



HAL
open science

Domain Optimization Problem for Stationary Heat Equation

Antoine Henrot, Werner Horn, Jan Sokolowski

► **To cite this version:**

Antoine Henrot, Werner Horn, Jan Sokolowski. Domain Optimization Problem for Stationary Heat Equation. [Research Report] RR-2788, INRIA. 1996, pp.26. inria-00073903

HAL Id: inria-00073903

<https://inria.hal.science/inria-00073903>

Submitted on 24 May 2006

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

INSTITUT NATIONAL DE RECHERCHE EN INFORMATIQUE ET EN AUTOMATIQUE

***Domain Optimization Problem for Stationary
Heat Equation***

Antoine Henrot, Werner Horn, Jan Sokolowski,

N° 2788

Janvier 1996

PROGRAMME 6



*R*apport
de recherche



Domain Optimization Problem for Stationary Heat Equation

Antoine Henrot, Werner Horn, Jan Sokolowski *

Programme 6 — Calcul scientifique, modélisation et logiciel numérique
Projet Numath

Rapport de recherche n2788 — Janvier 1996 — 26 pages

Abstract: In the paper the support of a Radon measure is selected in an optimal way. The solution of the second order elliptic equation depends on the measure via the mixed type boundary conditions. The existence of a solution for a class of domain optimization problems is shown. The relaxed formulation of the optimization problems is proposed. The first order necessary optimality conditions are derived.

Key-words: shape optimization, generalized solution, Radon measure, optimality conditions, heat equation, optimal design

(Résumé : tsvp)

*sokol@iecn.u-nancy.fr

Unité de recherche INRIA Lorraine
Technopôle de Nancy-Brabois, Campus scientifique,
615 rue de Jardin Botanique, BP 101, 54600 VILLERS LÈS NANCY (France)
Téléphone : (33) 83 59 30 30 – Télécopie : (33) 83 27 83 19
Antenne de Metz, technopôle de Metz 2000, 4 rue Marconi, 55070 METZ
Téléphone : (33) 87 20 35 00 – Télécopie : (33) 87 76 39 77

Un probleme d'optimisation de forme pour l'equation de la chaleur stationnaire

Résumé : Dans cet article, nous nous intéressons à l'optimisation de la forme d'un arc chauffant. Nous modélisons tout d'abord le probleme en faisant intervenir une mesure de Radon, dont nous allons chercher, en particulier, à optimiser le support. Nous montrons l'existence d'une solution optimale, d'abord en un sens classique, puis en un sens généralisé en relaxant le probleme. Nous écrivons aussi les conditions d'optimalité du premier ordre.

Mots-clé : l'optimisation de la forme, arc chauffant, mesure de Radon, l'existence d'une solution optimale, les conditions d'optimalité du premier ordre

DOMAIN OPTIMIZATION PROBLEM FOR STATIONARY HEAT EQUATION

A. HENROT*, W. HORN†
J. SOKOŁOWSKI‡

1. Introduction

In this paper we will consider a problem related to the following. Given a flat piece of material – a pane of glass in a window for example – we attach a heating wire to one surface of this material. This wire is modelled as a continuous curve connecting to fixed points A and B . We want to investigate which curve would optimize the temperature distribution on the opposite surface at a given time?

To elaborate the problem we use the following mathematical model:

Let Ω be a region in 2-dimensional Euclidean space. Define

$$\Sigma = \Omega \times (0, d), \quad d > 0,$$

and

$$\Omega_0 = \Omega \times \{0\}, \quad \Omega_1 = \Omega \times \{d\},$$

and

$$\Gamma = \partial\Omega \times (0, d)$$

the “vertical” boundary of Σ . Let u_0 and u_1 be positive functions on Ω_0 and Ω_1 , $A = (x_0, y_0)$ and $B = (x_1, y_1)$ two distinct points in Ω_0 , and $\gamma : [0, 1] \rightarrow \Omega_0$ a continuous curve of finite length in Ω_0 . Let U be the solution to the stationary heat equation

$$-\Delta U = 0,$$

on Σ , with boundary values

$$\partial_n U|_{\Gamma} = 0,$$

* Equipe de Mathématique, Université de Franche-Comté, 25030 Besancon Cedex, France; e-mail: henrot@vega.univ-fcomte.fr

† Department of Mathematics, California State University Northridge, 18111 Nordhoff St, Northridge, CA 91330, USA; e-mail: whorn@huey.csun.edu

‡ Institut Elie Cartan, Laboratoire de Mathématiques, Université Henri Poincaré Nancy I, B.P. 239, 54506 Vandoeuvre lès Nancy Cedex, France and Systems Research Institute of the Polish Academy of Sciences, ul. Newelska 6, 01-447 Warszawa, Poland; e-mail: sokol@iecn.u-nancy.fr

$$\partial_n U|_{\Omega_1} = U - u_1,$$

and

$$\partial_n U|_{\Omega_0} = U - u_0 - f_\gamma,$$

where f_γ is a positive function concentrated along the curve γ . One could think as f_γ to be an approximation of a delta-function at γ . The problem can now be stated as follows: Given a target function U^* on Ω_1 find γ such that

$$\|U|_{\Omega_1} - U^*\|_X^2$$

becomes minimal in a suitable Banach space X .

The crux of the matter is to find a suitable admissible set for the curves γ , as well as a convenient metric on this set of curves. First of all, it would be tempting to replace any curve by its parametrization in order to have a Banach structure on the set of curves. However, it is obvious that this point of view is not convenient, since a parametrization is not “intrinsic” enough to measure distances of two curves, as the following example illustrates. Let

$$\gamma_1 : \begin{cases} x(t) &= t^4 \\ y(t) &= 0 \end{cases}$$

and

$$\gamma_2 : \begin{cases} x(t) &= t \\ y(t) &= 0 \end{cases}$$

Both of these parametrizations give the same curve, but

$$\int_0^1 |\gamma_1(t) - \gamma_2(t)|^2 dt > 0,$$

showing that this integral does not define a metric on the set of curves.

A more classical idea is to work with the Hausdorff metric. For two curves parametrized by $\gamma_1(t)$, $t \in [0, 1]$ and $\gamma_2(t)$, $t \in [0, 1]$ this distance is defined by

$$d(\gamma_1, \gamma_2) = \max \left\{ \max_{t \in [0, 1]} \left(\min_{s \in [0, 1]} |\gamma_1(t) - \gamma_2(s)| \right), \max_{s \in [0, 1]} \left(\min_{t \in [0, 1]} |\gamma_1(t) - \gamma_2(s)| \right) \right\}.$$

This metric has good compactness properties. For example, for any sequence of compact sets K_n which are included in some large ball, there is a subsequence which converges in the Hausdorff metric to a compact set K . Unfortunately, for a sequence

of curves K_n the K does not have to be a curve, as the following example illustrates. The curves $K_n = \{y = \sin(nx), x \in [0, \pi]\}$ converge to the set $K = [0, \pi] \times [-1, 1]$. It therefore seems natural that one has to impose some additional constraints on the set of curves considered. These constraints could be on the length or on the Hausdorff measure of the curves. The following sections will elaborate on these ideas.

It is well-known (see, for example, Ziemer) that if γ_n is a sequence of continuous curves whose Hausdorff measure is uniformly bounded by a number M , and if γ_n converge to γ in the Hausdorff metric, then γ is also a continuous curve. However, it is not generally true that the Dirac-measures δ_{γ_n} converge weak* to δ_γ (see section 3). But this is exactly the kind of convergence necessary to prove the continuity of the solution of the problem above with respect to curves.

We are faced with a classical situation in shape optimization: the Hausdorff distance has very good compactness properties, but is not strong enough to ensure that the cost functional is lower semi-continuous.

Finally, we want to point out, that the results of this paper also hold if the Laplacian is replaced by more general uniform elliptic operators.

