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*Mathematical and Numerical Studies of  
1D Non Linear Ferromagnetic Materials*

Patrick Joly, Olivier Vacus

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*R*apport  
de recherche





# Mathematical and Numerical Studies of 1D Non Linear Ferromagnetic Materials

Patrick Joly, Olivier Vacus \*

Thème 4 — Simulation et optimisation  
de systèmes complexes  
Projet Ondes

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**Abstract:** In this paper we are interested in the numerical modeling of absorbing ferromagnetic materials obeying the non-linear Landau-Lifchitz-Gilbert law with respect to the propagation and scattering of electromagnetic waves. In this work we consider the 1D problem. We first show that the corresponding Cauchy problem has a unique global solution. We then derive a numerical scheme based on an appropriate modification of Yee's scheme, that we show to preserve some important properties of the continuous model such as the conservation of the norm of the magnetization and the decay of the electromagnetic energy. The stability is proved under a suitable CFL condition. Eventually some numerical results corresponding to the 1D model are presented.

**Key-words:** Maxwell's equations - Landau-Lifchitz-Gilbert law - Hille-Yoshida theorem - finite difference scheme

(Résumé : *tsvp*)

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# Etude mathématique et numérique de matériaux ferromagnétiques non linéaires 1D

**Résumé :** On s'intéresse dans ce rapport à la modélisation numérique dans le cas monodimensionnel de matériaux ferromagnétiques absorbants obéissants à la loi non linéaire de Landau-Lifchitz-Gilbert, et ce relativement à la propagation et à la diffraction des ondes électromagnétiques. Nous montrons d'abord que le problème de Cauchy associé admet une unique solution puis nous construisons un schéma numérique adapté, dérivé du classique schéma de Yee, dont nous établissons la stabilité. Quelques résultats d'expériences numériques sont finalement présentés.

**Mots-clé :** Equations de Maxwell - loi de Landau-Lifchitz-Gilbert - théorème de Hille-Yoshida - schéma aux différences finies

# Chapter 1

## Introduction

### 1.1 Physical introduction

Without entering into the details of the delicate theory of ferromagnetism (see for this [8], [13] or [16]), we can briefly describe the physics of our problem.

#### The ferromagnetic zone

What we call a “ferromagnetic material” may be simply defined as one that possesses a spontaneous magnetization  $\mathbf{M}$ . (Note that we always use bold letters for vectors.) These materials like ferrites or garnets are widely used in the microwave domain as shifters or circulators for instance. In this paper we would particularly like to study the propagation of electromagnetic waves inside such materials which are said to be absorbing materials and whose behavior is expected to be strongly dependent on the frequency. We will see what sense we can give to these assertions.

A ferromagnetic zone is a region in which  $\mathbf{M} \neq 0$ , and we denote by  $\Omega$  all such regions in space. Such a region will be examined in this paper on a scale large enough to permit the use of continuous magnetization vector  $\mathbf{M}(\mathbf{x})$  at any time ; so we do not consider *a priori* any domain or wall concept ([7], [2]). The magnetization  $\mathbf{M}$  is assumed to obey the Landau-Lifchitz-Gilbert equation (LLG) :

$$(1.1) \quad \frac{\partial \mathbf{M}}{\partial t} = |\gamma| \mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M} + \frac{\alpha}{|\mathbf{M}|} \mathbf{M} \times \frac{\partial \mathbf{M}}{\partial t}$$

where  $\times$  denotes the usual vector product,  $|\cdot|$  is the usual Euclidean norm,  $\gamma$  is the gyromagnetic factor ( $\gamma < 0$ ) and in the last term,  $\alpha$  is called the damping constant, as we shall explain later.

It is first straightforward to show from the (LLG) equation – at least formally – the conservation in time of the norm of the magnetization at any point :

$$\frac{\partial \mathbf{M}}{\partial t} \cdot \mathbf{M} = \frac{\partial}{\partial t} \{ |\mathbf{M}|^2 \} = 0$$

Consequently, the LLG equation can be simplified to

$$(1.2) \quad \frac{\partial \mathbf{M}}{\partial t} = |\gamma| \mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M} + \frac{\alpha}{|\mathbf{M}_0|} \mathbf{M} \times \frac{\partial \mathbf{M}}{\partial t} \quad \text{where} \quad \mathbf{M}_0(\mathbf{x}) = \mathbf{M}(\mathbf{x}, t = 0)$$

#### The total magnetic field $\mathbf{H}_T(\mathbf{H}, \mathbf{M})$

The total magnetic field  $\mathbf{H}_T(\mathbf{H}, \mathbf{M})$  is defined as a sum :

$$(1.3) \quad \mathbf{H}_T(\mathbf{H}, \mathbf{M}) = \mathbf{H} + \mathbf{H}_{eff}(\mathbf{M})$$

where  $\mathbf{H}$  is the magnetic field appearing in Maxwell’s equations

$$(1.4) \quad \begin{cases} \frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E} \\ \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t} = \nabla \times \mathbf{H} \end{cases}$$

the relation between  $\mathbf{H}$ ,  $\mathbf{B}$  and  $\mathbf{M}$  being

$$(1.5) \quad \mathbf{H} = \frac{\mathbf{B}}{\mu_0} - \mathbf{M} \text{ or } \mathbf{B} = \mu_0(\mathbf{H} + \mathbf{M})$$

In (1.4) and (1.5), we denote by :

- $\mathbf{B}$  the magnetic induction,
- $\mathbf{E}$  the electric field ,
- $\varepsilon_0$  and  $\mu_0$ , the electric permittivity and the magnetic permeability in the vacuum.

Moreover, we have preferred to present in this paper the non conductive case but no new difficulty would arise if the conductivity  $\sigma$  was different from 0.

It is important to note that the (LLG) equation is a non-linear differential equation in time, whereas people often study linear materials in the frequency domain where it is possible to define the magnetic susceptibility  $\chi$  and to write that  $\mathbf{M} = \chi\mathbf{H}$ . In our case, the relation which links  $\mathbf{M}$  and  $\mathbf{H}$  is implicitly defined by the coupled system (LLG) equation–Maxwell’s equations.

Last of all, the effective field  $\mathbf{H}_{eff}(\mathbf{M})$  is the result of several contributions. In the mode we are interested in, we retain two of them :

$$(1.6) \quad \mathbf{H}_{eff}(\mathbf{M}) = \mathbf{H}_s + \mathbf{H}_a(\mathbf{M})$$

where  $\mathbf{H}_s$  is a static field and  $\mathbf{H}_a(\mathbf{M})$  is a field of anisotropy, a field derived from an energy of anisotropy. We define these fields in a ferromagnetic zone  $\Omega$  at section 2.1.

## The mathematical model

As a model of insulating non-linear ferrites, we get the following Cauchy’s problem from (1.2), (1.4) and (1.5) :

$$(1.7) \quad \left\{ \begin{array}{l} \frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E} \\ \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t} = \nabla \times \mathbf{H} \\ \mathbf{H} = \frac{\mathbf{B}}{\mu_0} - \mathbf{M} \\ \frac{\partial \mathbf{M}}{\partial t} = |\gamma| \mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M} + \frac{\alpha}{|\mathbf{M}|} \mathbf{M} \times \frac{\partial \mathbf{M}}{\partial t} \end{array} \right. \quad \text{with} \quad \left\{ \begin{array}{l} \mathbf{B}(\mathbf{x}, t = 0) = \mathbf{B}_0(\mathbf{x}) \\ \mathbf{E}(\mathbf{x}, t = 0) = \mathbf{E}_0(\mathbf{x}) \\ \mathbf{M}(\mathbf{x}, t = 0) = \mathbf{M}_0(\mathbf{x}) \end{array} \right.$$

or, if we simplify the system by eliminating the magnetic induction  $\mathbf{B}$ ,

$$(1.8) \quad \left\{ \begin{array}{l} \mu_0 \frac{\partial \mathbf{H}}{\partial t} + \mu_0 \frac{\partial \mathbf{M}}{\partial t} = -\nabla \times \mathbf{E} \\ \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t} = \nabla \times \mathbf{H} \\ \frac{\partial \mathbf{M}}{\partial t} = |\gamma| \mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M} + \frac{\alpha}{|\mathbf{M}|} \mathbf{M} \times \frac{\partial \mathbf{M}}{\partial t} \end{array} \right. \quad \text{with} \quad \left\{ \begin{array}{l} \mathbf{H}(\mathbf{x}, t = 0) = \mathbf{H}_0(\mathbf{x}) \\ \mathbf{E}(\mathbf{x}, t = 0) = \mathbf{E}_0(\mathbf{x}) \\ \mathbf{M}(\mathbf{x}, t = 0) = \mathbf{M}_0(\mathbf{x}) \end{array} \right.$$

## 1.2 Presentation of the paper

Two main mathematical issues naturally arise : 1. Is it possible to show the existence, and then the uniqueness, of a solution to this non-linear problem ? , 2. How can we discretize the equations of this problem to get a numerical modeling of ferromagnetic materials which respects the main properties of the continuous problem ?

These are the two questions we intend to address in this paper. Concerning the first point, as far as we know, the only paper where a partial answer is given is due to Visintin (see [17]) who has established the existence of weak solutions to a problem very close to (1.7) ; however he considers a new field in addition in  $\mathbf{H}_{eff}(\mathbf{M})$ , derived for an exchange energy, and this assumption ensures more regularity for  $\mathbf{M}$ . His existence result and some secondary properties are proved for the corresponding variational formulation by using a Galerkin method. But he can not allow  $\alpha$  to be 0, and it is also said in [17] that “the possible uniqueness of the solution is an open question”.

Concerning the second point, the situation is even clearer : we have not found in the literature any satisfying answer to the question of the discretization of the complete non-linear problem (1.7). All the works on the subject mention at least one of the two following assumptions :

1. The damping is neglected and  $\alpha$  is taken equal to 0 ([9], [6], [5], [19]) ;
2. the material is supposed to be saturated and either the internal field is *a priori* known, or a small signal approximation is made, which leads to a linearization of the equations ([4], [3], [14], [9], [6], [5]).

In the present paper, we establish more complete results concerning both aspects of the problem in the one dimensional case – which is of course a limitation. However all that we do concerning the numerical analysis can be generalized to more space dimensions ([11]).

We shall show by applying the theory of semi-groups the existence of a unique maximal local strong solution to (1.7) and then we establish suitable *a priori* estimates to guarantee that this solution is indeed a maximal global strong solution (the two cases  $\alpha = 0$  and  $\alpha \neq 0$  are treated separately). This theoretical result requiring quite technical proofs, they must be seen as a complementary result to the theorem presented in [12], about existence and uniqueness of weak solutions to (1.7). (In [12], solutions are looked for in  $L^2(\mathbb{R}) \cap L^\infty(\mathbb{R})$  instead of  $H^1(\mathbb{R})$  in the present paper.) In particular, in the case  $\alpha \neq 0$ , uniqueness will be ensured by results of [12].

In the discrete case, we propose an approach based on the Yee's scheme ([18]) and an original way to solve numerically the discrete (LLG) equation. We show that this method ensures the conservation of the norm of the magnetization. We also establish the discrete equivalent of the *a priori* estimates of the continuous case which guarantees in particular the stability of the scheme under a classical Courant-Friedrichs-Levy condition.

Our paper is organized as follows. In section 2, we develop the machinery of semi-groups to study the existence and the uniqueness of solutions to the problem (1.7). In section 3, we give both time and space discretizations of the equations and their main properties. In section 4, a representative collection of numerical experiments is presented.



## Chapter 2

# Mathematical results

In this section we mainly work from the formulation (1.8) involving the three vectorial unknowns  $\mathbf{E}$ ,  $\mathbf{H}$  and  $\mathbf{M}$ .

### 2.1 About the effective field

In this paragraph we make precise the mathematical definition of the two local fields we consider in our work.

1. The static field  $\mathbf{H}_s$  is a vector field that is constant in time. For technical reason,  $\mathbf{H}_s$  is chosen so that :

$$(2.1) \quad \mathbf{H}_s \in H^1(\mathbf{R})$$

which does not constitute an actual restriction.

2. The field of anisotropy is derived as follows from the energy of anisotropy  $\mathcal{E}_a$  – we assume  $\mathcal{E}_a \in C^1(L^2(\Omega); \mathbf{R}_+)$ . For all  $\delta\mathbf{M} \in L^2(\Omega)$ , we have

$$(2.2) \quad \langle \mathcal{E}'_a(\mathbf{M}), \delta\mathbf{M} \rangle = -\mu_0 (\mathbf{H}_a(\mathbf{M}), \delta\mathbf{M})_{L^2}$$

which ensures that  $\mathbf{H}_a(\mathbf{M}) \in C^0(L^2(\Omega)^3; L^2(\Omega)^3)$ . For simplicity in this paper we consider a special case (see [7]) where

$$(2.3) \quad \mathcal{E}_a = \mu_0 \frac{K}{2} |P(\mathbf{M})|^2$$

where  $K$  is a positive constant depending on the material and  $P(\mathbf{M})$  the projection of  $\mathbf{M}$  on the plane perpendicular to the “easy axis” which is simply a privileged direction of the material, not depending on time. Consequently

$$(2.4) \quad \mathbf{H}_a(\mathbf{M}) = -KP(\mathbf{M})$$

**Remark 1** *If we introduce  $\mathbf{p}$ , a unit vector along the easy axis, we see that*

$$(2.5) \quad \mathbf{H}_a(\mathbf{M}) = -(\mathbf{M} - K(\mathbf{p} \cdot \mathbf{M})\mathbf{p})$$

*This new form will be particularly useful in the numerical part because it leads to*

$$(2.6) \quad \mathbf{H}_a(\mathbf{M}) \times \mathbf{M} = K(\mathbf{p} \cdot \mathbf{M})\mathbf{p} \times \mathbf{M}$$

*It can be seen in (2.6) that  $\mathbf{H}_a(\mathbf{M})$  can be also taken equal to  $K(\mathbf{p} \cdot \mathbf{M})\mathbf{p}$ .*

As a consequence of these definitions an important property of  $\mathbf{H}_{eff}(\mathbf{M})$  is to be an affine continuous function of  $\mathbf{M}$  – That is to say :  $[\mathbf{H}_{eff}(\mathbf{M}) - \mathbf{H}_0]$  is linear in  $\mathbf{M}$ .

### 2.2 Preliminary results independent on the space dimension

The results that we give in this paragraph are obtained **regardless** both the dimension of the problem and the geometry of the ferromagnetic zone  $\Omega$ . We denote by  $\mathbf{x} = (x, y, z)$  the current point of  $\mathbf{R}^3$ .

### 2.2.1 A new form of the (LLG) equation

For later use in the mathematical analysis, it would be useful to express  $\frac{\partial \mathbf{M}}{\partial t}$  as a function of  $\mathbf{M}$  and  $\mathbf{H}$ . We shall use the :

**Lemma 1** *Given any vector  $\mathbf{a}$ , function  $A : \mathbf{x} \mapsto \mathbf{x} + \mathbf{a} \times \mathbf{x}$  is a homeomorphism from  $\mathbb{R}^3$  to  $\mathbb{R}^3$  and*

$$(2.7) \quad A^{-1}(\mathbf{y}) = \frac{\mathbf{y} + (\mathbf{a} \cdot \mathbf{y})\mathbf{a} - \mathbf{a} \times \mathbf{y}}{1 + |\mathbf{a}|^2}$$

Moreover,  $A^{-1}$  is a contraction.

