

Improved coupling formulations and time integration strategies for ocean-atmosphere interactions

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Context: ocean - atmosphere coupling



Representing their complex interactions mostly means

- estimating OA fluxes
- imposing these fluxes through the OA interface
- → Numerous physical / mathematical / numerical / algorithmic issues



Context: ocean - atmosphere coupling



climate modeling



seasonal forecasts



short term predictions

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Some mathematical and numerical inputs

- Modeling aspects: provide a general framework for the formulation of the coupled OA problem
- Algorithmic aspects: revisit the coupling strategies presently used in such systems



Air-sea interactions: 3 different domains

with
$$\boldsymbol{U}^{\mathsf{a}} = \begin{pmatrix} \boldsymbol{u}_{h}^{\mathsf{a}} \\ T^{\mathsf{a}} \end{pmatrix} \quad \boldsymbol{U}^{\mathsf{o}} = \begin{pmatrix} \boldsymbol{u}_{h}^{\mathsf{o}} \\ T^{\mathsf{o}} \end{pmatrix}$$





E. Blayo et al - Coupled problems 2021, online, 14-16 June 2021

Grenoble Alnes

Air-sea interactions: 3 different domains Exchange conditions:

$$\begin{array}{lll} \text{momentum} & \rho^{a} \mathcal{K}_{m}^{a} \, \frac{\partial \boldsymbol{u}_{h}^{a}}{\partial z}(\boldsymbol{z}_{a}^{1},t) & = & \rho^{o} \mathcal{K}_{m}^{o} \, \frac{\partial \boldsymbol{u}_{h}^{o}}{\partial z}(\boldsymbol{z}_{o}^{1},t) & = & \boldsymbol{\tau}(t) \\ \text{heat flux} & \rho^{a} \boldsymbol{c}_{p}^{a} \mathcal{K}_{T}^{a} \, \frac{\partial T^{a}}{\partial z}(\boldsymbol{z}_{a}^{1},t) & = & \rho^{o} \boldsymbol{c}_{p}^{o} \mathcal{K}_{T}^{o} \, \frac{\partial T^{o}}{\partial z}(\boldsymbol{z}_{o}^{1},t) & = & \boldsymbol{Q}_{S}(t) + \mathcal{R}(t) \end{array}$$

where $\tau(t)$ and $Q_{S}(t)$ are obtained via a bulk closure.



Monin-Obukhov similarity theory (1954): a generalized law of the wall for stratified fluids

- Key parameters: turbulent scales (u_a^*, θ_a^*)
- Ingredients: constant flux hypothesis, stability functions, roughness lengths...

with

 Φ_m, Φ_h

•
$$L_O^a = \frac{(u_a^*)^2}{g\kappa} \frac{\theta(z_a^1)}{\theta_a^*}$$
 Obukhov length

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empirical stability functions

Usual practice in coupled OA models: constant profiles in AVSL and OSL

$$\frac{\partial u}{\partial z} = \frac{\partial \theta}{\partial z} = 0 \quad z \in (z_o^1, z_x^{a,r})$$



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By integration: nonlinear system of equations for (u_a^*, θ_a^*)

$$\begin{bmatrix} u \end{bmatrix}_{z_0^1}^{z_1^1} = \frac{u_a^*}{\kappa} \left[\ln\left(\frac{z_a^1}{z_u^{a,r}(u_a^*,\theta_a^*)}\right) - \psi_m\left(\frac{z_a^1}{L_O^a(u_a^*,\theta_a^*)}\right) \right]$$

$$\begin{bmatrix} \theta \end{bmatrix}_{z_0^1}^{z_1^1} = \frac{\theta_a^*}{\kappa} \left[\ln\left(\frac{z_a^1}{z_\theta^{a,r}(u_a^*,\theta_a^*)}\right) - \psi_h\left(\frac{z_a^1}{L_O^a(u_a^*,\theta_a^*)}\right) \right]$$

- ► These so-called "bulk formulas" are solved by fixed-point iterations.
- Hence the fluxes: $\tau = \rho_a (u_a^*)^2$ $Q_S = \rho_a c_p^a u_a^* \theta_a^*$
- Several such formulas have been developed (e.g. Fairall et al., 2002; Large, 2006; ECMWF, 2016), corresponding to different parameterizations for $z^{a,r}$, (ψ_m, ψ_h) , and the addition of other physical effects (gustiness, warm layers...). Calibration is done w.r.t. field experiment data.





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Classical bulk closures have been developed in the framework of atmosphere-only simulations:

- No consideration of the ocean surface layer
- The flux depends on the ocean model vertical discretization



Goal: building a framework more adapted to the context of coupled air-sea simulations (symmetry of the problem, regularity of the solution profile...)







