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# Stability criteria for singularly perturbed linear switching systems 

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#### Abstract

We present some results on the uniform exponential stability of singularly perturbed linear systems undergoing switching. In absence of dwell-time constraints, the switching parameter can evolve on the same time scale as the fast variables, or even faster. We investigate the effect of switching laws evolving at a time scale comparable with the fast variables, describing the corresponding asymptotic effect on the slow variable. Based on this analysis, we propose stability criteria for the overall system, uniform for small values of the parameter of singular perturbation.


## 1 Introduction

Singularly perturbed linear switching systems of the type

$$
\begin{aligned}
\dot{x}(t) & =A(t) x(t)+B(t) y(t), \\
\varepsilon \dot{y}(t) & =C(t) x(t)+D(t) y(t),
\end{aligned}
$$

where $\varepsilon$ denotes the small singular perturbation parameter, appear in several industrial and engineering applications (see, e.g., $[3,6,7]$ ). It has been noticed that, differently from the unswitching case [4], their

[^0]stability properties cannot be deduced from those of their formal slow and fast reductions [5, 1].

We give in this paper a new necessary condition for the stability of the singularly perturbed linear switched system for small values of $\varepsilon$. This condition is obtained by considering switching laws which evolve on a time scale of order $\varepsilon$, and by accounting for their asymptotic effect, as $\varepsilon$ goes to zero, on the slow variable $x$.

More precisely, let us consider a signal $t \mapsto$ $(A(t), B(t), C(t), D(t))=: \sigma(t)$ defined on an interval $[0, T]$ and all its reparameterizations $\sigma_{\varepsilon}: t \mapsto \sigma(t / \varepsilon)$ defined on $[0, \varepsilon T]$. The flow of the switching system, for the value of the singular perturbation parameter set to $\varepsilon$, corresponding to $\sigma_{\varepsilon}$ and evaluated at time $\varepsilon T$, turns out to have an effect of order $O(\varepsilon)$ on the coordinate $x$ (since the velocities are of order 1 and the length of the time interval is of order $\varepsilon$ ) and of order $O(1)$ on the coordinate $y$, provided that the switching system $\dot{y}(t)=D(t) y(t)$ is exponentially stable. Suitable computations show that the evolution of $x$ can actually be described as

$$
x \mapsto x+\varepsilon T \Lambda(T, \sigma) x+O\left(\varepsilon^{2}\right)
$$

where the term $\Lambda(T, \sigma)$ does not depend on the initial condition of the variable $y$. We then propose an auxiliary switching system, denoted $\check{\Sigma}$, having as possible modes all the matrices of the type $\Lambda(T, \sigma)$. We provide an explicit expression for such matrices and we prove that, indeed, if $\Sigma$ is exponentially unstable, then the same is true for the original singularly perturbed switching system, for all small enough values of the singular perturbation parameter $\varepsilon$.

### 1.1 Notations

By $\mathbb{R}$ we denote the set of real numbers, by $|\cdot|$ the Euclidean norm of a real vector, and by $\|\cdot\|$ the induced matrix norm. We write $\lfloor x\rfloor$ to denote the evaluation of the floor function at a real number $x$. By $M_{n}(\mathbb{R})$ we denote the set of $n \times n$ real matrices. The $n \times n$ identity matrix is denoted by $I_{n}$. We use $\rho(M)$ to denote the spectral radius of a matrix $M \in$ $M_{n}(\mathbb{R})$, defined as the largest modulus among the eigenvalues of $M$. Given a subset $\mathcal{N}$ of $M_{n}(\mathbb{R})$, we denote by $\mathcal{S}_{\mathcal{N}}$ the set of all measurable functions from $[0,+\infty)$ to $\mathcal{N}$.

## 2 Problem statement

Fix $n, m \in \mathbb{N}$ and a compact set of matrices $\mathcal{M} \subset$ $M_{n+m}(\mathbb{R})$. Let $\Sigma=\left(\Sigma_{\varepsilon}\right)_{\varepsilon>0}$ where, for every $\varepsilon>0$, $\Sigma_{\varepsilon}$ denotes the linear switching system

$$
\begin{aligned}
& \begin{array}{rrr} 
& \Sigma_{\varepsilon}: & \dot{x}(t)
\end{array} \quad A(t) x(t)+B(t) y(t), \\
& \varepsilon \dot{y}(t)=C(t) x(t)+D(t) y(t),
\end{aligned}
$$

where

$$
t \mapsto\left(\begin{array}{ll}
A(t) & B(t) \\
C(t) & D(t)
\end{array}\right)
$$

is an arbitrary element of the set $\mathcal{S}_{\mathcal{M}}$ of measurable functions from $[0,+\infty)$ to $\mathcal{M}$. In particular, piecewise-constant functions (that is, with locally finite set of jumps) with values in $\mathcal{M}$ are admissible signals, but also continuous functions or functions with accumulating sequences of jumps.

For a given $\varepsilon>0$, the usual stability notions, recalled in the following definition, apply to the linear switching system $\Sigma_{\varepsilon}$.

Definition 1. Let $d \in \mathbb{N}$ and $\mathcal{N}$ be a bounded subset of $M_{d}(\mathbb{R})$. Consider the linear switching system

$$
\begin{equation*}
\Sigma_{\mathcal{N}}: \quad \dot{x}(t)=N(t) x(t), \quad N \in \mathcal{S}_{\mathcal{N}} \tag{1}
\end{equation*}
$$

and denote by $\Phi_{N}(t, 0)$ the flow from time 0 to time $t$ of $\Sigma_{\mathcal{N}}$ associated with the switching signal $N$. Then $\Sigma_{\mathcal{N}}$ is said to be

1. exponentially stable if there exist $C>0$ and $\delta>$ 0 such that

$$
\left\|\Phi_{N}(t, 0)\right\| \leq C e^{-\delta t}, \quad \forall t \geq 0, \forall N \in \mathcal{S}_{\mathcal{N}}
$$

2. exponentially unstable if there exist $C>0, \delta>$ 0 , and a nonzero trajectory $t \mapsto x(t)$ of $\Sigma_{\mathcal{N}}$ such that

$$
|x(t)| \geq C e^{\delta t}|x(0)|, \quad \forall t \geq 0
$$

The maximal Lyapunov exponent of $\Sigma_{\mathcal{N}}$ is defined as

$$
\lambda\left(\Sigma_{\mathcal{N}}\right)=\limsup _{t \rightarrow+\infty} \frac{1}{t} \sup _{N \in \mathcal{S}_{\mathcal{N}}} \log \left\|\Phi_{N}(t, 0)\right\|
$$