*Proof.* Continuity of  $A$  is clear. Let  $(\mathbf{x}, \mathbf{y}) \in (\mathbb{R}^3)^2$  be such that  $A(\mathbf{x}) = \mathbf{y}$ . Taking the scalar product by  $\mathbf{a}$ , we get  $(\mathbf{a} \cdot \mathbf{x}) = (\mathbf{a} \cdot \mathbf{y})$ , and taking the vector product by  $\mathbf{a}$ ,

$$\mathbf{a} \cdot \mathbf{x} + \mathbf{a} \times (\mathbf{a} \times \mathbf{x}) = \mathbf{a} \cdot \mathbf{x} - |\mathbf{a}|^2 \mathbf{x} + (\mathbf{a} \cdot \mathbf{x})\mathbf{a} = \mathbf{a} \times \mathbf{y}$$

Then we can consider

$$\begin{cases} \mathbf{x} + \mathbf{a} \times \mathbf{x} = \mathbf{y} \\ -|\mathbf{a}|^2 \mathbf{x} + \mathbf{a} \times \mathbf{x} = \mathbf{a} \times \mathbf{y} - (\mathbf{a} \cdot \mathbf{y})\mathbf{a} \end{cases}$$

This system is linear in  $(\mathbf{x}, \mathbf{a} \times \mathbf{x})$  and can be solved for  $\mathbf{x}$  which leads to (2.7). Finally, taking the scalar product of  $A(\mathbf{x}) = \mathbf{y}$  by  $\mathbf{x}$  gives  $|\mathbf{x}|^2 \leq \mathbf{x} \cdot \mathbf{y}$ . Thus  $\forall \mathbf{y}$ ,  $|A^{-1}(\mathbf{y})| \leq |\mathbf{y}|$ .  $\square$

With  $\mathbf{a} = -\frac{\alpha}{|\mathbf{M}|}\mathbf{M}$ , the (LLG) equation can be written  $A\left(\frac{\partial \mathbf{M}}{\partial t}\right) = |\gamma|\mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M}$ . Since  $\mathbf{M}$  is orthogonal to  $\frac{\partial \mathbf{M}}{\partial t}$ , applying the lemma 1 leads to

$$(2.8) \quad \frac{\partial \mathbf{M}}{\partial t} = \frac{|\gamma|}{1 + \alpha^2} \left[ \mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M} + \frac{\alpha}{|\mathbf{M}|}\mathbf{M} \times (\mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M}) \right]$$

This new form of the (LLG) equation shows that  $\frac{\partial \mathbf{M}}{\partial t}$  is a continuous function of  $\mathbf{H}$  and  $\mathbf{M}$  that we shall denote by  $\mathbf{L}_G^\alpha$  :

$$(2.9) \quad \mathbf{L}_G^\alpha(\mathbf{M}, \mathbf{H}) = \mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M} + \frac{\alpha}{|\mathbf{M}|}\mathbf{M} \times (\mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M})$$

**Remark 2** *Using this new form of the (LLG) equation and the Pythagorean theorem (given that  $\mathbf{H}_T \times \mathbf{M}$  and  $\mathbf{M} \times (\mathbf{H}_T \times \mathbf{M})$  are orthogonal vectors), it is easy to see that*

$$(2.10) \quad \left| \frac{\partial \mathbf{M}}{\partial t} \right|^2 = \frac{\gamma^2}{1 + \alpha^2} |\mathbf{H}_T \times \mathbf{M}|^2$$

We shall also use, in the part devoted to the numerical method, the

**Lemma 2**  $\mathbf{L}_G^\alpha$  is linear with respect to  $\mathbf{H}$  and  $\mathbf{L}_G^\alpha(\mathbf{M}, \mathbf{H}) = \mathbf{L}_G^\alpha\left(\mathbf{M}, \frac{\mathbf{B}}{\mu_0}\right)$ .

*Proof.* It is straightforward.  $\square$

### 2.2.2 A priori estimates

We assume in this paragraph that there exists a local solution  $(\mathbf{E}, \mathbf{H}, \mathbf{M})$  to problem (1.8) in  $C^1(0, T; L^2(\mathbb{R}^3))^3 \cap C^0(0, T; H^1(\mathbb{R}^3))^3$ . We assume that the system (1.8) is satisfied in the classical sense which means that the three equations of (1.8) are equalities in  $C^0(0, T; L^2(\mathbb{R}^3))$ .

- Conservation of the norm of  $\mathbf{M}$ .

It is first straightforward to show – at least formally in our case – from the (LLG) equation the conservation in time of the norm of the magnetization at any point :

$$\frac{\partial \mathbf{M}}{\partial t} \cdot \mathbf{M} = \frac{\partial}{\partial t} \{ |\mathbf{M}|^2 \} = 0$$

This computation is valid if we assume that for each  $\mathbf{x}$ , the function  $t \mapsto \mathbf{M}(\mathbf{x}, t)$  is of class  $C^1$ , which is not necessary the case under our assumptions. Rigorously we can multiply the (LLG) equation in  $L^2(\mathbb{R})^3$  by  $\varphi \mathbf{M}$ , with  $\varphi \in \mathcal{D}(\mathbb{R})$ , to get :

$$(2.11) \quad \frac{d}{dt} \left( \int_{\mathbb{R}} \varphi(x) |\mathbf{M}(x, t)|^2 dx \right) = 0$$

that is to say

$$(2.12) \quad \int_{\mathbb{R}} \left( |\mathbf{M}(x, t)|^2 - |\mathbf{M}_0(x)|^2 \right) \varphi(x) dx = 0$$

Thus

$$(2.13) \quad |\mathbf{M}(x, t)|^2 = |\mathbf{M}_0(x)|^2, \text{ a.e. } x \in \mathbb{R}$$

In particular  $\|\mathbf{M}(t)\|_{L^2} = \|\mathbf{M}_0(t)\|_{L^2}$  but more generally, if  $\mathbf{M}_0$  belongs to  $L^p(\mathbb{R})$ , with  $1 \leq p \leq +\infty$ , then  $\mathbf{M}(t)$  belongs to  $L^p(\mathbb{R})$  for any  $t$  and  $\|\mathbf{M}(t)\|_{L^p} = \|\mathbf{M}_0(t)\|_{L^p}$ .

• **Decay in time of the electromagnetic energy.**

The second estimate that we can get in any case consists in the decay in time of the total electromagnetic energy. Multiplying the Maxwell's equations by  $\mathbf{H}$  and  $\mathbf{E}$ , we get

$$(2.14) \quad \mathbf{E} \cdot (\nabla \times \mathbf{H}) - \mathbf{H} \cdot (\nabla \times \mathbf{E}) = \varepsilon_0 \mathbf{E} \cdot \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{H} \cdot \frac{\partial \mathbf{H}}{\partial t} + \mu_0 \mathbf{H} \cdot \frac{\partial \mathbf{M}}{\partial t}$$

We then integrate in space and, using the Green's formula, we get

$$(2.15) \quad \int_{\mathbb{R}^3} \nabla \cdot (\mathbf{E} \times \mathbf{H}) d\mathbf{x} = 0 = \int_{\mathbb{R}^3} \left( \varepsilon_0 \mathbf{E} \cdot \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{H} \cdot \frac{\partial \mathbf{H}}{\partial t} \right) d\mathbf{x} + \mu_0 \int_{\mathbb{R}^3} \mathbf{H} \cdot \frac{\partial \mathbf{M}}{\partial t} d\mathbf{x}$$

By adding  $\mathbf{H}_{eff}(\mathbf{M})$  to  $\mathbf{H}$ , we obtain

$$(2.16) \quad \frac{d}{dt} \left[ \frac{\varepsilon_0}{2} \|\mathbf{E}\|_{L^2}^2 + \frac{\mu_0}{2} \|\mathbf{H}\|_{L^2}^2 \right] - \mu_0 \int_{\mathbb{R}^3} \mathbf{H}_{eff}(\mathbf{M}) \cdot \frac{\partial \mathbf{M}}{\partial t} d\mathbf{x} = -\mu_0 \int_{\mathbb{R}^3} \mathbf{H}_T(\mathbf{H}, \mathbf{M}) \cdot \frac{\partial \mathbf{M}}{\partial t} d\mathbf{x}$$

Let us transform the left hand side of (2.16). First, concerning the field of anisotropy, we have

$$(2.17) \quad -\mu_0 \int_{\mathbb{R}^3} (-KP(\mathbf{M})) \cdot \frac{\partial \mathbf{M}}{\partial t} d\mathbf{x} = \frac{d}{dt} \left[ \frac{\mu_0 K}{2} \int_{\mathbb{R}^3} |P(\mathbf{M})|^2 \right]$$

(In the general case, we would have

$$\mu_0 \int_{\mathbb{R}^3} \mathbf{H}_a(\mathbf{M}) \cdot \frac{\partial \mathbf{M}}{\partial t} d\mathbf{x} = \mu_0 \left( \mathbf{H}_a, \frac{\partial \mathbf{M}}{\partial t} \right)_{L^2} = -\langle \mathcal{E}'_a(\mathbf{M}), \frac{\partial \mathbf{M}}{\partial t} \rangle = -\frac{d}{dt} (\mathcal{E}_a(\mathbf{M}))$$

Finally, we obtain, using (1.6) and (2.17) :

$$(2.18) \quad \frac{d}{dt} \left[ \frac{\mu_0 K}{2} \|P(\mathbf{M})\|_{L^2}^2 - \mu_0 (\mathbf{H}_s, \mathbf{M})_{L^2} \right] = -\mu_0 \int_{\mathbb{R}} (\mathbf{H}_{eff}(\mathbf{M}) \cdot \mathbf{M}) d\mathbf{x}$$

or, by adding some terms constant in time,

$$(2.19) \quad \frac{d}{dt} \left[ \frac{\mu_0 K}{2} \|P(\mathbf{M})\|_{L^2}^2 + \frac{\mu_0}{2} \|\mathbf{H}_s - \mathbf{M}\|_{L^2}^2 \right] = -\mu_0 \int_{\mathbb{R}} (\mathbf{H}_{eff}(\mathbf{M}) \cdot \mathbf{M}) d\mathbf{x}$$

Besides, we have, taking the vector product of the (LLG) equation by  $\frac{\partial \mathbf{M}}{\partial t}$  :

$$(2.20) \quad \frac{\partial \mathbf{M}}{\partial t} \times \frac{\partial \mathbf{M}}{\partial t} = 0 = |\gamma| \left( \mathbf{H}_T(\mathbf{H}, \mathbf{M}) \cdot \frac{\partial \mathbf{M}}{\partial t} \right) \mathbf{M} - \frac{\alpha}{|\mathbf{M}|} \left| \frac{\partial \mathbf{M}}{\partial t} \right|^2 \mathbf{M}$$

that yields to, everywhere  $\mathbf{M}_0$  is different from 0,

$$(2.21) \quad \mathbf{H}_T(\mathbf{H}, \mathbf{M}) \cdot \frac{\partial \mathbf{M}}{\partial t} = -\frac{\alpha}{|\gamma|} \frac{1}{|\mathbf{M}|} \left| \frac{\partial \mathbf{M}}{\partial t} \right|^2$$

On the other hand, everywhere  $\mathbf{M}_0 = 0$ , this equality reduces to  $0 = 0$ , since (2.10) implies that

$$(2.22) \quad \lim_{|\mathbf{M}| \rightarrow 0} \left( \frac{1}{|\mathbf{M}|} \left| \frac{\partial \mathbf{M}}{\partial t} \right|^2 \right) = 0$$

Finally, from (2.16), (2.19) and (2.21), we deduce that

$$(2.23) \quad \frac{d}{dt} \left[ \frac{\varepsilon_0}{2} \|\mathbf{E}\|_{L^2}^2 + \frac{\mu_0}{2} \|\mathbf{H}\|_{L^2}^2 + \frac{\mu_0 K}{2} \|P(\mathbf{M})\|_{L^2}^2 + \frac{\mu_0}{2} \|\mathbf{H}_s - \mathbf{M}\|_{L^2}^2 \right] = -\mu_0 \frac{\alpha}{|\gamma|} \int_{\mathbb{R}} \frac{1}{|\mathbf{M}|} \left| \frac{\partial \mathbf{M}}{\partial t} \right|^2 dx$$

It can be understood with this equality, where the sign of  $\alpha$  is crucial, why the ferromagnetic materials are expected to be absorbing since it shows that the electromagnetic energy

$$(2.24) \quad \mathcal{E}_T(\mathbf{E}, \mathbf{H}, \mathbf{M}) = \frac{\varepsilon_0}{2} \|\mathbf{E}\|_{L^2}^2 + \frac{\mu_0}{2} \|\mathbf{H}\|_{L^2}^2 + \frac{\mu_0 K}{2} \|P(\mathbf{M})\|_{L^2}^2 + \frac{\mu_0}{2} \|\mathbf{H}_s - \mathbf{M}\|_{L^2}^2$$

is decreasing in time. From (2.23), we get

$$(2.25) \quad \begin{aligned} \varepsilon_0 \|\mathbf{E}(t)\|_{L^2}^2 &\leq \left( \varepsilon_0 \|\mathbf{E}_0\|_{L^2}^2 + \mu_0 \|\mathbf{H}_0\|_{L^2}^2 + (1+K)\mu_0 \|\mathbf{M}_0\|_{L^2}^2 + \mu_0 \|\mathbf{H}_s\|_{L^2}^2 \right) \\ \mu_0 \|\mathbf{H}(t)\|_{L^2}^2 &\leq \left( \varepsilon_0 \|\mathbf{E}_0\|_{L^2}^2 + \mu_0 \|\mathbf{H}_0\|_{L^2}^2 + (1+K)\mu_0 \|\mathbf{M}_0\|_{L^2}^2 + \mu_0 \|\mathbf{H}_s\|_{L^2}^2 \right) \end{aligned}$$

These estimates show that  $\mathbf{E} \in L^2(\mathbb{R}^3)^3$  and  $\mathbf{H} \in L^2(\mathbb{R}^3)^3$  as soon as  $\mathbf{E}_0 \in L^2(\mathbb{R}^3)^3$  and  $\mathbf{H}_0 \in L^2(\mathbb{R}^3)^3$ .

### 2.3 The 1D model

In order to shorten the forthcoming formulae in our mathematical developments, henceforth all the positive physical constants **but**  $\alpha - \varepsilon_0, \mu_0, |\gamma|$  and  $1 + \alpha^2$  – will be replaced by one.

Moreover we consider from now on a mono-dimensional problem : we assume that all the unknowns in the following problem depend only on  $x$ , the first space variable, and  $t$ , the time variable. In other words we are interesting in plane waves propagating along the  $x$ -direction. We shall denote by  $(\mathbf{e}_x, \mathbf{e}_y, \mathbf{e}_z)$  the direct orthonormal basis of  $\mathbb{R}^3$  where  $\mathbf{e}_x$  is parallel to the  $x$ -direction.

The first consequence of this 1D hypothesis is to simplify the definition of the curl operator. Given any field  $\mathbf{A}(A_x, A_y, A_z)$  depending only on the  $x$  variable, we simply have

$$\nabla \times \mathbf{A} = \left( 0, -\frac{\partial A_z}{\partial x}, \frac{\partial A_y}{\partial x} \right)$$

The interest of this remark appears if you notice that taking the projection of (1.4) on the propagation axis leads to

$$(2.26) \quad \frac{dB_x}{dt} = 0 \quad \text{and} \quad \frac{dE_x}{dt} = 0$$

and so, we get that the longitudinal components  $E_x$  and  $B_x$  are both constant in our 1D problem. (Concerning the magnetic field, we simply have  $H_x = -M_x$  up to a constant.) However we would like to underline the fact that, contrary to what happens in the vacuum, it is not possible – **even** in the 1D case – to consider the propagation of a transverse electric mode on the one hand and the propagation of a transverse magnetic mode on the other hand because all the components of  $\mathbf{H}$  are coupled by the (LLG) equation.

We shall also use in the sequel the two following properties of  $\mathbf{L}_G^\alpha$  in the 1D case :

**Lemma 3** *Assuming that  $\mathbf{H}_s$  belongs to  $L^2(\mathbb{R})^3$ ,  $\mathbf{L}_G^\alpha$  maps  $L^2(\mathbb{R})^3 \cap L^\infty(\mathbb{R})^3 \times L^2(\mathbb{R})^3$  into  $L^2(\mathbb{R})^3$ .*

*Proof.* We can deduce from (1.3) that

$$(2.27) \quad \|\mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M}\|_{L^2} \leq \|\mathbf{M}\|_{L^\infty} \|\mathbf{H}\|_{L^2} + C \|\mathbf{M}\|_{L^\infty} \|\mathbf{M}\|_{L^2}$$

and then

$$(2.28) \quad \|\mathbf{L}_G^\alpha(\mathbf{M}, \mathbf{H})\|_{L^2} \leq (1 + \alpha) \|\mathbf{M}\|_{L^\infty} (\|\mathbf{H}\|_{L^2} + \|\mathbf{H}_s\|_{L^2} + C \|\mathbf{M}\|_{L^2})$$

□

**Lemma 4** *Assuming that  $\mathbf{H}_s$  belongs to  $H^1(\mathbb{R})^3$ ,  $\mathbf{L}_G^\alpha$  is locally Lipschitz from  $H^1(\mathbb{R})^3 \times H^1(\mathbb{R})^3$  into  $H^1(\mathbb{R})^3$  if and only if  $\alpha = 0$ .*

*Proof.* The case when  $\mathbf{L}_G^\alpha$  is not locally Lipschitz is not important for our paper and so we shall only prove the case when  $\alpha = 0$ . (To understand what happens if  $\alpha \neq 0$ , the reader can check that the “absolute value” map is not locally Lipschitz from  $H^1(I)$  to  $H^1(I)$ , for any open interval  $I$ , by considering the two functions  $u_n = 2 + \sin(nx)$  and  $v_n = -2 + \sin(nx)$ .)

