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## - To cite this version:

Alain Krinik, Gerardo Rubino, David Beecher, Heba Ayeda. Steady-state Distributions of Somewhat Stochastic Matrices. 2023 - 7th Conference on Stochastic Models, Institute of Mathematics, Polish
Academy of Sciences, May 2023, Będlewo, Poland. pp.1-66. hal-04389349

HAL Id: hal-04389349
https://inria.hal.science/hal-04389349
Submitted on 11 Jan 2024

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## Steady-state Distributions of Somewhat Stochastic Matrices

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Dedication: To Pawel Lorek

Stochastic Models VII
Bẹdlewo, Poland
Monday, May 29 - Friday, June 2, 2023

## Sylvester's Eigenvalue Expansion Formula for $M^{k}$

Assume $M$ is a real $n \times n$ matrix that has $n$ known distinct eigenvalues $\omega_{1}, \omega_{2}, \ldots, \omega_{n}$. Then

$$
\begin{equation*}
M^{k}=A_{1} \omega_{1}^{k}+A_{2} \omega_{2}^{k}+A_{3} \omega_{3}^{k}+\ldots+A_{n} \omega_{n}^{k} \tag{1}
\end{equation*}
$$

where $k=0,1,2, \ldots$ and $A_{1}, A_{2}, \ldots, A_{n}$ are constant $n \times n$ coefficient matrices called spectral projectors (Meyer) or Frobenius covariants.

$$
M^{k}(i, j)=A_{1}(i, j) \omega_{1}^{k}+A_{2}(i, j) \omega_{2}^{k}+A_{3}(i, j) \omega_{3}^{k}+\ldots+A_{n}(i, j) \omega_{n}^{k}
$$

where $0 \leq i, j \leq n-1$ for $s=1,2,3, \ldots, n$.
Note: $A_{s}$ is the matrix outer product, $c^{2} \vec{R}_{\vec{L}} \vec{L}_{s}$, where $\vec{R}_{s}$ and $\vec{L}_{s}$ are right \& left eigenvectors of matrix $M$ and where $\vec{L}_{s}$ and $\vec{R}_{s}$ are scaled by a constant $c$ so that their dot product $c^{2} \vec{L}_{s} \cdot \vec{R}_{s}=1$.

## Extension of the Sylvester Eigenvalue Expansion

When the eigenvalues of our matrix $M$ are not distinct, then a more general form of Sylvester's Eigenvalue Expansion Formula must be used

$$
\begin{equation*}
M^{k}=\sum_{i=1}^{n} A_{i} \sum_{j=0}^{m_{i}-1}\binom{k}{j}\left(\omega_{i}\right)^{k-j}\left(M-\omega_{i}\right)^{j} \tag{2}
\end{equation*}
$$

where $n$ is the dimension of $M$ and $m_{i}$ is the multiplicity of the $i$ th eigenvalue. [Wik19b].

## Definitions

All matrices that appear in the following definitions are assumed to be real square matrices. Definitions $1 \& 2$ are conventional. Definitions $3 \& 4$ are comparatively new.

Definition 1. A sequence of matrices $A_{k}$ converge to a matrix $L$ if and only if $\left[A_{k}\right]_{j j} \rightarrow L_{i j}$ as $k \rightarrow \infty$ for all $i, j$.

Definition 2. A Markov chain with one-step transition probability matrix $P$ has a conventional steady-state distribution $\vec{\pi}$ if and only if $\lim _{k \rightarrow \infty} P^{k}=L$, where $L$ is a stochastic matrix and the rows of $L$ are identically $\vec{\pi}$.

Definition 3. A matrix $P$ is defined to be somewhat stochastic if the sum of each row of $P$ is 1 . Note that the entries of $P$ are permitted to be negative. [ĆJ15]

Definition 4. A somewhat stochastic matrix $P$ has a steady state distribution $\vec{\pi}$ if $\lim _{k \rightarrow \infty} P^{k}=L$ where the rows of $L$ are identically $\vec{\pi}$. Note that while the entries of $\vec{\pi}$ are permitted to be negative, the sum of the entries of $\vec{\pi}$ still equal one.

When does $\lim _{k \rightarrow \infty} A^{k}$ exists in Terms of the Eigenvalues of Matrix $A$ ?

Let $B=\{\omega \in \mathbb{C}:|\omega|<1$ or $\omega=1\}$, which is the open unit ball centered at the origin and the number 1 .

Theorem 1. Let $A$ be a square matrix with complex entries. Then $\lim _{k \rightarrow \infty} A^{k}$ exists if and only if both of the following conditions hold:

- Every eigenvalue of $A$ is an element of $B$.
- If 1 is an eigenvalue of $A$, then the dimension of the eigenspace corresponding to 1 equals the algebraic multiplicity of 1 .

Remark 1. Friedberg, Insel, and Spence discuss different proofs of Theorem 1 on page 284 of [FIS18].

## Somewhat Stochastic Matrices

Theorem 2. Suppose $P$ is a finite-dimensional, one-step transition probability matrix of a Markov chain. The Markov chain has a steady state distribution if and only if every eigenvalue of $P$ is an element of $B$ and the eigenvalue $\omega=1$ of $P$ has algebraic multiplicity 1 .

Theorem 3. Suppose $S$ is a finite-dimensional, somewhat stochastic matrix. $S$ has a steady-state distribution if and only if every eigenvalue of $S$ is an element of $B$ and the eigenvalue $\omega=1$ of $S$ has algebraic multiplicity 1.

Theorem 4 (Classical). Suppose $P$ is a finite-dimensional, one-step transition probability matrix of an irreducible Markov chain. The Markov chain has a steady-state distribution $\vec{\pi}$ if and only if the Markov chain has a stationary distribution $\vec{\pi}$ and is aperiodic [HPS72] [KS76].

## Sketch of the Converse for Theorem 2

For simplicity, we assume $P$ is $4 \times 4$ and satisfies Sylvester's eigenvalue expansion with distinct eigenvalues $\omega_{1}, \omega_{2}, \omega_{3}, \omega_{4} \in B$ where $\omega_{1}=1$ and $B=\{\omega \in \mathbb{C}:|\omega|<1$ or $\omega=1\}$ and $\left|\omega_{2}\right|,\left|\omega_{3}\right|,\left|\omega_{4}\right|<1$ :

$$
P^{k}=A_{1} \omega_{1}^{k}+A_{2} \omega_{2}^{k}+A_{3} \omega_{3}^{k}+A_{4} \omega_{4}^{k}
$$

