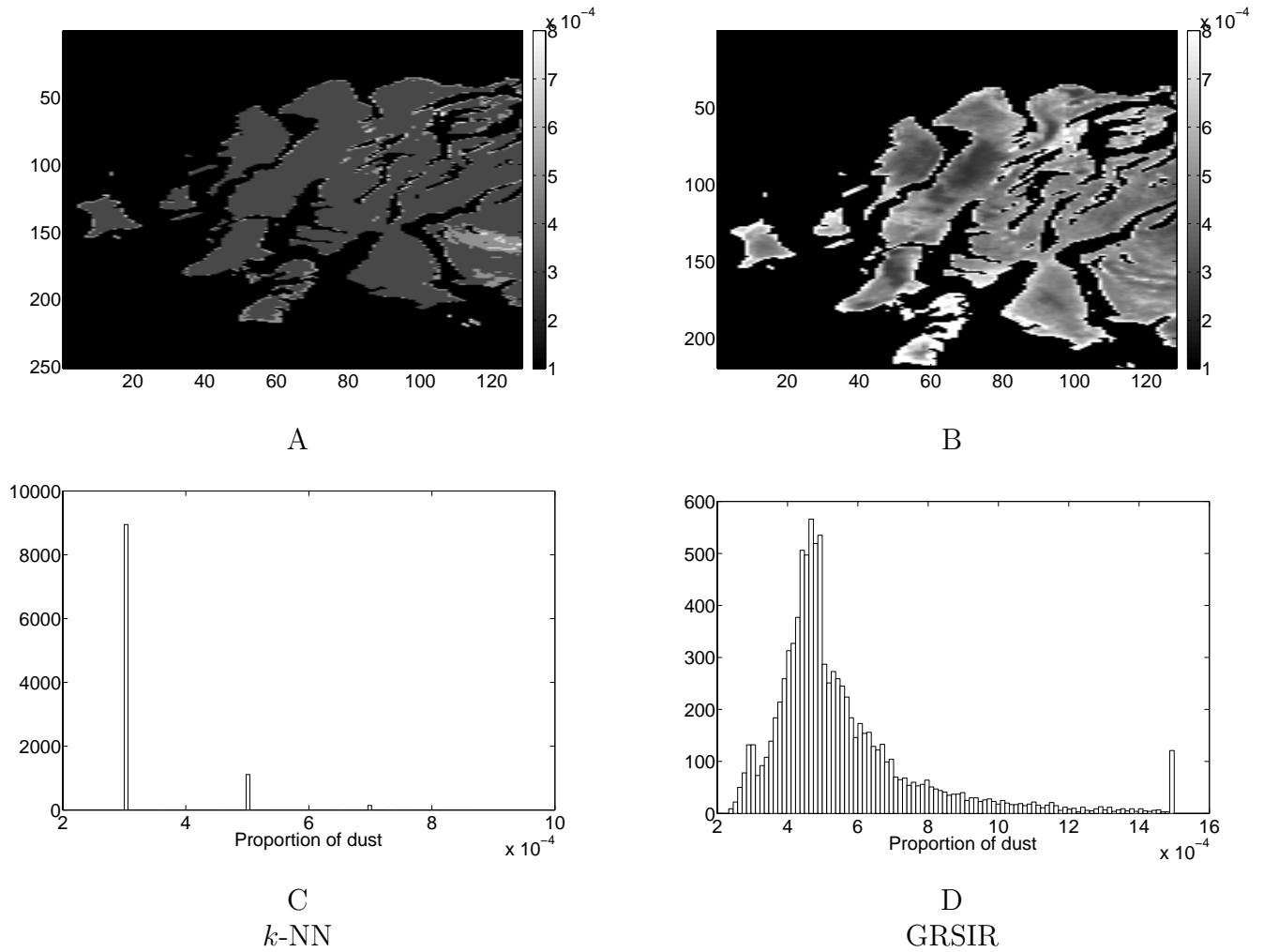


Figure 9. Proportion of dust estimated by k -NN (A) and GRSIR (B) from the hyperspectral image observed from orbit 103. The histograms of the estimations for the entire image are presented in (C) for k -NN and (D) for GRSIR.

