# Efficient a posteriori estimates for the wave equation 

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# Efficient a posteriori estimates for the wave equation 

T. Chaumont-Frelet ${ }^{\star}$

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1 The problem and available solutions

- The problem
- The goal
- The available solutions

2 The proposed error estimates

- The estimator
- The norm

■ The estimates

- A numerical example

3 A glance at the proofs

- A quick detour: the Poisson problem
- The Helmholtz problem
- The wave equation


## The problem and available solutions

The problem and available solutions
The problem

The wave equation

## The wave equation

Given $f: \mathbb{R}_{+} \times \Omega: \rightarrow \mathbb{R}$, find $u: \mathbb{R}_{+} \times \Omega \rightarrow \mathbb{R}$ such that

$$
\left\{\begin{aligned}
\ddot{u}(t)-\Delta u(t) & =f(t) & & \text { in } \Omega, \\
u(t) & =0 & & \text { on } \Gamma,
\end{aligned}\right.
$$

for $t>0$, together with initial conditions

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u(0)=\dot{u}(0)=0 \text { in } \Omega .
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For simplicity, $f$ is piecewise polynomial in space and $\Omega$ is convex.

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For simplicity, $f$ is piecewise polynomial in space and $\Omega$ is convex.

Importantly, we assume that $f$ is smooth in time.

The problem and available solutions The goal

We (semi) discretize the problem with finite element in space. This yields discrete solution $u_{h}: \mathbb{R}_{+} \rightarrow V_{h}$ smooth in time.

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We want to find a suitable norm $\|\mid \cdot\|$ such that estimates like

$$
\int_{0}^{+\infty}\left\|\left(u-u_{h}\right)(t)\right\|_{\Omega}^{2} \leq \int_{0}^{+\infty} \eta(t)^{2} d t
$$

and

$$
\eta_{K}(t) \lesssim\left\|\left(u-u_{h}\right)(t)\right\|_{\tilde{K}}
$$

hold true for a fully-computable estimator $\eta$.