Let us assume that  $\alpha = 0$ . We want to show that the map  $((\mathbf{M}, \mathbf{H}) \rightarrow \mathbf{H} \times \mathbf{M})$  is locally Lipschitz from  $H^1(\mathbb{R})^3 \times H^1(\mathbb{R})^3$  to  $H^1(\mathbb{R})^3$ . It obviously suffices to show that the map  $(u, v) \rightarrow uv$  is locally Lipschitz from  $H^1(\mathbb{R}) \times H^1(\mathbb{R})$  to  $H^1(\mathbb{R})$ .

Since  $H^1(\mathbb{R})$  is an algebra, we know that

$$u \in H^1(\mathbb{R}), v \in H^1(\mathbb{R}) \implies uv \in H^1(\mathbb{R})$$

and so  $H^1(\mathbb{R}) \times H^1(\mathbb{R})$  is mapped to  $H^1(\mathbb{R})$ . Then we can compute

$$(2.29) \quad \|u_1 v_1 - u_2 v_2\|_{L^2}^2 \leq \|u_1 (v_1 - v_2)\|_{L^2}^2 + \|(u_1 - u_2) v_2\|_{L^2}^2$$

hence, since  $\|u\|_{L^\infty(\mathbb{R})}^2 \leq \|u\|_{H^1(\mathbb{R})}^2$ ,

$$(2.30) \quad \|u_1 v_1 - u_2 v_2\|_{L^2}^2 \leq \|v_1 - v_2\|_{H^1}^2 \|u_1\|_{L^2}^2 + \|u_1 - u_2\|_{H^1}^2 \|v_2\|_{L^2}^2$$

We eventually get from

$$\frac{d}{dx} (u_1 v_1 - u_2 v_2) = \frac{du_1}{dx} v_1 - \frac{du_2}{dx} v_2 + \frac{dv_1}{dx} u_1 - \frac{dv_2}{dx} u_2$$

that

$$(2.31) \quad \left\| \frac{d}{dx} (u_1 v_1 - u_2 v_2) \right\|_{L^2}^2 = \left\| \frac{du_1}{dx} - \frac{du_2}{dx} \right\|_{L^2}^2 \|v_1\|_{H^1}^2 + \left\| \frac{du_2}{dx} \right\|_{L^2}^2 \|v_1 - v_2\|_{H^1}^2 + \left\| \frac{dv_1}{dx} - \frac{dv_2}{dx} \right\|_{L^2}^2 \|u_1\|_{H^1}^2$$

Inequalities (2.30) and (2.31) show that the map  $(u, v) \mapsto uv$  is locally Lipschitz from  $H^1(\mathbb{R}) \times H^1(\mathbb{R})$  to  $H^1(\mathbb{R})$ . □

**Remark 3** *The same argument also holds for the map  $(u, v, w) \mapsto uvw$ , which is locally Lipschitz from  $H^1(\mathbb{R}) \times H^1(\mathbb{R}) \times H^1(\mathbb{R})$  to  $H^1(\mathbb{R})$  as a composite of Lipschitz functions.*

## 2.4 Local results

The aim of this paragraph is to show the existence of a unique maximal local strong solution by applying the following theorem (see [10] or [15]) which applies to semi-linear abstract evolution equations :

**Theorem 1** *Let  $\mathcal{H}$  be a Banach space and  $A$  a maximal monotone operator (see [10]) with domain  $D(A)$ . Let  $F : \mathcal{H} \rightarrow \mathcal{H}$  be a locally Lipschitz function on  $D(A)$ . Then, for all  $u_0 \in D(A)$ , there exists a unique solution to*

$$\begin{cases} \frac{\partial u}{\partial t} = Au + F(u) \\ u(t=0) = u_0 \end{cases}$$

with  $u \in C^1(0, T_{max}; D(A)) \cap C^0(0, T_{max}; \mathcal{H})$  where

$$\text{either } (T_{max} = +\infty) \quad \text{or} \quad \left( T_{max} < +\infty \text{ and } \lim_{t \nearrow T_{max}} \|u\|_{\mathcal{H}} = +\infty \right)$$

In order to apply this theorem, it is natural to try to divide our problem into the linear part – the operator  $A$  – and the non-linear part – the function  $F$ , and then to check the required assumptions. It is exactly what is done in the next paragraph in the case  $\alpha = 0$ . But in the general case ( $\alpha \neq 0$ ), the function  $F$  (indeed the function  $\mathbf{L}_G^\alpha$ ) is no longer locally Lipschitz as mentioned in lemma 4 and we have to adapt our proof.

**Remark 4** *It would have been possible to treat the general case the same manner as the case  $\alpha = 0$  if we had considered (2.32), the same equation as in [17], where the term  $1/|\mathbf{M}|$  is omitted :*

$$(2.32) \quad \frac{\partial \mathbf{M}}{\partial t} = |\gamma| \mathbf{H} \times \mathbf{M} + \alpha \mathbf{M} \times (\mathbf{H} \times \mathbf{M})$$

*But this change in the LLG equation is valid only if the norm of the magnetization can be assumed to be constant in space, which is impossible when you consider the interface between a vacuum and a ferromagnetic material.*

We establish below local results applying theorem 1. We consider separately the two cases  $\alpha = 0$  and  $\alpha \neq 0$ , that is to say with and without damping. In each case, we shall decompose our proof into three steps. First we define the appropriate functional framework  $\mathcal{H}$  as well as the operator  $A$  and the non-linear part  $F$ . Then we establish that  $A$ , and  $F$ , have the required properties.

### 2.4.1 The case $\alpha = 0$

The (LLG) equation is

$$(2.33) \quad \frac{\partial \mathbf{M}}{\partial t} = \mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M}$$

- **Mathematical framework**

We introduce  $\mathcal{H} = L^2(\mathbb{R})^3 \times L^2(\mathbb{R})^3 \times H^1(\mathbb{R})^3$ , which is clearly a Hilbert space, and take as our unknown the vector  $\mathbf{u} = (\mathbf{E}, \mathbf{H}, \mathbf{M}) \in \mathcal{H}$ . We also introduce the Hilbert space  $H(\text{curl}, \mathbb{R}) = \{v \in L^2(\mathbb{R})^3; \text{curl}(v) \in L^2(\mathbb{R})^3\}$ , which is in the 1D case nothing other than  $L^2(\mathbb{R}) \times H^1(\mathbb{R}) \times H^1(\mathbb{R})$ . It is now possible to rewrite the three equations of (1.8) as :

$$(2.34) \quad \begin{cases} \frac{\partial \mathbf{E}}{\partial t} = \text{curl } \mathbf{H} \\ \frac{\partial \mathbf{H}}{\partial t} + \frac{\partial \mathbf{M}}{\partial t} = -\text{curl } \mathbf{E} \\ \frac{\partial \mathbf{M}}{\partial t} = \mathbf{H}_T \times \mathbf{M} \end{cases} \quad (\text{with } \mathbf{H}_T = \mathbf{H} + \mathbf{H}_s + \mathbf{H}_\alpha(\mathbf{M}))$$

that leads to the following system of evolution equations :

$$(2.35) \quad \begin{cases} \frac{\partial \mathbf{E}}{\partial t} = \text{curl } \mathbf{H} \\ \frac{\partial \mathbf{H}}{\partial t} = -\text{curl } \mathbf{E} - \mathbf{H} \times \mathbf{M} \\ \frac{\partial \mathbf{M}}{\partial t} = \mathbf{H} \times \mathbf{M} \end{cases}$$

It is now natural to introduce the unbounded operator

$$(2.36) \quad A(\mathbf{u}) = A(\mathbf{E}, \mathbf{H}, \mathbf{M}) = (\text{curl } \mathbf{H}, -\text{curl } \mathbf{E}, 0)$$

with domain  $D(A) = H(\text{curl}, \mathbb{R}) \times H(\text{curl}, \mathbb{R}) \times H^1(\mathbb{R})^3$ , and the function :

$$(2.37) \quad F(\mathbf{u}) = F(\mathbf{E}, \mathbf{H}, \mathbf{M}) = (0, -\mathbf{H} \times \mathbf{M}, \mathbf{H} \times \mathbf{M})$$

With these notations, the problem (1.8) is equivalent to :

$$(2.38) \quad \begin{cases} \frac{\partial \mathbf{u}}{\partial t} = A\mathbf{u} + F(\mathbf{u}) \\ \mathbf{u}(t=0) = \mathbf{u}_0 \end{cases}$$

where of course  $\mathbf{u}_0 = (\mathbf{E}_0, \mathbf{H}_0, \mathbf{M}_0)$ . It remains to show that  $A$  is maximal monotone and that  $F$  is locally Lipschitz from  $D(A)$  to  $D(A)$ .

• **A is a maximal monotone operator**

This result is not surprising since  $A$  is nothing other than the operator associated to the Maxwell's equations. First using the Green's formula :

$$(2.39) \quad \int_{\mathbb{R}} (\text{curl } \mathbf{E} \cdot \mathbf{H} - \text{curl } \mathbf{H} \cdot \mathbf{E}) dx = 0$$

for all  $(\mathbf{E}, \mathbf{H})$  in  $H(\text{curl}, \mathbb{R}) \times H(\text{curl}, \mathbb{R})$  yields to :

$$(2.40) \quad (A\mathbf{u}, \mathbf{u})_{\mathcal{H}} = 0, \quad \forall \mathbf{u} \in D(A)$$

Then, to get the surjectivity of  $A + I$  from  $D(A)$  in  $\mathcal{H}$ , we have to show that, for all  $(\mathbf{e}, \mathbf{h}, \mathbf{m})$  in  $\mathcal{H}$ , there exists  $(\mathbf{E}, \mathbf{H}, \mathbf{M})$  in  $D(A)$  such that :

$$(2.41) \quad \begin{cases} -\text{curl } \mathbf{H} + \mathbf{E} &= \mathbf{e} \\ \text{curl } \mathbf{E} + \mathbf{H} &= \mathbf{h} \\ \mathbf{M} &= \mathbf{m} \end{cases}$$

Reminding that the longitudinal components are constant, this leads to solve for the transverse components

$$(2.42) \quad \text{curl} (\text{curl } \mathbf{E}) + \mathbf{E} = \mathbf{e} + \text{curl } \mathbf{h}$$

The Lax-Milgram theorem ensures that there exists a unique solution  $H^1(\mathbb{R})^2$  – here are considered only the transverse components – for a right hand side in  $L^2(\mathbb{R})^2$ . This implies in our case that  $\mathbf{E}$  is in  $H(\text{curl}, \mathbb{R})$  for  $(\mathbf{e}, \mathbf{h})$  in  $L^2(\mathbb{R})^3$ . We can then conclude with  $\mathbf{H} = \mathbf{h} - \text{curl } \mathbf{E}$ .

•  **$F$  is a locally Lipschitz continuous function from  $D(A)$  to  $D(A)$**

It is clear by lemma 4.

This was the last point of our proof. Eventually we can state the

**Theorem 2** *Given  $(\mathbf{E}_0, \mathbf{H}_0, \mathbf{M}_0) \in H(\text{curl}, \mathbb{R}) \times H(\text{curl}, \mathbb{R}) \times H^1(\mathbb{R})^3$ , there exists a unique local maximal strong solution of (1.8) so that :*

$$(2.43) \quad \begin{cases} \mathbf{E} \in C^1(0, T_{max}; L^2(\mathbb{R})^3) \cap C^0(0, T_{max}; H(\text{curl}, \mathbb{R})) \\ \mathbf{H} \in C^1(0, T_{max}; L^2(\mathbb{R})^3) \cap C^0(0, T_{max}; H(\text{curl}, \mathbb{R})) \\ \mathbf{M} \in C^1(0, T_{max}; H^1(\mathbb{R})^3) \end{cases}$$

with moreover the following alternative : either  $T_{max} = +\infty$ ,

$$(2.44) \quad \text{or } \lim_{t \nearrow T_{max}} \left\{ \|\mathbf{E}(t)\|_{H(\text{curl})}^2 + \|\mathbf{H}(t)\|_{H(\text{curl})}^2 + \|\mathbf{M}(t)\|_{H^1}^2 \right\} = +\infty$$

**The case  $\alpha \neq 0$**

The (LLG) equation is

$$(2.45) \quad \frac{\partial \mathbf{M}}{\partial t} = \mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M} + \frac{\alpha}{|\mathbf{M}|} \mathbf{M} \times (\mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{M}) = \mathbf{L}_G^\alpha(\mathbf{M}, \mathbf{H}_T(\mathbf{H}, \mathbf{M}))$$

As said in the introduction, in the case  $\alpha \neq 0$ , we shall only present an existence result but, since

$$(2.46) \quad H^1(\mathbb{R}) \subset L^2(\mathbb{R}) \cap L^\infty(\mathbb{R})$$

uniqueness will be ensured by the theorem about weak solutions presented in [12]. (For this reason, this paragraph can be seen as complementary to [12].)

First of all it must be said that this existence result is obtained under a new and purely technical assumption on the initial distribution of magnetization  $\mathbf{M}_0$ , namely

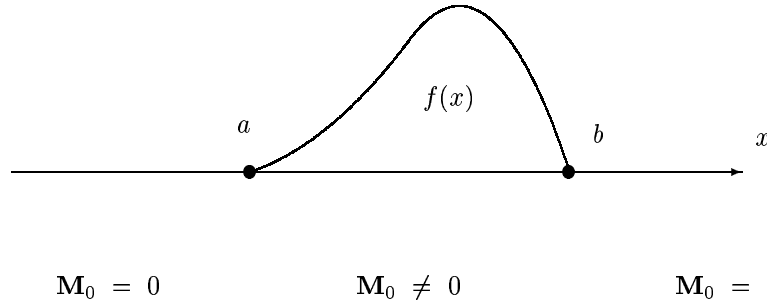
$$(2.47) \quad \frac{\mathbf{M}_0}{|\mathbf{M}_0|} \in H^1(\mathbb{R})$$

or, to be more rigorous,

$$(2.48) \quad \exists \mathbf{m}_0 \in H^1(\mathbb{R}) \text{ such that } : \forall x \in \mathbb{R}, |\mathbf{M}_0(x)| \neq 0 \Rightarrow \frac{\mathbf{M}_0(x)}{|\mathbf{M}_0(x)|} = \mathbf{m}_0(x)$$

However it is important to understand that (2.48) represents no serious restriction, other than  $\mathbf{M}_0$  must be of bounded support.

**Remark 5** *Without studying the most general case, let us see briefly that there exists a large class of admissible initial distributions. For instance it suffices that  $\mathbf{M}_0 = f(x)\mathbf{u}$  with  $\mathbf{u}$  a fixed vector and  $f \in H^1(\mathbb{R})$  a positive function with bounded support : it is easy in this case to construct  $\mathbf{m}_0 \in H^1(\mathbb{R})$  equal to  $\mathbf{u}$  on the support of  $f$ .*



From this first example, the reader can easily check that more generally, one may consider any ferromagnetic zone  $[a, b]$  such that

$$(2.49) \quad \begin{cases} \exists \varepsilon_a > 0, \exists \mathbf{u}_a \in \mathbb{R}^3, \forall x \in ]a, a + \varepsilon_a[, \mathbf{M}_0 = f_a(x)\mathbf{u}_a & (f_a(x) \geq 0) \\ \exists \varepsilon_b > 0, \exists \mathbf{u}_b \in \mathbb{R}^3, \forall x \in ]b - \varepsilon_b, b[, \mathbf{M}_0 = f_b(x)\mathbf{u}_b & (f_b(x) \geq 0) \end{cases}$$

In this case, we have

$$\left( \frac{\mathbf{M}_0}{|\mathbf{M}_0|} \right)_{]a, a + \varepsilon_a[} \in H^1(]a, a + \varepsilon_a[)^3 \text{ and } \left( \frac{\mathbf{M}_0}{|\mathbf{M}_0|} \right)_{]b - \varepsilon_b, b[} \in H^1(]b - \varepsilon_b, b[)^3$$

and is easy to check, because  $\exists C_\varepsilon > 0$  such that  $|\mathbf{M}_0| \geq C_\varepsilon$  on  $]a + \varepsilon_a, b - \varepsilon_b[$ , that

$$\frac{\mathbf{M}_0}{|\mathbf{M}_0|}_{]a + \varepsilon_a, b - \varepsilon_b[} \in H^1(]a + \varepsilon_a, b - \varepsilon_b[)^3$$

Hence the result since  $\mathbf{M}_0$  is continuous for  $x = a + \varepsilon_a$  and  $x = b - \varepsilon_b$ . And of course, the ferromagnetic zone can be defined as the union of such disjoint intervals :

$$(2.50) \quad \Omega = \bigcup_i [a_i, b_i]$$

which leads to a large class of admissible initial distributions.