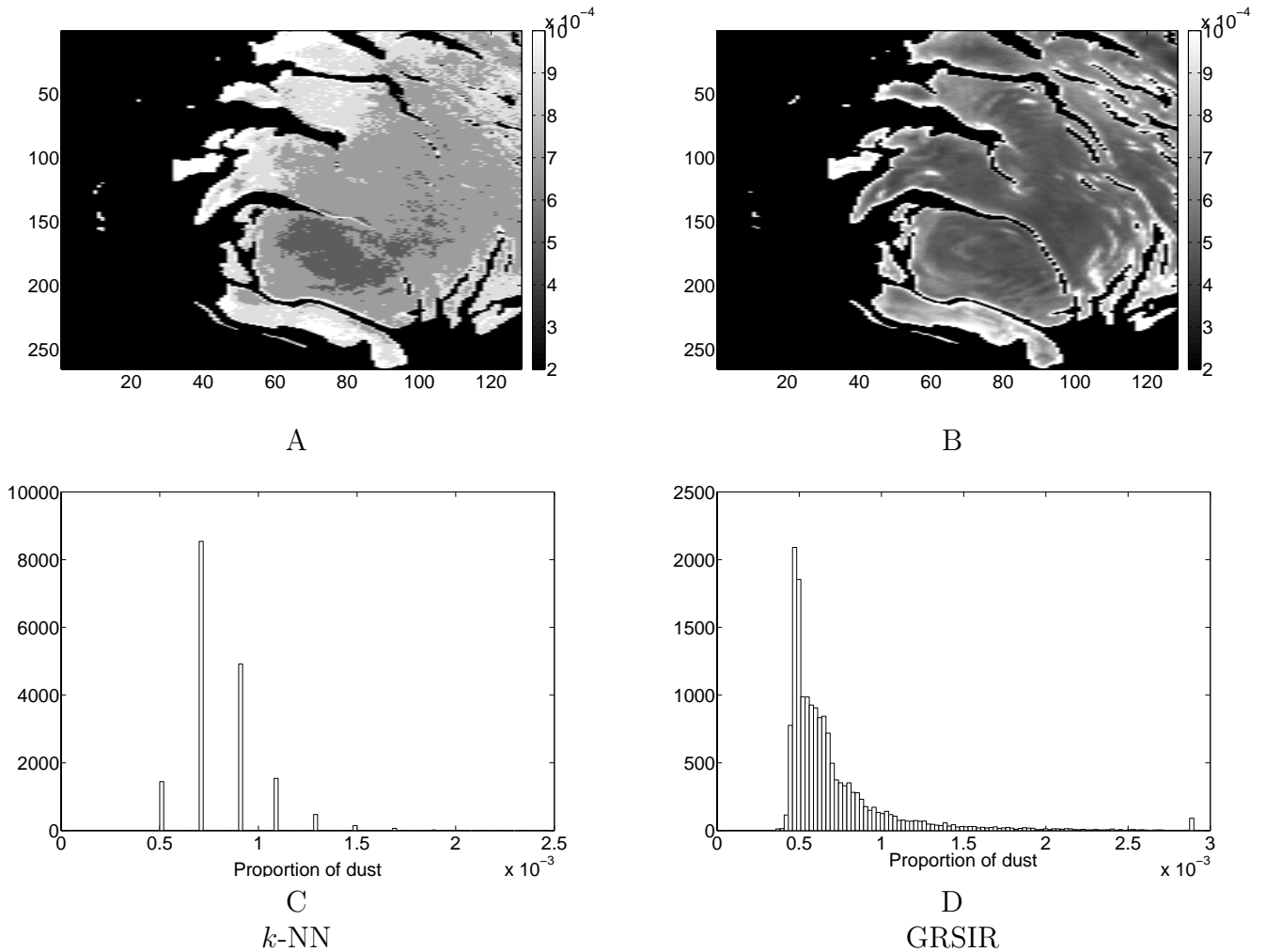


Figure 10. Proportion of dust estimated by k -NN (A) and GRSIR (B) from the hyperspectral image observed from orbit 41. The histograms of the estimations for the entire image are presented in (C) for k -NN and (D) for GRSIR.