2. Existence of a classical solution

We assume that Ω is a simply connected domain in \mathbb{R}^2 and let $\Sigma = \Omega \times (0, d)$. We denote $\Omega_0 = \Omega \times \{0\}$, $\Omega_1 = \Omega \times \{d\}$ and $\Gamma = \partial\Omega \times (0, d)$. Therefore

$$\partial\Sigma = \Omega_0 \cup \Omega_1 \cup \Gamma$$

Given a curve $\gamma \subset \Omega_0$ parametrized by $s \in [0, 1]$, we assume that $A = \gamma(0)$ and $B = \gamma(1)$ are fixed points in Ω_0 . For the stationary heat equation γ is the heat source. Let us consider the following elliptic equation.

$$(\mathcal{P}_1(\gamma)) \quad \begin{cases} -\Delta u = 0 & \text{in } \Omega \\ \frac{\partial u}{\partial n} = 0 & \text{on } \Gamma \\ -\frac{\partial u}{\partial n} = u - u_1 & \text{on } \Omega_1 \\ -\frac{\partial u}{\partial n} = u - u_0 - \delta_\gamma & \text{on } \Omega_0 \end{cases}$$

where u_0, u_1 are given L^2 functions, and δ_γ is a Dirac measure supported on the curve γ .

The variational formulation of the stationary heat equation is given as follows.

$$\text{Find } u \in H^1(\Sigma) \text{ such that for all functions } v \in H^1(\Sigma)$$

$$(\mathcal{P}_2(\gamma)) \quad a(u, v) = L(v)$$

where

$$a(u, v) = \int_{\Sigma} \nabla u \cdot \nabla v dx + \int_{\Omega_1} u v d\sigma + \int_{\Omega_0} u v d\sigma \quad (1)$$

$$L(v) = \int_{\Omega_1} u_1 v d\sigma + \int_{\Omega_0} u_0 v d\sigma + \int_{\gamma} v d\gamma \quad (2)$$

In order to have the problem well defined it is sufficient to show that the linear form

$$\delta_{\gamma} : v \mapsto \int_{\gamma} v d\gamma$$

is continuous on the space $H^1(\Sigma)$. We are going to define the set of admissible curves γ in such a way that the linear form is continuous.

To this end we denote by Q the cube $Q = (0, 1) \times (0, 1)$, by $I \subset Q$ the interval $I = \left[-\frac{1}{2}, \frac{1}{2}\right] \times \{0\}$.

Definition 1. A given curve γ is called admissible if there exists a one-to-one mapping $F : Q \mapsto \mathcal{O}$, where \mathcal{O} denotes an open neighbourhood of γ in Ω_0 such that

$$F(Q) = \mathcal{O} \quad F(I) = \gamma \quad (3)$$

$$\|F\|_{W^{1,\infty}(Q)} \leq L_1 \quad \|F^{-1}\|_{W^{1,\infty}(\mathcal{O})} \leq L_2 \quad (4)$$

Prescribing uniform bounds $L = L_1 = L_2 > 0$ and assuming that the following compactness condition is satisfied

(\mathcal{H}) Given a sequence F_n which satisfies uniformly the latter bounds, there exists a subsequence, still denoted by F_n such that

$$|F'_n(\cdot, 0)| \rightarrow |F'(\cdot, 0)| \quad \text{weakly in } L^2\left(-\frac{1}{2}, \frac{1}{2}\right). \quad (5)$$

we define an admissible family

$$\mathcal{F}_L = \{\gamma \text{ is admissible} \mid (\mathcal{H}) \text{ is satisfied, } \|F\|_{W^{1,\infty}(Q)} \leq L \text{ and } \|F^{-1}\|_{W^{1,\infty}(\mathcal{O})} \leq L\}$$

where $L > 0$ is a given constant.

Remark 1. Without the assumption (\mathcal{H}) on the family \mathcal{F}_L we cannot expect that for any sequence $\{\gamma_n\} \subset \mathcal{F}_L$, there exists a subsequence, still denoted by $\{\gamma_n\}$ such that

$$\delta_{\gamma_n} \rightarrow \delta_{\gamma} \quad \text{weak-}(\ast) \text{ in the space } (H^1(\Sigma))' .$$

A counterexample can be constructed using $F_n(x, y) = \{x, y + \frac{1}{n} \sin(nx)\}$.

Our problem becomes now:

Minimize the cost functional

$$J(\gamma) = \|u_\gamma - u_d\|$$

where u_γ denotes a solution to the stationary heat equation for any $\gamma \in \mathcal{F}_L$ and the Dirac measure δ_γ in the boundary conditions, u_d is a given function, and $\|\cdot\|$ is a norm, or a seminorm on the space $H^1(\Sigma)$ which will be specified later.

Remark 2. We use the above definition of a set of admissible curves \mathcal{F}_L , since we want to apply an appropriate trace theorem on γ . Such a definition is better suited for our applications than the simple definition of curves parametrized over an interval.

Remark 3. We can replace Definition 1 by a more general notion of a Lipschitzian manifold, where the existence of a global parametrization is not required. We prefer to work with the global parametrization for the sake of simplicity. The same result can be obtained for the more general setting of a Lipschitzian manifold, provided that the uniform bounds are prescribed with the same Lipschitz constant for any collection of charts. Using a partition of unity the problem can be localized in a standard way.

Remark 4. Some classes of admissible curves in the plane are introduced by I.I. Daniliuk (Daniliuk, 1975) in the framework of integral equations in non-smooth domains.

On the other hand, it seems to be possible to use some families of admissible curves defined by using capacity type constraints, which probably assure the existence of a solution in a slightly wider class. But this approach is rather complicated and it is not evident that such families of admissible curves can be of any interest for the numerical methods. We refer the reader to the monograph (Ziemer, 1989) for the definition and properties of capacity, and to (Bucur, Zolesio, to appear) for some results in the case of admissible domains with capacity constraints for homogeneous Dirichlet problems. In the present paper we rather use the notion of a generalized solution to the problem defined in section 3.

Proposition 1. For any admissible curve $\gamma \in \mathcal{F}_L$ the linear form

$$\langle \delta_\gamma, \varphi \rangle = \int_\gamma \varphi d\gamma$$

is continuous with the norm in the dual space bounded

$$\|\delta_\gamma\|_* \leq C_\gamma P(\gamma)$$

where $C_\gamma = C_\gamma(L, \Sigma)$ and $P(\gamma) = \int_\gamma d\gamma$ is the length of γ .

Proof. For an element $\varphi \in H^1(\Sigma)$ the trace on Ω_0 is also denoted by φ and the trace satisfies $\varphi \in H^{\frac{1}{2}}(\Omega_0)$, we refer to (Adams, 1975) and (Lions, Magenes, 1968) for a proof. A first important question is, whether or not it is possible to define a trace on γ for any element of the space $H^{\frac{1}{2}}(\Omega_0)$. The positive answer is obtained by applying the theorem of Besov–Uspienskii (Adams, Thm. 7.58):

The injection of the space $H^{\frac{1}{2}}(\mathbb{R}^2)$ into $L^2(\mathbb{R})$ is continuous.

Let us show that if u is in $H^{\frac{1}{2}}(\Omega_0)$ the function $\hat{u} = u \circ F$ defined on Q belongs to the space $H^{\frac{1}{2}}(Q)$. Hence, it is sufficient to prove that the following integral is finite

$$\begin{aligned} I &= \int_Q \int_Q \left(\frac{|\hat{u}(x) - \hat{u}(y)|}{|x - y|^{\frac{3}{2}}} \right)^2 dx dy \\ &= \int_Q \int_Q \frac{|u(F(x)) - u(F(y))|^2}{|x - y|^3} dx dy \end{aligned} \quad (6)$$

We denote $x_1 = F(x)$, $y_1 = F(y)$, $|DF^{-1}|$ the determinant of Jacobian DF^{-1} , hence

$$I = \int_{\mathcal{O}} \int_{\mathcal{O}} \frac{|u(x_1) - u(y_1)|^2}{|F^{-1}(x_1) - F^{-1}(y_1)|^3} |DF^{-1}| dx_1 dy_1 \quad (7)$$

Since the mapping F is Lipschitz,

$$L_1 |F^{-1}(x_1) - F^{-1}(y_1)| = L_1 |x - y| \geq |F(x) - F(y)| = |x_1 - y_1| \quad (8)$$

$$\text{and} \quad |DF^{-1}| \leq C = 2L_2^2 \quad (9)$$

thus

$$I \leq 2L_2^2 L_1^3 \int_{\mathcal{O}} \int_{\mathcal{O}} \frac{|u(x_1) - u(y_1)|^2}{|x_1 - y_1|^3} dx_1 dy_1 < \infty$$

since $u \in H^{\frac{1}{2}}(\Omega_0)$.