On the contrary  $\mathbf{M}_0 = e^{-x^2} \sin(x)\mathbf{u}$  is not a admissible initial data.

As a conclusion let us simply say that assumption (2.48) only implies a kind of control on the way  $\mathbf{M}_0$  vanishes to zero.

- **A new formulation**



We know that it is possible to write the three equations of (1.8) as the following system of evolution equations :

$$(2.51) \quad \begin{cases} \frac{\partial \mathbf{E}}{\partial t} = \operatorname{curl} \mathbf{H} \\ \frac{\partial \mathbf{H}}{\partial t} = -\operatorname{curl} \mathbf{E} - \mathbf{L}_G^\alpha(\mathbf{M}, \mathbf{H}) \\ \frac{\partial \mathbf{M}}{\partial t} = \mathbf{L}_G^\alpha(\mathbf{M}, \mathbf{H}) \end{cases}$$

associated to initial data  $(\mathbf{E}_0, \mathbf{H}_0, \mathbf{M}_0)$ . The idea of this section is to consider a normalized magnetization vector  $\mathbf{m}$  ; more exactly, we define the following problem

$$(2.52) \quad \begin{cases} \frac{\partial \mathbf{E}}{\partial t} = \operatorname{curl} \mathbf{H} \\ \frac{\partial \mathbf{H}}{\partial t} = -\operatorname{curl} \mathbf{E} - |\mathbf{M}_0| \mathbf{L}^\alpha(\mathbf{m}, \mathbf{H}) \\ \frac{\partial \mathbf{m}}{\partial t} = \mathbf{L}^\alpha(\mathbf{m}, \mathbf{H}) \end{cases}$$

where

$$(2.53) \quad \mathbf{L}^\alpha(\mathbf{m}, \mathbf{H}) = \frac{|\gamma|}{1 + \alpha^2} [\mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{m} + \alpha \mathbf{m} \times (\mathbf{H}_T(\mathbf{H}, \mathbf{M}) \times \mathbf{m})]$$

associated to initial data  $(\mathbf{E}_0, \mathbf{H}_0, \mathbf{m}_0)$  with  $\mathbf{m}_0 = \frac{\mathbf{M}_0}{|\mathbf{M}_0|} \in H^1(\mathbb{R})$  as indicated in assumption (2.48). Now let  $(\mathbf{E}, \mathbf{H}, \mathbf{m})$  be a solution to problem (2.52) and let  $\mathbf{M}(x, t) = |\mathbf{M}_0(x)| \mathbf{m}(x, t)$ . We have  $\mathbf{M} \in H^1(\Omega)$ , and also, because of the conservation of the norm of  $\mathbf{M}$  (and  $\mathbf{m}$ ) :

$$(2.54) \quad \mathbf{L}_G^\alpha(\mathbf{M}, \mathbf{H}) = |\mathbf{M}_0| \mathbf{L}^\alpha(\mathbf{m}, \mathbf{H})$$

and

$$(2.55) \quad \frac{\partial \mathbf{M}}{\partial t} = \mathbf{L}_G^\alpha(\mathbf{M}, \mathbf{H}) \Leftrightarrow \frac{\partial \mathbf{m}}{\partial t} = \mathbf{L}^\alpha(\mathbf{m}, \mathbf{H})$$

In other words every solution to problem (2.52) is clearly a solution to problem (2.51), and it only remains now to conclude to apply theorem 1 to problem (2.52).

#### • Mathematical framework

The notations are the same as in the case  $\alpha = 0$ . We take as unknown the vector  $\mathbf{u} = (\mathbf{E}, \mathbf{H}, \mathbf{m}) \in \mathcal{H}$ . This time again we introduce the same operator

$$(2.56) \quad A(\mathbf{u}) = A(\mathbf{E}, \mathbf{H}, \mathbf{M}) = (\operatorname{curl} \mathbf{H}, -\operatorname{curl} \mathbf{E}, 0)$$

with domain  $D(A) = H_0(\operatorname{curl}, \Omega) \times H(\operatorname{curl}, \Omega) \times H^1(\Omega)^3$ , and the function :

$$(2.57) \quad F_{\mathbf{M}_0}(\mathbf{u}) = F(\mathbf{E}, \mathbf{H}, \mathbf{M}) = (0, -|\mathbf{M}_0| \mathbf{L}^\alpha(\mathbf{m}, \mathbf{H}), \mathbf{L}^\alpha(\mathbf{m}, \mathbf{H}))$$

With these notations, the problem (2.52) is equivalent to :

$$(2.58) \quad \begin{cases} \frac{\partial \mathbf{u}}{\partial t} = A\mathbf{u} + F_{\mathbf{M}_0}(\mathbf{u}) \\ \mathbf{u}(t=0) = \mathbf{u}_0 \end{cases}$$

with  $\mathbf{u}_0 = (\mathbf{E}_0, \mathbf{H}_0, \mathbf{m}_0)$ . It remains to see that

#### • $A$ is a maximal monotone operator

It is clear since this operator is exactly the same as in the previous section.

#### • $F_{\mathbf{M}_0}$ is a locally Lipschitz continuous function from $D(A)$ to $D(A)$

First it is obvious that  $F_{\mathbf{M}_0}$  maps  $D(A)$  to  $D(A)$  since  $H^1(\Omega)^3$  is an algebra. Moreover  $F_{\mathbf{M}_0}$  is locally Lipschitz by remark 3.

Applying theorem 1, we obtain a local solution  $(\mathbf{E}, \mathbf{H}, \mathbf{m})$  to the problem (2.52) ; it is then possible to define  $\mathbf{M} = |\mathbf{M}_0| \mathbf{m}$  a solution to problem (2.51) and finally we can state the

**Theorem 3** Given  $(\mathbf{E}_0, \mathbf{H}_0, \mathbf{M}_0) \in H_0(\text{curl}, \Omega) \times H(\text{curl}, \Omega) \times H^1(\Omega)^3$ , there exists a unique local maximal strong solution of (1.8) so that :

$$(2.59) \quad \begin{cases} \mathbf{E} \in C^1(0, T_{max}; L^2(\Omega)^3) \cap C^0(0, T_{max}; H(\text{curl}, \Omega)) \\ \mathbf{H} \in C^1(0, T_{max}; L^2(\Omega)^3) \cap C^0(0, T_{max}; H(\text{curl}, \Omega)) \\ \mathbf{M} \in C^1(0, T_{max}; H^1(\Omega)^3) \end{cases}$$

with moreover the following alternative : either  $T_{max} = +\infty$ ,

$$(2.60) \quad \text{or } \lim_{t \nearrow T_{max}} \left\{ \|\mathbf{E}(t)\|_{H(\text{curl})}^2 + \|\mathbf{H}(t)\|_{H(\text{curl})}^2 + \|\mathbf{M}(t)\|_{H^1}^2 \right\} = +\infty$$

It remains now to show that the local solutions of theorem 2 and theorem 3 are indeed global. Rather than expliciting two very similar demonstrations, we now end this mathematical part of our work in the most general case : no assumption is made neither on  $\alpha$  nor on  $\mathbf{M}_0$ . This way we can eventually conclude in each case.

## 2.5 A new estimate for $\mathbf{H}$

We have now to verify that no local solution can blow up in a finite time. With (2.13) and (2.25) which are valid in our case for any  $T < T_{max}$ , only the  $L^2$  norm of the local solutions provided by theorem 2 is controlled. So, to prove global existence we have to find  $H^1$  estimates for these solutions depending on  $t$  in a locally bounded manner. More precisely we first establish in this paragraph that  $\mathbf{H} \in H(\text{curl}, \mathbb{R})$ .

Before beginning, let us introduce the following technical lemma.

**Lemma 5** Concerning the infinity norm of the magnetic field, we have

$$(2.61) \quad \|\mathbf{H}\|_{L^\infty}^2 \leq \|\mathbf{M}\|_{L^\infty}^2 + 2 \|\mathbf{H}\|_{L^2} \|\text{curl } \mathbf{H}\|_{L^2}$$

*Proof.* First let us recall that for all  $u \in H^1(\mathbb{R})$ , we have

$$(2.62) \quad \|u\|_{L^\infty}^2 \leq 2 \|u\|_{L^2} \left\| \frac{\partial u}{\partial x} \right\|_{L^2}$$

as it can be shown by density of  $\mathcal{D}(\mathbb{R})$  in  $H^1(\mathbb{R})$ . Then it suffices to apply (2.62) to the transverse components of  $\mathbf{H}$  and to recall, for the longitudinal component, that  $H_x = -M_x$ .  $\square$

Now we can give the three steps of our proof.

- **step 1** : Estimates on  $\mathbf{M}$

As far as we know at this point,  $\mathbf{M} \in C^1(0, T_{max}; H^1(\mathbb{R})^3)$  but it is not difficult to prove that  $\mathbf{M} \in C^2(0, T_{max}; H^1(\mathbb{R})^3)$ , and moreover that  $\|\mathbf{M}(t)\|_{L^\infty} = \|\mathbf{M}_0\|_{L^\infty}$ . Indeed

$$(2.63) \quad \frac{\partial \mathbf{M}}{\partial t} = \mathbf{H}_T \times \mathbf{M} + \frac{\alpha}{|\mathbf{M}|} \mathbf{M} \times (\mathbf{H}_T \times \mathbf{M})$$

and as we know that  $\mathbf{H} \in C^1(0, T_{max}; L^2(\mathbb{R})^3)$ , it is clear that

$$\mathbf{H}_T = [\mathbf{H} + \mathbf{H}_s + \mathbf{H}_a(\mathbf{M})] \in C^1(0, T_{max}; L^2(\mathbb{R})^3)$$

since  $\mathbf{H}_a(\mathbf{M})$  linearly depends on  $\mathbf{M}$ . Moreover we know that

$$\mathbf{M} \in C^1(0, T_{max}; H^1(\mathbb{R})^3) \hookrightarrow C^1(0, T_{max}; L^\infty(\mathbb{R})^3)$$

We can conclude that  $\mathbf{H}_T \times \mathbf{M}$  and  $\mathbf{M} \times (\mathbf{H}_T \times \mathbf{M})$  are in  $C^1(0, T_{max}; L^2(\mathbb{R})^3)$ . Thus  $\mathbf{M} \in C^2(0, T_{max}; H^1(\mathbb{R})^3)$ .

- **step 2** : A new energy-type identity.

The idea is to do the same kind of computations we made to get the decay in time of the electromagnetic energy but for the time derivatives of the fields  $\mathbf{E}$  and  $\mathbf{H}$ . Let us give a formal proof assuming that  $\mathbf{E}$  and  $\mathbf{H}$  are smooth enough. We start from

$$(2.64) \quad \begin{cases} \frac{\partial^2 \mathbf{H}}{\partial t^2} + \frac{\partial^2 \mathbf{M}}{\partial t^2} = -\nabla \times \frac{\partial \mathbf{E}}{\partial t} \\ \frac{\partial^2 \mathbf{E}}{\partial t^2} = \nabla \times \frac{\partial \mathbf{H}}{\partial t} \end{cases}$$

Then we multiply by  $\frac{\partial \mathbf{H}}{\partial t}$  and  $\frac{\partial \mathbf{E}}{\partial t}$  and integrate in space to obtain

$$(2.65) \quad \frac{d}{dt} \left[ \int_{\mathbb{R}} \left( \left| \frac{\partial \mathbf{E}}{\partial t} \right|^2 + \left| \frac{\partial \mathbf{H}}{\partial t} \right|^2 \right) dx \right] = - \int_{\mathbb{R}} \frac{\partial^2 \mathbf{M}}{\partial t^2} \cdot \frac{\partial \mathbf{H}}{\partial t} dx$$

or, since  $\frac{\partial \mathbf{E}}{\partial t} = \text{curl } \mathbf{H}$ ,

$$(2.66) \quad \frac{d}{dt} \left[ \int_{\mathbb{R}} \left( |\text{curl } \mathbf{H}|^2 + \left| \frac{\partial \mathbf{H}}{\partial t} \right|^2 \right) dx \right] = - \int_{\mathbb{R}} \frac{\partial^2 \mathbf{M}}{\partial t^2} \cdot \frac{\partial \mathbf{H}}{\partial t} dx$$

That yields to,  $\forall t \leq T_{max}$ ,

$$(2.67) \quad \frac{1}{2} \int_{\mathbb{R}} \left( |\text{curl } \mathbf{H}(t)|^2 + \left| \frac{\partial \mathbf{H}}{\partial t}(t) \right|^2 \right) dx = \frac{1}{2} \int_{\mathbb{R}} \left( |\text{curl } \mathbf{H}(0)|^2 + \left| \frac{\partial \mathbf{H}}{\partial t}(0) \right|^2 \right) dx - \int_0^t \int_{\mathbb{R}} \frac{\partial^2 \mathbf{M}}{\partial t^2} \cdot \frac{\partial \mathbf{H}}{\partial t} dx ds$$

As we know that  $\forall t \in [0, T_{max}]$

$$(2.68) \quad \frac{\partial \mathbf{H}}{\partial t} = \text{curl } \mathbf{E} - \frac{\partial \mathbf{M}}{\partial t} = \text{curl } \mathbf{E} - \mathbf{L}_G^\alpha(\mathbf{M}, \mathbf{H})$$

we introduce  $\Lambda_0 = \text{curl } \mathbf{E}_0 - \mathbf{L}_G^\alpha(\mathbf{M}_0, \mathbf{H}_0) \in L^2$ . ( $\Lambda_0$  is well defined since  $\mathbf{E}_0$  and  $\mathbf{H}_0$  are taken in  $H(\text{curl}, \mathbb{R})$  and  $\mathbf{M}_0$  is taken in  $H^1(\mathbb{R})^3$ .) Then (2.67) is equivalent to

$$(2.69) \quad \frac{1}{2} \left( \|\text{curl } \mathbf{H}(t)\|_{L^2}^2 + \left\| \frac{\partial \mathbf{H}}{\partial t}(t) \right\|_{L^2}^2 \right) = \frac{1}{2} \left( \|\text{curl } \mathbf{H}_0\|_{L^2}^2 + \|\Lambda_0\|_{L^2}^2 \right) - \int_0^t \int_{\mathbb{R}} \frac{\partial^2 \mathbf{M}}{\partial t^2} \cdot \frac{\partial \mathbf{H}}{\partial t} dx ds$$

Now, to give a rigorous proof of (2.69), since  $\mathbf{E}$  and  $\mathbf{H}$  are not regular enough, we have to use a discrete differentiation. More precisely, for any Banach space  $X$  and any  $u \in C^0([0, T^*]; X)$ , we introduce

$$\forall T < T^*, \forall h \in [0; T^* - T], \quad D_h u(t) = \frac{u(t+h) - u(t)}{h} \in C^1(0, T; X)$$

Moreover, if  $u \in C^1([0, T^*]; X)$ , then, when  $h \rightarrow 0$ ,

$$D_h u(t) \longrightarrow \frac{\partial u}{\partial t}(t) \quad \text{in } C^0(0, T; X)$$

Applying  $D_h$  to Maxwell's equations – with  $T^* = T_{max}$  –, we get by linearity

$$(2.70) \quad \begin{cases} \frac{\partial}{\partial t} (D_h \mathbf{E}) - \text{curl} (D_h \mathbf{H}) = 0 \\ \frac{\partial}{\partial t} (D_h \mathbf{H}) + \text{curl} (D_h \mathbf{E}) = -\frac{\partial}{\partial t} (D_h \mathbf{M}) \end{cases}$$

We multiply the first equation by  $D_h \mathbf{E}$  and the second one by  $D_h \mathbf{H}$  and then integrate the sum to get

$$(2.71) \quad \frac{1}{2} \frac{d}{dt} \left( \|D_h \mathbf{E}\|_{L^2}^2 + \|D_h \mathbf{H}\|_{L^2}^2 \right) = - \int_{\mathbb{R}} D_h \frac{\partial \mathbf{M}}{\partial t} \cdot D_h \mathbf{H} dx$$

We now integrate in time, for any  $t \in [0, T_{max}]$ ,

$$(2.72) \quad \frac{1}{2} \left( \|D_h \mathbf{E}(t)\|_{L^2}^2 + \|D_h \mathbf{H}(t)\|_{L^2}^2 \right) = \frac{1}{2} \left( \|D_h \mathbf{E}(0)\|_{L^2}^2 + \|D_h \mathbf{H}(0)\|_{L^2}^2 \right) - \int_0^t \int_{\mathbb{R}} D_h \frac{\partial \mathbf{M}}{\partial t} \cdot D_h \mathbf{H} dx dt$$

As  $\mathbf{E} \in C^1(0, T_{max}; L^2(\mathbb{R}))^3$  and  $\mathbf{H} \in C^1(0, T_{max}; L^2(\mathbb{R}))^3$ , we have

$$(2.73) \quad \begin{cases} \lim_{h \rightarrow 0} \|D_h \mathbf{E}(t)\|_{L^2}^2 = \left\| \frac{\partial \mathbf{E}}{\partial t}(t) \right\|_{L^2}^2 = \|\operatorname{curl} \mathbf{H}(t)\|_{L^2}^2 \\ \lim_{h \rightarrow 0} \|D_h \mathbf{H}(t)\|_{L^2}^2 = \left\| \frac{\partial \mathbf{H}}{\partial t}(t) \right\|_{L^2}^2 \end{cases}$$

Moreover, as  $\mathbf{H} \in C^1(0, T; L^2(\mathbb{R}))^3$  and  $\mathbf{M} \in C^2(0, T; L^2(\mathbb{R}))^3$ , we have,  $\forall t \in [0, T_{max}]$ ,

$$(2.74) \quad \lim_{h \rightarrow 0} \int_0^t \int_{\mathbb{R}} D_h \frac{\partial \mathbf{M}}{\partial t} \cdot D_h \mathbf{H} dx dt = \int_0^t \int_{\mathbb{R}} \frac{\partial^2 \mathbf{M}}{\partial t^2} \cdot \frac{\partial \mathbf{H}}{\partial t} dx dt$$

It is thus possible to take the limit when  $h \rightarrow 0$  in (2.72) and then to get (2.69).