Remark 2. Notice that, for every $t>0$ and every $N \in \mathcal{S}_{\mathcal{N}}$, by Gelfand's formula,
$\rho\left(\Phi_{N}(t, 0)\right)=\lim _{k \rightarrow+\infty}\left\|\Phi_{N}(t, 0)^{k}\right\|^{\frac{1}{k}}=\lim _{k \rightarrow+\infty}\left\|\Phi_{\tilde{N}}(k t, 0)\right\|^{\frac{1}{k}}$
where $\tilde{N}$ denotes the $t$-periodic signal obtained by periodization of $\left.N\right|_{[0, t]}$. As a consequence,

$$
\begin{equation*}
\rho\left(\Phi_{N}(t, 0)\right) \leq e^{t \lambda\left(\Sigma_{\mathcal{N}}\right)} \tag{2}
\end{equation*}
$$

Actually, it is well known (see, e.g., [8]) that an equivalent definition of the maximal Lyapunov exponent is

$$
\begin{equation*}
\lambda\left(\Sigma_{\mathcal{N}}\right)=\limsup _{t \rightarrow+\infty} \frac{1}{t} \sup _{N \in \mathcal{S}_{\mathcal{N}}} \log \rho\left(\Phi_{N}(t, 0)\right) \tag{3}
\end{equation*}
$$

Remark 3. It follows from the definitions that $\Sigma_{\mathcal{N}}$ is exponentially stable if and only if $\lambda\left(\Sigma_{\mathcal{N}}\right)<0$. Also, by (3), we have that $\Sigma_{\mathcal{N}}$ is exponentially unstable if and only if $\lambda\left(\Sigma_{\mathcal{N}}\right)>0$. Moreover, the properties of exponential stability/instability and the value of $\lambda\left(\Sigma_{\mathcal{N}}\right)$ keep unchanged if we replace $\mathcal{S}_{\mathcal{N}}$ by the class of piecewise-constant functions from $[0,+\infty)$ to $\mathcal{N}$.

In the following we write $\Phi_{M}^{\varepsilon}(t, 0)$ to denote the flow from time 0 to time $t$ of $\Sigma_{\varepsilon}$ associated with a signal $M \in \mathcal{S}_{\mathcal{M}}$ and we introduce the following stability notions for the 1-parameter family $\Sigma$ of switching systems.

Definition 4. We say that $\Sigma=\left(\Sigma_{e}\right)_{\varepsilon>0}$ is

1. $\varepsilon$-uniformly exponentially stable if there exist $\varepsilon^{\star}>0, C>0$, and $\delta>0$ such that

$$
\begin{equation*}
\left\|\Phi_{M}^{\varepsilon}(t, 0)\right\| \leq C e^{-\delta t}, \forall t \geq 0, \forall M \in \mathcal{S}_{\mathcal{M}}, \forall \varepsilon \in\left(0, \varepsilon^{\star}\right) \tag{4}
\end{equation*}
$$

2. $\varepsilon$-uniformly exponentially unstable if there exist $\delta>0, C>0$, and $\varepsilon^{\star}>0$ such that for every $\varepsilon \in\left(0, \varepsilon^{\star}\right)$ there exist $M \in \mathcal{S}_{\mathcal{M}}$ and $z_{0} \neq 0$ for which

$$
\left|\Phi_{M}^{\varepsilon}(t, 0) z_{0}\right| \geq C e^{\delta t}\left|z_{0}\right|, \quad \forall t \geq 0
$$

Remark 5. The $\varepsilon$-uniform exponential stability of $\Sigma$ is stronger than the property of global uniform asymptotic stability introduced in [1], since it not only guarantees that $\lambda\left(\Sigma_{\varepsilon}\right)<0$ for every $\varepsilon$ small enough, but also that the stability is uniform with respect to $\varepsilon$, i.e., $\lambda\left(\Sigma_{\varepsilon}\right)$ is smaller than the negative number $-\delta$, independent of $\varepsilon$ small, and the constant $C$ appearing in (4) is independent of $\varepsilon$ small. Similar considerations concern the $\varepsilon$-uniform exponential instability.

## 3 Main results

In order to study the $\varepsilon$-uniform exponential stability of the two-time-scales system $\Sigma$, we are going to consider single-time-scale systems whose stability properties give information on the asymptotic behavior of $\Sigma_{\varepsilon}$ for $\varepsilon>0$ small.

Let us first discuss the system $\Sigma_{D}$ corresponding to the evolution of the fast variable $y$ with $x=0$. More precisely, let

$$
\mathcal{M}_{D}=\left\{D \left\lvert\,\left(\begin{array}{cc}
A & B \\
C & D
\end{array}\right) \in \mathcal{M}\right.\right\}
$$

and consider the linear switching system

$$
\Sigma_{D}: \quad \dot{y}(t)=D(t) y(t), \quad D \in \mathcal{S}_{\mathcal{M}_{D}}
$$

We first provide a useful result for the asymptotic behavior of $\Sigma$.

Lemma 6. Assume that $\Sigma_{D}$ is exponentially unstable. Then $\Sigma$ is $\varepsilon$-uniformly exponentially unstable.

Proof. Consider the family $\left(\Sigma_{\mathcal{N}_{\varepsilon}}\right)_{\varepsilon \geq 0}$ of linear switching systems, where, for every $\varepsilon \geq 0$, the compact set of matrices $\mathcal{N}_{\varepsilon} \subset M_{n+m}(\mathbb{R})$ is defined by

$$
\mathcal{N}_{\varepsilon}=\left\{\left.\left(\begin{array}{cc}
\varepsilon A & \varepsilon B \\
C & D
\end{array}\right) \right\rvert\,\left(\begin{array}{ll}
A & B \\
C & D
\end{array}\right) \in \mathcal{M}\right\} .
$$

Then one notices that, for every $\varepsilon>0$, the time rescaling $t \mapsto \varepsilon t$ yields

$$
\phi_{N_{\varepsilon}}(t, 0)=\phi_{M}^{\varepsilon}(\varepsilon t, 0),
$$

where $M$ is an arbitrary signal in $\mathcal{M}$ and the signal $N_{\varepsilon} \in \mathcal{N}_{\varepsilon}$ is defined as

$$
N_{\varepsilon}(t)=\left(\begin{array}{cc}
\varepsilon A(t) & \varepsilon B(t) \\
C(t) & D(t)
\end{array}\right), \quad t \geq 0
$$

In particular,

$$
\lim _{\varepsilon \rightarrow 0} \phi_{M}^{\varepsilon}(\varepsilon t, 0)=\phi_{N_{0}}(t, 0)=\left(\begin{array}{cc}
I_{n} & 0 \\
* & \Phi_{D}(t, 0)
\end{array}\right) .
$$

Since $\Sigma_{D}$ is exponentially unstable, we can assume that the signal $M$ is $T$-periodic for some $T>0$ and such that $\rho\left(\phi_{N_{0}}(T, 0)\right)=\rho\left(\Phi_{D}(T, 0)\right)>1$.