The trace operator maps $H^{\frac{1}{2}}(Q)$ into $L^2(Q_0)$ by the theorem of Besov and Uspienskii, is defined by means of the mapping F , as a trace operator for the space $H^{\frac{1}{2}}(\mathcal{O})$ into the space $L^2(\gamma)$, furthermore,

$$\|\varphi\|_{L^2(\gamma)} \leq C \|\varphi\|_{H^{\frac{1}{2}}(\mathcal{O})} \leq C \|\varphi\|_{H^{\frac{1}{2}}(\Omega_0)}$$

with $C = C(L_1, L_2)$, L_1, L_2 are the Lipschitz constants of F and F^{-1} . In particular $L = L_1 = L_2$ for $\gamma \in \mathcal{F}_L$. In view of the continuity of the trace operator $H^1(\Sigma) \mapsto H^{\frac{1}{2}}(\Omega_0)$ it follows that

$$\|\varphi\|_{L^2(\gamma)} \leq C\|\varphi\|_{H^1(\Sigma)} \quad (10)$$

with a constant $C = C(L_1, L_2, \Sigma)$. Therefore

$$\left| \int_{\gamma} \varphi d\gamma \right| \leq P(\gamma) \left(\int_{\gamma} \varphi^2 d\gamma \right)^{\frac{1}{2}} \leq CP(\gamma)\|\varphi\|_{H^1(\Sigma)}$$

which completes the proof. \blacksquare

An admissible curve is defined in the parametric form

$$\begin{cases} x(t) = F_1(t, 0) \\ y(t) = F_2(t, 0) \end{cases} \quad t \in \left[-\frac{1}{2}, \frac{1}{2}\right]$$

where $F = (F_1, F_2)$ is bi-Lipschitz mapping. For $\gamma \in \mathcal{F}_L$ it follows that

$$P(\gamma) = \int_{-\frac{1}{2}}^{\frac{1}{2}} \sqrt{x'^2(t) + y'^2(t)} dt = \int_{-\frac{1}{2}}^{\frac{1}{2}} \left(\frac{\partial F_1^2}{\partial x}(t, 0) + \frac{\partial F_2^2}{\partial x}(t, 0) \right)^{\frac{1}{2}} dt \leq \sqrt{2}L$$

therefore the length of admissible curves in the set \mathcal{F}_L is uniformly bounded, but the uniform boundedness of the length is a weaker condition for a curve than the condition to be a member of \mathcal{F}_L .

The class \mathcal{F}_L is sufficiently small to obtain an existence result for the problem under considerations.

Proposition 2. *Given a sequence of curves γ_n in \mathcal{F}_L , there exists a curve $\gamma \in \mathcal{F}_L$ and a subsequence γ_{n_k} such that*

$$\delta_{\gamma_{n_k}} \rightarrow \delta_{\gamma} \quad \text{weak-} (*) \text{ in the space } (H^1(\Sigma))'$$

i.e.,

$$\langle \delta_{\gamma_{n_k}}, \varphi \rangle \rightarrow \langle \delta_{\gamma}, \varphi \rangle \quad \text{for all } \varphi \in H^1(\Sigma)$$

Proof. Given $\gamma_n = F_n(Q_0) \in \mathcal{F}_L$, we have

$$\|F_n\|_{W^{1,\infty}} \leq L \quad \text{and} \quad \|F_n^{-1}\|_{W^{1,\infty}} \leq L$$

By the theorem of Ascoli there exists a function F which is continuous over Q such that for a subsequence F_{n_k}

$$F_{n_k}(x) \rightarrow F(x) \quad \text{uniformly over } \overline{Q} .$$

The functions F_{n_k} are uniformly Lipschitz continuous with the constant L , the same remains valid for F , thus $F \in W^{1,\infty}(Q)$ with $\|F\|_{W^{1,\infty}} \leq L$. We denote $\gamma = F(Q_0)$.

Furthermore, the inequality $\|F_n^{-1}\|_{W^{1,\infty}} \leq L$ implies that

$$|F_n(x) - F_n(y)| \geq \frac{1}{L}|x - y| \quad \forall x, y \in Q \quad (11)$$

hence taking the limit it follows that

$$|F(x) - F(y)| \geq \frac{1}{L}|x - y| \quad \forall x, y \in Q \quad (12)$$

which shows that F is one-to-one. We denote $\mathcal{O} = F(Q)$, thus there exists the inverse mapping $F^{-1} : \mathcal{O} \mapsto Q$, F^{-1} being Lipschitz continuous with the constant L in view of the latter inequality. Therefore $\gamma \in \mathcal{F}_L$.

For the sake of simplicity we denote by γ_n the subsequence γ_{n_k} .

We are going to show that δ_{γ_n} converges to δ_γ . To this end we assume that there is given a continuous on $\overline{\Sigma}$ function φ , henceforth

$$\langle \delta_{\gamma_n}, \varphi \rangle = \int_{\gamma_n} \varphi d\gamma_n = \int_{-\frac{1}{2}}^{\frac{1}{2}} \varphi(F_n(t, 0)) |F'_n(t, 0)| dt$$

The sequence F_n satisfies uniformly (4), using the assumption (\mathcal{H}) it follows that

$$|F'_n(\cdot, 0)| \rightarrow |F'(\cdot, 0)| \quad \text{weakly in } L^2\left(-\frac{1}{2}, \frac{1}{2}\right) .$$

Since φ is continuous, hence uniformly continuous on $\overline{\Omega}_0$,

$$\varphi(F_n(\cdot, 0)) \rightarrow \varphi(F(\cdot, 0)) \quad \text{in } L^\infty\left(-\frac{1}{2}, \frac{1}{2}\right)$$

thus

$$\langle \delta_{\gamma_n}, \varphi \rangle \rightarrow \int_{-\frac{1}{2}}^{\frac{1}{2}} \varphi(F(t, 0)) |F'(t, 0)| dt = \langle \delta_\gamma, \varphi \rangle$$

The same result can be obtained for an arbitrary $\varphi \in H^1(\Sigma)$ since $C(\overline{\Sigma}) \cap H^1(\Sigma)$ is dense in $H^1(\Sigma)$. ■

Let us consider a sequence of admissible curves γ_n and the admissible curve γ such that δ_{γ_n} converges to δ_γ weakly in $(H^1(\Sigma))'$. We denote by u_n, u solutions to (\mathcal{P}_1) or (\mathcal{P}_2) for the boundary data $\delta_{\gamma_n}, \delta_\gamma$, respectively. We are interested in the convergence $u_n \rightarrow u$.

Proposition 3. *Let $\{\gamma_n\}, \gamma \in \mathcal{F}_L$ be given, such that $\delta_{\gamma_n} \rightarrow \delta_\gamma$ weakly in $(H^1(\Sigma))'$. Then*

$$u_n \rightarrow u \quad \text{in } H^1(\Sigma) \text{ weakly and in } L^2(\Sigma) \text{ strongly} .$$

Proof. u_n is the unique solution to the following variational problem,

$$a(u_n, v) = \mathcal{L}_n(v) \quad \forall v \in H^1(\Sigma) \quad (13)$$

where

$$\mathcal{L}_n(v) = \int_{\Omega_1} u_1 v d\sigma + \int_{\Omega_0} u_0 v d\sigma + \int_{\gamma_n} v d\gamma_n$$

From the proposition 1 it follows that $\|\mathcal{L}_n\|_* \leq C$, where the constant C is independent of $n = 1, 2, \dots$. The bilinear form $a(\cdot, \cdot)$ being coercive by the inequality of Friedrichs–Poincaré, we obtain directly from the variational formulation that

$$\alpha \|u_n\|_{H^1(\Sigma)}^2 \leq a(u_n, u_n) = \mathcal{L}_n(u_n) \leq C \|u_n\|_{H^1(\Sigma)}$$

therefore the sequence $u_n, n = 1, 2, \dots$, is bounded in $H^1(\Sigma)$. There exists a subsequence of the sequence u_n , still denoted u_n , such that

$$u_n \rightarrow u^* \quad \text{weakly in } H^1(\Sigma) \text{ and strongly in } L^2(\Sigma)$$

the strong convergence follows by an application of the Rellich theorem. We show that $u^* = u$.