So $\lim _{k \rightarrow \infty} P^{k}=A_{1}=c^{2} \overrightarrow{R_{1}} \overrightarrow{L_{1}}$ where $\overrightarrow{R_{1}}=\overrightarrow{1}$ and $\overrightarrow{L_{1}}$ are the right and left eigenvectors of $P$ corresponding to $\omega_{1}=1$, then

$$
A_{1}=c^{2} \vec{R}_{1} \vec{L}_{1}=c^{2}\left[\begin{array}{l}
1 \\
1 \\
1 \\
1
\end{array}\right]\left[\begin{array}{llll}
I_{1} & I_{2} & I_{3} & I_{4}
\end{array}\right]=c^{2}\left[\begin{array}{llll}
I_{1} & I_{2} & I_{3} & I_{4} \\
I_{1} & I_{2} & I_{3} & I_{4} \\
I_{1} & I_{2} & I_{3} & I_{4} \\
I_{1} & I_{2} & I_{3} & I_{4}
\end{array}\right]
$$

## Sketch of the Converse for Theorem 2 and Theorem 3

Since $P$ is stochastic, it follows that $P^{k}$ is also stochastic and this implies that the $\lim _{k \rightarrow \infty} P^{k}$ is also stochastic. Therefore, $A_{1}$ is stochastic and the rows form the steady-state distribution $\vec{\pi}$ of $P$ by Definition 2. This conclusion still holds for non-distinct eigenvalues using equation (2) in place of equation (1).

Remark 2. The proof of the converse of Theorem 3 for somewhat stochastic matrices is essentially the same proof as the converse of Theorem 2 for stochastic matrices.

## Sketch of the Original Implication for Theorem 2 and 3

Suppose $P$ has a double eigenvalue root 1. Using equation (2)

$$
\begin{aligned}
P^{k}= & \sum_{i=1}^{n} A_{i} \sum_{j=0}^{m_{i}-1}\binom{k}{j}\left(\omega_{i}\right)^{k-j}\left(P-\omega_{i} I\right)^{j} \\
= & A_{1}[I+k(P-I)]+A_{1^{\prime}}[I+k(P-I)] \\
& +\sum_{i=3}^{n} A_{i} \sum_{j=0}^{m_{i}-1}\binom{k}{j}\left(\omega_{i}\right)^{k-j}\left(P-\omega_{i} I\right)^{j} \\
= & \left(A_{1}+A_{1^{\prime}}\right)[I+k(P-I)]+\sum_{i=3}^{n} A_{i} \sum_{j=0}^{m_{i}-1}\binom{k}{j}\left(\omega_{i}\right)^{k-j}\left(P-\omega_{i} I\right)^{j}
\end{aligned}
$$

SO

$$
\lim _{k \rightarrow \infty} P^{k}=\left(A_{1}+A_{1^{\prime}}\right)
$$

## Sketch of the Original Implication for Theorem 2 and 3

$$
\lim _{k \rightarrow \infty} P^{k}=\left(A_{1}+A_{1^{\prime}}\right)
$$

This leads to a contradiction because if $P$ converges in distribution by Definition 2, then its limit is a matrix of rank 1 . However, $A_{1}+A_{1^{\prime}}$ has rank 2 by [rh], which is the contradiction. This argument also applies to show that if $P$ has a steady-state distribution by Definition 2, then $P$ can only have eigenvalue 1 with algebraic multiplicity 1.

Remark 3. Using Theorem 3, we will be able to develop an alternative method to find Gambler's Ruin probabilities. This will be explained later as a consequence of algebraic Duality Theory. For now, we illustrate steady-state distributions (Definition 4) for somewhat stochastic matrices that are $2 \times 2$.

## Example 1: General $2 \times 2$ Real Matrix

Consider the following transition diagram where $P$ is a stochastic matrix:


$$
P=\left[\begin{array}{cc}
1-p & p \\
q & 1-q
\end{array}\right]
$$

where $0<p, q<1$
$P$ has eigenvalues $\omega_{1}=1$ and $\omega_{2}=1-p-q$, and the Sylvester expansion:

$$
\begin{aligned}
P^{k} & =A_{1} \omega_{1}^{k}+A_{2} \omega_{2}^{k} \\
& =\left[\begin{array}{cc}
\frac{q}{p+q} & \frac{p}{p+q} \\
\frac{q}{p+q} & \frac{p}{p+q}
\end{array}\right](1)^{k}+\left[\begin{array}{cc}
\frac{p}{p+q} & -\frac{p}{p+q} \\
-\frac{q}{p+q} & \frac{q}{p+q}
\end{array}\right](1-p-q)^{k}
\end{aligned}
$$

## New Assumption $-1<1-p-q<1$ Instead of $0<p, q<1$

For illustration, consider the following somewhat stochastic transition matrix: $P=\left[\begin{array}{rr}-\frac{1}{2} & \frac{3}{2} \\ \frac{1}{3} & \frac{2}{3}\end{array}\right]$ where $p=\frac{3}{2}, q=\frac{1}{3}$, and $p+q=\frac{3}{2}+\frac{1}{3}=\frac{11}{6}$.

As before, the eigenvalues are $\omega_{1}=1$ and $\omega_{2}=1-p-q=1-\frac{3}{2}-\frac{1}{3}=-\frac{1}{6}$. By letting $k \rightarrow \infty$ in Sylvester's Eigenvalue Expansion, we obtain the following:

$$
\lim _{k \rightarrow \infty} P^{k}=\lim _{k \rightarrow \infty}\left(A_{1} \omega_{1}^{k}+A_{2} \omega_{2}^{k}\right)=A_{1}=\left[\begin{array}{cc}
\frac{q}{p+q} & \frac{p}{p+q} \\
\frac{q}{p+q} & \frac{p}{p+q}
\end{array}\right]=\left[\begin{array}{cc}
\frac{2}{11} & \frac{9}{11} \\
\frac{2}{11} & \frac{9}{11}
\end{array}\right]=L
$$

So by Definition 4, $P$ has a steady state distribution $\vec{\pi}=\left[\begin{array}{cc}\frac{2}{11} & \frac{9}{11}\end{array}\right]$. When does this behavior happen?

We need $-1<\omega_{2}<1$, which means $-1<1-p-q<1$. This implies we are in the following slanted strip:


The Steady State Distribution of $P=\left[\begin{array}{cc}1-p & p \\ q & 1-q\end{array}\right]$ is $\vec{\pi}=\left[\begin{array}{cc}\frac{q}{p+q} & \frac{p}{p+q}\end{array}\right]$ when $-1<1-p-q<1$.