We claim that there exist $\delta, C>0$ such that, for $\varepsilon$ small enough there exists $w_{\varepsilon} \in \mathbb{R}^{n+m} \backslash\{0\}$ with

$$
\begin{equation*}
\left|\phi_{N_{\varepsilon}}(t, 0) w_{\varepsilon}\right| \geq C e^{\delta t}\left|w_{\varepsilon}\right|, \quad \forall t \geq 0 \tag{5}
\end{equation*}
$$

Let $H_{\varepsilon}=\Phi_{N_{\varepsilon}}(T, 0)$ for $\varepsilon \geq 0$. Denote by $\pi$ the projector on a generalized eigenspace corresponding to an eigenvalue of maximal module of $\phi_{N_{0}}(T, 0)$. Since $\lim _{\varepsilon \rightarrow 0} H_{\varepsilon}=H_{0}$, by classical results (see [2, Theorem 5.1, Chapter II]) there exists a sum $\pi_{\varepsilon}$ of projectors on generalized eigenspaces of $H_{\varepsilon}$, satisfying $\lim _{\varepsilon \rightarrow 0} \pi_{\varepsilon}=\pi$, and the corresponding eigenvalues also converge. Let $v_{\varepsilon}$ be a possibly complex eigenvector of the restriction of $H_{\varepsilon}$ to the image of $\pi_{\varepsilon}$, associated with an eigenvalue $\alpha_{\varepsilon}$. If $\alpha_{\varepsilon}$ is real then $v_{\varepsilon}$ can be taken real as well, otherwise we assume without loss of generality that $v_{\varepsilon}$ satisfies $\left|\operatorname{Re}\left(v_{\varepsilon}\right)\right|=\min _{\theta \in \mathbb{R}}\left|\operatorname{Re}\left(e^{i \theta} v_{\varepsilon}\right)\right|$. In particular

$$
\left|\operatorname{Re}\left(\beta v_{\varepsilon}\right)\right| \geq|\beta|\left|\operatorname{Re}\left(v_{\varepsilon}\right)\right|, \quad \forall \beta \in \mathbb{C}
$$

Note that, for every positive integer $k$, one has that $H_{\varepsilon}^{k} v_{\varepsilon}=\alpha_{\varepsilon}^{k} v_{\varepsilon}$ and $H_{\varepsilon}^{k} \bar{v}_{\varepsilon}=\bar{\alpha}_{\varepsilon}^{k} \bar{v}_{\varepsilon}$ which implies that

$$
H_{\varepsilon}^{k} \operatorname{Re}\left(v_{\varepsilon}\right)=\operatorname{Re}\left(\alpha_{\varepsilon}^{k} v_{\varepsilon}\right)
$$

Set

$$
w_{\varepsilon}=\operatorname{Re}\left(v_{\varepsilon}\right)
$$

and consider the trajectory $t \mapsto z_{\varepsilon}(t):=\Phi_{N_{\varepsilon}}(t, 0) w_{\varepsilon}$. For every nonnegative integer $h$ and $\varepsilon$ small enough, we have

$$
\left|z_{\varepsilon}(h T)\right| \geq\left|\alpha_{\varepsilon}\right|^{h}\left|w_{\varepsilon}\right| \geq \bar{\alpha}^{h}\left|w_{\varepsilon}\right|,
$$

where $\bar{\alpha}$ is some fixed real value in the open interval $\left(1, \rho\left(\phi_{N_{0}}(T, 0)\right)\right)$, independent of $\varepsilon$. Setting $\kappa=2 \max _{M \in \mathcal{M}}\|M\|$ it follows that

$$
\left|z_{\varepsilon}(h T+\tau)\right| \geq\left|z_{\varepsilon}(h T)\right| e^{-\kappa \tau}
$$

for every $\tau \in[0, T]$. This completes the proof of Equation (5). We then conclude by noticing that $\left|\phi_{M}^{\varepsilon}(t, 0) w_{\varepsilon}\right| \geq C e^{\delta t / \varepsilon}\left|w_{\varepsilon}\right| \geq C e^{\delta t}\left|w_{\varepsilon}\right|$.

The above result motivates the following assumption.

Assumption 7. The switching system $\Sigma_{D}$ is exponentially stable.

In particular, all matrices in $\mathcal{M}_{D}$ are Hurwitz (and then invertible).

Lut us now introduce an auxiliary switching system, denoted by $\check{\Sigma}$, which could be thought of as the slow dynamics corresponding to signals whose switching occurs at rate exactly $1 / \varepsilon$. In order to define $\check{\Sigma}$, let us associate with every $T>0$ and every $\sigma=\left(\begin{array}{ll}A & B \\ C & B\end{array}\right) \in \mathcal{S}_{\mathcal{M}}$ the matrices

$$
\begin{aligned}
& \Lambda_{0}(T, \sigma)=\int_{0}^{T} \Phi_{D}(T, s) C(s) d s \\
& \Lambda_{1}(T, \sigma)=\int_{0}^{T}\left(A(s)+B(s) \Lambda_{0}(s, \sigma)\right) d s, \\
& \Lambda_{2}(T, \sigma)=\int_{0}^{T} B(s) \Phi_{D}(s, 0) d s,
\end{aligned}
$$

and

$$
\begin{aligned}
& \Lambda(T, \sigma)= \\
& \frac{\Lambda_{1}(T, \sigma)+\Lambda_{2}(T, \sigma)\left(I_{m}-\Phi_{D}(T, 0)\right)^{-1} \Lambda_{0}(T, \sigma)}{T} .
\end{aligned}
$$

The following property holds true.

Lemma 8. Under Assumption 7 the set

$$
\check{\mathcal{M}}:=\left\{\Lambda(T, \sigma) \mid T>0, \sigma \in \mathcal{S}_{\mathcal{M}}\right\}
$$

is bounded.

Proof. By Assumption 7 we have that

$$
\begin{equation*}
\left\|\Phi_{D}\left(t_{2}, t_{1}\right)\right\| \leq c e^{-\alpha\left(t_{2}-t_{1}\right)} \tag{6}
\end{equation*}
$$

for every $t_{2} \geq t_{1}$, for some constants $\alpha>0$ and $c \geq 1$ independent of the switching law $D \in \mathcal{S}_{\mathcal{M}_{D}}$. As a consequence of (6) and the boundedness of $\mathcal{M}$, we have $\left\|\Lambda_{0}(T, \sigma)\right\| \leq C_{1} \min \{1, T\}$ and $\left\|\Lambda_{2}(T, \sigma)\right\| \leq$ $C_{1} \min \{1, T\}$ for some $C_{1}>0$, and $\left\|\Lambda_{1}(T, \sigma)\right\| \leq C_{2} T$ for some $C_{2}>0$. Moreover since $\rho\left(\Phi_{D}(T, 0)\right)<1$ for every $T>0$ and $D \in \mathcal{S}_{\mathcal{M}_{D}}$ according to (2), $\left(I_{m}-\Phi_{D}(T, 0)\right)^{-1}$ is well defined and expandable as an absolutely convergent power series in $\Phi_{D}(T, 0)$. Letting $\bar{T}=\frac{\log (2 c)}{\alpha}$ so that $\left\|\Phi_{D}(\tau, 0)\right\| \leq \frac{1}{2}$ for every $\tau \geq \bar{T}$, we can write