- **step 3** : The estimate on  $\left( \frac{\partial \mathbf{H}}{\partial t} \right)$  and  $(\operatorname{curl} \mathbf{H})$ .

To shorten the proof given below we shall consider here the case ( $\alpha = 0, \mathbf{H}_{eff}(\mathbf{M}) = 0$ ) which presents all the difficulties of the case ( $\alpha \neq 0, \mathbf{H}_{eff}(\mathbf{M}) \neq 0$ ) that we shall detail in appendix A. In other words we assume for a while that

$$(2.75) \quad \frac{\partial \mathbf{M}}{\partial t} = \mathbf{H} \times \mathbf{M}$$

while the case of the complete (LLG) is considered in the appendix.

In the case ( $\alpha = 0, \mathbf{H}_{eff}(\mathbf{M}) = 0$ ), it is easy to compute at each point :

$$\frac{\partial^2 \mathbf{M}}{\partial t^2} \cdot \frac{\partial \mathbf{H}}{\partial t} = \left[ \frac{\partial \mathbf{H}}{\partial t} \times \mathbf{M} + \mathbf{H} \times \frac{\partial \mathbf{M}}{\partial t} \right] \cdot \frac{\partial \mathbf{H}}{\partial t} = [(\mathbf{H} \times \mathbf{M}) \times \mathbf{H}] \cdot \frac{\partial \mathbf{H}}{\partial t}$$

that yields to

$$(2.76) \quad \left| \left( \frac{\partial^2 \mathbf{M}}{\partial t^2} \cdot \frac{\partial \mathbf{H}}{\partial t} \right) \right| \leq |\mathbf{M}| |\mathbf{H}|^2 \left| \frac{\partial \mathbf{H}}{\partial t} \right|$$

Thus, plugging (2.76) into (2.69), we obtain

$$(2.77) \quad \frac{1}{2} \left( \left\| \frac{\partial \mathbf{H}}{\partial t} \right\|_{L^2}^2 + \|\operatorname{curl} \mathbf{H}\|_{L^2}^2 \right) \leq \frac{1}{2} \left( \|\operatorname{curl} \mathbf{H}_0\|_{L^2}^2 + \|\Lambda_0\|_{L^2}^2 \right) + \|\mathbf{M}\|_{L^\infty} \int_0^t \int_{\mathbb{R}} |\mathbf{H}|^2 \left| \frac{\partial \mathbf{H}}{\partial t} \right| dx ds$$

We then note that

$$(2.78) \quad \int_{\mathbb{R}} |\mathbf{H}|^2 \left| \frac{\partial \mathbf{H}}{\partial t} \right| dx \leq \|\mathbf{H}\|_{L^\infty} \left\| \frac{\partial \mathbf{H}}{\partial t} \right\|_{L^2} \|\mathbf{H}\|_{L^2}$$

that is to say, applying lemma 5,

$$(2.79) \quad \int_{\mathbb{R}} |\mathbf{H}|^2 \left| \frac{\partial \mathbf{H}}{\partial t} \right| dx \leq 2 \|\mathbf{H}\|_{L^2}^{\frac{3}{2}} \|\operatorname{curl} \mathbf{H}\|_{L^2}^{\frac{1}{2}} \left\| \frac{\partial \mathbf{H}}{\partial t} \right\|_{L^2} + \|\mathbf{M}\|_{L^\infty}^2 \|\mathbf{H}\|_{L^2} \left\| \frac{\partial \mathbf{H}}{\partial t} \right\|_{L^2}$$

In the case ( $\alpha = 0, \mathbf{H}_{eff}(\mathbf{M}) = 0$ ), (2.23) gives :

$$(2.80) \quad \|\mathbf{E}\|_{L^2}^2 + \|\mathbf{H}\|_{L^2}^2 = \|\mathbf{E}_0\|_{L^2}^2 + \|\mathbf{H}_0\|_{L^2}^2$$

Introducing

$$(2.81) \quad \mathcal{E}_0 = \frac{1}{2} \left( \|\mathbf{E}_0\|_{L^2}^2 + \|\mathbf{H}_0\|_{L^2}^2 \right)$$

we get

$$(2.82) \quad \int_{\mathbb{R}} |\mathbf{H}|^2 \left| \frac{\partial \mathbf{H}}{\partial t} \right| dx \leq 2 (2\mathcal{E}_0)^{\frac{3}{4}} \|\operatorname{curl} \mathbf{H}\|_{L^2}^{\frac{1}{2}} \left\| \frac{\partial \mathbf{H}}{\partial t} \right\|_{L^2} + \mathcal{E}_0^{\frac{1}{2}} \|\mathbf{M}_0\|_{L^\infty}^2 \left\| \frac{\partial \mathbf{H}}{\partial t} \right\|_{L^2}$$

Therefore, with  $C = \text{Max} \left( 2 \|\mathbf{M}_0\|_{L^\infty} (2\mathcal{E}_0)^{\frac{3}{4}}, \|\mathbf{M}_0\|_{L^\infty}^3 \mathcal{E}_0^{\frac{1}{2}} \right)$ , we get from (2.77) and (2.82)

$$(2.83) \quad \frac{1}{2} \left( \left\| \frac{\partial \mathbf{H}}{\partial t} \right\|_{L^2}^2 + \|\text{curl } \mathbf{H}\|_{L^2}^2 \right) \leq \frac{1}{2} \left( \|\text{curl } \mathbf{H}_0\|_{L^2}^2 + \|\Lambda_0\|_{L^2}^2 \right) + C \int_0^t \|\text{curl } \mathbf{H}\|_{L^2}^{\frac{1}{2}} \left\| \frac{\partial \mathbf{H}}{\partial t} \right\|_{L^2} ds$$

and so

$$(2.84) \quad \frac{1}{2} \left( \left\| \frac{\partial \mathbf{H}}{\partial t} \right\|_{L^2}^2 + \|\text{curl } \mathbf{H}\|_{L^2}^2 \right) \leq \frac{1}{2} \left( \|\text{curl } \mathbf{H}_0\|_{L^2}^2 + \|\Lambda_0\|_{L^2}^2 \right) + \frac{C}{2} \int_0^t \left( \|\text{curl } \mathbf{H}\|_{L^2} + \left\| \frac{\partial \mathbf{H}}{\partial t} \right\|_{L^2} \right)^2 ds$$

Introducing the function  $F(t) = \frac{1}{2} \left( \|\text{curl } \mathbf{H}\|_{L^2}^2 + \left\| \frac{\partial \mathbf{H}}{\partial t} \right\|_{L^2}^2 \right)$  and noticing that

$$(2.85) \quad \|\text{curl } \mathbf{H}\|_{L^2} \leq 1 + \|\text{curl } \mathbf{H}\|_{L^2}^2$$

(2.84) is equivalent to

$$(2.86) \quad F(t) \leq F(0) + C \int_0^t \left( \frac{1}{2} + F(s) \right) ds$$

We can then apply Gronwall's lemma to the function  $\left( \frac{1}{2} + F(t) \right)$  to get

$$(2.87) \quad \left( \frac{1}{2} + F(t) \right) \leq \left( \frac{1}{2} + F(0) \right) \exp(Ct)$$

that yields to

$$(2.88) \quad \left( \|\text{curl } \mathbf{H}\|_{L^2}^2 + \left\| \frac{\partial \mathbf{H}}{\partial t} \right\|_{L^2}^2 \right) \leq \left( \|\text{curl } \mathbf{H}_0\|_{L^2}^2 + \|\Lambda_0\|_{L^2}^2 \right) \exp(Ct) + \frac{1}{2} (\exp(Ct) - 1)$$

which is exactly the kind of estimate that we look for.

## 2.6 The global existence and uniqueness results

We can conclude from estimates (2.25) and (2.88) that for all  $T > 0$ , there exists  $C_H(T) > 0$  such that

$$(2.89) \quad \forall t \in [0, T], \quad \|\mathbf{H}(t)\|_{H^1} < C_H(T)$$

Besides, we know that  $\text{curl } \mathbf{E} = \frac{\partial \mathbf{H}}{\partial t} + \frac{\partial \mathbf{M}}{\partial t}$ . We recall that  $\frac{\partial \mathbf{M}}{\partial t} \in L^2(\mathbb{R})$  since  $\mathbf{M} \in L^\infty(\mathbb{R})$  and  $\mathbf{H} \in L^2(\mathbb{R})$ . Thus (2.25) and (2.88) yields to : for all  $T > 0$ , there exists  $C_E(T) > 0$  such that

$$(2.90) \quad \forall t \in [0, T], \quad \|\mathbf{E}(t)\|_{H^1} < C_E(T)$$

It remains to see that we can get the same conclusion for  $\mathbf{M}$  by taking the derivative in space of the (LLG) equation :

$$(2.91) \quad \begin{aligned} \frac{\partial}{\partial t} \left( \frac{\partial \mathbf{M}}{\partial x} \right) &= \frac{\partial \mathbf{H}}{\partial x} \times \mathbf{M} + \frac{\partial \mathbf{H}_{eff}}{\partial x} \times \mathbf{M} + \mathbf{H} \times \frac{\partial \mathbf{M}}{\partial x} \\ &+ \frac{\alpha}{|\mathbf{M}|} \frac{\partial \mathbf{M}}{\partial x} \times [(\mathbf{H} + \mathbf{H}_{eff}) \times \mathbf{M}] + \frac{\alpha}{|\mathbf{M}|} \mathbf{M} \times \left[ \left( \frac{\partial \mathbf{H}}{\partial x} + \frac{\partial \mathbf{H}_{eff}}{\partial x} \right) \times \mathbf{M} \right] \\ &+ \frac{\alpha}{|\mathbf{M}|} \mathbf{M} \times \left[ (\mathbf{H} + \mathbf{H}_{eff}) \times \frac{\partial \mathbf{M}}{\partial x} \right] - \alpha \left( \frac{\mathbf{M}}{|\mathbf{M}|^3} \cdot \frac{\partial \mathbf{M}}{\partial x} \right) \mathbf{M} \times \left[ (\mathbf{H} + \mathbf{H}_{eff}) \times \frac{\partial \mathbf{M}}{\partial x} \right] \end{aligned}$$

Then, multiplying by  $\frac{\partial \mathbf{M}}{\partial x}$  and integrating in space, we get :

$$(2.92) \quad \begin{aligned} \frac{\partial}{\partial t} \left( \int_{\mathbb{R}} \left| \frac{\partial \mathbf{M}}{\partial x} \right|^2 dx \right) &= \int_{\mathbb{R}} \left( \frac{\partial \mathbf{H}}{\partial x}, \mathbf{M}, \frac{\partial \mathbf{M}}{\partial x} \right) dx + \int_{\mathbb{R}} \left( \frac{\partial \mathbf{H}_{eff}}{\partial x}, \mathbf{M}, \frac{\partial \mathbf{M}}{\partial x} \right) dx \\ &- \int_{\mathbb{R}} \frac{\alpha}{|\mathbf{M}|} \left( \mathbf{M} \times \frac{\partial \mathbf{M}}{\partial x}, \left( \frac{\partial \mathbf{H}}{\partial x} + \frac{\partial \mathbf{H}_{eff}}{\partial x} \right), \mathbf{M} \right) dx \\ &- \int_{\mathbb{R}} \frac{\alpha}{|\mathbf{M}|} \left( \mathbf{M} \times \frac{\partial \mathbf{M}}{\partial x}, (\mathbf{H} + \mathbf{H}_{eff}), \left( \frac{\partial \mathbf{M}}{\partial x} - \frac{\mathbf{M}}{|\mathbf{M}|^2} \cdot \frac{\partial \mathbf{M}}{\partial x} \mathbf{M} \right) \right) dx \end{aligned}$$

where we denote by  $(\mathbf{a}, \mathbf{b}, \mathbf{c})$  the product  $(\mathbf{a} \times \mathbf{b}) \cdot \mathbf{c}$ . It is then easy to see that (2.92) yields to

$$(2.93) \quad \begin{aligned} \frac{d}{dt} \left( \left\| \frac{\partial \mathbf{M}}{\partial x} \right\|_{L^2}^2 \right) &\leq \|\mathbf{M}\|_{L^\infty} \|\mathbf{H}\|_{H^1} \left\| \frac{\partial \mathbf{M}}{\partial x} \right\|_{L^2} + \|\mathbf{M}\|_{L^\infty} \|\mathbf{H}_{eff}\|_{H^1} \left\| \frac{\partial \mathbf{M}}{\partial x} \right\|_{L^2} \\ &+ \alpha \|\mathbf{M}\|_{L^\infty} \left\| \frac{\partial \mathbf{H}}{\partial x} \right\|_{L^2} \left\| \frac{\partial \mathbf{M}}{\partial x} \right\|_{L^2} + \alpha \|\mathbf{M}\|_{L^\infty} \left\| \frac{\partial \mathbf{H}_{eff}}{\partial x} \right\|_{L^2} \left\| \frac{\partial \mathbf{M}}{\partial x} \right\|_{L^2} \\ &+ 2\alpha \|\mathbf{H}\|_{L^2} \left\| \frac{\partial \mathbf{M}}{\partial x} \right\|_{L^2}^2 + 2\alpha \|\mathbf{H}_{eff}\|_{L^2} \left\| \frac{\partial \mathbf{M}}{\partial x} \right\|_{L^2}^2 \end{aligned}$$

Let us recall that we know that all the terms in the right hand side of (2.93), but maybe  $\|\mathbf{H}_{eff}\|_{H^1}$  and  $\left\| \frac{\partial \mathbf{M}}{\partial x} \right\|_{L^2}$ , are bounded. In what concerns  $\mathbf{H}_{eff}$ , we use the fact that  $\mathbf{H}_{eff}(\mathbf{M})$  is affine with respect to  $\mathbf{M}$ . Then, using Cauchy-Schwartz inequalities, it is straightforward to see that (2.93) leads to : there exists  $C_1(T) > 0$  and  $C_1(T) > 0$  such that, for all  $t \leq T$ ,

$$(2.94) \quad \frac{d}{dt} \left( \left\| \frac{\partial \mathbf{M}}{\partial x} \right\|_{L^2}^2 \right) \leq C_1(T) + C_2(T) \left\| \frac{\partial \mathbf{M}}{\partial x} \right\|_{L^2}^2$$

We can then conclude with the Gronwall's lemma that for all  $T > 0$ , there exists  $C_M(T) > 0$  such that

$$(2.95) \quad \forall t \in [0, T], \quad \|\mathbf{M}(t)\|_{H^1} < C_M(T)$$

Finally, we can state the

**Theorem 4** *Given  $(\mathbf{E}_0, \mathbf{H}_0, \mathbf{M}_0) \in H(\text{curl}, \mathbb{R}) \times H(\text{curl}, \mathbb{R}) \times H^1(\mathbb{R})^3$ , there exists a unique global strong solution of (1.8) so that :*

$$(2.96) \quad \begin{cases} \mathbf{E} \in C^1(\mathbb{R}_+; L^2(\mathbb{R})^3) \cap C^0(\mathbb{R}_+; H^1(\mathbb{R})^2) \\ \mathbf{H} \in C^1(\mathbb{R}_+; L^2(\mathbb{R})^3) \cap C^0(\mathbb{R}_+; H^1(\mathbb{R})^3) \\ \mathbf{M} \in C^2(\mathbb{R}_+; H^1(\mathbb{R})^3) \end{cases}$$

## 2.7 About finite velocity of propagation

Using an energetic method, it is possible to show that the velocity of propagation of electromagnetic waves in ferromagnetic media is finite. The idea of our proof consists in evaluating the electromagnetic energy outside a bounded domain but whose width increases with time. Note that only little changes are required to obtain the same result regardless the dimension of the problem.