- Introduction of oceanic turbulent scales (u_o^*, θ_o^*)
- Symmetric point of view between ocean and atmosphere
- Mathematical regularity of the vertical profiles





- ▶ Introduction of oceanic turbulent scales (u_o^*, θ_o^*)
- Symmetric point of view between ocean and atmosphere
- Mathematical regularity of the vertical profiles

$$\begin{split} \llbracket u \rrbracket_{z_{o}^{1}}^{z_{o}^{1}} &= \frac{u_{a}^{*}}{\kappa} \left[\ln \left(\frac{z_{a}^{1}}{z_{u}^{a,r} \left(u_{a}^{*}, \theta_{a}^{*} \right)} \right) - \psi_{m}^{a} \left(\frac{z_{a}^{1}}{L_{O}^{o} \left(u_{a}^{*}, \theta_{a}^{*} \right)} \right) \right. \\ &+ \lambda \ln \left(\frac{z_{o}^{1}}{z_{u}^{o,r} \left(u_{o}^{*}, \theta_{o}^{*} \right)} \right) - \lambda \psi_{m}^{o} \left(\frac{-z_{o}^{1}}{L_{O}^{o} \left(u_{o}^{*}, \theta_{o}^{*} \right)} \right) \right] \\ \llbracket \theta \rrbracket_{z_{o}^{1}}^{z_{o}^{1}} &= \frac{\theta_{a}^{*}}{\kappa} \left[\ln \left(\frac{z_{a}^{1}}{z_{\theta}^{a,r} \left(u_{a}^{*}, \theta_{a}^{*} \right)} \right) - \psi_{h}^{a} \left(\frac{z_{a}^{1}}{L_{O}^{o} \left(u_{a}^{*}, \theta_{a}^{*} \right)} \right) \right. \\ &+ \lambda \ln \left(\frac{-z_{o}^{1}}{z_{u}^{o,r} \left(u_{o}^{*}, \theta_{o}^{*} \right)} \right) - \lambda \psi_{h}^{o} \left(\frac{-z_{o}^{1}}{L_{O}^{o} \left(u_{o}^{*}, \theta_{o}^{*} \right)} \right) \right] \end{split}$$



Impact on turbulent fluxes

- Systematic decrease of the wind stress, slight increase (on average) of the heat flux
- Indicates that the ocean part is significant. But must not be understood as an error: need for recalibration in the context of two-sided closures



- ► A generic modeling framework, with a clear view of the hypotheses, and avoiding non physical properties of the vertical profiles
- Allows for an easy introduction of new parameterizations and/or additional effects: waves, radiative flux...

More details in:

Pelletier C., F. Lemarié, E. Blayo, M.-N. Bouin and J.-L. Redelsperger, 2021: Two-sided turbulent surface layer parameterizations for computing air-sea fluxes. *Quart. J. Roy. Meteorol. Soc.* http://dx.doi.org/10.1002/qj.3991



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- Modeling aspects: provide a general framework for the representation of the coupled surface layers
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Algorithmic issues

One has to solve:

 $\left\{ \begin{array}{ll} \mathcal{L}_{atm} \ \boldsymbol{\textit{U}}^{a} = f_{atm} & \text{ in } \Omega_{atm} \times [0, T] \\ \mathcal{L}_{oce} \ \boldsymbol{\textit{U}}^{o} = f_{oce} & \text{ in } \Omega_{oce} \times [0, T] \\ \mathcal{F}_{atm} \ \boldsymbol{\textit{U}}^{a} = \mathcal{F}_{oce} \ \boldsymbol{\textit{U}}^{o} = F_{OA}(\boldsymbol{\textit{U}}^{a}, \boldsymbol{\textit{U}}^{o}, \mathcal{R}) & \text{ on } \Gamma \times [0, T] \end{array} \right.$





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Two usual approaches:

- Asynchronous coupling by time windows (averaged fluxes)
- Synchronous coupling at the time step (instantaneous fluxes)



Asynchronous coupling

▶ Split the simulation into time windows $[0, T] = \bigcup_{i=1}^{M} [t_i, t_{i+1}]$

Example: atmosphere-first sequential algorithm



- ▶ same mean fluxes over each time window [t_i, t_{i+1}]
- synchrony issue: time lag between model(s) and interface cond.

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Synchronous coupling at the time step



- much smaller synchrony issue
- difficult to implement efficiently
- physical validity issue (w.r.t. Monin-Obukhov hypotheses)
- stability issues with the coupling algorithm (Lemarié et al., 2015; Beljaars et al., 2017; Zhang et al., 2020)

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Algorithmic issues

- Current coupling methods are ad-hoc algorithms which ensure that fluxes are balanced and are computationally cheap.
- These methods raise a number of issues.