$$
\begin{aligned}
& \left(I_{m}-\Phi_{D}(T, 0)\right)^{-1}=\sum_{k \geq 0} \Phi_{D}(T, 0)^{k} \\
& =\sum_{h=0}^{\left\lfloor\frac{\bar{T}}{T}\right\rfloor} \Phi_{D}(T, 0)^{h}+\Phi_{D}(T, 0)^{\left\lfloor\frac{\bar{T}}{T}\right\rfloor+1}\left(\sum_{h=0}^{\left\lfloor\frac{T}{T}\right\rfloor} \Phi_{D}(T, 0)^{h}\right) \\
& +\Phi_{D}(T, 0)^{2\left(\left\lfloor\frac{\bar{T}}{T}\right\rfloor+1\right)}\left(\sum_{h=0}^{\left\lfloor\frac{\bar{T}}{T}\right\rfloor} \Phi_{D}(T, 0)^{h}\right)+\ldots \\
& =\left(\sum_{k \geq 0} \Phi_{D}(T, 0)^{\left\lfloor\left\lfloor\left\lfloor\frac{\bar{T}}{T}\right\rfloor+1\right) k\right.}\right)\left(\sum_{h=0}^{\left\lfloor\frac{\bar{T}}{T}\right\rfloor} \Phi_{D}(T, 0)^{h}\right) .
\end{aligned}
$$

As $\Phi_{D}(T, 0)^{h}$ corresponds to the flow of $\Sigma_{D}$ for a $T$ periodic signal at time $h T$ we have $\left\|\Phi_{D}(T, 0)^{h}\right\| \leq c$ by (6). Hence

$$
\left\|\sum_{h=0}^{\left\lfloor\frac{\bar{T}}{T}\right\rfloor} \Phi_{D}(T, 0)^{h}\right\| \leq\left(1+\left\lfloor\frac{\bar{T}}{T}\right\rfloor\right) c \leq c+\frac{c \bar{T}}{T} .
$$

Moreover, as $\left(\left\lfloor\frac{\bar{T}}{T}\right\rfloor+1\right) T \geq \bar{T}$,

$$
\begin{aligned}
\left\|\sum_{k \geq 0} \Phi_{D}(T, 0)^{\left(\left\lfloor\frac{\bar{T}}{T}\right\rfloor+1\right) k}\right\| & \leq \sum_{k \geq 0}\left\|\Phi_{D}(T, 0)^{\left(\left\lfloor\frac{\bar{T}}{T}\right\rfloor+1\right)}\right\|^{k} \\
& \leq \sum_{k \geq 0} \frac{1}{2^{k}} \\
& =2
\end{aligned}
$$

Summing up

$$
\begin{aligned}
\|\Lambda(T, \sigma)\| & \leq C_{2}+\frac{C_{1}^{2} \min \{1, T\}^{2}(2 c+2 c \bar{T} / T)}{T} \\
& =C_{2}+2 c C_{1}^{2} \min \left\{\frac{1}{T}+\frac{\bar{T}}{T^{2}}, T+\bar{T}\right\} \\
& \leq C_{2}+2 c C_{1}^{2}(1+\bar{T})
\end{aligned}
$$

concluding the proof of the lemma.
The switching system $\Sigma \check{\Sigma}$ is then defined as

$$
\begin{equation*}
\check{\Sigma}: \quad \dot{\tilde{x}}(t)=M(t) \check{x}(t), \quad M \in \mathcal{S}_{\check{\mathcal{M}}} \tag{7}
\end{equation*}
$$

We can now state the main result of this contribution.

Theorem 9. Suppose that Assumption 7 holds and that system $\dot{\Sigma}$ is exponentially unstable. Then system $\Sigma$ is $\varepsilon$-exponentially unstable.

Therefore, under Assumption 7, a necessary condition for $\Sigma$ to be stable for small values of $\varepsilon$ is that $\check{\Sigma}$ has no trajectory whose norm diverges exponentially.

## 4 Proof of Theorem 9

Before entering in the core of the proof of the theorem, let us prove a useful technical result.

Lemma 10. Let $T>0, \sigma=\left(\begin{array}{ll}A & B \\ C & B\end{array}\right) \in \mathcal{S}_{\mathcal{M}}$, and $\mu \in \mathbb{R}$. For every $\varepsilon>0$ denote by $\mathscr{M}(\varepsilon)$ the flow at time $\varepsilon T$ of $\Sigma_{\varepsilon}$ corresponding to the signal $\sigma(\cdot / \varepsilon)$. Then there exists $P(\varepsilon)$ of the form

$$
P(\varepsilon)=\left(\begin{array}{cc}
I_{n} & 0  \tag{8}\\
Q(\varepsilon) & I_{m}
\end{array}\right)
$$

such that

$$
\begin{align*}
& P(\varepsilon)^{-1} \mathscr{M}(\varepsilon) P(\varepsilon)= \\
& \left(\begin{array}{cc}
I_{n}+\varepsilon T \Lambda(T, \sigma)+O\left(\varepsilon^{2}\right) & O(\varepsilon) \\
0 & \Phi_{D}(T, 0)+O(\varepsilon)
\end{array}\right) \tag{9}
\end{align*}
$$

where

$$
\begin{equation*}
Q(\varepsilon)=\left(I_{m}-\Phi_{D}(T, 0)\right)^{-1} \Lambda_{0}(T, \sigma)+O(\varepsilon) \tag{10}
\end{equation*}
$$

and the functions $O\left(\varepsilon^{k}\right)$ are such that $\left\|O\left(\varepsilon^{k}\right)\right\| \leq C \varepsilon^{k}$ for $\varepsilon$ small, for some positive constant $C$ independent of $\varepsilon$ and $\sigma$.

Proof. We first prove that $\varepsilon \mapsto \mathscr{M}(\varepsilon)$ admits a first order expansion

$$
\begin{equation*}
\mathscr{M}(\varepsilon)=\mathscr{M}_{0}+\varepsilon \mathscr{M}_{1}+O\left(\varepsilon^{2}\right) \tag{11}
\end{equation*}
$$

for some matrices $\mathscr{M}_{0}, \mathscr{M}_{1}$ to be computed. To see that, we first apply the time rescaling $\tau=t / \varepsilon$ and we have that $\mathscr{M}(\varepsilon)$ is equal to $\Phi_{N_{0}+\varepsilon N_{1}}(T, 0)$, where the signals $N_{0}$ and $N_{1}$ are defined as

$$
\begin{aligned}
& N_{0}(\tau)=\left(\begin{array}{cc}
0 & 0 \\
C(\tau) & D(\tau)
\end{array}\right), \\
& N_{1}(\tau)=\left(\begin{array}{cc}
A(\tau) & B(\tau) \\
0 & 0
\end{array}\right)
\end{aligned}
$$