## Cool Alternative Algorithm to Find Unique Stationary Distributions

For convenience, we describe this result for a 4-state irreducible one-step transition probability matrix $P$. Assume $\vec{\pi}=\left[\begin{array}{llll}\pi_{0} & \pi_{1} & \pi_{2} & \pi_{3}\end{array}\right]$ is its unique stationary distribution and suppose $\widetilde{P}$ is as defined below:

$$
P=\left[\begin{array}{llll}
P_{00} & P_{01} & P_{02} & P_{03} \\
P_{10} & P_{11} & P_{12} & P_{13} \\
P_{20} & P_{21} & P_{22} & P_{23} \\
P_{30} & P_{31} & P_{32} & P_{33}
\end{array}\right] \quad \widetilde{P}=\left[\begin{array}{cccc}
P_{00}-1 & P_{01} & P_{02} & 1 \\
P_{10} & P_{11}-1 & P_{12} & 1 \\
P_{20} & P_{21} & P_{22}-1 & 1 \\
P_{30} & P_{31} & P_{32} & 1
\end{array}\right]
$$

Then $\widetilde{P}$ is invertible and $\vec{\pi}=\left[\begin{array}{llll}0 & 0 & 0 & 1\end{array}\right] \widetilde{P}^{-1}$. Moreover

$$
\pi_{0}=\frac{N_{0}}{D} \quad \pi_{1}=\frac{N_{1}}{D} \quad \pi_{2}=\frac{N_{2}}{D} \quad \pi_{3}=\frac{N_{3}}{D}
$$

where $N_{i}=\widetilde{M}_{i 3}, \widetilde{M}_{i j}$ is the minor of the $i j$ entry of $\widetilde{P}$, and $D$ is the determinant of $\widetilde{P}$. The result is due to Mihoc (1934) and can be found in [Pyk03]; special thanks to Stewart Ethier for this reference.

## Example 2: The Dual Matrix of a Birth-Death Matrix



$$
P^{*}=\left[\begin{array}{cccc}
1 & 0 & 0 & 0 \\
0.3 & 0.5 & 0.2 & 0 \\
0 & 0.2 & 0.1 & 0.7 \\
0 & 0 & 0 & 1
\end{array}\right]
$$



## Formal Definition of the Dual Matrix

Assume $\mathrm{P}=\left[P_{i, j}\right]$ for $i, j$ going from $0,1, \ldots, H$ is a $(H+1) \times(H+1)$ stochastic matrix. Suppose that the $(H+2) \times(H+2)$ matrix $P^{*}=\left[P_{i, j}^{*}\right]$ where:

$$
P_{i, j}^{*}=\sum_{s=i+1}^{H}\left[P_{j+1, s}-P_{j, s}\right]
$$

is also a stochastic matrix under the following conventions:

$$
\begin{aligned}
& \text { - } P_{-1, s}=0 \text { if }-1<s \leq H \\
& \text { - } P_{H, H}^{*}=1 \\
& P_{i, H}^{*}=1-\sum_{s=0}^{H-1} P_{i, s}^{*}
\end{aligned}
$$

Then $P^{*}$ is called the dual matrix of $P$. Note that $P^{*}$ may not exist as a stochastic matrix. $P$ and $P^{*}$ have the same set of eigenvalues, see [Lyc18].

## Solving $P^{(k)}$ Using the Duality Theorem

Once we solve the P* we can use the Duality Theorem to calculate each element of our P matrix

$$
\begin{equation*}
P_{i, j}^{*(k)}=\sum_{s=i+1}^{H}\left[P_{j+1, s}^{(k)}-P_{j, s}^{(k)}\right] \quad P_{i, j}^{(k)}=\sum_{s=i}^{H}\left[P_{j, s}^{*(k)}-P_{j-1, s}^{*(k)}\right] \tag{3}
\end{equation*}
$$

for $n \geq 0$ and for all states $i, j=0,1,2, \ldots, H$ with the conventions:

$$
P_{-1, s}^{(k)}=0 \text { if } s>-1
$$

$$
P_{H, H}^{*(k)}=1
$$

In particular, the finite time Gambler's Ruin probability is:

$$
\begin{equation*}
P_{i,-1}^{*(k)}=\sum_{s=i+1}^{H} P_{0, s}^{(k)} \quad \text { where } i=0,1,2, \cdots, H-1 . \tag{4}
\end{equation*}
$$

Assume $P$ is the 1 -step transition probability matrix of an irreducible finite Markov chain, then we know (see [HPS72]) there exists a steady state distribution:

$$
\vec{\pi}=[\pi(0), \pi(1), \pi(2), \ldots, \pi(H)]
$$

And then

$$
\lim _{k \rightarrow \infty} \sum_{s=i+1}^{H} P_{0, s}^{(k)}=\sum_{s=i+1}^{H} \lim _{k \rightarrow \infty} P_{0, s}^{(k)}=\sum_{s=i+1}^{H} P_{0, s}^{(\infty)}=\sum_{s=i+1}^{H} \pi(s)=\lim _{k \rightarrow \infty} P_{i,-1}^{*(k)}=P_{i,-1}^{*(\infty)}
$$

where $P_{i,-1}^{*(\infty)}$ is the Gambler's Ruin probability starting at $i$ and ending at -1 in infinite time. So

$$
\begin{equation*}
P_{i,-1}^{*(\infty)}=\sum_{s=i+1}^{H} \pi(s) \tag{5}
\end{equation*}
$$

## Example 3: Gambler's Ruin Probabilities Using Duality



By our Algorithm, $P$ has the following steady-state distribution
$P=\left[\begin{array}{cccc}0.75 & 0.05 & 0.05 & 0.15 \\ 0.5 & 0.1 & 0.05 & 0.35 \\ 0.4 & 0.1 & 0.05 & 0.45 \\ 0.2 & 0.05 & 0.05 & 0.7\end{array}\right] \quad \begin{array}{ll} & \pi(0)=0.504 \\ \pi(1)=0.055 \\ & \pi(2)=0.05\end{array} \pi(3)=0.391$

## $\infty$-Time Gambler's Ruin Probabilities Using Duality Theory



By the Duality Theorem, it follows
\(P^{*}=\left[\begin{array}{ccccc}1 \& 0 \& 0 \& 0 \& 0 <br>
0.25 \& 0.25 \& 0.1 \& 0.2 \& 0.2 <br>
0.2 \& 0.2 \& 0.1 \& 0.25 \& 0.25 <br>
0.15 \& 0.2 \& 0.1 \& 0.25 \& 0.3 <br>

0 \& 0 \& 0 \& 0 \& 1\end{array}\right] \quad\)| $P_{0,-1}^{*(\infty)}(p)=\pi(1)+\pi(2)+\pi(3)=0.496$ |
| :--- |
| $P_{1,-1}^{*(\infty)}(p)=\pi(2)+\pi(3)=0.441$ |
| $P_{2,-1}^{*(\infty)}(p)=\pi(3)=0.391$ |

## Example 4: Starting with a Dual Markov Chain

Consider the following transition diagram of a conventional five-state Markov chain having two absorbing states.