Questions

- Can the coupling algorithm be improved ?
- Does it improve the physics of the coupled solution ?
- Can this be done at a reasonable CPU cost ?

















 Present ocean-atmosphere coupling algorithms correspond to one single iteration of a Schwarz-like coupling method



Toward iterative ocean-atmosphere coupling

 Present ocean-atmosphere coupling algorithms correspond to one single iteration of a Schwarz-like coupling method

$$\begin{cases} \mathcal{L}_{\mathsf{atm}} \boldsymbol{U}^{\mathsf{a}} &= f_{\mathsf{atm}} & \text{ in } \Omega_{\mathsf{atm}} \times [t_i, t_{i+1}] \\ \mathcal{F}_{\mathsf{atm}} \boldsymbol{U}^{\mathsf{a}} &= F_{\mathsf{OA}}(\boldsymbol{U}^{\mathsf{o}}, \boldsymbol{U}^{\mathsf{a}}, \mathcal{R}) & \text{ on } \Gamma \times [t_i, t_{i+1}] \end{cases}$$

$$\begin{cases} \mathcal{L}_{\text{oce}} \boldsymbol{U}^{\text{o}} &= f_{\text{oce}} & \text{in } \Omega_{\text{oce}} \times [t_i, t_{i+1}] \\ \mathcal{F}_{\text{oce}} \boldsymbol{U}^{\text{o}} &= \mathcal{F}_{\text{atm}} \boldsymbol{U}^{\text{a}} & \text{on } \Gamma \times [t_i, t_{i+1}] \end{cases}$$



Toward iterative ocean-atmosphere coupling

 Present ocean-atmosphere coupling algorithms correspond to one single iteration of a Schwarz-like coupling method

Iterate until convergence

$$\begin{cases} \mathcal{L}_{\mathsf{atm}} \boldsymbol{U}_{k+1}^{\mathsf{a}} = f_{\mathsf{atm}} & \text{in } \Omega_{\mathsf{atm}} \times [t_i, t_{i+1}] \\ \mathcal{F}_{\mathsf{atm}} \boldsymbol{U}_{k+1}^{\mathsf{a}} = F_{\mathsf{OA}}(\boldsymbol{U}_k^{\mathsf{o}}, \boldsymbol{U}_{k+1}^{\mathsf{a}}, \mathcal{R}_{k+1}) & \text{on } \Gamma \times [t_i, t_{i+1}] \\ \end{cases}$$
then

$$\begin{cases} \mathcal{L}_{\text{oce}} \boldsymbol{U}_{k+1}^{\text{o}} &= f_{\text{oce}} & \text{in } \Omega_{\text{oce}} \times [t_i, t_{i+1}] \\ \mathcal{F}_{\text{oce}} \boldsymbol{U}_{k+1}^{\text{o}} &= \mathcal{F}_{\text{atm}} \boldsymbol{U}_{k+1}^{\text{a}} & \text{on } \Gamma \times [t_i, t_{i+1}] \end{cases}$$

 Should lead (if convergence) to the exact solution of the coupled problem



Does it improve the physics of the coupled solution ?

A major difficulty: no idealized coupled ocean-atmosphere testcase with a known reference solution.

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Lemarié et al., 2014: ensemble simulations of a tropical cyclone with a coupled ROMS-WRF system.



30%-40% decrease of the uncertainty on trajectory and intensity





Does it improve the physics of the coupled solution ?

24h simulations with the 1-D version of the climate model of Météo-France (S. Valcke and A. Voldoire)



Resynchronization of the diurnal cycle

See also next talk by Olivier Marti with the IPSL Climate Model.

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Analysis of the iterative process







Analysis of the iterative process



Systems satisfied by the errors $e_i^k = u_i^k - u_i$:

ſ	$L_1 e_1^{k+1}$	= 0	$\Omega_1 imes [0, T]$	$\int L_2 e_2^{k+1}$	= 0	$\Omega_2 imes [0, T]$
J	e_1^{k+1}	= 0	at $t = 0$	e_2^{k+1}	= 0	at $t = 0$
)	$B_1e_1^{\overline{k}+1}$	= 0	$\partial \Omega_1^{ext} imes [0, T]$	$B_2 e_2^{\overline{k}+1}$	= 0	$\partial \Omega_2^{\text{ext}} \times [0, T]$
l	Ce_1^{k+1}	$= C' e_2^k$	$\Gamma \times [0, T]$	be_2^{k+1}	$= D' e_1^k$	$\Gamma \times [0, T]$

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convergence factor

 $\rho_k = \frac{\|\boldsymbol{e}_j^{k+1}\|}{\|\boldsymbol{e}_i^k\|}$

- analytical expression ?
- convergence: $\rho < 1$?
 - minimization of ρ ?