By the weak convergence of the sequence $\{u_n\}$ in $H^1(\Sigma)$, since the trace mapping is linear and continuous, we have the following convergence of the traces

$$u_n \rightarrow u^* \quad \text{in } L^2(\Omega_0) \text{ and in } L^2(\Omega_1)$$

hence for any fixed test function $v \in H^1(\Sigma)$

$$a(u_n, v) \rightarrow a(u^*, v)$$

and with our assumptions

$$\int_{\gamma_n} v d\gamma_n \rightarrow \int_{\gamma} v d\gamma$$

whence

$$\mathcal{L}_n(v) \rightarrow \mathcal{L}(v) .$$

We obtain

$$a(u^*, v) = \mathcal{L}(v),$$

and the solution to the problem $(\mathcal{P}_2(\gamma))$ being unique, it follows that $u^* = u$ which completes the proof. \blacksquare

Remark 5. In order to show that $u_n \rightarrow u$ strongly in $H^1(\Sigma)$ it is sufficient to have the following convergence

$$\int_{\gamma_n} u_n d\gamma_n \rightarrow \int_{\gamma} u d\gamma \quad (14)$$

since, using the variational formulation of the problem (\mathcal{P}_2) , we obtain

$$\int_{\Sigma} |\nabla u_n|^2 dx \rightarrow \int_{\Sigma} |\nabla u|^2 dx \quad (15)$$

Using the above results we are in position to prove an existence result for the optimization problem under considerations. Assume that there is given a functional $J(\cdot)$ continuous with respect to $u = u(\gamma)$ in the norm topology of the space $L^2(\Sigma)$ or weakly lower semicontinuous on $H^1(\Sigma)$. Let us consider, as an example, the following cost functional

$$J(\gamma) = \int_{\Sigma} (u(\gamma) - u_d)^2 dx + \int_{\Sigma} |\nabla u(\gamma) - \nabla u_d|^2 dx \quad (16)$$

Theorem 1. *There exists a solution to the minimization problem*

$$\inf_{\gamma \in \mathcal{F}_L} J(\gamma) \quad (17)$$

Proof. Let $\{\gamma_n\}$ denote a minimizing sequence, then for a subsequence, still denoted by $\{\gamma_n\}$ we have

$$u(\gamma_n) \rightarrow u(\gamma) \quad \text{in } H^1(\Sigma) \text{ weakly} \quad (18)$$

hence

$$\liminf J(\gamma_n) \geq J(\gamma)$$

which completes the proof of theorem. ■

Let us present another formulation of the problem which furnishes a smooth solution. We denote by $D \subset \mathbb{R}^2 = \mathbb{C}$ the unit disk. We denote by \mathfrak{D} the collection of holomorphic functions defined on D with values in Ω_0 ,

$$\Phi(z) = \sum_{n=0}^{\infty} a_n z^n \quad a_n \in \mathbb{C} \quad \text{the serie converges in } D, \quad \Phi(z) \in \Omega_0$$

The curve γ is defined by the following parametrization

$$z(t) = \sum_{n=0}^{\infty} a_n t^n \quad t \in \left[-\frac{1}{2}, \frac{1}{2}\right].$$

In fact, by the theorem of Stone–Weierstrass, any curve located in Ω_0 can be approximated by the curve of this form.

Using the set \mathfrak{D} we obtain the existence of the solution for our problem since if the sequence $\{\Phi_p\}_{p \in \mathbb{N}}$ is a minimizing sequence, using the theorem of Montel it follows, since Ω_0 is bounded and $\Phi_p(z) \in \Omega_0$, that there exists a subsequence which converges uniformly on any compact, along with all derivatives. In this case the assumption (\mathcal{H}) is satisfied. In particular Φ'_p converges to the limit Φ' on the interval $[-\frac{1}{2}, \frac{1}{2}]$ which implies that $\delta_{\gamma_p} \rightarrow \delta_{\gamma}$ weakly in $(H^1(\Sigma))'$.

3. Generalized solutions to domain optimization problem

We start with the classical definition of a solution $u \in W^{1,p}(\Sigma)$, $1 \leq p < \frac{3}{2}$, to the system (\mathcal{P}_1) in the form

$$(\mathcal{P}_1(\mu)) \quad \begin{cases} -\Delta u = 0 & \text{in } \Omega \\ \frac{\partial u}{\partial n} = 0 & \text{on } \Gamma \\ -\frac{\partial u}{\partial n} = u - u_1 & \text{on } \Omega_1 \\ -\frac{\partial u}{\partial n} = u - u_0 - \mu & \text{on } \Omega_0 \end{cases}$$

where μ is a Radon measure supported on Ω_0 .

We are going to prove the existence and uniqueness of the solution to $\mathcal{P}_1(\mu)$. First, we recall the Friedrich’s type inequality related to our problem.

Lemma 1. *Let $\Omega \subset \mathbb{R}^N$ be a bounded simply connected domain with smooth boundary $\Gamma = \partial\Omega$, $\tilde{\gamma} \subset \Gamma$ a given set with $|\tilde{\gamma}| = \int_{\tilde{\gamma}} d\Gamma(x) > 0$. Then there exists a constant*

$C = C(\Omega, \tilde{\gamma}, p)$ such that

$$\|v\|_{L^p(\Omega)} \leq C \left[\sum_{i=1}^N \int_{\Omega} \left| \frac{\partial v}{\partial x_i} \right|^p dx + \int_{\tilde{\gamma}} |v|^p d\Gamma(x) \right]^{\frac{1}{p}} \quad \forall v \in W^{1,p}(\Omega)$$

From lemma 1 it follows that the equivalent norm in the space $W^{1,p}(\Omega)$ can be defined by

$$\|v\|_{1,p} = \left[\sum_{i=1}^N \int_{\Omega} \left| \frac{\partial v}{\partial x_i} \right|^p dx + \int_{\tilde{\gamma}} |v|^p d\Gamma(x) \right]^{\frac{1}{p}}. \quad (19)$$

We can evaluate the norm $\|\cdot\|_{1,p}$ by duality, where q satisfies $\frac{1}{p} + \frac{1}{q} = 1$,

$$\|v\|_{1,p} = \sup_{\substack{\xi_1, \dots, \xi_N \in L^q(\Omega) \\ \xi_0 \in L^q(\tilde{\gamma}) \\ \sum_{i=1}^N \|\xi_i\|_{L^q} = 1}} \left\{ \sum_{i=1}^N \int_{\Omega} \xi_i \frac{\partial v}{\partial x_i} dx + \int_{\tilde{\gamma}} v \xi_0 d\Gamma(x) \right\} \quad (20)$$

We denote by $\|\cdot\|_{\mathcal{M}_b(\Gamma_0)}$ the norm defined by

$$\|\nu\|_{\mathcal{M}_b(\Gamma_0)} = \sup_{\substack{\varphi \in C^0(\Gamma_0) \\ \varphi \neq 0}} \frac{|\int \varphi d\nu|}{\|\varphi\|_{\infty}}$$

Proposition 4. *Let $\Omega \subset \mathbb{R}^N$, $N \geq 2$ be a bounded simply connected domain with smooth boundary $\Gamma = \Gamma_0 \cup \Gamma_1$, $|\Gamma_0| > 0$, μ a bounded Radon measure supported on Γ_0 . There exists the unique solution $u \in W^{1,p}(\Omega)$, for all $p \in \left[1, \frac{N}{N-1}\right)$ to the problem*

$$(\mathcal{P}_{\mu}) \quad \begin{cases} -\Delta u & = 0 & \text{in } \Omega \\ \frac{\partial u}{\partial n} & = 0 & \text{on } \Gamma_1 \\ \frac{\partial u}{\partial n} & = \mu - u & \text{on } \Gamma_0 \end{cases}$$

moreover

$$\|u\|_{W^{1,p}(\Omega)} \leq \|\mu\|_{\mathcal{M}_b(\Gamma_0)}$$

Proof. Let us first consider $d\nu = \psi d\Gamma(x)$, $\psi \in C(\Gamma_0)$, and afterwards we use the density argument.

