

Linear Response in Neuronal Network Models Bruno Cessac

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Linear Response in Neuronal Network Models

Bruno Cessac

Université Côte d'Azur, INRIA, France, Biovision team and Neuromod Institute

06-12-2021

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2 From firing rate neurons dynamics to linear response.

3 From spiking neurons dynamics to linear response.



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 Introduction
 From firing rate neurons dynamics to linear response.
 From spiking neurons dynamics to linear response.
 Conclus

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Neuronal networks



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Neurons are electrically active cells where electric currents are produced by ions transfer.



Figure: Neuron. From https://en.wikipedia.org/wiki/Neuron

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Neurons are electrically active cells where electric currents are produced by ions transfer.

Most neurons (but not all of them) communicate by action potentials ("spikes").



Figure: Spike. From http://faculty.washington.edu/chudler/ap.html

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Figure: Spike propagation. From https://openbooks.lib.msu.edu/neuroscience/back-matter/images-of-animations/

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Neurons are electrically active cells where electric currents are produced by ions transfer.

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Figure: Spike train.

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Neurons

Neurons are electrically active cells where electric currents are produced by ions transfer.

Most neurons (but not all of them) communicate by action potentials ("spikes"). Spike trains contain "a" "neural code".

Two main descriptions:

I Firing rates (frequency of spikes within a small time interval) ⇒ first part of the talk.

Individual spikes ⇒ second part of the talk.

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Neuronal networks are usually submitted to external, non stationary, entries (sensory stimuli, interactions between brain areas, ...)

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How to compute the response of a neuronal network to a time dependent stimulus ?

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How to compute the response of a neuronal network to a time dependent stimulus ?



How to compute the response of a neuronal network dynamical system to a time dependent stimulus perturbation ?



From firing rate neurons dynamics to linear response.

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Amari-Wilson-Cowan Model

Amari, 1972; Wilson-Cowan, 1972 Continuous time.

$$\frac{dV_i}{dt} = -V_i + \sum_{j=1}^N J_{ij}f(V_j(t)) + \varepsilon S_i(t); \quad i = 1...N.$$

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- f(x), transfer function (sigmoid) with gain g, e.g.

$$f(x) = \frac{1 + \tanh(gx))}{2};$$

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S:(t), time dependent stimulus applied to neuron i:

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$$f(x) = \frac{1 + \tanh(gx))}{2};$$

- $S_i(t)$, time dependent stimulus applied to neuron *i*;
- ε > 0, small parameter.

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Amari-Wilson-Cowan Model

Amari, 1972; Wilson-Cowan, 1972 Discrete time.

$$V_i(t+1) = \sum_{j=1}^N J_{ij}f(V_j(t)) + \varepsilon S_i(t); \quad i = 1...N.$$

- V:(t), membrane potential of neuron i at time t:
- J_{ii}, (oriented) synaptic weight from presynaptic neuron j to post synaptic neuron i (does not evolve in time);
- f(x), transfer function (sigmoid) with gain g, e.g.

$$f(x) = \frac{1 + \tanh(gx))}{2};$$

- $S_i(t)$, time dependent stimulus applied to neuron *i*;
- ε > 0, small parameter.

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Vector form.

• $\vec{V} = (V_i) \sum_{i=1}^{N}$, state vector; • $\vec{S} = (S_i) \sum_{i=1}^{N}$, stimulus vector; • $\mathscr{J} = (J_{ij}) \sum_{i,j=1}^{N}$, matrix of synaptic weights • $f(\vec{V})$, vector $(f(V_i)) \sum_{i=1}^{N}$.

Continuous time.

$$\frac{d\vec{V}}{dt} = \underbrace{-\vec{V} + \mathscr{J}.f(\vec{V})}_{\vec{F}(\vec{V})} + \varepsilon\vec{S}(t),$$

Discrete time.

$$\vec{V}(t+1) = \underbrace{\mathscr{J}.f(\vec{V}(t))}_{\vec{G}(\vec{V}(t))} + \varepsilon \vec{S}(t).$$

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The 3 "regions" of the sigmoid transfer function

Effect of $f(x) = \frac{1 + \tanh(g x)}{2}$, g > 1, on a time varying signal.