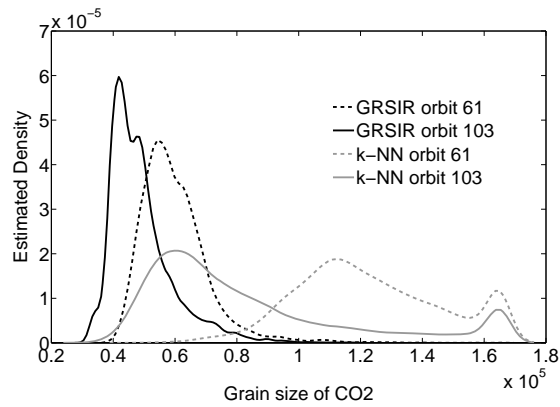


Figure 11. Comparison between the densities of the estimated grain sizes of CO₂ ice by *k*-NN and GRSIR in images from orbit 61 and 103.

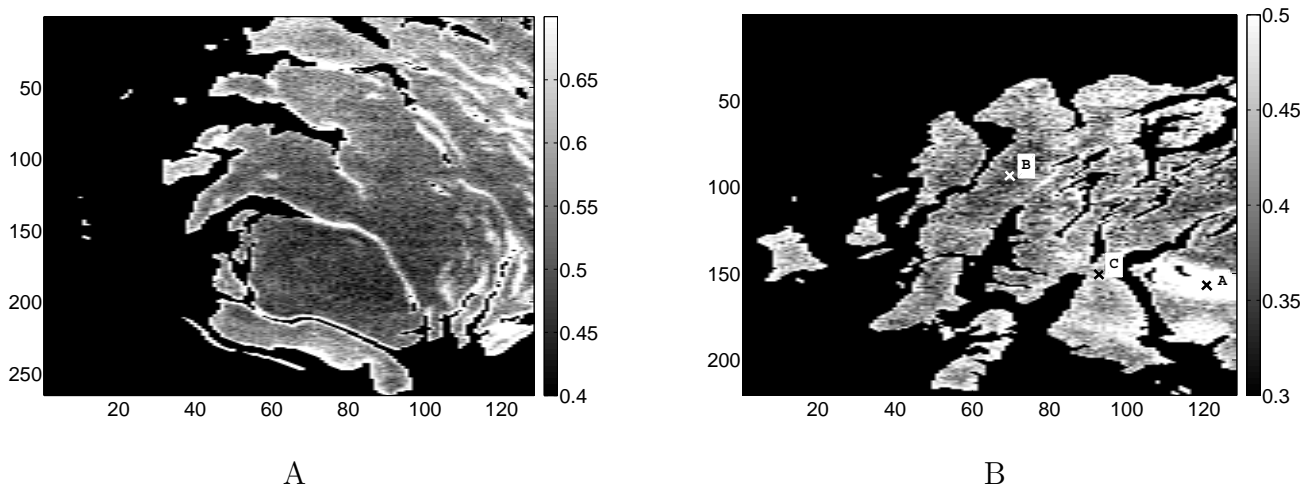


Figure 12. Cosinus of the Wavanglet angle between the observed spectra and a reference spectrum of martian dust for the hyperspectral images acquired during orbit 41 (A) and during orbit 103 (B).

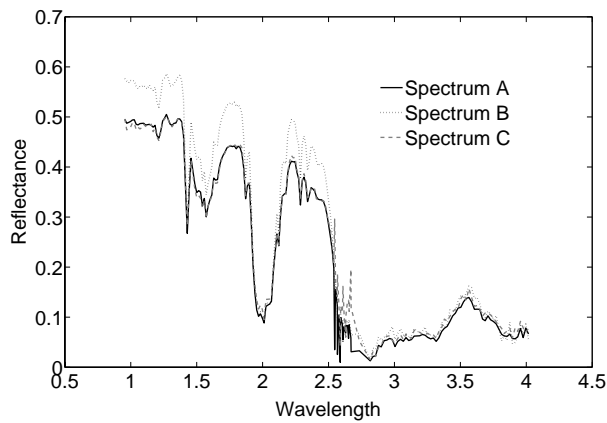


Figure 13. Spectra extracted from the locations labelled A, B and C in the image 103 of figure 12 (B)