The problem and available solutions
The available solutions

## The available solutions

S. Adjerid, 2002
hierarchical estimator
$\checkmark$ asymptotically exact
$X$ requires Cartesian grids and smooth solutionsC. Bernardi and E. Süli, 2005 residual-based estimator
$\checkmark$ upper bound and lower bounds, fully discrete problem
$\boldsymbol{X}$ weaker norm in the upper bound that in the lower boundM. Picasso, 2010
elliptic reconstructions
$\checkmark$ upper bound in $L^{2}\left(H^{1}\right)$
$X$ no lower bound
E.H. Georgoulis, O. Lakkis and C. Makridakis, 2013 elliptic reconstructions

```
\checkmark ~ u p p e r ~ b o u n d ~ i n ~ L ' \infty ~ ( ~ L ~ ' ~ ) , ~ f u l l y ~ d i s c r e t e ~ p r o b l e m ~
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```

To the best of my knwoledge, no rigorous efficiency proofs.

The proposed error estimates

The proposed error estimates
The estimator

## An instantaneous estimator

## An instantaneous estimator

The estimator is designed to control the "instantaneous residual"

$$
\left|\left(\ddot{u}_{h}, v\right)_{\Omega}+\left(\boldsymbol{\nabla} u_{h}(t), \boldsymbol{\nabla} v\right)_{\Omega}-(f(t), v)_{\Omega}\right| \leq \eta(t)\|\boldsymbol{\nabla} v\|_{\Omega} \quad \forall v \in \boldsymbol{H}_{0}^{1}(\Omega),
$$

at each $t \in \mathbb{R}_{+}$.

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at each $t \in \mathbb{R}_{+}$.
The associated residual-based estimator would be

$$
\eta_{K}(t):=h_{K}\left\|f(t)-\ddot{u}_{h}(t)+\Delta u_{h}(t)\right\|_{K}+h_{K}^{1 / 2}\left\|\llbracket \nabla u_{h}(t) \rrbracket \cdot \boldsymbol{n}\right\|_{\partial K \backslash\ulcorner } .
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$$

Here, we use an equilibrated flux $\sigma(t)$ with $\nabla \cdot \sigma(t)=f(t)-\ddot{u}_{h}(t)$, and

$$
\eta_{K}(t):=\left\|\sigma(t)+\nabla u_{h}(t)\right\|_{\kappa} .
$$

The proposed error estimates
The norm

## The norm

Fix a time-scale $T>0$.

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Let $\xi_{h}:=u-u_{h}$. The "instantaneous" error measure is

$$
\left\|\xi_{h}(t)\right\|_{\Omega}^{2}:=\underbrace{T^{2}\left\|\ddot{\xi}_{h}(t)\right\|_{\Omega}^{2}}_{\text {optional }}+\underbrace{\left\|\dot{\xi}_{h}(t)\right\|_{\Omega}^{2}+\left\|\nabla \xi_{h}(t)\right\|_{\Omega}^{2}}_{\text {physical energy }}
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$$

The exponantial weight means that we essentially control up to $t \sim T$.

The proposed error estimates
The estimates

The estimator is space-time locally efficient:

$$
\eta_{\kappa}(t) \lesssim\left\|\xi_{h}(t)\right\|_{\tilde{\kappa}} .
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The estimator is also reliable:

$$
\int_{0}^{+\infty}\left\|\xi_{h}(t)\right\| \|_{\Omega}^{2} e^{-2 t / T} d t \leq \underbrace{\left(1+o\left(h^{1 / 2}\right)\right)^{2}}_{\text {pre factor }} \int_{0}^{+\infty} \eta(t)^{2} e^{-2 t / T} d t+\underbrace{\operatorname{osc}^{2}(f)}_{\text {data oscillation }}
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$$

$$
\operatorname{osc}(f)=h^{p+1} \times \text { "high power of } T^{\prime \prime} \times \text { "high order time derivatives of } f \text { ". }
$$

The proposed error estimates
A numerical example

Reflection by a wall with $p=2$


## Convergence and effiency




A glance at the proofs

## A glance at the proofs

A quick detour: the Poisson problem

## The Poisson problem

Consider the Poisson problem

$$
(\nabla u, \nabla v)_{\Omega}=(f, v)_{\Omega} \quad \forall v \in H_{0}^{1}(\Omega)
$$

and its conforming FEM discretization $u_{h}$.

With constructions for $\eta$ similar to what I presented, we have

$$
\left|\left(\boldsymbol{\nabla} \xi_{h}, \boldsymbol{\nabla} v\right)_{\Omega}\right| \leq \eta\|\boldsymbol{\nabla} v\|_{\Omega} \quad \forall v \in \boldsymbol{H}_{0}^{1}(\Omega) .
$$

Because $(\boldsymbol{\nabla} \cdot, \boldsymbol{\nabla} \cdot)_{\Omega}$ is actually an energy norm, we have