More precisely, let us assume that there exists two reals  $a$  and  $b$  such that, at time  $t = 0$ ,

1.  $\mathbf{E}_0$  and  $\mathbf{H}_0$  are equal to 0 outside the interval  $[a, b]$  ;
2. the initial distribution  $\mathbf{M}_0$ , outside  $[a, b]$ , corresponds to a stable steady state ; that means that  $\mathbf{M}_0$  is parallel to  $(\mathbf{H}_s + \mathbf{H}_a(\mathbf{M}))$  outside  $[a, b]$  and that the function

$$(2.97) \quad U(\mathbf{M}) = K|P(\mathbf{M})|^2 + |\mathbf{H}_s - \mathbf{M}|^2$$

reaches a global minimum, almost everywhere outside  $[a, b]$ , for  $\mathbf{M} = \mathbf{M}_0$ .

In this case, it is possible to multiply the Maxwell's equations by  $\mathbf{H}$  and  $\mathbf{E}$  and to integrate the result in space as follows :

$$(2.98) \quad \int_{b+vt}^{+\infty} \left( \varepsilon_0 \mathbf{E} \cdot \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{H} \cdot \frac{\partial \mathbf{H}}{\partial t} \right) dx + \mu_0 \int_{b+vt}^{+\infty} \mathbf{H} \cdot \frac{\partial \mathbf{M}}{\partial t} dx + \int_{b+vt}^{+\infty} \nabla \cdot (\mathbf{E} \times \mathbf{H}) dx = 0$$

Our aim is to estimate the electromagnetic energy for all  $x$  greater than  $b + vt$ , the positive constant  $v$  being defined later. By integrating the first term in time, we get

$$(2.99) \quad \int_{b+vt}^{+\infty} \frac{d}{dt} \left[ \varepsilon_0 |\mathbf{E}|^2 + \mu_0 |\mathbf{H}|^2 + K(x) \right] dx + \mu_0 \int_{b+vt}^{+\infty} \mathbf{H} \cdot \frac{\partial \mathbf{M}}{\partial t} dx + \int_{b+vt}^{+\infty} \nabla \cdot (\mathbf{E} \times \mathbf{H}) dx = 0$$

In this expression, the function  $K(x)$  must be taken constant in time. Our choice is consists in  $K(x) = -U(\mathbf{M}_0)$ . Then, from (2.99), the computations unfolds like in the general case : by adding  $\mathbf{H}_{eff}(\mathbf{M})$  to  $\mathbf{H}$ , we obtain

$$(2.100) \quad \int_{b+vt}^{+\infty} \frac{d}{dt} \left[ \varepsilon_0 |\mathbf{E}|^2 + \mu_0 |\mathbf{H}|^2 - U(\mathbf{M}_0) \right] dx + \mu_0 \int_{b+vt}^{+\infty} \mathbf{H}_T(\mathbf{H}, \mathbf{M}) \cdot \frac{\partial \mathbf{M}}{\partial t} dx - (\mathbf{E} \times \mathbf{H})(b+vt) = 0$$

where we denote by  $(\mathbf{E} \times \mathbf{H})(x)$  the scalar function  $(Ey(x)Hz(x) - Ez(x)Hy(x))$ . Using (2.19) and (2.21) yields to

$$(2.101) \quad \int_{b+vt}^{+\infty} \frac{d}{dt} \left[ \varepsilon_0 |\mathbf{E}|^2 + \mu_0 |\mathbf{H}|^2 + U(\mathbf{M}) - U(\mathbf{M}_0) \right] dx \\ = -\mu_0 \frac{\alpha}{|\gamma|} \int_{b+vt}^{+\infty} \frac{1}{|\mathbf{M}|} \left| \frac{\partial \mathbf{M}}{\partial t} \right|^2 dr + (\mathbf{E} \times \mathbf{H})(b+vt)$$

We deduce that

$$(2.102) \quad \int_{b+vt}^{+\infty} \frac{d}{dt} \left[ \frac{\varepsilon_0}{2} |\mathbf{E}|^2 + \frac{\mu_0}{2} |\mathbf{H}|^2 + U(\mathbf{M}) - U(\mathbf{M}_0) \right] dr \\ \leq -\mu_0 \frac{\alpha}{|\gamma|} \int_{b+vt}^{+\infty} \frac{1}{|\mathbf{M}|} \left| \frac{\partial \mathbf{M}}{\partial t} \right|^2 dr + (|\mathbf{E}| |\mathbf{H}|)(b+vt)$$

As we know that

$$(2.103) \quad \int_{b+vt}^{+\infty} \frac{d}{dt} f(s) ds = \frac{d}{dt} \left[ \int_{b+vt}^{+\infty} f(s) ds \right] + \int_{\Sigma_t} f(s) ds$$

we can finally conclude that

$$(2.104) \quad \frac{d}{dt} \left[ \int_{b+vt}^{+\infty} \left( \frac{\varepsilon_0}{2} |\mathbf{E}|^2 + \frac{\mu_0}{2} |\mathbf{H}|^2 + U(\mathbf{M}) - U(\mathbf{M}_0) \right) dr \right] \leq -\mu_0 \frac{\alpha}{|\gamma|} \int_{b+vt}^{+\infty} \frac{1}{|\mathbf{M}|} \left| \frac{\partial \mathbf{M}}{\partial t} \right|^2 dr \\ - \left[ v \left( \frac{\varepsilon_0}{2} |\mathbf{E}|^2 + \frac{\mu_0}{2} |\mathbf{H}|^2 + U(\mathbf{M}) - U(\mathbf{M}_0) \right) - |\mathbf{E}| |\mathbf{H}| \right] (b+vt)$$

Because of our choice  $K(x) = -U(\mathbf{M}_0)$ , the function of the left hand side, which is the total electromagnetic energy for  $x$  greater than  $b$ , is always positive and equal to 0 at time  $t = 0$ . To show that the speed of propagation is smaller than  $v$ , it suffices to show that this energy can not increase in time, that is to say that the derivative in time can not be positive. Since  $(U(\mathbf{M}) - U(\mathbf{M}_0))(b+vt) \geq 0$ , it remains to choose  $v$  such that

$$(2.105) \quad \left( v \frac{\varepsilon_0}{2} |\mathbf{E}|^2 + v \frac{\mu_0}{2} |\mathbf{H}|^2 - |\mathbf{E}| |\mathbf{H}| \right) (b+vt) \geq 0$$

The reader can easily check that a possible choice is  $v = c_o = \left( \frac{1}{\varepsilon_o \mu_o} \right)^{\frac{1}{2}}$ .

Of course it would remain to do the same for  $x$  smaller than  $a$ .

## Chapter 3

# The numerical method

We present in this third part the discretization of our problem, first in space and then in time. The physical constants are not taken equal to 1 any longer.

For these numerical results we shall work in a bounded interval  $[0, L]$  of  $\mathbb{R}$  with Dirichlet conditions for the electric field  $\mathbf{E}$ . The reader will easily check that the extension to this case of existence and uniqueness theorems of section 2 is straightforward. We introduce the space  $H^0(\text{curl}, \mathbb{R})$  :

$$(3.1) \quad H^0(\text{curl}, \mathbb{R}) = \{\psi \in H(\text{curl}, \mathbb{R}), \psi(0) \times \mathbf{e}_x = 0, \psi(L) \times \mathbf{e}_x = 0\}$$

For the ‘‘Maxwell part’’ of our problem, our choice consists in working from the following variational formulation ([1]) :

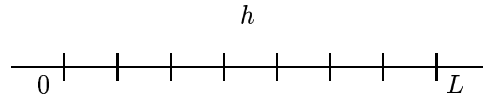
$$(3.2) \quad \begin{cases} \left( \frac{\partial \mathbf{B}}{\partial t}, \phi \right) + (\text{curl } \mathbf{E}, \phi) = 0 & , \quad \forall \phi \in L^2 \\ \varepsilon_0 \left( \frac{\partial \mathbf{E}}{\partial t}, \psi \right) - (\mathbf{H}, \text{curl } \psi) = 0 & , \quad \forall \psi \in H^1 \\ (\mathbf{H}, \phi) = \left( \frac{\mathbf{B}}{\mu_0}, \phi \right) - (\mathbf{M}, \phi) & , \quad \forall \phi \in L^2 \end{cases}$$

It is important to see that the curl term in the second equation have been integrated by parts. For this reason functions  $\phi$  can be taken in  $L^2$  rather than in  $H(\text{curl})$ . The third equation looks like a trivality but yet it indicates that it would be better to look for  $\mathbf{H}$ ,  $\mathbf{B}$  and  $\mathbf{M}$  in the same functional space.

### 3.1 Semi-discretization in space

#### 3.1.1 Construction of the approximate

We assume that  $L = Nh$ , where  $h > 0$  is our space step :



We take our discrete electric field with piecewise linear functions (that we denote as the space  $V_h$ ) and the discrete fields  $\mathbf{H}_h$ ,  $\mathbf{B}_h$  and  $\mathbf{M}_h$  with piece constant functions (that we denote as the space  $X_h$ ). We introduce  $\{\psi_i \in P_1\}_{i=1}^{N-1}$  the canonical basis functions of  $V_h$  to define  $\mathbf{E}_h$  :

$$(3.3) \quad \mathbf{E}_h = \sum_{i=1}^{N-1} (E_y(i)\psi_i(x)\mathbf{e}_y + E_z(i)\psi_i(x)\mathbf{e}_z)$$

and  $\{\phi_i \in P_0\}_{i=1}^N$  the canonical basis functions of  $X_h$  to define  $\mathbf{H}_h$ ,  $\mathbf{B}_h$  and  $\mathbf{M}_h$  :

$$(3.4) \quad \mathbf{B}_h = \sum_{i=1}^N (B_y(i - \frac{1}{2})\phi_i(x)\mathbf{e}_y + B_z(i - \frac{1}{2})\phi_i(x)\mathbf{e}_z)$$



$$(3.5) \quad \mathbf{H}_h = \sum_{i=1}^N \left( -M_x(i - \frac{1}{2})\phi_i(x)\mathbf{e}_x + H_y(i - \frac{1}{2})\phi_i(x)\mathbf{e}_y + H_z(i - \frac{1}{2})\phi_i(x)\mathbf{e}_z \right)$$

$$(3.6) \quad \mathbf{M}_h = \sum_{i=1}^N \left( M_x(i - \frac{1}{2})\phi_i(x)\mathbf{e}_x + M_y(i - \frac{1}{2})\phi_i(x)\mathbf{e}_y + M_z(i - \frac{1}{2})\phi_i(x)\mathbf{e}_z \right)$$

It is now possible to write a discrete variational formulation of (3.2) :

$$(3.7) \quad \begin{cases} \left( \frac{\partial \mathbf{B}_h}{\partial t}, \phi_h \right) + (\text{curl } \mathbf{E}_h, \phi_h) = 0 & , \forall \phi_h \in (P_0)^2 \\ \varepsilon_0 \left( \frac{\partial \mathbf{E}_h}{\partial t}, \psi_h \right) - (\mathbf{H}_h, \text{curl } \psi_h) = 0 & , \forall \psi_h \in (P_1)^2 \end{cases}$$

This leads to the following scalar system for Maxwell's equations :

$$(3.8) \quad 1 \leq i \leq N, \quad \begin{cases} \frac{dB_y(i - \frac{1}{2})}{dt} = \frac{E_z(i) - E_z(i - 1)}{\Delta x} \\ \frac{dB_z(i - \frac{1}{2})}{dt} = -\frac{E_y(i) - E_y(i - 1)}{\Delta x} \end{cases}$$

(where  $E_y(0) = 0 = E_z(0)$  and  $E_y(N) = 0 = E_z(N)$ ), and

$$(3.9) \quad 1 \leq i \leq N - 1, \quad \begin{cases} \frac{dE_y(i)}{dt} = -\frac{H_z(i + \frac{1}{2}) - H_z(i - \frac{1}{2})}{\Delta x} \\ \frac{dE_z(i)}{dt} = \frac{H_y(i + \frac{1}{2}) - H_y(i - \frac{1}{2})}{\Delta x} \end{cases}$$

With this approach a discrete curl operator is defined for the space  $H_h$  : In the 1D case it is exactly the usual operator corresponding to the Yee's scheme ([18]), but this would be no longer true in the 2D case (see [11]). This discrete curl operator is defined with centered finite difference which guarantees order 2 in space for our scheme.

The situation is of course different for the (LLG) equation and the magnetic coupling since no derivative occurs in these equations. We simply have (see (2.6)) :

$$(3.10) \quad \frac{\partial \mathbf{M}_h}{\partial t} = |\gamma| (\mathbf{H}_{Th}(\mathbf{M}) + \mathbf{H}_s) \times \mathbf{M}_h + \frac{\alpha}{|\mathbf{M}|} \mathbf{M}_h \times \frac{\partial \mathbf{M}_h}{\partial t}$$

with  $\mathbf{H}_{Th}(\mathbf{M}) = \mathbf{H}_h + \mathbf{H}_s - KP(\mathbf{M}_h)$  and  $\mathbf{H}_h = \frac{\mathbf{B}_h}{\mu_0} - \mathbf{M}_h$ .

### 3.1.2 Existence of solutions and stability analysis

It can be seen that the problem defined from (3.7) and (3.10) consists in solving the following system :

$$(3.11) \quad \begin{cases} \frac{\partial \mathbf{B}_h}{\partial t} + \text{curl } \mathbf{E}_h = 0 \\ \mathbf{B}_h = \mu_0 (\mathbf{H}_h + \mathbf{M}_h) \\ \varepsilon_0 \frac{\partial \mathbf{E}_h}{\partial t} - \text{curl}_h \mathbf{H}_h = 0 \\ \frac{\partial \mathbf{M}_h}{\partial t} = \mathbf{L}_G^\alpha(\mathbf{M}_h, \mathbf{H}_h) \end{cases}$$

where  $\text{curl}_h$  is the discrete curl operator defined on  $(P_0)^2$  by :

$$(3.12) \quad \forall \phi_h \in (P_0)^2, \forall \psi_h \in (P_1)^2, (\text{curl}_h \phi_h, \psi_h) = (\phi_h, \text{curl } \psi_h)$$

We see that (3.11) is a system of  $(2 + 2 + 3)N$  ordinary differential equations. Thus Cauchy-Lipschitz' theorem applies and ensures that there exists a unique local solution  $(\mathbf{E}_h, \mathbf{H}_h, \mathbf{M}_h)$ .