## Original System Has Both Negative Transitions \& $\vec{\pi}$ Entries



By our Algorithm, $P$ has the unconventional "steady-state" distribution

$$
P=\left[\begin{array}{cccc}
0.75 & 0.1 & -0.05 & 0.2 \\
0.5 & 0.15 & -0.05 & 0.4 \\
0.4 & 0.15 & -0.05 & 0.5 \\
0.2 & 0.1 & -0.05 & 0.75
\end{array}\right]
$$

$$
\begin{array}{ll}
\pi(0)=0.491 & \pi(1)=0.103 \\
\pi(2)=-0.05 & \pi(3)=0.456
\end{array}
$$

Eigenvalues:

$$
\begin{array}{ll}
\omega_{1}=1 & \omega_{2}=1 / 20 \\
\omega_{3}=11 / 20 & \omega_{4}=0
\end{array}
$$

Infinite Time Gambler's Ruin Probabilities Still Found as in Example 3


By the Duality Theorem, it follows

$$
\begin{aligned}
& P_{0,-1}^{*(\infty)}(p)=\pi(1)+\pi(2)+\pi(3)=0.103-0.05+0.457=0.509 \\
& P_{1,-1}^{*(\infty)}(p)=\pi(2)+\pi(3)=-0.05+0.457=0.407 \\
& P_{2,-1}^{*(\infty)}(p)=\pi(3)=0.457
\end{aligned}
$$

## General Gambler's Ruin Problem Solved Using Duality

The Gambler's Ruin diagram can be thought of as the dual transition diagram of a general $n$-state chain with one-step transition probability matrix $P$. We let the figure below represent a Markov chain having 2 absorbing states and $n-2$ transient states.


The Dual transition diagram on the previous slide has the corresponding one-step transition probability matrix
$P^{*}=\left[\begin{array}{cccccc}1 & 0 & 0 & 0 & 0 & 0 \\ p_{0,-1}^{*} \\ p_{1,-1}^{*} \\ p_{2,-1}^{*} & {\left[\begin{array}{c}p_{0,0}^{*} \\ p_{1,0}^{*} \\ p_{3,-1}^{*} \\ p_{3,-1}^{*}\end{array}\right.} & p_{1,1}^{*} & p_{0,2}^{*} & p_{1,2}^{*} & p_{0,3}^{*} \\ 0 & 0 & p_{1,3}^{*} \\ p_{3,0}^{*} & p_{3,1}^{*} & p_{3,2}^{*} & p_{2,3}^{*} & p_{0,4}^{*} \\ p_{1,4}^{*} \\ p_{3,3}^{*}\end{array}\right]$

The $k$ th power of $P^{*}$ can be represented as:

$$
\left(P^{*}\right)^{k}=\left[\begin{array}{cccccc}
1 & 0 & 0 & 0 & 0 & 0 \\
p_{0,-1}^{*(k)} & {\left[\begin{array}{c}
p_{0,0}^{*(k)}
\end{array}\right.} & p_{0,1}^{*(k)} & p_{0,2}^{*(k)} & p_{0,3}^{*(k)} \\
p_{1,-1}^{*(k)} & p_{1,0}^{*(k)} & p_{1,1}^{*(k)} & p_{1,2}^{*(k)} & p_{1,3}^{*(k)} \\
p_{2,-1}^{*(k)} & p_{0,4}^{*(k)} \\
p_{2,0}^{*(k)} & p_{2,1}^{*(k)} & p_{2,2}^{*(k)} & p_{2,3}^{*(k)} \\
p_{3,-1}^{*(k)} & {\left[\begin{array}{c}
\text { (k) } \\
p_{3,0}^{*(k)} \\
0
\end{array}\right.} & p_{3,1}^{*(k)} & p_{3,2}^{*(k)} & p_{3,3}^{*(k)} & p_{2,4}^{*(k)} \\
0 & 0 & 0 & 0 & p_{3,4}^{*(k)} \\
0
\end{array}\right]
$$

where $p_{i, j}^{*(k)}$ are the $k$-step transition probabilities.
the center matrix of $P^{*}$ is $C^{*}$ as shown below:

$$
C^{*}=\left[\begin{array}{llll}
p_{0,0}^{*} & p_{0,1}^{*} & p_{0,2}^{*} & p_{0,3}^{*} \\
p_{1,0}^{*} & p_{1,1}^{*} & p_{1,2}^{*} & p_{1,3}^{*} \\
p_{2,0}^{*} & p_{2,1}^{*} & p_{2,2}^{*} & p_{2,3}^{*} \\
p_{3,0}^{*} & p_{3,1}^{*} & p_{3,2}^{*} & p_{3,3}^{*}
\end{array}\right]
$$

and

$$
\left(C^{*}\right)^{k}=\left[\begin{array}{llll}
p_{0,0}^{*(k)} & p_{0}^{*(k)} & p_{0,2}^{*(k)} & p_{0,3}^{*(k)} \\
p_{1,0}^{*(k)} & p_{1,1}^{*(k)} & p_{1,2}^{*(k)} & p_{1,3}^{*(k)} \\
p_{2,0}^{*(k)} & p_{2,1}^{*(k)} & p_{2,2}^{*(k)} & p_{2,3}^{*(k)} \\
p_{3,0}^{*(k)} & p_{3,1}^{*(k)} & p_{3,2}^{*(k)} & p_{3,3}^{*(k)}
\end{array}\right]
$$

Note that

$$
\lim _{k \rightarrow \infty}\left(C^{*}\right)^{k}=\left[\begin{array}{llll}
0 & 0 & 0 & 0  \tag{6}\\
0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0
\end{array}\right]
$$

because states $0,1,2$, and 3 are all transient states. Suppose $C^{*}$ has distinct eigenvalues $\omega_{0}, \omega_{1}, \omega_{2}, \omega_{3}$, then by Sylvester's eigenvalue expansion:

$$
\left(C^{*}\right)^{k}=C_{0}^{*} \omega_{0}^{k}+C_{1}^{*} \omega_{1}^{k}+C_{2}^{*} \omega_{2}^{k}+C_{3}^{*} \omega_{3}^{k}
$$

where $C_{0}^{*}, C_{1}^{*}, C_{2}^{*}$, and $C_{3}^{*}$ are the spectral projectors of $C^{*}$. By (6), it follows that $\left|\omega_{0}\right|,\left|\omega_{1}\right|,\left|\omega_{2}\right|,\left|\omega_{3}\right|<1$.

We assume the original system has one-step transition probability matrix $P$ where $P$ is somewhat stochastic.