By the variation of constants formula, one has

$$
\begin{aligned}
\Phi_{N_{0}+\varepsilon N_{1}}(T, 0) & =\Phi_{N_{0}}(T, 0) \\
& +\varepsilon \int_{0}^{T} \Phi_{N_{0}}(T, \tau) N_{1}(\tau) \Phi_{N_{0}+\varepsilon N_{1}}(\tau, 0) d \tau
\end{aligned}
$$

One deduces that (11) holds true with

$$
\begin{aligned}
\mathscr{M}_{0} & =\Phi_{N_{0}}(T, 0), \\
\mathscr{M}_{1} & =\int_{0}^{T} \Phi_{N_{0}}(T, \tau) N_{1}(\tau) \Phi_{N_{0}}(\tau, 0) d \tau .
\end{aligned}
$$

It is easy to get that

$$
\begin{align*}
\mathscr{M}_{0} & =\left(\begin{array}{cc}
I_{n} & 0 \\
\Lambda_{0}(T, \sigma) & \Phi_{D}(T, 0)
\end{array}\right)  \tag{12}\\
\mathscr{M}_{1} & =\left(\begin{array}{ll}
\Lambda_{1}(T, \sigma) & \Lambda_{2}(T, \sigma) \\
\Lambda_{3}(T, \sigma) & \Lambda_{4}(T, \sigma)
\end{array}\right) \tag{13}
\end{align*}
$$

where

$$
\begin{aligned}
& \Lambda_{3}(T, \sigma)=\Lambda_{0}(T, \sigma) \Lambda_{1}(T, \sigma) \\
& \quad-\int_{0}^{T} \Phi_{D}(T, \tau) \Lambda_{0}(\tau, \sigma)\left(A(\tau)+B(\tau) \Lambda_{0}(\tau, \sigma)\right) d \tau
\end{aligned}
$$

and

$$
\begin{aligned}
\Lambda_{4}(T, \sigma)= & \Lambda_{0}(T, \sigma) \Lambda_{2}(T, \sigma) \\
& -\int_{0}^{T} \Phi_{D}(T, \tau) \Lambda_{0}(\tau, \sigma) B(\tau) \Phi_{D}(\tau, 0) d \tau
\end{aligned}
$$

According to (11), (12), and (13), it follows that (9) holds true with $P(\varepsilon)$ and $Q(\varepsilon)$ as in (8) and (10). This concludes the proof of the lemma.

We can now provide a proof for Theorem 9.
Proof of Theorem 9. Since $\check{\Sigma}$ is exponentially unstable then, as recalled in Remark 3, $\lambda(\check{\Sigma})>0$. Then, according to (3), there exist $\ell \in \mathbb{N}$, $\ell$ matrices $\Lambda\left(T_{1}, \sigma_{1}\right), \ldots, \Lambda\left(T_{\ell}, \sigma_{\ell}\right) \in \mathcal{M}$, and $\ell$ positive times $t_{1}, \ldots, t_{\ell}$ so that

$$
\begin{equation*}
\rho\left(e^{t_{\ell} \Lambda\left(T_{\ell}, \sigma_{\ell}\right)} \ldots e^{t_{1} \Lambda\left(T_{1}, \sigma_{1}\right)}\right)>1 \tag{14}
\end{equation*}
$$

Let $\varepsilon$ be sufficiently small so that $\varepsilon T_{k}<t_{k}$ for every $k=1, \ldots, \ell$, and denote by $N_{k}=\left\lfloor\frac{t_{k}}{\varepsilon T_{k}}\right\rfloor$ the number of intervals of length $\varepsilon T_{k}$ contained in $\left[0, t_{k}\right]$.

Then, for every $k=1, \ldots, \ell$, consider the flow $\mathscr{M}_{k}(\varepsilon)$ of system $\Sigma_{\varepsilon}$ corresponding to the signal $\sigma_{k}(\cdot / \varepsilon)$, evaluated at time $\varepsilon T_{k}$. Thanks to Lemma 10, there exists $P_{k}(\varepsilon)$ given by $P_{k}(\varepsilon)=\left(\begin{array}{cc}I_{n} & 0 \\ Q_{k}(\varepsilon) & I_{m}\end{array}\right)$ such that

$$
\mathcal{T}_{k}(\varepsilon)=P_{k}(\varepsilon)^{-1} \mathscr{M}_{k}(\varepsilon) P_{k}(\varepsilon)
$$

satisfies

$$
\begin{aligned}
& \mathcal{T}_{k}(\varepsilon)= \\
& \left(\begin{array}{cc}
I_{n}+\varepsilon T_{k}\left(\Lambda\left(T_{k}, \sigma_{k}\right)\right)+O\left(\varepsilon^{2}\right) & O(\varepsilon) \\
0 & \Phi_{D}\left(T_{k}, 0\right)+O(\varepsilon)
\end{array}\right)
\end{aligned}
$$

We repeat $N_{k}$ times $\sigma_{k}(\cdot / \varepsilon)$ to get a signal on $\left[0, \varepsilon T_{k} N_{k}\right]$ and the corresponding flow of $\Sigma_{\varepsilon}$ at time $\varepsilon T_{k} N_{k}$ is given by

$$
\begin{equation*}
\mathscr{M}_{k}(\varepsilon)^{N_{k}}=P_{k}(\varepsilon) \mathcal{T}_{k}(\varepsilon)^{N_{k}} P_{k}(\varepsilon)^{-1} \tag{15}
\end{equation*}
$$

We claim that

$$
\mathcal{T}_{k}(\varepsilon)^{N_{k}}=\left(\begin{array}{cc}
e^{t_{k} \Lambda\left(T_{k}, \sigma_{k}\right)}+O(\varepsilon) & O(\varepsilon)  \tag{16}\\
0 & O(\varepsilon)
\end{array}\right)
$$

This follows from the general formula

$$
\left(\begin{array}{cc}
A_{11} & A_{12} \\
0 & A_{22}
\end{array}\right)^{N}=\left(\begin{array}{cc}
A_{11}^{N} & W \\
0 & A_{22}^{N}
\end{array}\right)
$$

with

$$
W=\sum_{j=0}^{N-1} A_{11}^{j} A_{12} A_{22}^{N-j-1}
$$

applied to $N=N_{k}$,

$$
\begin{aligned}
& A_{11}=I+B_{11} / N_{k}+O\left(\varepsilon^{2}\right)=e^{\frac{B_{11}}{N_{k}}}+O\left(\varepsilon^{2}\right) \\
& A_{22}=\Phi_{D}\left(T_{k}, 0\right)+O(\varepsilon)
\end{aligned}
$$

where $B_{11}=t_{k} \Lambda\left(T_{k}, \sigma_{k}\right)$. Then, for $\varepsilon$ small enough and $1 \leq j \leq N_{k}$ one gets that

$$
\begin{align*}
\left\|A_{11}^{j}\right\| & \leq\left(1+\frac{2\left\|B_{11}\right\|}{j}\right)^{j} \leq e^{2\left\|B_{11}\right\|}  \tag{17}\\
A_{11}^{N_{k}} & =e^{B_{11}}+O(\varepsilon)
\end{align*}
$$