Concerning this local solution, we see first that there is nothing new to say about the conservation at each point of the norm of  $\mathbf{M}_h$ , since the (LLG) equation is exactly the same as in the continuous case :

$$(3.13) \quad |\mathbf{M}_h(x, t)|^2 = |\mathbf{M}_{0h}(x)|^2, \quad \forall x \in [0, L]$$

To get the complete control of the local solution  $(\mathbf{E}_h, \mathbf{H}_h, \mathbf{M}_h)$ , it suffices then to prove again the decay in time of the electromagnetic energy. If we take  $\phi_h = \mathbf{H}_h$  and  $\psi_h = \mathbf{E}_h$  in (3.7) and then integrate, it comes by adding

$$(3.14) \quad \int_0^L \left\{ \mu_0 \left( \frac{\partial \mathbf{H}_h}{\partial t}, \mathbf{H}_h \right) + \varepsilon_0 \left( \frac{\partial \mathbf{E}_h}{\partial t}, \mathbf{E}_h \right) \right\} dx = \mu_0 \int_0^L \left( \mathbf{H}_h \cdot \frac{\partial \mathbf{M}_h}{\partial t} \right) dx$$

We do the same computations but with the 1D assumptions :

$$(3.15) \quad -\mu_0 \int_0^L \left( \mathbf{H}_{ah} \cdot \frac{\partial \mathbf{M}_h}{\partial t} \right) dx = \mu_0 \int_0^L \left( KP(\mathbf{M}_h) \cdot \frac{\partial P(\mathbf{M}_h)}{\partial t} \right) dx = \frac{1}{2} \frac{d}{dt} [\mathcal{E}_a(\mathbf{M}_h)]$$

Since the (LLG) equation has not been modified, the other results are unchanged and lead to

$$(3.16) \quad \begin{aligned} \frac{d}{dt} \left[ \frac{\varepsilon_0}{2} \|\mathbf{E}_h\|_{L^2}^2 + \frac{\mu_0}{2} \|\mathbf{H}_h\|_{L^2}^2 + \frac{\mu_0 K}{2} \|P(\mathbf{M}_h)\|_{L^2}^2 + \frac{\mu_0}{2} \|\mathbf{H}_s - \mathbf{M}_h\|_{L^2}^2 \right] \\ = -\mu_0 \frac{\alpha}{|\gamma|} \int_0^L \frac{1}{|\mathbf{M}|} \left| \frac{\partial \mathbf{M}_h}{\partial t} \right|^2 dx \end{aligned}$$

Thus, as in the continuous case, we conclude that

$$(3.17) \quad \begin{aligned} \varepsilon_0 \|\mathbf{E}_h(t)\|_{L^2}^2 &\leq \left( \varepsilon_0 \|\mathbf{E}_{h0}\|_{L^2}^2 + \mu_0 \|\mathbf{H}_{h0}\|_{L^2}^2 + (1 + K)\mu_0 \|\mathbf{M}_{h0}\|_{L^2}^2 + \mu_0 \|\mathbf{H}_s\|_{L^2}^2 \right) \\ \mu_0 \|\mathbf{H}_h(t)\|_{L^2}^2 &\leq \left( \varepsilon_0 \|\mathbf{E}_{h0}\|_{L^2}^2 + \mu_0 \|\mathbf{H}_{h0}\|_{L^2}^2 + (1 + K)\mu_0 \|\mathbf{M}_{h0}\|_{L^2}^2 + \mu_0 \|\mathbf{H}_s\|_{L^2}^2 \right) \end{aligned}$$

These estimations ensure that the local solution is indeed a global solution. Moreover, they do not depend on  $h$ .

## 3.2 Time stepping

Our semi-discrete equations must be now discretized in time.

### 3.2.1 Construction of the approximate

We shall use to discretize (3.8) and (3.9) in time the standard leapfrog technique : if  $\Delta t > 0$  is our time step, the electric field is discretized at times  $t^n = n\Delta t$  :

$$(3.18) \quad \mathbf{E}_h^n \simeq \mathbf{E}_h(t = n\Delta t)$$

while the magnetic field, the magnetic induction and the magnetization are discretized at time  $t^{n+\frac{1}{2}} = (n + \frac{1}{2}) \Delta t$  :

$$(3.19) \quad \mathbf{B}_h^{n+\frac{1}{2}} \simeq \mathbf{B}_h \left( (n + \frac{1}{2}) \Delta t \right), \quad \mathbf{M}_h^{n+\frac{1}{2}} \simeq \mathbf{M}_h \left( (n + \frac{1}{2}) \Delta t \right), \quad \mathbf{H}_h^{n+\frac{1}{2}} \simeq \mathbf{H}_h \left( (n + \frac{1}{2}) \Delta t \right)$$

Besides we denote by  $\mathbf{M}_h^n$  (resp.  $\mathbf{H}_h^n$  and  $\mathbf{B}_h^n$ ) the half sum of  $\mathbf{M}_h^{n-\frac{1}{2}}$  and  $\mathbf{M}_h^{n+\frac{1}{2}}$  (resp. of  $\mathbf{H}_h^{n-\frac{1}{2}}$  and  $\mathbf{H}_h^{n+\frac{1}{2}}$ , and of  $\mathbf{B}_h^{n-\frac{1}{2}}$  and  $\mathbf{B}_h^{n+\frac{1}{2}}$ ).

Discretizing the system (3.11), we get the following scheme : on the one hand, for the Maxwell's equations,

$$(3.20) \quad \frac{\mathbf{B}_h^{n+\frac{1}{2}} - \mathbf{B}_h^{n-\frac{1}{2}}}{\Delta t} = -\text{curl}_h \mathbf{E}_h^n$$

and

$$(3.21) \quad \frac{\mathbf{E}_h^{n+1} - \mathbf{E}_h^n}{\Delta t} = \text{curl}_h \mathbf{H}_h^{n+\frac{1}{2}}, \text{ with } \mathbf{H}_h^{n+\frac{1}{2}} = \frac{\mathbf{B}_h^{n+\frac{1}{2}}}{\mu_0} - \mathbf{M}_h^{n+\frac{1}{2}}$$

and on the other hand, for the (LLG) equation,

$$(3.22) \quad \frac{\mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}}}{\Delta t} = |\gamma| \mathbf{H}_{Th}^n \times \frac{\mathbf{M}_h^{n+\frac{1}{2}} + \mathbf{M}_h^{n-\frac{1}{2}}}{2} + \alpha \frac{\mathbf{M}_h^{n+\frac{1}{2}} + \mathbf{M}_h^{n-\frac{1}{2}}}{2} \times \frac{\mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}}}{\Delta t}$$

with  $\mathbf{H}_{Th}^n = \frac{1}{\mu_0} \mathbf{B}_h^n + \mathbf{H}_s - KP(\mathbf{M}_h^n)$  and

$$|\mathbf{M}_h|^n = \frac{|\mathbf{M}_h^{n-\frac{1}{2}}| + |\mathbf{M}_h^{n+\frac{1}{2}}|}{2}.$$

When  $\mathbf{B}_h^{n-\frac{1}{2}}$ ,  $\mathbf{M}_h^{n-\frac{1}{2}}$  and  $\mathbf{E}_h^n$  are assumed to be known, the computations at the following step unfold like this :

1.  $\mathbf{B}_h^{n+\frac{1}{2}}$  is calculated by (3.20).
2.  $\mathbf{M}_h^{n+\frac{1}{2}}$  is calculated by (3.22), hence  $\mathbf{H}_{Th}^n$  can be computed.
3.  $\mathbf{E}_h^{n+1}$  is calculated by (3.21).

Moreover it can be seen that our time stepping is centered which ensures an approximation of order 2 in time.

### 3.2.2 Existence and actual computation of the discrete solution

The calculus of  $\mathbf{B}_h^{n+\frac{1}{2}}$ ,  $\mathbf{H}_h^{n+\frac{1}{2}}$  and  $\mathbf{E}_h^{n+1}$  from (3.20) and (3.21) is obvious but we have to make precise the way we determinate  $\mathbf{M}_h^{n+\frac{1}{2}}$  from (3.22). Moreover, because of non-linearity of this equation, we can't be sure *a priori* neither of existence of a solution, nor of uniqueness. Also we shall see that although the scheme is implicit, the computation of  $\mathbf{M}_h^{n+1}$  can be done explicitly.

- **A new form of the discrete (LLG) equation.**

First of all, we can "decrease" the non-linearity of (3.22) with the help of three simple mathematical remarks :

1. The conservation of the norm of the magnetization : one only has to take the scalar product of (3.22) by  $\mathbf{M}_h^n$  to find

$$(3.23) \quad \left( \frac{\mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}}}{\Delta t} \right) \cdot \left( \frac{\mathbf{M}_h^{n+\frac{1}{2}} + \mathbf{M}_h^{n-\frac{1}{2}}}{2} \right) = |\mathbf{M}_h^{n+\frac{1}{2}}|^2 - |\mathbf{M}_h^{n-\frac{1}{2}}|^2 = 0$$

This means that  $|\mathbf{M}_h|^n$  is a actual constant in (3.22) and the superscript  $n$  can be omitted.

2. A change in the second vector product : we have

$$(3.24) \quad \frac{\mathbf{M}_h^{n+\frac{1}{2}} + \mathbf{M}_h^{n-\frac{1}{2}}}{2} \times \frac{\mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}}}{\Delta t} = \frac{\mathbf{M}_h^{n+\frac{1}{2}} + \mathbf{M}_h^{n-\frac{1}{2}}}{2} \times \frac{\left( \mathbf{M}_h^{n+\frac{1}{2}} + \mathbf{M}_h^{n-\frac{1}{2}} \right) - 2\mathbf{M}_h^{n-\frac{1}{2}}}{\Delta t}$$

Then, (3.22) becomes

$$(3.25) \quad \mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}} = \left[ \frac{|\gamma| \Delta t}{2} \mathbf{H}_{Th}^n + \alpha \mathbf{M}_h^{n-\frac{1}{2}} \right] \times \left( \mathbf{M}_h^{n+\frac{1}{2}} + \mathbf{M}_h^{n-\frac{1}{2}} \right)$$

3. The remark (2.6) about  $\mathbf{H}_a(\mathbf{M})$  : we have

$$(3.26) \quad -KP(\mathbf{M}_h^n) \times \mathbf{M}_h^n = K(\mathbf{p} \cdot \mathbf{M}) \mathbf{p} \times \mathbf{M}$$

Finally we have shown that the (LLG) equation is equivalent to

$$(3.27) \quad \mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}} = \left[ \frac{|\gamma|\Delta t}{2} (\mathbf{B}_h^n + \mathbf{H}_s) + \alpha \mathbf{M}_h^{n-\frac{1}{2}} - \frac{K|\gamma|\Delta t}{2} (\mathbf{p} \cdot \mathbf{M}_h^n) \mathbf{p} \right] \times (2\mathbf{M}_h^n)$$

Introducing the following notations :

- $\mathbf{f} = 2\mathbf{M}_h^{n-\frac{1}{2}}$ ,  $\mathbf{a} = -\left(\frac{|\gamma|\Delta t}{2} (\mathbf{B}_h^n + \mathbf{H}_0) + \alpha \mathbf{M}_h^{n-\frac{1}{2}}\right)$  and  $\lambda = \frac{K|\gamma|\Delta t}{4}$  ;
- $\mathbf{x} = \mathbf{M}_h^{n+\frac{1}{2}} + \mathbf{M}_h^{n-\frac{1}{2}} = 2\mathbf{M}_h^n$ ,

we get the following non-linear vectorial equation :

$$(3.28) \quad \mathbf{x} + \mathbf{a} \times \mathbf{x} + \lambda(\mathbf{p} \cdot \mathbf{x}) \mathbf{p} \times \mathbf{x} = \mathbf{f}$$

where  $\mathbf{f}$ ,  $\mathbf{a}$  and  $\lambda$  are data and  $\mathbf{x}$  is our new unknown.

**Remark 6** If  $K$  was equal to 0, we would get  $\mathbf{x} + \mathbf{a} \times \mathbf{x} = \mathbf{f}$  ; in this case, lemma 1 applies, and we conclude that

$$(3.29) \quad \mathbf{x} = \frac{\mathbf{f} + (\mathbf{a} \cdot \mathbf{f})\mathbf{a} - \mathbf{a} \times \mathbf{f}}{1 + |\mathbf{a}|^2}$$

- **Existence and uniqueness of the solution.**

We first consider the linear function of lemma 1

$$(3.30) \quad A(\mathbf{x}) = \mathbf{x} + \mathbf{a} \times \mathbf{x} = \mathbf{y} \Leftrightarrow \mathbf{x} = A^{-1} \mathbf{y}$$

We recall that  $A^{-1}$  is a contraction. Let  $\mathbf{x}$  be now a solution to (3.28) ; we introduce

$$(3.31) \quad \mathbf{x} = A^{-1}(\mathbf{f} - \lambda(\mathbf{p} \cdot \mathbf{x}) \mathbf{p} \times \mathbf{x}) = F(\mathbf{x})$$

It is not difficult to see that  $F : B(0, 1 + |\mathbf{f}|) \mapsto B(0, 1 + |\mathbf{f}|)$  as soon as  $\lambda \leq \frac{1}{(1 + |\mathbf{f}|)^2}$  :

$$(3.32) \quad |F(\mathbf{x})| \leq |\mathbf{f}| + \lambda|\mathbf{x}|^2 \leq |\mathbf{f}| + \lambda(1 + |\mathbf{f}|)$$

Then we compute

$$(3.33) \quad \begin{aligned} |F(\mathbf{x}) - F(\mathbf{y})| &\leq \lambda |(\mathbf{p} \cdot \mathbf{x}) \mathbf{p} \times \mathbf{x} - (\mathbf{p} \cdot \mathbf{y}) \mathbf{p} \times \mathbf{y}| \\ &\leq \lambda |(\mathbf{p} \cdot \mathbf{x})(\mathbf{p} \times \mathbf{x} - \mathbf{p} \times \mathbf{y}) - (\mathbf{p} \cdot \mathbf{x} - \mathbf{p} \cdot \mathbf{y}) \mathbf{p} \times \mathbf{y}| \\ &\leq \lambda |(\mathbf{p} \cdot \mathbf{x}) \mathbf{p} \times (\mathbf{x} - \mathbf{y}) - \mathbf{p} \cdot (\mathbf{x} - \mathbf{y}) \mathbf{p} \times \mathbf{y}| \end{aligned}$$

Hence

$$(3.34) \quad |F(\mathbf{x}) - F(\mathbf{y})| \leq 2 \left(1 + \frac{K}{\mu_0}\right) (1 + |\mathbf{f}|) |\mathbf{x} - \mathbf{y}|$$

We conclude that  $F$  is a contraction from  $B(0, 1 + |\mathbf{f}|)$  to  $B(0, 1 + |\mathbf{f}|)$  as soon as

$$(3.35) \quad \lambda \leq \text{Max} \left( \left[ 2 \left(1 + \frac{K}{\mu_0}\right) (1 + |\mathbf{f}|) \right]^{-1}, (1 + |\mathbf{f}|)^{-2} \right)$$

Therefore, for  $\Delta t$  small enough, the contraction theorem applies and there exists a unique solution to (3.28).

• **Actual computation.**

If we try to expand (3.28), a quadratic vectorial equation, it may lead to solve a scalar polynomial of degree 6. We should use in this case some iterative method and give up the hope of an explicit scheme.

To avoid this difficulty, our idea consists in projecting – when it is possible – the vectorial unknown  $\mathbf{x}$  in the basis  $(\mathbf{a}, \mathbf{p}, \mathbf{a} \times \mathbf{p})$ . We distinguish two cases

1.  $|\mathbf{a} \times \mathbf{p}| = 0$ .

Then  $(\mathbf{a}, \mathbf{p}, \mathbf{a} \times \mathbf{p})$  is not a basis, but there exists  $\theta \in \mathbf{R}$  such that  $\mathbf{a} = \theta \mathbf{p}$ . Thus equation (3.28) is equivalent to

$$(3.36) \quad \mathbf{x} + [\theta + \lambda(\mathbf{p} \cdot \mathbf{x})] \mathbf{p} \times \mathbf{x} = \mathbf{f}$$

Taking the scalar product by  $\mathbf{p}$  yields to  $(\mathbf{p} \cdot \mathbf{x}) = (\mathbf{p} \cdot \mathbf{f})$ , so that (3.28) is equivalent to

$$(3.37) \quad \mathbf{x} + [\theta + \lambda(\mathbf{p} \cdot \mathbf{f})] \mathbf{p} \times \mathbf{x} = \mathbf{f}$$

and lemma 1 applies. The resolution is explicit.

2.  $|\mathbf{a} \times \mathbf{p}| \neq 0$ .