Amplification/Expansion



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The 3 "regions" of the sigmoid transfer function

Effect of $f(x) = \frac{1 + \tanh(gx)}{2}$, g > 1, on a time varying signal.

Saturation/Contraction



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Contractive regime

Spontaneous dynamics. If g is small enough, spontaneous dynamics is contractive \Rightarrow there exists a unique stable fixed point \vec{V}^* attracting all trajectories.

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Contractive regime

Spontaneous dynamics. If g is small enough, spontaneous dynamics is contractive \Rightarrow there exists a unique stable fixed point \vec{V}^* attracting all trajectories.

Time dependent perturbation $\vec{V} = \vec{V}^* + \vec{\delta}V$.

Continuous time.

$$\begin{split} \frac{d\vec{V}}{dt} &= \vec{F}(\vec{V}) + \varepsilon \vec{S}(t) \quad \Rightarrow \quad \frac{d\vec{\delta}V}{dt} = D\vec{F}_{\vec{V}^*} \cdot \vec{\delta}V + \varepsilon \vec{S}(t) + O\left(\varepsilon^2\right), \\ \vec{S}V(t) &= e^{D\vec{F}} \vec{V}^* \cdot (t-t_0) \cdot \vec{\delta}V(t_0) + \varepsilon \int_{t_0}^t e^{D\vec{F}} \vec{V}^* \cdot (t-s) \cdot \vec{S}(s) ds + O\left(\varepsilon^2\right), \\ \vec{\delta}V(t) &\sim \varepsilon \int_{-\infty}^t e^{D\vec{F}} \vec{V}^* \cdot (t-s) \cdot \vec{S}(s) ds = \varepsilon \left[\chi * S\right](t) \\ \hline \chi(t) &= e^{D\vec{F}} \vec{V}^* t \end{split}$$

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Contractive regime

Spontaneous dynamics. If g is small enough, spontaneous dynamics is contractive \Rightarrow there exists a unique stable fixed point \vec{V}^* attracting all trajectories.

Time dependent perturbation $\vec{V} = \vec{V}^* + \vec{\delta} V$.

Discrete time.

$$\vec{\delta}V(t) \sim \varepsilon \sum_{\tau=-\infty}^{t-1} DG_{\vec{V}^*}^{t-\tau-1} \cdot \vec{S}(\tau) = \varepsilon \left[\chi * S\right](t)$$

$$\chi(t) = DG_{\vec{V}^*}^t$$

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Susceptibility and resonances

Linear response

 $\vec{V} = \vec{V}^* + \vec{\delta}V \Rightarrow$

$$\vec{\delta}V(t) \sim \varepsilon[\chi * S](t)$$

Fourier transform.

$$\delta \hat{V}(\omega) = \varepsilon \hat{\chi}(\omega) . \hat{S}(\omega)$$

Complex susceptibility

Susceptibility and resonances



Resonances 40 35 20 15 10 35 30 $|\chi(\omega)|_{20}^{25}$

Susceptibility and resonances



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Susceptibility and resonances

Linear response

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Fourier transform.

$$\delta \hat{V}(\omega) = \varepsilon \hat{\chi}(\omega) . \hat{S}(\omega)$$

Complex susceptibility

Complex poles of the susceptibility give resonances in the real frequency domain

- Real part of the pole \Rightarrow maximum of the ۰ resonance.
- Imaginary part of the pole \Rightarrow width of the • resonance.

Resonances



Bruno Cessac

Linear Response in Neuronal Network Models

The "neural" interpretation of resonances

Discrete time model

$$\begin{split} \vec{V}(t+1) &= \underbrace{\mathscr{J}.f(\vec{V}(t))}_{\vec{G}(\vec{V}(t))} + \varepsilon \vec{S}(t). \\ \chi(t) &= D \vec{G}_{\vec{V}^*}^t \end{split}$$

$$f(x) &= \frac{1 + \tanh(g x))}{2}; \quad f'(x) = \frac{g}{2} \left(1 - \tanh^2(g x)\right) \\ D \vec{G}_{\vec{V}^*} &= g \mathscr{J}.\Lambda(\vec{V}^*); \quad \Lambda_{ij}(\vec{V}^*) = f'(V_i^*) \delta_{ij} \\ \chi_{i,j}(t) &= \sum_{\gamma_{ij}(t)} \prod_{l=1}^t J_{k_l k_{l-1}} \prod_{l=1}^t f'(V_{k_{l-1}}^*), \end{split}$$

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The "neural" interpretation of resonances





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k_2 .t=2

 $k_1, t=4$

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The "neural" interpretation of resonances



The "neural" interpretation of resonances



dependent neural pathways.