Since it is known that the matrix of the original system, $P$, and $P^{*}$ have the same set of eigenvalues, see [Lyc18], it follows that the eigenvalues of $P$ are $1, \omega_{0}, \omega_{1}, \omega_{2}$, and $\omega_{3}$. Assuming these eigenvalues are distinct, the Sylvester eigenvalue expansion of $P^{k}$ is

$$
P^{k}=A_{-1}+A_{0} \omega_{0}^{k}+A_{1} \omega_{1}^{k}+A_{2} \omega_{2}^{k}+A_{3} \omega_{3}^{k}
$$

And so

$$
\begin{equation*}
\lim _{k \rightarrow \infty} P^{k}=A_{-1} \tag{7}
\end{equation*}
$$

because $\left|\omega_{0}\right|,\left|\omega_{1}\right|,\left|\omega_{2}\right|,\left|\omega_{3}\right|<1$. In other words, from the equation (7) and Definition 4, $P$ has a steady-state distribution, even though some of the entries may be negative.

## Remarks

Remark 4. Finding the preceding steady-state distribution of $P$ as shown
in Equation (7) provides an alternative approach to find infinite time
Gambler's Ruin probabilities using Equation (5) even when some entries of $P$ may be negative.

Remark 5. Another classical method for determining infinite time
Gambler's Ruin probabilities uses the Fundamental Matrix approach given in [KS76].

Remark 6. Gambler's Ruin probabilities have historically been by determined by solving a system of linear recurrences, see [HPS72].

## Xiaoxiao Eigenvalues



## Summary

(1) Suppose $P^{*}$ is any conventional $n$-state Markov chain having exactly two absorbing states and $n-2$ transient states, then $P$ is a somewhat stochastic matrix that has a single root of unity in its characteristic polynomial and all the remaining eigenvalues lie within the open unit disk on the complex plane. By Theorem 3, $P$ has a steady-state distribution $\vec{\pi}$. $\vec{\pi}$ may be found by determining the spectral projector of $P$ corresponding to eigenvalue $\omega=1$. This implies a solution of the infinite gambler's ruin probability for $P^{*}$ using the duality equation (5).
(2) If the eigenvalues and the eigenvectors of $P^{*}$ are also known, then by the Sylvester eigenvalue expansion (1) or the generalized Sylvester eigenvalue expansion (2), $P_{i, j}^{*(k)}$ is explicitly known. It follows by the Duality Theorem (3), $P_{i, j}^{(k)}$ is also known. So by equation (4), the finite-time gambler's ruin probabilities are explicitly known.

## Summary

(3) Therefore we have studied and have some results:

- when $P^{*}$ is a tridiagonal Toeplitz matrix
- when $P^{*}$ is a tridiagonal matrix with alternating entries along the diagonals
- when $P^{*}$ is a symmetric, pentadiagonal, Toeplitz matrix
- when $P^{*}$ is a pentadiagonal, Toeplitz matrix with zeros along the first subdiagonal and first superdiagonal.
(4) An interesting element of our duality solution today is the inclusion of somewhat stochastic matrices. The extension of our approach to finding the absorption probability in an infinite state Markov chain setting is still in progress.


## Example 5: Negative Probabilities - Infinite State Space

The next example extends our example on negative probabilities to an infinite number of states. Consider the following diagram for the dual matrix $P^{*}$


## Example 5: Negative Probabilities - Infinite State Space

We can regard the conventional Markov Chain on the previous slide to be the dual of another Markov system. Using the definition of Duality, the transition probabilities of the original Markov system include some negative probabilities as shown in the figure below


## Example 5: Negative Probabilities - Infinite State Space

The one-step transition probability matrix $P$ of the preceding figure looks like this

$$
P=\left[\begin{array}{ccccccc}
1 & -q / 2 & q / 2 & 0 & 0 & 0 & \ldots \\
1 & -q & 0 & q & 0 & 0 & \ldots \\
q & p-q & 0 & q & 0 & 0 & \ldots \\
0 & 0 & p & 0 & q & 0 & \ldots \\
0 & 0 & 0 & p & 0 & q & \ldots \\
0 & 0 & 0 & 0 & p & 0 & \ldots \\
\vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots
\end{array}\right]
$$

To find the stationary distribution $\vec{\pi}$ of $P$, we assume the following,

$$
\begin{equation*}
1=\pi(0)+\pi(1)+\pi(2)+\pi(3)+\pi(4)+\ldots \quad \text { where } q<p \tag{*}
\end{equation*}
$$

Solving $\vec{\pi}=\vec{\pi} P$ gives us the following algebraic relations
$\pi(1)=-\frac{q^{2}}{2\left(q^{2}+p\right)} \pi(0) \quad \pi(2)=\frac{q}{2\left(q^{2}+p\right)} \pi(0) \quad \pi(k)=\frac{q^{2}}{2\left(q^{2}+p\right)}\left(\frac{q}{p}\right)^{k-3} \pi(0)$ for $k \geq 3$. To solve for $\pi(0)$, substitute the preceding equations into (*)

We find that

$$
\begin{aligned}
1 & =\left[1-\frac{q^{2}}{2\left(q^{2}+p\right)}+\frac{q}{2\left(q^{2}+p\right)}\right] \pi(0)+\left[1+\frac{q}{p}+\left(\frac{q}{p}\right)^{2}+\ldots\right] \frac{q\left(1-q^{2}-p\right)}{2 p\left(q^{2}+p\right)} \pi(0) \\
& =\left[1-\frac{q^{2}}{2\left(q^{2}+p\right)}+\frac{q}{2\left(q^{2}+p\right)}\right] \pi(0)+\left(\frac{1}{1-q / p}\right) \frac{q^{2}}{2\left(q^{2}+p\right)} \pi(0) \\
& =\left[1-\frac{q^{2}}{2\left(q^{2}+p\right)}+\frac{q}{2\left(q^{2}+p\right)}+\left(\frac{1}{1-q / p}\right) \frac{q^{2}}{2\left(q^{2}+p\right)}\right] \pi(0) \\
& \text { and so }
\end{aligned}
$$

$\pi(0)=\left[1-\frac{q^{2}}{2\left(q^{2}+p\right)}+\frac{q}{2\left(q^{2}+p\right)}+\left(\frac{1}{1-q / p}\right) \frac{q^{2}}{2\left(q^{2}+p\right)}\right]^{-1}$
Simplifying and substituting $p$ in terms of $q$ by $p=1-q$

$$
\pi(0)=\frac{4 q^{3}-6 q^{2}+6 q-2}{3 q^{3}-4 q^{2}+5 q-2} \quad \pi(1)=-\frac{q^{2}}{2\left(q^{2}-q+1\right)} \pi(0) \quad \pi(2)=\frac{q}{2\left(q^{2}-q+1\right)} \pi(0)
$$

and for $k \geq 3$

$$
\pi(k)=\frac{q^{2}}{2\left(q^{2}-q+1\right)}\left(\frac{q}{1-q}\right)^{k-3} \pi(0)
$$