Moreover, since $\rho\left(A_{22}\right)<1$ for $\varepsilon$ small, it follows that

$$
\begin{equation*}
\left\|A_{22}^{j}\right\| \leq K \lambda^{j}, \quad 1 \leq j \leq N_{k} \tag{18}
\end{equation*}
$$

for some $K>0$ and $\lambda \in(0,1)$ independent of $\varepsilon$ small enough and $k \in\{1, \ldots \ell\}$. In particular,

$$
\left\|A_{22}^{N_{k}}\right\| \leq K \lambda^{\frac{t_{k}}{\varepsilon T_{k}}-1}=O(\varepsilon)
$$

Using now (17), (18), and recalling that $K$ can be taken so that $\left\|A_{12}\right\| \leq K \varepsilon$, one deduces that

$$
\begin{aligned}
\|W\| & \leq \sum_{j=0}^{N_{k}-1}\left\|A_{11}^{j}\right\|\left\|A_{12}\right\|\left\|A_{22}^{N_{k}-j-1}\right\| \\
& \leq K^{2} e^{2\left\|B_{11}\right\|} \varepsilon \sum_{j=0}^{N_{k}-1} \lambda^{N_{k}-1-j} \\
& =C \varepsilon
\end{aligned}
$$

for some $C>0$ independent of $\varepsilon$ small enough and $k \in\{1, \ldots, \ell\}$. This concludes the proof of (16).

We next use (10) and (16) in (15) to deduce that

$$
\mathscr{M}_{k}(\varepsilon)^{N_{k}}=\left(\begin{array}{cc}
e^{t_{k} \Lambda\left(T_{k}, \sigma_{k}\right)} & 0  \tag{19}\\
r_{k} & 0
\end{array}\right)+O(\varepsilon)
$$

where

$$
r_{k}=\left(I_{m}-\Phi_{D_{\sigma_{k}}}\left(T_{k}, 0\right)\right)^{-1} \Lambda_{0}\left(T_{k}, \sigma_{k}\right) e^{t_{k} \Lambda\left(T_{k}, \sigma_{k}\right)}
$$

Set $t_{\varepsilon}=\varepsilon \sum_{k=1}^{\ell} N_{k} T_{k}$ and notice that $t_{\varepsilon}$ tends to $t_{1}+\cdots+t_{\ell}$ as $\varepsilon$ tends to zero. We concatenate the $N_{k}$ times repetitions of $\sigma_{k}(\cdot / \varepsilon)$ for $k=1, \ldots, \ell$ to get a signal on $\left[0, t_{\varepsilon}\right]$ and the corresponding flow of $\Sigma_{\varepsilon}$ at time $t_{\varepsilon}$ is given by the matrix product

$$
\Upsilon_{\varepsilon}=\mathscr{M}_{\ell}(\varepsilon)^{N_{\ell}} \mathscr{M}_{\ell-1}(\varepsilon)^{N_{\ell-1}} \cdots \mathscr{M}_{1}(\varepsilon)^{N_{1}} .
$$

Using (19), one deduces that

$$
\Upsilon_{\varepsilon}=\left(\begin{array}{cc}
e^{t_{\ell} \Lambda\left(T_{\ell}, \sigma_{\ell}\right)} \cdots e^{t_{1} \Lambda\left(T_{1}, \sigma_{1}\right)} & 0 \\
r & 0
\end{array}\right)+O(\varepsilon)
$$

where the matrix $r$ does not depend on $\varepsilon$.
Using (14), one gets that $\rho\left(\Upsilon_{\varepsilon}\right)>1$ for $\varepsilon$ small enough, yielding that $\lambda\left(\Sigma_{\varepsilon}\right)>0$. In particular $\liminf _{\varepsilon \rightarrow 0^{+}} \lambda\left(\Sigma_{\varepsilon}\right) \geq 0$.

In order to show that $\Sigma$ is $\varepsilon$-exponentially unstable, observe that the flow $\Upsilon_{\varepsilon, t}$ of $\Sigma_{\varepsilon}$ at time $t \in\left[0, t_{\varepsilon}\right]$ corresponding to the signal defined above satisfies

$$
\begin{align*}
& \Upsilon_{\varepsilon, t}= \\
& \left(\begin{array}{cc}
e^{\left(t-\sum_{h=1}^{k-1} t_{h}\right) \Lambda\left(T_{k}, \sigma_{k}\right)} \prod_{\ell=1}^{k-1} e^{t_{\ell} \Lambda\left(T_{\ell}, \sigma_{\ell}\right)}+O(\varepsilon) & O(\varepsilon) \\
r(\varepsilon, t) & q(\varepsilon, t)
\end{array}\right) \tag{20}
\end{align*}
$$

whenever $t \in\left[\sum_{h=1}^{k-1} t_{h}, \sum_{h=1}^{k} t_{h}\right]$, for some matrix functions $r, q$, where the terms $O(\varepsilon)$ are uniform with respect to $t \in\left[0, t_{\varepsilon}\right]$. The matrix $\Upsilon_{\varepsilon, t_{\varepsilon}}$ converges, as $\varepsilon$ goes to zero, to

$$
\bar{\Upsilon}=\left(\begin{array}{cc}
e^{t_{\ell} \Lambda\left(T_{\ell}, \sigma_{\ell}\right)} \cdots e^{t_{1} \Lambda\left(T_{1}, \sigma_{1}\right)} & 0 \\
\bar{r} & 0
\end{array}\right)
$$

for some matrix $\bar{r}$. For $\varepsilon$ small enough we construct a trajectory $z_{\varepsilon}(t)=\left(x_{\varepsilon}(t), y_{\varepsilon}(t)\right)$ of $\Sigma_{\varepsilon}$ satisfying $\left|z_{\varepsilon}(t)\right| \geq \hat{C} e^{\hat{\lambda} t}\left|z_{\varepsilon}(0)\right|>0$ for every $t \geq 0$, with $\hat{\lambda} \in(0, \lambda(\check{\Sigma}))$ and $\hat{C}>0$ independent of $\varepsilon$.