There exists the unique solution $u \in H^1(\Omega)$, satisfying $u \in W^{1,p}(\Omega)$, to the problem \mathcal{P}_ψ in the variational form

$$\int_{\Omega} \nabla u \cdot \nabla v dx + \int_{\Gamma_0} u v d\Gamma(x) = \int_{\Gamma_0} \psi v d\Gamma(x) \quad \forall v \in H^1(\Omega)$$

Given $\xi_0 \in C^\infty(\Gamma_0)$ and $\{\xi_1, \dots, \xi_N\} \in (\mathcal{D}(\Omega))^N$, we denote by v the unique solution to the following elliptic equation

$$(\mathcal{P}'_\psi) \quad \begin{cases} -\Delta v &= \sum_{i=1}^N \frac{\partial \xi_i}{\partial x_i} & \text{in } \Omega \\ \frac{\partial v}{\partial n} &= 0 & \text{on } \Gamma_1 \\ \frac{\partial v}{\partial n} &= \psi - v & \text{on } \Gamma_0 \end{cases}$$

Let $q > N$ be fixed, by the injection theorem of Sobolev, taking into account the variational formulation of (\mathcal{P}'_ψ) and the definition of the norm by duality, it follows that

$$\|v\|_{L^\infty(\Omega)} \leq C \|v\|_{W^{1,p}(\Omega)} \leq C \left(\sum_{i=1}^N \|\xi_i\|_{L^q(\Omega)} + \|\xi_0\|_{L^q(\Gamma_0)} \right).$$

Furthermore

$$\begin{aligned} \sum_{i=1}^N \int_{\Omega} \xi_i \frac{\partial u}{\partial x_i} dx + \int_{\Gamma_0} \xi_0 u d\Gamma(x) &= - \int_{\Omega} u \Delta v dx + \int_{\Gamma_0} \xi_0 u d\Gamma(x) \\ &= \int_{\Omega} \nabla u \cdot \nabla v dx + \int_{\Gamma_0} u v d\Gamma(x) = \int_{\Gamma_0} \left(\frac{\partial u}{\partial n} + u \right) v d\Gamma(x) \end{aligned}$$

hence for all ξ_0, \dots, ξ_N ,

$$\begin{aligned} \left| \sum_{i=1}^N \int_{\Omega} \xi_i \frac{\partial u}{\partial x_i} dx + \int_{\Gamma_0} \xi_0 u d\Gamma(x) \right| &\leq \left\| \frac{\partial u}{\partial n} + u \right\|_{L^1(\Gamma_0)} \|v\|_{L^\infty(\Gamma_0)} \\ &\leq C \left\| \frac{\partial u}{\partial n} + u \right\|_{L^1(\Gamma_0)} \left(\sum_{i=1}^N \|\xi_i\|_{L^q(\Omega)} + \|\xi_0\|_{L^q(\Gamma_0)} \right). \end{aligned}$$

For p the conjugate of q , $p < \frac{N}{N-1}$ since $q > N$, by the definition of the norm by duality, it follows that

$$\|u\|_{W^{1,p}(\Omega)} \leq C \left\| \frac{\partial u}{\partial n} + u \right\|_{L^1(\Gamma_0)} = C \|\psi\|_{L^1(\Gamma_0)} = C \|\psi\|_{\mathcal{M}_b(\Gamma_0)}.$$

Since the space $C(\Gamma_0)$ is dense in the space $\mathcal{M}_b(\Gamma_0)$, this completes the proof of proposition. \blacksquare

Furthermore, we have the following variational formulation

$$\int_{\Omega} \nabla u \cdot \nabla \xi dx + \int_{\Gamma_0} u \xi d\Gamma(x) = \int \xi d\mu \quad \forall \xi \in \mathcal{D}(\mathbb{R}^N) \quad (21)$$

For $N = 3$ the latter formulation remains valid, by density, for an arbitrary test function $\xi \in H^2(\Omega)$, since by the Sobolev imbedding theorem $H^2(\Omega) \subset C(\Omega)$ and the integral $\int \xi d\mu$ is well defined.

From proposition 4 we obtain the following result.

Proposition 5. *Given a sequence $\{\mu_n\}$ of Radon measures supported on Γ_0 , $\|\mu_n\|_{\mathcal{M}_b(\Gamma_0)} \leq C$, there exists a subsequence, still denoted $\{\mu_n\}$ and a Radon measure $\mu \in \mathcal{M}_b(\Gamma_0)$ such that*

$$\begin{aligned} \mu_n &\rightarrow \mu && \text{in } \mathcal{M}_b(\Gamma_0) \text{ weak-} (*), \\ u_n &\rightarrow u && \text{in } W^{1,p}(\Omega) \text{ weak-} (*), \\ u_n &\rightarrow u && \text{in } L^p(\Omega), \quad 1 < p < \frac{N}{N-1}, \end{aligned}$$

where by u_n we denote a solution to the problem $\mathcal{P}_1(\mu_n)$.

The proof of Proposition 5 is omitted here, it uses the theorem of Banach–Alaoglu and the same argument as in the proof of Proposition 3.

We shall consider the admissible measures of the form

$$\mu = \psi \delta_{\gamma}$$

with some regularity properties imposed on the density $\psi \in L^\infty(\gamma)$ and on the curve $\gamma = \text{supp} \mu$. The reason to consider such a class is the following, it is easy to construct a sequence γ_n such that the length of the curve γ_n is uniformly bounded, and the following convergence takes place $\delta_{\gamma_n} \rightarrow \psi \delta_{\gamma}$ weak-(*).

Example 1. Let us consider the family of curves, $n = 1, 2, \dots$,

$$\gamma_n = \{x_n(t), y_n(t)\} \quad t \in [0, 1]$$

where

$$x_n(t) = t, \quad y_n(t) = \begin{cases} 0, & t \in \left[0, \frac{1}{3}\right] \cup \left[\frac{2}{3}, 1\right]; \\ \frac{1}{n} \sin 3n\pi x, & t \in \left(\frac{1}{3}, \frac{2}{3}\right). \end{cases}$$

It can be shown that

$$\delta_{\gamma_n} \rightarrow \psi \delta_\gamma \quad \text{weak-}(*)$$

where

$$\gamma = \begin{cases} x = x(t) = t, & t \in [0, 1]; \\ y = y(t) = 0 & t \in [0, 1]. \end{cases}$$

and

$$\psi(t) = \begin{cases} 1, & t \in \left[0, \frac{1}{3}\right] \cup \left[\frac{2}{3}, 1\right]; \\ L, & t \in \left(\frac{1}{3}, \frac{2}{3}\right). \end{cases}$$

with $L = \int_0^1 \sqrt{1 + \cos^2 \pi t} dt > 1$. ■

Let us recall that for a sequence of Radon measures $\{\mu_n\}$ such that

$$\mu_n \rightarrow \mu \quad \text{weak-}(*)$$

we have in general the only information on the support of the limit

$$\text{supp} \mu \subset \limsup_{n \rightarrow \infty} \text{supp} \mu_n$$

where the lim sup is taken in the sense of Kuratowski.

For the choice we make $\mu = \psi \delta_\gamma$, since

$$\int \varphi d\mu = \int_\gamma \varphi(\sigma) \psi(\sigma) d\sigma$$

it follows that

$$\|\mu\|_{\mathcal{M}_b} \leq L(\gamma) \|\psi\|_{L^\infty}, \quad \text{where } L(\gamma) = \int_\gamma d\sigma$$

i.e., $L(\gamma)$ denotes the length of the curve γ .

Now let $\alpha > \frac{1}{2}$ and $M > 0$ be given constants, we introduce the set of admissible Radon measures of the following form

$$\begin{aligned} \mathfrak{M}_{\alpha, M} = & \{ \mu = \psi \delta_\gamma \mid \gamma = \{x(\cdot), y(\cdot)\} \in [W^{1, 2\alpha}(0, 1)]^2 \\ & \int_0^1 (x'^2(t) + y'^2(t))^\alpha dt \leq M, \int_0^1 |\psi(x(t), y(t))| \sqrt{x'^2(t) + y'^2(t)} dt \leq M \} \end{aligned}$$

Theorem 2. *Given a sequence $\mu_n \in \mathfrak{M}_{\alpha, M}$, there exists a subsequence, still denoted by μ_n , a measure $\mu \in \mathcal{M}_b(\Omega_0)$ and a curve γ such that*

$$\mu_n \rightarrow \mu \quad \text{in } \mathcal{M}_b(\Omega_0) \text{ weak-}(*)$$

where

$$\gamma = \text{supp}\mu = \{x(\cdot), y(\cdot)\} \in \left[W^{1,2\alpha}(0,1) \right]^2 .$$

Furthermore, if the following conditions are satisfied,

$$\|\psi_n\|_{L^p(\gamma_n)} \leq C \quad \text{for some } p > 1, \quad \text{with } C \leq M^{\frac{p+1}{2p}}$$

$$\text{and } \sqrt{x'^2(t) + y'^2(t)} \geq \beta > 0 \quad \text{for a.e. } t \in (0,1)$$

then there exists a function $\psi \in L^1(\gamma)$ such that

$$\mu = \psi \delta_\gamma$$

Proof. First, since $\{x_n\}, \{y_n\}$ are the bounded sequences in the Sobolev space $W^{1,2\alpha}(0,1)$ which is a reflexive Banach space, by the theorem of Rellich compactly imbedded in the space of continuous functions, it follows that there exist elements $x, y \in W^{1,2\alpha}(0,1)$ such that for subsequences, still denoted by $\{x_n\}, \{y_n\}$,

$$x_n \rightarrow x \quad \text{uniformly in } [0,1] \text{ and weakly in } L^{2\alpha}(0,1) \quad (22)$$

$$y_n \rightarrow y \quad \text{uniformly in } [0,1] \text{ and weakly in } L^{2\alpha}(0,1) \quad (23)$$

By lower semicontinuity of the norm we obtain

$$\int_0^1 (x'^2(t) + y'^2(t))^\alpha dt \leq \liminf \int_0^1 (x_n'^2(t) + y_n'^2(t))^\alpha dt \leq M$$

thus the curve $\gamma = \{x, y\}$ is the admissible support for the measure we are going to construct.