Substituting $\pi(0)=\left(4 q^{3}-6 q^{2}+6 q-2\right) /\left(3 q^{3}-4 q^{2}+5 q-2\right)$ gives:

$$
\begin{aligned}
& P_{0,-1}^{*(\infty)}(q)=1-\pi(0)=1-\frac{4 q^{3}-6 q^{2}+6 q-2}{3 q^{3}-4 q^{2}+5 q-2} \\
& \begin{aligned}
P_{1,-1}^{*(\infty)}(q) & =1-\pi(0)-\pi(1)=1-\left(1-\frac{q^{2}}{2\left(q^{2}-q+1\right)}\right) \frac{4 q^{3}-6 q^{2}+6 q-2}{3 q^{3}-4 q^{2}+5 q-2} \\
P_{2,-1}^{*(\infty)}(q) & =1-\pi(0)-\pi(1)-\pi(2) \\
& =1-\left(1-\frac{q^{2}}{2\left(q^{2}-q+1\right)}+\frac{q}{2\left(q^{2}-q+1\right)}\right) \frac{4 q^{3}-6 q^{2}+6 q-2}{3 q^{3}-4 q^{2}+5 q-2} \\
P_{3,-1}^{*(\infty)}(q) & =1-\pi(0)-\pi(1)-\pi(2)-\pi(3)
\end{aligned}
\end{aligned}
$$

$$
=1-\left(1-\frac{q^{2}}{2\left(q^{2}-q+1\right)}+\frac{q}{2\left(q^{2}-q+1\right)}+\frac{q^{2}}{2\left(q^{2}-q+1\right)}\right) \frac{4 q^{3}-6 q^{2}+6 q-2}{3 q^{3}-4 q^{2}+5 q-2}
$$

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## Thank

 you!
## Numerical Example

For the case of $p=2 / 3$ and $q=1 / 3, \pi(0)$ and the stationary distribution become

$$
\begin{array}{lll}
\pi(0)=\frac{7}{9} & \pi(1)=-\frac{1}{18} & \pi(2)=\frac{1}{6} \\
\pi(3)=\frac{1}{18} & \pi(4)=\frac{1}{36} & \pi(k)=\frac{1}{18}\left(\frac{1}{2}\right)^{k-3}
\end{array}
$$



## Stationary Distribution



## Absorption Probabilities

## Friedberg Material

Theorem. Let $A \in M_{n \times n}(\mathbb{C})$ satisfy the following two conditions.

- Every eigenvalue of $A$ is contained in $S$.
- $A$ is diagonalizable.

Then $\lim _{m \rightarrow \infty} A^{m}$ exists. See page 284-285 of Friedberg's Linear Algebra [FIS18].

## Example 6

Consider the following dual transition diagram where $0<p, q<1$ and $p+q=1$

and corresponding dual transition matrix $P^{*}$

$$
P^{*}=\left[\begin{array}{cccccccc}
1 & 0 & 0 & 0 & 0 & 0 & 0 & \ldots \\
0 & 0 & p & q & 0 & 0 & 0 & \ldots \\
q / 2 & q / 2 & 0 & p & 0 & 0 & 0 & \ldots \\
0 & q & 0 & 0 & p & 0 & 0 & \ldots \\
0 & 0 & 0 & 0 & 0 & p & q & \ldots \\
0 & 0 & 0 & q / 2 & q / 2 & 0 & p & \ldots \\
0 & 0 & 0 & 0 & q & 0 & 0 & \ldots \\
\vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots
\end{array}\right]
$$

Taking the dual of the dual we obtain the following generalized diagram for the original system:

with the corresponding original matrix $P$

$$
P=\left[\begin{array}{cccccccc}
1 & -q / 2 & q / 2 & 0 & 0 & 0 & 0 & \ldots \\
1 & -q & 0 & q & 0 & 0 & 0 & \ldots \\
q & p-q & 0 & q & 0 & 0 & 0 & \ldots \\
0 & 0 & p & q & -q / 2 & q / 2 & 0 & \ldots \\
0 & 0 & 0 & 1 & -q & 0 & q & \ldots \\
0 & 0 & 0 & q & p-q & 0 & q & \ldots \\
0 & 0 & 0 & 0 & 0 & p & q & \ldots \\
0 & 0 & 0 & 0 & 0 & 0 & 1 & \ldots \\
\vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots
\end{array}\right]
$$

By solving the system $\vec{\pi} P=\vec{\pi}$ we obtain the following for the stationary distribution

$$
\pi(1)=-q C(q) \pi(0) \quad \pi(2)=C(q) \pi(0) \quad \pi(3)=q C(q) \pi(0)
$$

And in general, the following formulas will give us any distribution probability, for $k=0,1,2,3, \ldots$ :

$$
\begin{align*}
& \pi(3 k+1)=-[q C(q)]^{k+1} \pi(0)  \tag{8}\\
& \pi(3 k+2)=q^{k} C(q)^{k+1} \pi(0)  \tag{9}\\
& \pi(3 k+3)=[q C(q)]^{k+1} \pi(0) \tag{10}
\end{align*}
$$

Where

$$
\begin{equation*}
C(q)=\frac{q}{2\left(q^{2}-q+1\right)} \tag{11}
\end{equation*}
$$

Since the terms of the form $\pi(3 k+1)$ and $\pi(3 k+3)$ form a telescoping series, and knowing $\sum_{n=0}^{\infty} \pi(n)=1$, produces the following equation:

$$
1=\pi(0)+\sum_{k=0}^{\infty} \pi(3 k+2)=\pi(0)+\pi(0) \sum_{k=0}^{\infty} q^{k} C(q)^{k+1}=\pi(0)\left[1+\sum_{k=0}^{\infty} q^{k} C(q)^{k+1}\right]
$$

Therefore

$$
\pi(0)=\frac{1}{1+\sum_{k=0}^{\infty} q^{k} C(q)^{k+1}}=\frac{1}{1+C(q) \sum_{k=0}^{\infty} q^{k} C(q)^{k}}
$$

Using the formula for the sum of a geometric series, we obtain

$$
\begin{equation*}
\pi(0)=\frac{q^{2}-2 q+2}{q^{2}-q+2} \tag{12}
\end{equation*}
$$

In general, for $k=0,1,2,3, \cdots$,

$$
\begin{aligned}
& P_{0,-1}^{*(\infty)}(q)=1-\pi(0) \\
& P_{3 k+1,-1}^{*(\infty)}(q)=1-\pi(0)+[q C(q)]^{k+1} \pi(0)-\pi(0) C(q)\left[\frac{q^{k} C(q)^{k}-1}{q C(q)-1}\right] \\
& P_{3 k+2,-1}^{*(\infty)}(q)=1-\pi(0)+[q C(q)]^{k+1} \pi(0)-\pi(0) C(q)\left[\frac{q^{k+1} C(q)^{k+1}-1}{q C(q)-1}\right] \\
& P_{3 k+3,-1}^{*(\infty)}(q)=1-\pi(0)-\pi(0) C(q)\left[\frac{q^{k+1} C(q)^{k+1}-1}{q C(q)-1}\right]
\end{aligned}
$$