Let $\pi$ be the sum of the projectors on the generalized eigenspaces associated with the eigenvalues of $\bar{\Upsilon}$ of modulus $\rho\left(e^{t_{\ell} \Lambda\left(T_{\ell}, \sigma_{\ell}\right)} \cdots e^{t_{1} \Lambda\left(T_{1}, \sigma_{1}\right)}\right)>1$. Since $\Upsilon_{\varepsilon, t_{\varepsilon}}$ converges to $\bar{\Upsilon}$ as $\varepsilon$ goes to zero, by classical results (see [2, Theorem 5.1, Chapter II]) there exists $\pi_{\varepsilon}$, a sum of projectors on generalized eigenspaces of $\Upsilon_{\varepsilon, t_{\varepsilon}}$, satisfying $\lim _{\varepsilon \rightarrow 0} \pi_{\varepsilon}=\pi$, and the corresponding eigenvalues also converge. Let $v_{\varepsilon}$ be a possibly complex eigenvector of the restriction of $\Upsilon_{\varepsilon, t_{\varepsilon}}$ to the image of $\pi_{\varepsilon}$, associated with an eigenvalue $\alpha_{\varepsilon}$. If $\alpha_{\varepsilon}$ is real then $v_{\varepsilon}$ can be taken real as well, otherwise we assume without loss of generality that $v_{\varepsilon}$ satisfies $\left|\operatorname{Re}\left(v_{\varepsilon}\right)\right|=\min _{\theta \in \mathbb{R}}\left|\operatorname{Re}\left(e^{i \theta} v_{\varepsilon}\right)\right|$. In particular $\left|\operatorname{Re}\left(\beta v_{\varepsilon}\right)\right| \geq|\beta|\left|\operatorname{Re}\left(v_{\varepsilon}\right)\right|$ for every $\beta \in \mathbb{C}$. Note that, for every positive integer $k$, one has that $\left(\Upsilon_{\varepsilon, t_{\varepsilon}}\right)^{k} v_{\varepsilon}=\alpha_{\varepsilon}^{k} v_{\varepsilon}$ and $\left(\Upsilon_{\varepsilon, t_{\varepsilon}}\right)^{k} \bar{v}_{\varepsilon}=\bar{\alpha}_{\varepsilon}^{k} \bar{v}_{\varepsilon}$ which implies $\left(\Upsilon_{\varepsilon, t_{\varepsilon}}\right)^{k} \operatorname{Re}\left(v_{\varepsilon}\right)=\operatorname{Re}\left(\alpha_{\varepsilon}^{k} v_{\varepsilon}\right)$.

Consider the trajectory $z_{\varepsilon}(t)=\left(x_{\varepsilon}(t), y_{\varepsilon}(t)\right)$ of $\Sigma_{\varepsilon}$ obtained applying the flow $\Upsilon_{\varepsilon, t}$ to the initial condition $z_{\varepsilon}(0)=\operatorname{Re}\left(v_{\varepsilon}\right)$ and repeating periodically after time $t_{\varepsilon}$. Letting $\Pi_{x}$ be the projection of a vector of $\mathbb{R}^{n+m}$ onto its first $n$ components, it is easy to see that $\left|\Pi_{x} v\right| \geq C|v|$ for every $v$ in the image of $\bar{\Upsilon}$, where

$$
C=\left(1+\|\bar{r}\|^{2}\left\|\left(e^{t_{\ell} \Lambda\left(T_{\ell}, \sigma_{\ell}\right)} \cdots e^{t_{1} \Lambda\left(T_{1}, \sigma_{1}\right)}\right)^{-1}\right\|^{2}\right)^{-1 / 2}
$$

Hence, for every nonnegative integer $h$ and $\varepsilon$ small enough,

$$
\begin{align*}
\left|x_{\varepsilon}\left(h t_{\varepsilon}\right)\right| & =\left|\Pi_{x} z_{\varepsilon}\left(h t_{\varepsilon}\right)\right| \\
& =\left|\Pi_{x} \pi_{\varepsilon} z_{\varepsilon}\left(h t_{\varepsilon}\right)\right| \\
& \geq\left|\Pi_{x} \pi z_{\varepsilon}\left(h t_{\varepsilon}\right)\right|-\left\|\Pi_{x}\right\|\left\|\pi-\pi_{\varepsilon}\right\|\left|z_{\varepsilon}\left(h t_{\varepsilon}\right)\right| \\
& \geq\left|\Pi_{x} \pi z_{\varepsilon}\left(h t_{\varepsilon}\right)\right|-\left\|\pi-\pi_{\varepsilon}\right\|\left|z_{\varepsilon}\left(h t_{\varepsilon}\right)\right| \\
& \geq C\left|\pi z_{\varepsilon}\left(h t_{\varepsilon}\right)\right|-\left\|\pi-\pi_{\varepsilon}\right\|\left|z_{\varepsilon}\left(h t_{\varepsilon}\right)\right| \\
& \geq C\left|z_{\varepsilon}\left(h t_{\varepsilon}\right)\right|-(1+C)\left\|\pi-\pi_{\varepsilon}\right\|\left|z_{\varepsilon}\left(h t_{\varepsilon}\right)\right|  \tag{21}\\
& \geq \frac{C}{2}\left|z_{\varepsilon}\left(h t_{\varepsilon}\right)\right| .
\end{align*}
$$

As $z_{\varepsilon}\left(h t_{\varepsilon}\right)=\operatorname{Re}\left(\alpha_{\varepsilon}^{h} v_{\varepsilon}\right)$ we also have

$$
\begin{equation*}
\left|x_{\varepsilon}\left(h t_{\varepsilon}\right)\right| \geq \frac{C}{2}\left|\alpha_{\varepsilon}\right|^{h}\left|z_{\varepsilon}(0)\right| . \tag{22}
\end{equation*}
$$

By (20), (21), and setting $\kappa=\max _{\Lambda \in \mathscr{\mathcal { M }}}\|\Lambda\|$ it follows
that

$$
\begin{align*}
\left|x_{\varepsilon}\left(h t_{\varepsilon}+\tau\right)\right| & \geq\left|x_{\varepsilon}\left(h t_{\varepsilon}\right)\right| e^{-\kappa \tau}-\left|z_{\varepsilon}\left(h t_{\varepsilon}\right)\right| O(\varepsilon) \\
& \geq\left|x_{\varepsilon}\left(h t_{\varepsilon}\right)\right| e^{-\kappa \tau}-\frac{2}{C}\left|x_{\varepsilon}\left(h t_{\varepsilon}\right)\right| O(\varepsilon) \\
& \geq \frac{1}{2}\left|x_{\varepsilon}\left(h t_{\varepsilon}\right)\right| e^{-\kappa \tau} \tag{23}
\end{align*}
$$

for every $\tau \in\left[0, t_{\varepsilon}\right]$. Take $\varepsilon$ small enough in such a way that $1<\underline{\alpha} \leq\left|\alpha_{\varepsilon}\right| \leq \bar{\alpha}$ for some positive constants $\underline{\alpha}, \bar{\alpha}$, and $t_{\varepsilon} \leq \bar{t}$, where $\bar{t}=2 \sum_{h=1}^{\ell} t_{h}$. By (22) and (23) we thus get for every $t \geq 0$

$$
\begin{aligned}
\left|z_{\varepsilon}(t)\right| & \geq\left|x_{\varepsilon}(t)\right| \geq \frac{C}{4} e^{-\kappa t_{\varepsilon}}\left|\alpha_{\varepsilon}\right|^{\left\lfloor\frac{t}{t_{\varepsilon}}\right\rfloor}\left|z_{\varepsilon}(0)\right| \\
& \geq \frac{C e^{-\kappa t_{\varepsilon}}}{4\left|\alpha_{\varepsilon}\right|} e^{\frac{\log \left|\alpha_{\varepsilon}\right|}{t_{\varepsilon}} t}\left|z_{\varepsilon}(0)\right| \geq \hat{C} e^{\hat{\lambda} t}\left|z_{\varepsilon}(0)\right|,
\end{aligned}
$$

where $\hat{C}=\frac{C e^{-\kappa \bar{t}}}{4 \bar{\alpha}}$ and $\hat{\lambda}=\frac{\log \underline{\underline{\alpha}}}{t}$, which shows that $\Sigma$ is $\varepsilon$-exponentially unstable.