We can work in the basis  $(\mathbf{a}, \mathbf{p}, \mathbf{a} \times \mathbf{p})$  and get a system of 3 equations in scalar unknowns  $X = (\mathbf{a} \cdot \mathbf{x})$ ,  $Y = (\mathbf{p} \cdot \mathbf{x})$  and  $Z = (\mathbf{a} \times \mathbf{p}) \cdot \mathbf{x}$  by taking the scalar product of (3.28) by  $\mathbf{a}$ ,  $\mathbf{p}$  and  $\mathbf{a} \times \mathbf{p}$  :

$$(3.38) \quad \begin{cases} Y - Z = (\mathbf{u} \cdot \mathbf{f}) \\ X + \lambda Y Z = (\mathbf{a} \cdot \mathbf{f}) \\ Z - (\mathbf{a} \cdot \mathbf{u}) X + \|\mathbf{a}\|^2 Y + \lambda(\mathbf{a} \cdot \mathbf{u}) Y^2 - \lambda X Y = (\mathbf{a} \times \mathbf{u}) \cdot \mathbf{f} \end{cases}$$

This leads to

$$(3.39) \quad \begin{cases} Y = Z + (\mathbf{u} \cdot \mathbf{f}) \\ X = (\mathbf{a} \cdot \mathbf{f}) - \lambda Z (Z + (\mathbf{u} \cdot \mathbf{f})) \\ \lambda^2 Z^3 + 2\lambda((\mathbf{a} \cdot \mathbf{u}) + \lambda(\mathbf{u} \cdot \mathbf{f})) Z^2 \\ + \left(1 + \|\mathbf{a}\|^2 - \lambda(\mathbf{a} \cdot \mathbf{f}) + 3\lambda(\mathbf{a} \cdot \mathbf{u})(\mathbf{u} \cdot \mathbf{f}) + \lambda^2(\mathbf{u} \cdot \mathbf{f})^2\right) Z \\ - \lambda(\mathbf{a} \cdot \mathbf{f})(\mathbf{u} \cdot \mathbf{f}) - (\mathbf{a} \cdot \mathbf{u})(\mathbf{u} \cdot \mathbf{f}) + \lambda(\mathbf{a} \cdot \mathbf{u})(\mathbf{u} \cdot \mathbf{f})^2 + \|\mathbf{a}\|^2(\mathbf{u} \cdot \mathbf{f}) - (\mathbf{a} \times \mathbf{u}) \cdot \mathbf{f} = 0 \end{cases}$$

Solving this system is equivalent to find the – unique – real root of a polynomial of degree 3 in  $Z$ , while degree 6 was expected. This time again, the resolution is explicit. Knowing  $X$ ,  $Y$  and  $Z$ , we can determinate  $\mathbf{x}$  by :

$$(3.40) \quad \mathbf{x} = \frac{1}{|\mathbf{a} \times \mathbf{p}|^2} [(X - (\mathbf{a} \cdot \mathbf{p})Y) \mathbf{a} + (|\mathbf{a}|^2 Y - (\mathbf{a} \cdot \mathbf{p})X) \mathbf{p} + Z \mathbf{a} \times \mathbf{p}]$$

### 3.2.3 Stability analysis

We want to show the decay at each time step of a discrete electromagnetic energy by adapting the proof of the continuous case. Multiplying (3.20) by  $\Delta t \left( \mathbf{H}_h^{n+\frac{1}{2}} + \mathbf{H}_h^{n-\frac{1}{2}} \right)$ , we obtain

$$(3.41) \quad \mu_0 \|\mathbf{H}_h^{n+\frac{1}{2}}\|^2 - \mu_0 \|\mathbf{H}_h^{n-\frac{1}{2}}\|^2 + \mu_0 \left( \mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}}, \frac{\mathbf{B}_h^{n+\frac{1}{2}} + \mathbf{B}_h^{n-\frac{1}{2}}}{\mu_0} - (\mathbf{M}_h^{n+\frac{1}{2}} + \mathbf{M}_h^{n-\frac{1}{2}}) \right) \\ + \Delta t \left( \text{curl}_h \mathbf{E}_h^n, \mathbf{H}_h^{n+\frac{1}{2}} + \mathbf{H}_h^{n-\frac{1}{2}} \right) = 0$$

Using the conservation of  $|\mathbf{M}_h|$ , the fact that the operator  $curl_h$  as we defined it is a self-adjoint operator, and (3.21), this is equivalent to

$$(3.42) \quad \left[ \mu_0 \|\mathbf{H}_h^{n+\frac{1}{2}}\|^2 + \varepsilon_0(\mathbf{E}_h^n, \mathbf{E}_h^{n+1}) \right] - \left[ \mu_0 \|\mathbf{H}_h^{n-\frac{1}{2}}\|^2 + \varepsilon_0(\mathbf{E}_h^{n-1}, \mathbf{E}_h^n) \right] = -\mu_0 \left( \mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}}, \frac{2\mathbf{B}_h^n}{\mu_0} \right)$$

By adding  $\mathbf{H}_{eff}(\mathbf{M}_h^n)$  to  $\frac{2\mathbf{B}_h^n}{\mu_0}$ , we obtain

$$(3.43) \quad \left[ \mu_0 \|\mathbf{H}_h^{n+\frac{1}{2}}\|^2 + \varepsilon_0(\mathbf{E}_h^n, \mathbf{E}_h^{n+1}) \right] - \left[ \mu_0 \|\mathbf{H}_h^{n-\frac{1}{2}}\|^2 + \varepsilon_0(\mathbf{E}_h^{n-1}, \mathbf{E}_h^n) \right] - \mu_0 \left( \mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}}, \mathbf{H}_{eff}(\mathbf{M}_h^n) \right) = -\mu_0 \left( \mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}}, \mathbf{H}_{Th}^n \right)$$

Let us compute the left hand side ; we have

$$(3.44) \quad \mu_0 \left( 2KP(\mathbf{M}_h^n), \mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}} \right) = \left| 2KP(\mathbf{M}_h^{n+\frac{1}{2}}) \right|^2 - \left| 2KP(\mathbf{M}_h^{n-\frac{1}{2}}) \right|^2$$

Finally we obtain that  $-\mu_0 \left( \mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}}, \mathbf{H}_{eff}(\mathbf{M}_h^n) \right)$  is equal to

$$(3.45) \quad \left[ 2K\mu_0 \left| P(\mathbf{M}_h^{n+\frac{1}{2}}) \right|^2 - \mathbf{H}_s \cdot \mathbf{M}_h^{n+\frac{1}{2}} \right] - \left[ 2K\mu_0 \left| P(\mathbf{M}_h^{n-\frac{1}{2}}) \right|^2 - \mathbf{H}_s \cdot \mathbf{M}_h^{n-\frac{1}{2}} \right]$$

Besides one has only to take the vector product of (3.22) by  $\mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}}$  to get

$$(3.46) \quad -\mu_0 \left( \mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}}, \mathbf{H}_{Th}^n \right) = -\frac{\alpha}{\gamma} \mu_0 \left| \mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}} \right|^2$$

Thus, if we introduce,  $\forall n \in \mathbf{N}$ ,

$$(3.47) \quad U_h(n + \frac{1}{2}) = \mu_0 \|\mathbf{H}_h^{n+\frac{1}{2}}\|^2 + \varepsilon_0(\mathbf{E}_h^n, \mathbf{E}_h^{n+1}) + \mu_0 |2K\mu_0(\mathbf{M}_h^{n+\frac{1}{2}})|^2 + \mu_0 |\mathbf{H}_s - \mathbf{M}_h^{n+\frac{1}{2}}|^2$$

we have established that,  $\forall n \geq 1$ ,

$$(3.48) \quad U_h(n + \frac{1}{2}) - U_h(n - \frac{1}{2}) = -\frac{\alpha}{\gamma} \mu_0 \left| \mathbf{M}_h^{n+\frac{1}{2}} - \mathbf{M}_h^{n-\frac{1}{2}} \right|^2$$

This shows that the discrete quantity  $U_h(n + \frac{1}{2})$  decreases at each step of the computation. We know that

$\mu_0 |2K\mu_0(\mathbf{M}_h^{n+\frac{1}{2}})|^2 + \mu_0 |\mathbf{H}_s - \mathbf{M}_h^{n+\frac{1}{2}}|^2$  is a positive term. What can we conclude about the quantity  $\hat{U}_h(n + \frac{1}{2}) =$

$\mu_0 \left\| \mathbf{H}_h^{n+\frac{1}{2}} \right\|^2 + \varepsilon_0(\mathbf{E}_h^n, \mathbf{E}_h^{n+1})$  ? Using the fact that

$$(\mathbf{E}_h^n, \mathbf{E}_h^{n+1}) = \left\| \frac{\mathbf{E}_h^{n+1} + \mathbf{E}_h^n}{2} \right\|^2 - \frac{\Delta t^2}{4} \left\| \frac{\mathbf{E}_h^{n+1} - \mathbf{E}_h^n}{\Delta t} \right\|^2 = \left\| \frac{\mathbf{E}_h^{n+1} + \mathbf{E}_h^n}{2} \right\|^2 - \frac{\Delta t^2}{4} \left\| curl_h \mathbf{H}_h^{n+\frac{1}{2}} \right\|^2$$

and the fact that  $\|curl_h \mathbf{H}_h^{n+\frac{1}{2}}\| \leq \frac{4}{h^2} \|\mathbf{H}_h^{n+\frac{1}{2}}\|^2$ , we see that

$$(3.49) \quad \hat{U}_h(n + \frac{1}{2}) \geq \left( 1 - \frac{\Delta t^2}{h^2} \right) \|\mathbf{H}_h^{n+\frac{1}{2}}\|^2 + \left\| \frac{\mathbf{E}_h^n + \mathbf{E}_h^{n+1}}{2} \right\|^2$$

and we can conclude that  $\hat{U}_h(n + \frac{1}{2})$  is positive, and therefore that the scheme is stable, under the classical

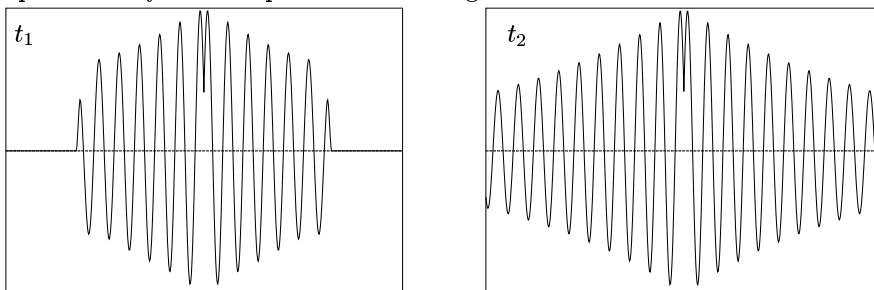
CFL condition :  $\frac{\Delta t}{h} \leq 1$ .

## Chapter 4

# Numerical experiments

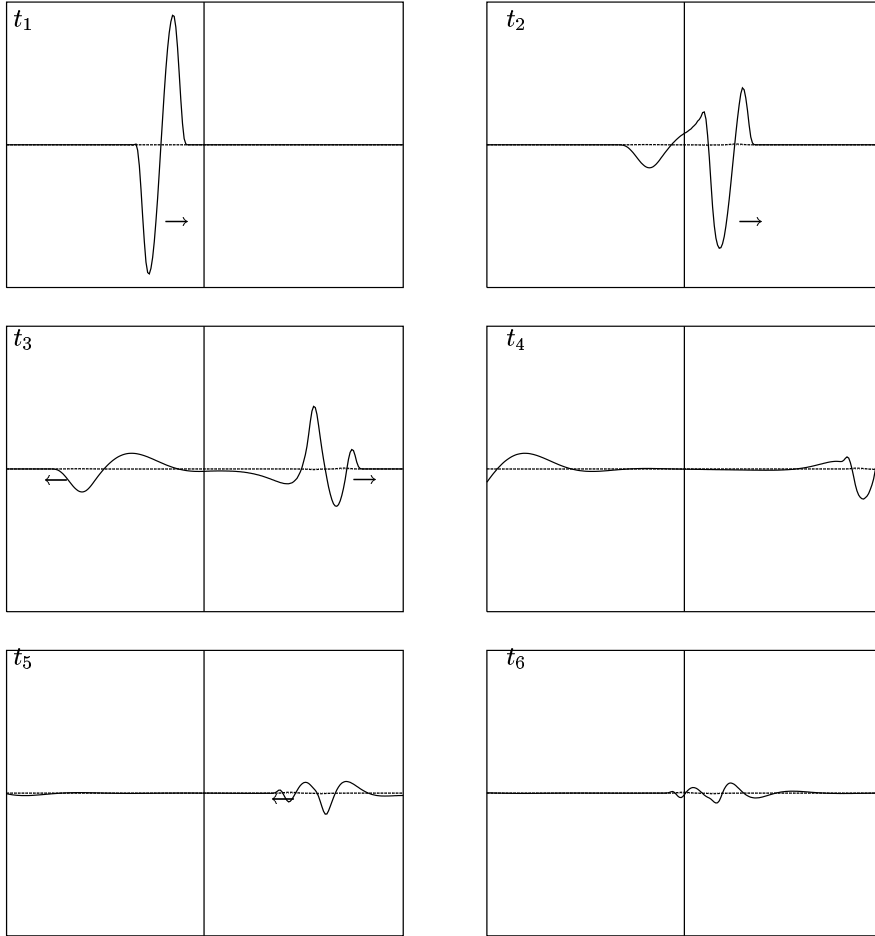
In this last part, we perform some numerical simulations. The results that we present are not really “discussed” ; our goal consists only in illustrating our numerical method.

First of all we consider the propagation of a sustained signal inside a homogeneous ferromagnetic domain. We represent only one component of the magnetic field  $\mathbf{H}$  in two different cases



On the first one, we see the decay of the propagating signal. On the second one, a steady state is reached.

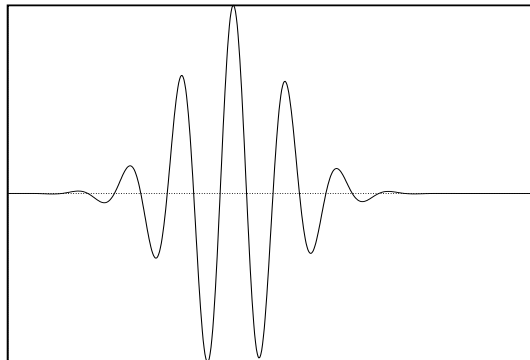
We show now some scattering experiments. The calculus domain is divided in two parts : on the left side we consider a vacuum ( $\mathbf{M} = 0$ ), on the right side a ferromagnetic material. Six snapshots of one component of the magnetic field  $\mathbf{H}$  are represented.



At time  $t_2$  and  $t_3$ , we see a reflected part of the signal due to the interface propagating to the left, while a transmitted part is decreasing inside the ferromagnetic media. At time  $t_4$ , the signal is totally reflected on the right because of a Dirichlet boundary condition.

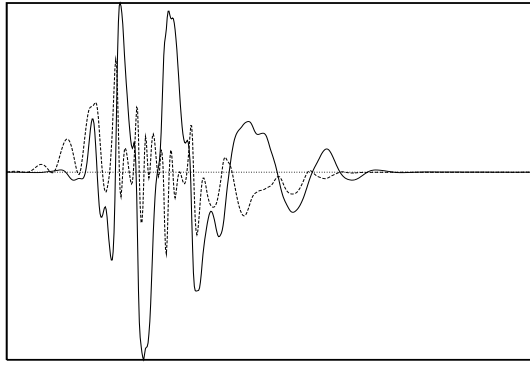
To end this final part, we would like to show some amazing nonlinear effects on a electromagnetic signal crossing a ferromagnetic layer. Rather than snapshots of the signal, we study the fields in time at two different points, respectively before and after the layer.

For the following incident signal  $\mathbf{H}(t) = h(t)\mathbf{e}_z$ ,

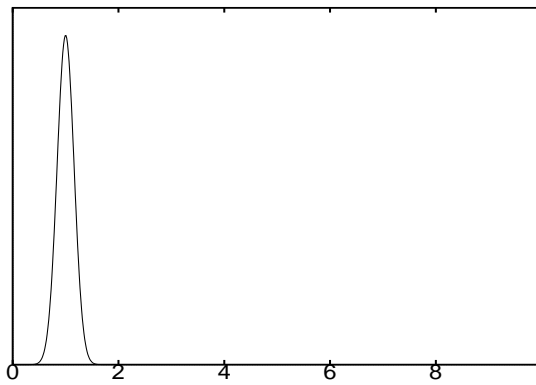




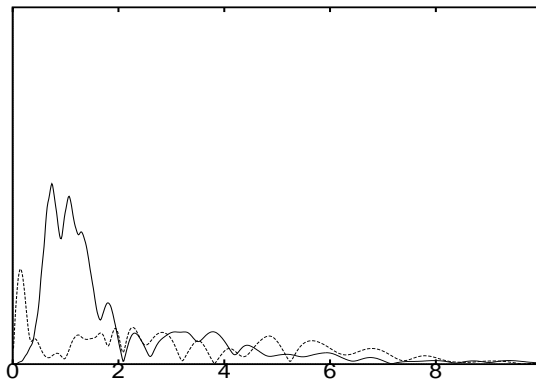
we obtain, after the ferromagnetic layer, the following transmitted signal



It is important to see that both the nature and the polarization of the signal have been modified. A part of the incident signal has rotated. To end this part, we give below the spectrum, obtained by Fourier transform, of the incident signal



and of the transmitted signal



As a typically nonlinear effect of our model, we see here that a lot of new frequencies are created.

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