Substituting $C(q)=\frac{q}{2\left(q^{2}-q+1\right)}$ from equation (?) and $\pi(0)=\frac{q^{2}-2 q+2}{q^{2}-q+2}$ from equation (5), we obtain the first 4 absorption probabilities:

$$
\begin{aligned}
& P_{0,-1}^{*(\infty)}(q)=1-\pi(0)=\frac{q}{q^{2}-q+2} \\
& P_{1,-1}^{*(\infty)}(q)=1-\pi(0)-\pi(1)=\frac{q^{4}+2 q}{2\left(q^{2}-q+1\right)\left(q^{2}-q+2\right)} \\
& P_{2,-1}^{*(\infty)}(q)=1-\pi(0)-\pi(1)-\pi(2)=\frac{q^{4}-q^{3}+2 q^{2}}{2\left(q^{2}-q+1\right)\left(q^{2}-q+2\right)} \\
& P_{3,-1}^{*(\infty)}(q)=1-\pi(0)-\pi(1)-\pi(2)-\pi(3)=\frac{q^{3}}{2\left(q^{2}-q+1\right)\left(q^{2}-q+2\right)}
\end{aligned}
$$

For $q=1 / 3$, we obtain the following numerical example for the dual chain


With absorption probabilities

$$
\begin{array}{llll}
P_{0,-1}^{*(\infty)}=\frac{3}{16} & P_{1,-1}^{*(\infty)}=\frac{55}{224} & P_{2,-1}^{*(\infty)}=\frac{1}{14} & P_{3,-1}^{*(\infty)}=\frac{3}{224} \\
P_{3 k+1,-1}^{*(\infty)}=\left(\frac{55}{224}\right) \frac{1}{14^{k}} & P_{3 k+2,-1}^{*(\infty)}=\frac{1}{14^{k+1}} & P_{3 k+3,-1}^{*(\infty)}=\left(\frac{3}{16}\right) \frac{1}{14^{k+1}} & k=0,1,2,3, \ldots
\end{array}
$$

where the original chain is shown below


## Sylvester's Eigenvalue Expansion Formula for $M^{k}$

Assume $M$ is a real $n \times n$ matrix that has $n$ known distinct eigenvalues $\omega_{1}, \omega_{2}, \ldots, \omega_{n}$. Then

$$
\begin{equation*}
M^{k}=A_{1} \omega_{1}^{k}+A_{2} \omega_{2}^{k}+A_{3} \omega_{3}^{k}+\ldots+A_{n} \omega_{n}^{k} \tag{13}
\end{equation*}
$$

where $k=0,1,2, \ldots$ and $A_{1}, A_{2}, \ldots, A_{n}$ are constant $n \times n$ coefficient matrices called spectral projectors (Meyer) or Frobenius covariants.

$$
M^{k}(i, j)=A_{1}(i, j) \omega_{1}^{k}+A_{2}(i, j) \omega_{2}^{k}+A_{3}(i, j) \omega_{3}^{k}+\ldots+A_{n}(i, j) \omega_{n}^{k}
$$

where $0 \leq i, j \leq n-1$ for $s=1,2,3, \ldots, n$.
Note: $A_{s}$ is the matrix outer product, $c^{2} \vec{R}_{s} \vec{L}_{s}$, where $\vec{R}_{s}$ and $\vec{L}_{s}$ are right \& left eigenvectors of matrix $M$ and where $\vec{L}_{s}$ and $\vec{R}_{s}$ are scaled by a constant $c$ so that their dot product $c^{2} \vec{L}_{s} \cdot \vec{R}_{s}=1$.

## Example 7: Unconventional Infinite Markov Processes

Consider the $M^{x_{2}} / M^{y_{2}} / 1$ process shown below

with the $Q$ matrix of this Markov Process being

$$
Q=\left[\begin{array}{cccccc}
-(\lambda+\beta) & \lambda & \beta & 0 & 0 & \cdots \\
\mu & -(\lambda+\beta+\mu) & \lambda & \beta & 0 & \cdots \\
\gamma & \mu & -(\lambda+\beta+\mu+\gamma) & \lambda & \beta & \cdots \\
0 & \gamma & \mu & -(\lambda+\beta+\mu+\gamma) & \lambda & -(\lambda+\beta+\mu+\gamma) \\
0 & 0 & \vdots & \mu & \cdots \\
\vdots & \vdots & \vdots & \vdots & \vdots & \ddots
\end{array}\right]
$$

Constructing the rate matrix $Q^{*}$ for the dual process gives us

$$
Q^{*}=\left[\begin{array}{ccccccc}
0 & 0 & 0 & 0 & 0 & \cdots \\
\lambda+\beta & -(\lambda+\beta+\mu) & \mu-\gamma & 0 & \cdots & 0 & \cdots \\
\beta & \lambda & -(\gamma+\beta+\mu+\gamma) & \mu & \mu & \cdots \\
0 & \beta & \beta & \lambda & -(\gamma+\beta+\mu+\gamma) & -(\gamma+\beta+\mu+\gamma) & \cdots \\
0 & 0 & \vdots & \vdots & \vdots & \ddots & \\
\vdots & \vdots & \lambda & & & \\
0 & \vdots & & & & \\
\hline
\end{array}\right.
$$

Note that there is an off-diagonal entry in $Q^{*}$ that could theoretically be negative, $\mu-\gamma$. Thus we assume $\mu>\gamma$ so that the dual process exists and is a well-defined Markov process. Using the $Q^{*}$ matrix we may draw its state transition rate diagram as follows:


## Major Theorem

Consider the $M^{x_{2}} / M^{y_{2}} / 1$ process as shown in the previous two slides. This process has a steady-state distribution iff

$$
2 \beta+\lambda<2 \gamma+\mu
$$

When the steady state distribution exists, it may be represented as

$$
\pi(n)=c_{2} r_{2}^{n}+c_{3} r_{3}^{n} \quad n \geq 0
$$

where

$$
\rho_{n}^{*}=1-\pi(0)-\pi(1)-\cdots-\pi(n-1) \quad \text { or } \quad P_{n, 0}^{*(\infty)}=\sum_{k=n}^{\infty} \pi(k)
$$

where

$$
\begin{aligned}
& r_{2}=-\frac{a}{3}-2 \sqrt{U} \cos \left(\frac{\theta}{3}+\frac{4 \pi}{3}\right) \quad r_{3}=-\frac{a}{3}-2 \sqrt{U} \cos \left(\frac{\theta}{3}+\frac{2 \pi}{3}\right) \\
& c_{2}=\frac{\left(\frac{\beta}{r_{3}}-\gamma r_{3}\right)\left(1-r_{2}\right)\left(1-r_{3}\right)}{\left(r_{3}-r_{2}\right)\left[2 \gamma\left(r_{2}+r_{3}\right)+\mu\right]} \quad c_{3}=\frac{\left(\frac{\beta}{r_{2}}-\gamma r_{2}\right)\left(1-r_{2}\right)\left(1-r_{3}\right)}{\left(r_{2}-r_{3}\right)\left[2 \gamma\left(r_{2}+r_{3}\right)+\mu\right]}
\end{aligned}
$$