On the other hand, the sequence $\{\psi_n\}$ is bounded in $L^1(\gamma_n)$,

$$\int_0^1 |\psi_n(x_n(t), y_n(t))| \sqrt{x_n'^2(t) + y_n'^2(t)} dt \leq M$$

i.e., the function $t \mapsto \psi_n(x_n(t), y_n(t)) \sqrt{x_n'^2(t) + y_n'^2(t)}$ is bounded in $L^1(0,1)$. We denote by μ_n the measure defined in the following way,

$$\int v d\mu_n = \int_0^1 v(x_n(t), y_n(t)) \psi_n(x_n(t), y_n(t)) \sqrt{x_n'^2(t) + y_n'^2(t)} dt$$

for any $v \in C(\overline{\Omega}_0)$. Therefore, there exists a subsequence, still denoted by $\{\mu_n\}$, such that

$$\int v d\mu_n \rightarrow \int v d\mu \quad \text{for any } v \in C(\overline{\Omega}_0)$$

where the limit measure satisfies

$$\text{supp}\mu = \gamma$$

since

$$v(x_n(t), y_n(t)) \rightarrow v(x(t), y(t)) \quad \text{uniformly on } [0, 1].$$

Let us show the second part of the theorem. To this end we observe that by our assumptions the sequences $\{\psi_n(\cdot)\}$, $\left\{\sqrt{x_n'^2(\cdot) + y_n'^2(\cdot)}\right\}$ are bounded in $L^p(\gamma_n)$, $L^\alpha(0, 1)$, respectively, where $p > 1$ and $\alpha > \frac{1}{2}$.

Set $\beta = \frac{2\alpha p}{p+2\alpha-1}$, since $p > 1$ and $2\alpha > 1$ it follows that $\beta > 1$ and we verify that the sequence $\left\{\psi_n(x_n, y_n)\sqrt{x_n'^2 + y_n'^2}\right\}$ is bounded in $L^\beta(0, 1)$. Let

$$m = \frac{p+2\alpha-1}{2\alpha} \quad \text{and} \quad m^* = \frac{p+2\alpha-1}{p-1} \quad \text{so that} \quad \frac{1}{m} + \frac{1}{m^*} = 1,$$

by an application of the Hölder inequality it follows that

$$\begin{aligned} & \int_0^1 |\psi_n(x_n(t), y_n(t))|^\beta \left(x_n'^2(t) + y_n'^2(t)\right)^{\frac{\beta}{2}} dt \\ & \leq \left(\int_0^1 |\psi_n(x_n(t), y_n(t))|^{\beta m} \left(x_n'^2(t) + y_n'^2(t)\right)^{\frac{\alpha m}{p+2\alpha-1}} dt\right)^{\frac{1}{m}} \cdot \left(\int_0^1 \left(x_n'^2(t) + y_n'^2(t)\right)^{\frac{\alpha p - \alpha}{p+2\alpha-1} \cdot m^*} dt\right)^{\frac{1}{m^*}} \\ & = \left(\int_0^1 |\psi_n(x_n, y_n)|^p \sqrt{x_n'^2 + y_n'^2} dt\right)^{\frac{1}{m}} \cdot \left(\int_0^1 \left(x_n'^2 + y_n'^2\right)^\alpha dt\right)^{\frac{1}{m^*}} \\ & = \|\psi_n\|_{L^p(\gamma_n)}^{\frac{p}{m}} \cdot \|x_n'^2 + y_n'^2\|_{L^\alpha(0,1)}^{\frac{\alpha}{m^*}} \leq C^{\frac{p}{m}} M^{\frac{\alpha}{m^*}} \end{aligned}$$

the latter inequality follows by our assumptions, hence

$$\|\psi_n(x_n, y_n)\sqrt{x_n'^2 + y_n'^2}\|_{L^\beta(0,1)} \leq C^{\frac{p}{m\beta}} M^{\frac{\alpha}{m^*\beta}} = CM^{\frac{p-1}{2p}} \leq M.$$

Therefore, there exists an element $\varphi \in L^\beta(0, 1)$ such that

$$\psi_n(x_n, y_n)\sqrt{x_n'^2 + y_n'^2} \rightarrow \varphi \quad \text{weakly in } L^\beta(0, 1),$$

$$\|\varphi\|_{L^\beta(0,1)} \leq \liminf \|\psi_n(x_n, y_n) \sqrt{x_n'^2 + y_n'^2}\|_{L^\beta(0,1)} \leq M$$

and

$$\begin{aligned} \int v d\mu_n &= \int_0^1 v(x_n(t), y_n(t)) \psi_n(x_n(t), y_n(t)) \sqrt{x_n'^2(t) + y_n'^2(t)} dt \\ &\rightarrow \int v d\mu = \int_0^1 v(x(t), y(t)) \varphi(t) dt \end{aligned}$$

so that we define

$$\psi(x(t), y(t)) = \frac{\varphi(t)}{\sqrt{x'^2(t) + y'^2(t)}}$$

with $\psi \in L^1(\gamma)$ and

$$\|\psi\|_{L^1(\gamma)} = \int_0^1 |\psi(x(t), y(t))| \sqrt{x'^2(t) + y'^2(t)} dt = \int_0^1 |\varphi(t)| dt \leq \|\varphi\|_{L^\beta(0,1)} \leq M$$

thus

$$\mu = \psi \delta_\gamma \in \mathcal{M}_{\alpha, M} .$$

■

Conclusion. According to theorem 2, there exists a solution to the minimization problem

$$\min_{\mu \in \mathcal{M}_{\alpha, M}} J(\mu)$$

for a class of cost functionals, e.g. $J(\mu) = \int_\Sigma (u - \bar{u})^2 dx$, u being a solution to $\mathcal{P}_1(\mu)$.

4. Optimality conditions

We start with the auxiliary results on the differentiability of the following shape functional

$$\gamma \rightarrow \int_\gamma \mathcal{G} d\gamma .$$

We assume that the function $\mathcal{G} \in L^1(\gamma)$ may depend on the curve γ . We use the same approach as in the case of thin shell, where we consider a curve γ on the manifold, here Ω_0 is flat set.

Let the sufficiently smooth mapping $\mathcal{F}_s : \mathbb{R}^3 \mapsto \mathbb{R}^3$ be given, $s \in [0, \delta)$ is a parameter, such that $F_s = \mathcal{F}_s|_Q$ for any $s \in [0, \delta)$ satisfies the assumptions of definition 1, i.e.

$$\begin{aligned} F_s(Q) &= \mathcal{O} & F_s(I) &= \gamma \\ \|F_s\|_{W^{1,\infty}(Q)} &\leq L_1 & \|F_s^{-1}\|_{W^{1,\infty}(\mathcal{O})} &\leq L_2 \end{aligned}$$

Given parametrization $\{x_s(t), y_s(t)\}$, $t \in [0, 1]$, of the curve γ_s , we denote

$$j(s) = \int_{\gamma_s} \mathcal{G}_s d\gamma_s = \int_0^1 \mathcal{G}_s(x_s(t), y_s(t)) \sqrt{x_s'^2(t) + y_s'^2(t)} dt$$

The derivative takes the form

$$\begin{aligned} j'(s) &= \int_0^1 \left\{ \frac{\partial \mathcal{G}_s}{\partial s} + \nabla \mathcal{G}_s(x_s(t), y_s(t)) \cdot \xi_s(t) \right\} \sqrt{x_s'^2(t) + y_s'^2(t)} dt \\ &\quad + \int_0^1 \mathcal{G}_s(x_s(t), y_s(t)) \tau_s(t) \cdot \frac{d\xi_s}{dt}(t) dt \end{aligned}$$

where $\tau_s(t) = \frac{(x_s'(t), y_s'(t))}{\sqrt{x_s'^2(t) + y_s'^2(t)}}$ is the unit tangent vector to γ and $\xi_s(t) = \frac{d}{ds}(x_s(t), y_s(t))$.