$$
\left\|\boldsymbol{\nabla} \xi_{h}\right\|_{\Omega} \leq \eta .
$$

What makes the wave equation tough is the lack of inf-sup stability.

## A glance at the proofs

The Helmholtz problem

## The Helmholtz problem

We next consider the Helmholtz problem with complex frequency $s:=\nu+i \rho$

$$
-s^{2}(u, v)_{\Omega}+(\nabla u, \nabla v)_{\Omega}=(f, v)_{\Omega} \quad \forall v \in H_{0}^{1}(\Omega)
$$

The sesquilinear form is not an energy norm. We only have

$$
|s|^{2}\left\|\xi_{h}\right\|_{\Omega}^{2}+\left\|\nabla \xi_{h}\right\|_{\Omega}^{2} \leq\left(1+\frac{|\nu|^{2}}{\rho^{2}}\right) \eta^{2}
$$

Using duality techniques, this may be improved to

$$
|s|^{2}\left\|\xi_{h}\right\|_{\Omega}^{2}+\left\|\nabla \xi_{h}\right\|_{\Omega}^{2} \leq \underbrace{\left(1+C \frac{|\nu|}{\rho} \frac{|\nu| h}{c}\right)^{2}}_{\text {pre factor }} \eta^{2}
$$

$\square$ W. Dörfler and S. Sauter, 2013T. Chaumont-Frelet, A. Ern and M. Vohralík, 2021

## A glance at the proofs

The wave equation

## The wave equation

Key observation: letting $\rho:=1 / T$, we have

$$
\int_{0}^{+\infty}\|\xi(t)\|_{\Omega}^{2} e^{-2 t / T} d t=\int_{\rho-i \infty}^{\rho+i \infty}\left(\frac{|s|^{4}}{\rho^{2}}\|\widehat{\xi}(s)\|_{\Omega}^{2}+|s|^{2}\|\widehat{\xi}(s)\|_{\Omega}^{2}+\|\nabla \widehat{\xi}(s)\|_{\Omega}^{2}\right) d s .
$$

The Laplace transform $\widehat{u}(s)$ of $u(t)$ solves a Helmholtz problem:

$$
\frac{|s|^{4}}{\rho}\left\|\widehat{\xi}_{h}(s)\right\|_{\Omega}^{2}+|s|^{2}\left\|\widehat{\xi}_{h}(s)\right\|_{\Omega}^{2}+\left\|\nabla \widehat{\xi}_{h}(s)\right\|_{\Omega}^{2} \leq\left(1+C \frac{|\nu|^{2}}{\rho^{2}}|\nu| h\right)^{2} \widehat{\eta}(s)^{2} .
$$

Problem: the pre factor is not bounded as $\nu \rightarrow+\infty$.
Solution: use a cutoff frequency $\omega$ and do something else for $|\nu|>\omega$.

## Treatment of high-frequency content

For the high-frequency content, we simply use "plain" stability:

$$
\|\widehat{\xi}(s)\|_{s} \leq\|\widehat{u}(s)\|_{s}+\left\|\widehat{u}_{h}(s)\right\|_{s} \leq 4 \frac{|s|^{2}}{\rho^{2}}\|\widehat{f}(s)\|_{\Omega} \leq 4 \frac{1}{\rho^{2}}\left\|\widehat{f}^{(2)}(s)\right\|_{\Omega}
$$

Next, we use that

$$
\int_{\rho-i \infty}^{\rho+i \infty}\|\widehat{\phi}\|_{\Omega}^{2} \chi_{|s|>\omega} d s \leq \frac{1}{\omega^{2 r}} \int_{0}^{+\infty}\left\|\phi^{(r)}(t)\right\|_{\Omega}^{2} e^{-2 \rho t} d t
$$

It easily follows that term

$$
\int_{\rho-i \infty}^{\rho+i \infty}\left\|\widehat{\xi}_{h}\right\|_{\Omega}^{2} \chi_{|s|>\omega} d s \leq \frac{1}{\omega^{2 r}} \int_{0}^{+\infty}\left\|f^{(r+2)}(t)\right\|_{\Omega}^{2} e^{-2 \rho t} d t
$$

can be made small by approxiately picking $\omega$ and $r$.

## Regrouping the pieces

Reliability follows by regrouping the low- and high-frequency terms:

$$
\begin{gathered}
\int_{0}^{+\infty}\left\|\xi_{h}(t)\right\|^{2} e^{-2 \rho t} d t=\int_{\rho-i \infty}^{\rho+i \infty}\left\|\widehat{\xi}_{h}\right\|_{\Omega}^{2} \chi_{|s|<\omega} d s+\int_{\rho-i \infty}^{\rho+i \infty}\left\|\widehat{\xi}_{h}\right\|_{\Omega}^{2} \chi_{|s|>\omega} d s \\
\leq \\
\left(1+C \frac{\omega^{2}}{\rho^{2}} \frac{\omega h}{c}\right) \int_{0}^{+\infty} \eta(t)^{2} e^{-2 \rho t} d t+\frac{4}{\rho^{2}} \frac{1}{\omega^{2 r}} \int_{0}^{+\infty}\left\|f^{(r+1)}(t)\right\|_{\Omega}^{2} e^{-2 \rho t} d t
\end{gathered}
$$

The estimates I presented earlier simply rely on a clever choice of $\omega$ and $r$.

## Concluding remarks

## Concluding remarks

This is a new framework to derive a posteriori estimates for the wave equation.

It can handle space-varying coefficients and disspative terms (with $\dot{u}$ ).

With equilibration, it gives asymptotically constant-free and p-robust estimates.

Similar (slightly weaker) results are possible with residual-based estimators.

Future challenges include time-discretization and adaptivity.
$\square$ T. Chaumont-Frelet, arxiv:2205.13939, 2022.