## Major Theorem

and where

$$
\begin{array}{lll}
U=\frac{a^{2}-3 b}{9} & V=\frac{2 a^{3}-9 a b+27 c}{54} & \theta=\cos ^{-1}\left(\frac{V}{\sqrt{U^{3}}}\right) \\
a=\frac{\mu+\gamma}{\gamma} & b=-\frac{\beta+\lambda}{\gamma} & c=-\frac{\beta}{\gamma}
\end{array}
$$

Furthermore, whenever $(\lambda / \mu)^{2}=\beta / \gamma$, the steady state distribution exists iff

$$
\frac{\lambda}{\mu}<1 \quad \text { or } \quad \frac{\beta}{\gamma}<1
$$

In this special case, the steady state distribution may then be simplified to

$$
\pi(n)=\left(1-\frac{\lambda}{\mu}\right)\left(\frac{\lambda}{\mu}\right)^{n}
$$

$$
n \geq 0
$$

for the conventional case where $\mu-\gamma \geq 0$, from [KS11].

## First Numerical Example

If we let $\lambda=3, \beta=1, \mu=6, \gamma=4$, our values become

$$
\begin{array}{ll}
r_{2}=-0.1771 & r_{3}=0 \cdot \frac{1}{2} \\
c_{2}=0 & c_{3}=\frac{1}{2}
\end{array}
$$

and the values of the steady state distribution become

$$
\begin{array}{lll}
\pi(0)=\frac{1}{2} & \pi(1)=\frac{1}{4} & \pi(2)=\frac{1}{8} \\
\pi(3)=\frac{1}{16} & \pi(4)=\frac{1}{32} & \pi(5)=\frac{1}{64}
\end{array}
$$

and so the absorption probabilities become

$$
\begin{aligned}
& P_{1,0}^{*(\infty)}=\frac{1}{2} \\
& P_{3,0}^{*(\infty)}=\frac{1}{8}
\end{aligned}
$$

$$
\begin{aligned}
& P_{2,0}^{*(\infty)}=\frac{1}{4} \\
& P_{4,0}^{*(\infty)}=\frac{1}{16}
\end{aligned}
$$

## Transition Diagram



For $\lambda=3, \beta=1, \mu=6, \gamma=4$

## Second Numerical Example

If we let $\lambda=1, \beta=1, \mu=2, \gamma=4$, our values become

$$
\begin{array}{ll}
r_{2}=-0.2929 & r_{3}=\frac{1}{2} \\
c_{2}=0 & c_{3}=\frac{1}{2}
\end{array}
$$

and the values of the steady state distribution become

$$
\begin{array}{lll}
\pi(0)=\frac{1}{2} & \pi(1)=\frac{1}{4} & \pi(2)=\frac{1}{8} \\
\pi(3)=\frac{1}{16} & \pi(n)=\frac{1}{2^{n+1}} &
\end{array}
$$

and so the absorption probabilities become

$$
\begin{array}{ll}
P_{1,0}^{*(\infty)}=\frac{1}{2} & P_{2,0}^{*(\infty)}=\frac{1}{4} \\
P_{3,0}^{*(\infty)}=\frac{1}{8} & P_{n, 0}^{*(\infty)}=\frac{1}{2^{n}}
\end{array}
$$

## Transition Diagram



For $\lambda=1, \beta=1, \mu=2, \gamma=4$

## Summary

- Siegmund duality provides a precise way to calculate Gambler's Ruin probabilities and absorption probabilities (as opposed to simulations and matrix approximations)
- This methodology still "works" using algebraic duality as described in Rubino \& Krinik [RK21] and past JMM presentations when Siegmund duality doesn't exist. We demonstrated in two cases.
- Case 1: The dual chain or process is conventional and the original chain or process is "unconventional". Unconventional means that the chain or process has negative one-step transition probabilities (or rates) or the chain has a somewhat stochastic stationary distribution.

As the eigenvector corresponding to $\omega=1$ is $\vec{R}_{1}=\overrightarrow{1}$, when we construct the spectral projector corresponding to $\omega=1$ we find that the rows are identical:

$$
\begin{aligned}
A_{-1} & =c^{2} \vec{R}_{1} \vec{L}_{1}=c^{2}\left[\begin{array}{l}
r_{1} \\
r_{2} \\
r_{3} \\
r_{4}
\end{array}\right]\left[\begin{array}{llll}
l_{1} & l_{2} & l_{3} & l_{4}
\end{array}\right] \text { where }\left[\begin{array}{l}
r_{1} \\
r_{2} \\
r_{3} \\
r_{4}
\end{array}\right]=\left[\begin{array}{l}
1 \\
1 \\
1 \\
1
\end{array}\right] \\
& =c^{2}\left[\begin{array}{llll}
r_{1} l_{1} & r_{1} l_{2} & r_{1} l_{3} & r_{1} l_{4} \\
r_{2} l_{1} & r_{2} l_{2} & r_{2} /_{3} & r_{2} /_{4} \\
r_{3} l_{1} & r_{3} l_{2} & r_{3} /_{3} & r_{3} /_{4} \\
r_{4} l_{1} & r_{4} l_{2} & r_{4} /_{3} & r_{4} l_{4}
\end{array}\right]=c^{2}\left[\begin{array}{llll}
l_{1} & l_{2} & l_{3} & l_{4} \\
l_{1} & l_{2} & l_{3} & l_{4} \\
l_{1} & l_{2} & l_{3} & l_{4} \\
l_{1} & l_{2} & l_{3} & l_{4}
\end{array}\right]
\end{aligned}
$$

Since $P$ is somewhat stochastic, it follows that $P^{k}$ is also somewhat stochastic and this implies that the $\lim _{k \rightarrow \infty} P^{k}$ is also somewhat stochastic. Therefore, $A_{-1}$ is somewhat stochastic and the rows form the steady-state distribution $\vec{\pi}$ of $P$.

## Summary and Future Work

- Case 2: The original chain or process is conventional and the dual chain or process is unconventional.
- Gambler's Ruin probabilities and absorption probabilities are "correct" as checked by simulations or matrix approximations.
- Much more to do:
- Since negative probabilities are useful, we need to interpret them beyond being just signed measures.
- For example, we know positive recurrent conventional Markov chains have a unique stationary distribution. Does a similar result hold for somewhat stochastic matrices, see [ĆJ15] for some progress in this direction.