Under regularity assumptions, after integration by parts the latter integral can be rewritten in the following form

$$\begin{aligned} &\int_0^1 \mathcal{G}_s(x_s(t), y_s(t)) \tau_s(t) \cdot \frac{d\xi_s}{dt}(t) dt \\ &= - \int_0^1 \left\{ \nabla \mathcal{G}_s(x_s(t), y_s(t)) \cdot (x_s(t), y_s(t)) \tau_s(t) \cdot \xi_s(t) + \mathcal{G}_s(x_s(t), y_s(t)) \frac{d\tau_s}{dt}(t) \cdot \xi_s(t) \right\} dt \\ &\quad + \mathcal{G}_s(x_s(1), y_s(1)) \tau_s(1) \cdot \xi_s(1) - \mathcal{G}_s(x_s(0), y_s(0)) \tau_s(0) \cdot \xi_s(0) \end{aligned}$$

On the other hand, we can use the material derivative method to obtain the same derivative $j'(s)$. Namely, we introduce the vector field

$$V(s, x, y, z) = \left(\frac{\partial \mathcal{F}_s}{\partial s} \circ \mathcal{F}_s^{-1} \right) (x, y, z)$$

and assume that the support of the vector field is included in a small neighbourhood $\mathcal{O}(\gamma)$ of the curve γ in \mathbb{R}^3 . Furthermore, we assume that for $(x, y, z) \in \mathcal{O}(\gamma)$ and sufficiently small $z \in (-\varepsilon, \varepsilon)$, $\varepsilon > 0$, the field is of the following form

$$V(s, x, y, z) = \begin{pmatrix} V_1(s, x, y) \\ V_1(s, x, y) \\ 0 \end{pmatrix} = V(s, x, y, 0)$$

The shape functional we consider takes the form

$$J(\gamma) = \int_{\gamma} \mathcal{G} d\gamma .$$

With the vector field V we associate the mapping

$$T_s(V) : \mathbb{R}^3 \mapsto \mathbb{R}^3 ,$$

in particular, under our assumptions on the support of the field V , $\text{supp}V \subset \mathcal{O}(\gamma)$, it follows that $T_s(V) \equiv I$ on $\mathbb{R}^3 \setminus \mathcal{O}(\gamma)$, where I denotes the identity mapping.

Let us define the Eulerian semiderivative

$$dJ(\gamma; V) = \lim_{s \downarrow 0} \frac{1}{s} (J(T_s(\gamma)) - J(\gamma)) .$$

For

$$\gamma_s = T_s(\gamma), \quad s \in [0, \delta)$$

it follows that

$$j'(0^+) = dJ(\gamma; V)$$

and therefore, by an application of the structure theorem for the shape gradient, we obtain

$$\begin{aligned} dJ(\gamma; V) &= \int_0^1 \left\{ \frac{\partial \mathcal{G}_s}{\partial s} \Big|_{s=0} + \nabla \mathcal{G}(x(t), y(t)) \cdot \xi(t) \right\} \sqrt{x'^2(t) + y'^2(t)} dt \\ &\quad + \int_0^1 \mathcal{G}(x(t), y(t)) \frac{d\tau}{dt}(t) \cdot \xi(t) dt \\ &\quad + \mathcal{G}(x(1), y(1)) \tau(1^-) \cdot \xi(1) - \mathcal{G}(x(0), y(0)) \tau(0^+) \cdot \xi(0) \end{aligned}$$

since $V(s, x(t), y(t), 0) = (\xi_s(t), 0)$ for $t \in [0, 1]$, and the vector $\tau(t) \in \Omega_0$, $t \in (0, 1)$, is tangent to γ . If $\nu(t) \in \Omega_0$, $t \in (0, 1)$, denotes the normal vector field on γ , the equivalent form of the first integral reads

$$\begin{aligned} &\int_0^1 \left\{ \frac{\partial \mathcal{G}_s}{\partial s} \Big|_{s=0} + \nabla \mathcal{G}(x(t), y(t)) \cdot \xi(t) \right\} \sqrt{x'^2(t) + y'^2(t)} dt \\ &= \int_0^1 \left\{ \frac{\partial \mathcal{G}_s}{\partial s} \Big|_{s=0} + [\nabla \mathcal{G}(x(t), y(t)) \cdot \nu(t)] \xi(t) \cdot \nu(t) \right\} \sqrt{x'^2(t) + y'^2(t)} dt \end{aligned}$$

since the integral part of $dJ(\gamma; V)$, by the structure theorem, depends only on the normal component $V(0, x(t), y(t), 0) \cdot n = \xi(t) \cdot \nu(t)$, $t \in (0, 1)$, of the field $V(0, x(t), y(t), 0)$. We denote

$$\begin{aligned} \int_{\gamma} \dot{\mathcal{G}} d\gamma &= \int_0^1 \left\{ \frac{\partial \mathcal{G}_s}{\partial s} \Big|_{s=0} + \nabla \mathcal{G}(x(t), y(t)) \cdot \xi(t) \right\} \sqrt{x'^2(t) + y'^2(t)} dt \\ \int_{\gamma} \mathcal{G} \tau' \cdot V d\gamma &= \int_0^1 \mathcal{G}(x(t), y(t)) \frac{d\tau}{dt}(t) \cdot \xi(t) dt \\ (x(1), y(1)) &= (x_1, y_1), \quad (x(0), y(0)) = (x_0, y_0) \end{aligned}$$

Proposition 6. *The shape functional $J(\gamma) = \int_{\gamma} \mathcal{G} d\gamma$ is shape differentiable, the Eulerian semiderivative takes the following form*

$$\begin{aligned} dJ(\gamma; V) &= \int_{\gamma} \dot{\mathcal{G}} d\gamma + \int_{\gamma} \mathcal{G} \tau' \cdot V d\gamma + \mathcal{G}(x_1, y_1) \tau(x_1^-, y_1^-) \cdot V(0, x_1, y_1, 0) \\ &\quad - \mathcal{G}(x_0, y_0) \tau(x_0^+, y_0^+) \cdot V(0, x_0, y_0, 0) \end{aligned}$$

where $\dot{\mathcal{G}}$ denotes the material derivative of \mathcal{G} in the direction of the vector field V .

Remark 6. In particular for $\mathcal{G} = 1$, $J(\gamma) = |\gamma| = \int_{\gamma} d\gamma$,

$$dJ(\gamma; V) = \int_{\gamma} \tau \cdot DV \cdot \tau d\gamma .$$

We have the property

$$|\gamma_s| = |T_s(\gamma)| = \int_{\gamma_s} d\gamma_s = \int_{\gamma} d\gamma \quad \gamma_s = T_s(\gamma)$$

provided that the vector field V satisfies the equation

$$\int_{\gamma} \tau' \cdot V d\gamma + \tau(x_1^-, y_1^-) \cdot V(0, x_1, y_1, 0) - \tau(x_0^+, y_0^+) \cdot V(0, x_0, y_0, 0) = 0$$

■

Now, we are in the position to obtain the shape differentiability of solutions to the problem $\mathcal{P}_1(\gamma)$.

We denote $\Sigma_s = T_s(\Sigma)$, $u_s \in W^{1,p}(\Sigma_s)$ the unique solution to the following integral identity

$$\int_{\Sigma_s} \nabla u_s \cdot \nabla \varphi d\Sigma_s + \int_{\Omega_0^s} u_s \varphi d\sigma_s + \int_{\Omega_1^s} u_s \varphi d\sigma_s = \int_{\Omega_0^s} u_0 \varphi d\sigma_s + \int_{\Omega_1^s} u_1 \varphi d\sigma_s + \int_{\gamma_s} \varphi d\gamma_s$$

for all $\varphi \in W^{1,q}(\Sigma_s)$, where $\Omega_i^s = T_s(\Omega_i)$, $i = 0, 1$, $\gamma_s = T_s(\gamma)$.