Remark 11. The standard approach to singular perturbations in the non-switching case consists in identifying separated slow and fast variables where different controllers for different time-scale variables can be designed in order to lead the overall system to its desired performance (see, e.g., [4]). We could apply formally such an approach in our case by saying that $y$ tends instantaneously to $-D(t)^{-1} C(t) x(t)$ (the equilibrium of the equation for $y$ when $x \equiv x(t)$ ) and obtaining the slow-variable switching system

$$
\begin{equation*}
\dot{x}(t)=\left(A(t)-B(t) D(t)^{-1} C(t)\right) x(t) \tag{24}
\end{equation*}
$$

However, as it has already been observed in the literature, the switching system (24) may be exponentially stable even when, for every $\varepsilon>0, \Sigma_{\varepsilon}$ is unstable [5]. The modes of (24) are given by matrices in the set

$$
\overline{\mathcal{M}}:=\left\{A-B D^{-1} C \left\lvert\,\left(\begin{array}{ll}
A & B \\
D
\end{array}\right) \in \mathcal{M}\right.\right\} \subset M_{n}(\mathbb{R})
$$

It is not difficult to check that $\overline{\mathcal{M}} \subset \mathcal{M}$. Indeed, letting $T>0$ and $\bar{\sigma} \in \mathcal{S}_{\mathcal{M}}$ constantly equal to $M$ on $[0, T]$, it holds

$$
\Lambda(T, \bar{\sigma})=A-B D^{-1} C
$$

Example 12. Let us show by an example that $\mathcal{M}$ gives sharper necessary conditions for stability than the set $\overline{\mathcal{M}}$ introduced in the previous remark. Consider system $\Sigma$ with

$$
\mathcal{M}=\left\{M_{1}=\left(\begin{array}{cc}
-1 & 1 \\
0 & -0.1
\end{array}\right), M_{2}=\left(\begin{array}{cc}
-3 & 0 \\
2 & -0.1
\end{array}\right)\right\} .
$$

The stability of singularly perturbed planar switching systems is completely characterized in [1, Theorem 2] through some necessary and sufficient conditions. Based on this characterization (cf., in particular, Item (SP5) in [1, Theorem 2]) the condition
$\frac{1}{2}\left(\operatorname{tr}\left(M_{1}\right) \operatorname{tr}\left(M_{2}\right)-\operatorname{tr}\left(M_{1} M_{2}\right)\right)<-\sqrt{\operatorname{det}\left(M_{1}\right) \operatorname{det}\left(M_{2}\right)}$
implies that $\Sigma_{\varepsilon}$ is exponentially unstable for all $\varepsilon>0$. Condition (25) is satisfied in the case of this example since

$$
\frac{1}{2}\left(\operatorname{tr}\left(M_{1}\right) \operatorname{tr}\left(M_{2}\right)-\operatorname{tr}\left(M_{1} M_{2}\right)\right)=-0.8
$$

and

$$
\operatorname{det}\left(M_{1} M_{2}\right)=0.03
$$

Look now at systems $\bar{\Sigma}$ and $\check{\Sigma}$. We have $\overline{\mathcal{M}}=$ $\{-1,-3\}$ and then system (24) is exponentially stable. Concerning system $\check{\Sigma}$, let us consider the switching signal

$$
\sigma(t)=\alpha(t) M_{1}+(1-\alpha(t)) M_{2}
$$

associated with the 2-periodic function

$$
\alpha(t)= \begin{cases}1 & t \in[0,1] \\ 0 & t \in[1,2]\end{cases}
$$

and take $T=2$. For this choice of $\sigma$ and $T$ one can easily verify that

$$
\Lambda(T, \sigma)=-2+100\left(1-e^{-0.2}\right)^{-1}\left(1-e^{-0.1}\right)^{2}>0
$$

Then $\check{\Sigma}$ is exponentially unstable, as illustrated in Figure 1.


Figure 1: Fast and slow variables evolution of system $\Sigma_{\varepsilon}$ with signal $\sigma$ and $\varepsilon=0.1$, starting from $\left(x_{0}, y_{0}\right)=$ (1, 1).

## 5 Conclusions

The main goal of this contribution is to identify linear modes (the elements of $\check{\mathcal{M}}$ ) that can be approximated arbitrarily well, in the slow-varying variables $x$, by a singularly perturbed switching linear system for sufficiently small values of the singular perturbation parameter. The corresponding signals are periodic with a period of order $\varepsilon$. In this way, we are able to single out a switched system with a single time-scale (the system $\check{\Sigma}$ ), whose stability is a necessary condition for the stability of the original singularly perturbed switching system for small values of $\varepsilon$.

## References

[1] F. El Hachemi, M. Sigalotti, and J. Daafouz. Stability analysis of singularly perturbed switched linear systems. IEEE Trans. Automat. Control, 57(8):2116-2121, 2012.
[2] T. Kato. Perturbation theory for linear operators. Die Grundlehren der mathematischen Wissenschaften, Band 132. Springer-Verlag New York, Inc., New York, 1966.
[3] J. W. Kimball and P. T. Krein. Singular perturbation theory for DC-DC converters and application to PFC converters. In 2007 IEEE Power Electronics Specialists Conference, pages 882-887, 2007.
[4] P. Kokotović, H. K. Khalil, and J. O'Reilly. Singular perturbation methods in control, volume 25 of Classics in Applied Mathematics. Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 1999. Analysis and design, Corrected reprint of the 1986 original.
[5] I. Malloci, J. Daafouz, and C. Iung. Stabilization of continuous-time singularly perturbed switched systems. In Proceedings of the $48 h$ IEEE Conference on Decision and Control (CDC) held jointly with 2009 28th Chinese Control Conference, pages 6371-6376, 2009.
[6] I. Malloci, J. Daafouz, C. Iung, R. Bonidal, and P. Szczepanski. Switched system modeling and robust steering control of the tail end phase in a hot strip mill. Nonlinear Anal. Hybrid Syst., 3(3):239-250, 2009.
[7] J. Wang, Z. Huang, Z. Wu, J. Cao, and H. Shen. Extended dissipative control for singularly perturbed pdt switched systems and its application. IEEE Transactions on Circuits and Systems I: Regular Papers, 67(12):5281-5289, 2020.
[8] F. Wirth. The generalized spectral radius and extremal norms. Linear Algebra and its Applications, 342(1):17-40, 2002.


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