$$\begin{cases} \partial_{t} u_{a} - f v_{a} - \partial_{z} (\nu_{a}(z)\partial_{z}u_{a}) &= F_{a}^{x} \\ \partial_{t} v_{a} + f u_{a} - \partial_{z} (\nu_{a}(z)\partial_{z}v_{a}) &= F_{a}^{y} \end{cases}$$
$$\begin{cases} \partial_{t} u_{o} - f v_{o} - \partial_{z} (\nu_{o}(z)\partial_{z}u_{o}) &= F_{o}^{x} \\ \partial_{t} v_{o} + f u_{o} - \partial_{z} (\nu_{o}(z)\partial_{z}v_{o}) &= F_{o}^{y} \end{cases}$$

- + initial conditions
- + boundary conditions
 - influence of the Coriolis term
 - \triangleright v depends on z and is discontinuous at z = 0 $\nu(z) \in P_0$ (viscous layer), P_1 (log layer), P_2 , P_3 (KPP)
 - Interface conditions:

Robin-Robin

Dirichlet-Neumann $\begin{array}{l} \mathbf{U}_{a}(0) = \mathbf{U}_{o}(0) \\ \nu_{o}(0) \partial_{z} \mathbf{U}_{o}(0) = \nu_{a}(0) \partial_{z} \mathbf{U}_{a}(0) \\ \end{array} \begin{cases} \nu_{a}(0) \partial_{z} \mathbf{U}_{a}(0) + \lambda \mathbf{U}_{a}(0) = \nu_{o}(0) \partial_{z} \mathbf{U}_{o}(0) + \lambda \mathbf{U}_{o}(0) \\ \nu_{o}(0) \partial_{z} \mathbf{U}_{o}(0) + \mu \mathbf{U}_{o}(0) = \nu_{a}(0) \partial_{z} \mathbf{U}_{a}(0) + \mu \mathbf{U}_{a}(0) \end{cases} \end{cases}$

Bulk (linearized or not) $\nu_o(0) \partial_z \mathbf{U}_o(0) = \nu_a(0) \partial_z \mathbf{U}_a(0) = \alpha ||\mathbf{U}_a(0) - \mathbf{U}_o(0)|| (\mathbf{U}_a(0) - \mathbf{U}_o(0))|| (\mathbf{U}_a(0) - \mathbf{U}_o(0$ With or without an averaging of the fluxes $\nu_o(0) \partial_z \mathbf{U}_o(0,t) = \langle \nu_a(0) \partial_z \mathbf{U}_a(0) \rangle$





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 \blacktriangleright Analytical expressions of ρ for the different viscosity profiles and interface conditions







- Analytical expressions of ρ for the different viscosity profiles and interface conditions
- The convergence properties depend on the whole vertical profile ν(z). Strong influence of the turbulent zone.





- \blacktriangleright Analytical expressions of ρ for the different viscosity profiles and interface conditions
- The convergence properties depend on the whole vertical profile ν(z). Strong influence of the turbulent zone.
- Optimization results
- Averaging operator: solution has the correct average. Impact on ρ
- Impact of the nonlinearity of the bulk formula

References

- Thery S., C. Pelletier, F. Lemarié and E. Blayo, 2021: Analysis of Schwarz waveform relaxation for the coupled Ekman boundary layer problem with continuously variable coefficient. *Numerical Algorithms*, doi 10.1007/s11075-021-01149-y.
- Thery S., 2021: Étude numérique des algorithmes de couplage océan-atmosphère avec prise en compte des paramétrisations physiques de couches limites. PhD thesis, Univ. Grenoble Alpes.

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 Clement S., F. Lemarié, and E. Blayo, 2021: Discrete analysis of Schwarz Waveform Relaxation for a simplified air-sea coupling problem with nonlinear transmission conditions. Proceedings of the 26th conference on Domain Decomposition Methods.



Summary and future work

Modeling aspects

- A general framework for the formulation of the coupled OA problem: hypotheses are clearer, two-sided point of view, possible additional parameterizations...
- Still a lot of (difficult) theoretical work to do:
 - Well-posedness of the coupled surface layers problem (a unique fixed-point ?)
 - Coherence of the different parameterizations (O and A turbulent closures and bulk formula)





Summary and future work

Modeling aspects

- A general framework for the formulation of the coupled OA problem: hypotheses are clearer, two-sided point of view, possible additional parameterizations...
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 - Coherence of the different parameterizations (O and A turbulent closures and bulk formula)

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Algorithmic aspects

- Present OA coupling algorithms raise consistency and stability issues. Schwarz algorithms are a convenient framework for developing improved methods.
- Iterative coupling can lead to improved solutions.
- The CPU cost must be decreased: optimizing the interface conditions, imagining preconditioners...