The integral identity is transported to the fixed domain Σ , so we denote $u^s = u_s \circ T_s \in W^{1,p}(\Sigma)$, set $\varphi = v \circ T_s^{-1}$, and by standard change of variables it follows that u^s is the unique solution to the following integral identity,

$$\begin{aligned} & \int_{\Sigma} \langle A(s) \cdot \nabla u_s, \nabla v \rangle_{\mathbb{R}^3} d\Sigma + \int_{\Omega_0} u^s v \omega(s) d\sigma + \int_{\Omega_1} u^s v \omega(s) d\sigma \\ & = \int_{\Omega_0} u_0^s v \omega(s) d\sigma + \int_{\Omega_1} u_1^s v \omega(s) d\sigma + \int_{\gamma} v \rho(s) d\gamma \end{aligned}$$

for all $v \in W^{1,q}(\Sigma)$, where the matrix $A(s)$, the boundary terms $\omega(s), \rho(s)$ are given, sufficiently smooth functions of space variables and $s \in [0, \delta)$,

$$\begin{aligned} A(s) &= \det(DT_s) DT_s^{-1} \cdot * DT_s^{-1} \\ \omega(s) &= \|\det(DT_s)^* DT_s^{-1} \cdot n\|_{\mathbb{R}^3} \\ \rho(s) &= \left(\frac{x_s'^2(t) + y_s'^2(t)}{x'^2(t) + y'^2(t)} \right)^{\frac{1}{2}}, \quad (x(t), y(t)) \in \gamma, \quad \gamma_s = T_s(\gamma), \quad t \in (0, 1). \end{aligned}$$

By an application of the implicit function theorem for solutions of the latter integral identity we obtain the existence of the weak material derivative in $W^{1,p}(\Sigma)$, $1 < p < \frac{3}{2}$,

$$\dot{u} = \lim_{s \downarrow 0} \frac{1}{s} (u^s - u).$$

The material derivative $\dot{u} \in W^{1,p}(\Sigma)$ satisfies the following integral identity

$$\begin{aligned} & \int_{\Sigma} \nabla \dot{u} \cdot \nabla v d\Sigma + \int_{\Sigma} \langle A'(0) \cdot \nabla u, \nabla v \rangle_{\mathbb{R}^3} d\Sigma + \int_{\Omega_0} \dot{u} v d\sigma + \int_{\Omega_0} u v \omega'(0) d\sigma + \int_{\Omega_1} \dot{u} v d\sigma \\ & + \int_{\Omega_1} u v \omega'(0) d\sigma = \int_{\Omega_0} (\dot{u}_0 + u_0 \omega'(0)) v d\sigma + \int_{\Omega_1} (\dot{u}_1 + u_1 \omega'(0)) v d\sigma + \int_{\gamma} v \rho'(0) d\gamma, \end{aligned}$$

where we denote

$$\begin{aligned} A'(0) &= \operatorname{div} V(0) I - DV(0) - * DV(0) \\ \omega'(0) &= \operatorname{div} V(0) - \langle DV(0) \cdot n, n \rangle_{\mathbb{R}^3} \\ \rho'(0) &= \tau \cdot DV(0) \cdot \tau \end{aligned}$$

Finally, the weak shape derivative $u' = \dot{u} - \nabla u \cdot V$ in $L^p(\Sigma)$, $1 < p < \frac{3}{2}$, satisfies the following integral identity

$$\begin{aligned} - \int_{\Sigma} u' \Delta v dx &= - \int_{\gamma} v \tau' \cdot V d\gamma + v(x_1, y_1) \tau(x_1^-, y_1^-) \cdot V(0, x_1, y_1, 0) \\ &\quad - v(x_0, y_0) \tau(x_0^+, y_0^+) \cdot V(0, x_0, y_0, 0) \end{aligned}$$

for all test functions $v \in W^{2,q}(\Sigma)$, $\frac{\partial v}{\partial n} = 0$ on Γ , $\frac{\partial v}{\partial n} + v = 0$ on $\Omega_0 \cup \Omega_1$.

Theorem 3. *A solution to the minimization problem*

$$\inf_{\gamma \in \mathcal{F}_L} J(\gamma)$$

satisfies the first order necessary optimality conditions

$$dJ(\gamma; V) = 0$$

for all admissible vector fields V , where

$$\begin{aligned} dJ(\gamma; V) &= 2 \int_{\Sigma} (u(\gamma) - u_d) \dot{u} d\Sigma + 2 \int_{\Sigma} \langle \nabla u(\gamma) - \nabla u_d, \nabla \dot{u} \rangle_{\mathbb{R}^3} d\Sigma \\ &\quad - 2 \int_{\Sigma} \langle {}^* DV \cdot \nabla(u(\gamma) - u_d), \nabla(u(\gamma) - u_d) \rangle_{\mathbb{R}^3} d\Sigma \\ &\quad + \int_{\Sigma} \left(|\nabla(u(\gamma) - u_d)|^2 + |u(\gamma) - u_d|^2 \right) \operatorname{div} V d\Sigma \end{aligned}$$

Remark 7. For any vector field V such that

$$V(0) \cdot \nu = 0 \text{ on } \gamma, \quad V(0, A) = V(0, B) = 0$$

it follows that $dJ(\gamma; V) = 0$, therefore we obtain the following Green formula for such fields

$$\begin{aligned} 0 &= 2 \int_{\Sigma} (u(\gamma) - u_d) \nabla u \cdot V(0) d\Sigma \\ &\quad + 2 \int_{\Sigma} \langle \nabla u(\gamma) - \nabla u_d, \nabla(\nabla u \cdot V(0)) \rangle_{\mathbb{R}^3} d\Sigma \\ &\quad - 2 \int_{\Sigma} \langle {}^* DV \cdot \nabla(u(\gamma) - u_d), \nabla(u(\gamma) - u_d) \rangle_{\mathbb{R}^3} d\Sigma \\ &\quad + \int_{\Sigma} \left(|\nabla(u(\gamma) - u_d)|^2 + |u(\gamma) - u_d|^2 \right) \operatorname{div} V d\Sigma \end{aligned}$$

Remark 8.

In the particular case of the cost functional $I(\gamma) = \int_{\Omega_1} (u(\gamma) - u_d)^2 d\Omega$, it follows that

$$dI(\gamma; V) = 2 \int_{\Omega_1} (u(\gamma) - u_d) u'(\gamma; V) d\Omega.$$

References

- Adams R.A. (1975): *Sobolev Spaces*. — Academic Press, New York.
- Bucur, D. and Zolesio, J.P.: *N-dimensional shape optimization under the capacitary constraints* to appear in J. Diff. Eq.
- Daniliuk, I.I. (1975): *Nonsmooth boundary value problems in the plane*. — Nauka, Moscou (in Russian).
- Hoffmann, K.-H. and Sokolowski, J.(1991): *Domain optimization problem for parabolic equation*. — DFG Report No. 342, Augsburg, Germany.
- Hoffmann, K.-H. and Sokolowski, J.(1994): *Interface optimization problems for parabolic equations*. — Control and Cybernetics, v. 23, pp. 445–452.
- Lions J.L., Magenes E. (1968): *Problèmes aux limites non homogènes*. — Dunod, Paris.
- Sokolowski, J. and Zolesio, J.-P.(1992): *Introduction to Shape Optimization. Shape sensitivity analysis*. — Springer Verlag, New York.
- Ziemer P.W. (1989): *Weakly Differentiable Functions*. — Springer Verlag, New York.



Unité de recherche INRIA Lorraine, Technopôle de Nancy-Brabois, Campus scientifique,
615 rue du Jardin Botanique, BP 101, 54600 VILLERS LÈS NANCY
Unité de recherche INRIA Rennes, Irisa, Campus universitaire de Beaulieu, 35042 RENNES Cedex
Unité de recherche INRIA Rhône-Alpes, 46 avenue Félix Viallet, 38031 GRENOBLE Cedex 1
Unité de recherche INRIA Rocquencourt, Domaine de Voluceau, Rocquencourt, BP 105, 78153 LE CHESNAY Cedex
Unité de recherche INRIA Sophia-Antipolis, 2004 route des Lucioles, BP 93, 06902 SOPHIA-ANTIPOLIS Cedex

Éditeur
INRIA, Domaine de Voluceau, Rocquencourt, BP 105, 78153 LE CHESNAY Cedex (France)
ISSN 0249-6399