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Beyond the contractive regime

Discrete time model

$$\vec{V}(t+1) = \underbrace{\mathscr{J}.f(\vec{V}(t))}_{\vec{G}(\vec{V}(t))}$$

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<u>Beyond</u> the contractive regime

Discrete time model

$$ec{V}(t+1) = \underbrace{\mathscr{J}.f(ec{V}(t))}_{ec{G}(ec{V}(t))}$$

$$f(\mathbf{x}) = \frac{1 + \tanh(g\mathbf{x}))}{2}; \quad f'(\mathbf{x}) = \frac{g}{2} \left(1 - \tanh^2(g\mathbf{x})\right)$$
$$D\vec{G}_{1'\mathbf{x}} = g \not J \cdot \Lambda(\vec{\mathbf{x}}^*); \quad \Lambda_{ii}(\vec{\mathbf{y}}^*) = f'(\mathbf{y}_i^*) \,\delta_{ii}$$

The stability of \vec{V}^* is controlled by g and the spectrum of J.

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Beyond the contractive regime

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The stability of \vec{V}^* is controlled by g and the spectrum of J.

Take \mathscr{J} a random matrix, indep. entries, $\mathscr{N}(0, \frac{1}{N})$. (\mathscr{J} does not evolve in time.)

Discrete time model

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As g increases the stable fixed destabilizes (generically) by a Hopf bifurcation.



Beyond the contractive regime

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Transition to chaos

$$m(t) = \frac{1}{N} \sum_{j=1}^{N} V_j(t)$$



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Discrete time model

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As g increases the stable fixed destabilizes (generically) by a Hopf bifurcation.

When g increases, there is a transition to chaos by quasi-periodicity (Ruelle-Takens, 81).

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$$D\vec{G}_{\vec{U}^*} = g \not J \cdot \Lambda(\vec{V}^*); \quad \Lambda_{ii}(\vec{V}^*) = f'(V_i^*) \delta_{ii}$$

The stability of \vec{V}^* is controlled by g and the spectrum of J.

Take \mathscr{J} a random matrix, indep. entries, $\mathscr{N}(0, \frac{1}{N})$. (I does not evolve in time.)

The spectral density is know Girko, 1985; Edelman, 1993.

As g increases the stable fixed destabilizes (generically) by a Hopf bifurcation.

When g increases, there is a transition to chaos by quasi-periodicity (Ruelle-Takens, 81).

$$m(t) = \frac{1}{N} \sum_{j=1}^{N} V_j(t)$$



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Linear response in the chaotic regime

B. Cessac, J.A. Sepulchre, PRE (2004); Chaos (2006); Physica D (2006)

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Linear response in the chaotic regime

B. Cessac, J.A. Sepulchre, PRE (2004); Chaos (2006); Physica D (2006)

Time dependent perturbation

 $\vec{\delta}$

$$\begin{split} \vec{V}(t+1) &= \vec{G} \left(\vec{V}(t) \right), \quad \vec{V}'(t+1) = \vec{G} \left(\vec{V}'(t) \right) + \varepsilon \vec{S}(t) \\ \vec{V}(t_0) &= \vec{V}'(t_0); \quad \vec{\delta} V(t) = \vec{V}'(t) - \vec{V}(t), t \ge t_0 \\ \vec{\delta} V(t_0+1) &= \vec{V}'(t_0+1) - \vec{V}(t_0+1) = \varepsilon \vec{S}(t_0) \\ V(t_0+2) &= \vec{G} \left(\vec{V}(t_0+1) + \varepsilon \vec{S}(t_0) \right) + \varepsilon \vec{S}(t_0+1) - \vec{G}(\vec{V}(t_0+1)) \\ \vec{\delta} V(t_0+2) &= \varepsilon \left[DG_{\vec{V}(t_0+1)} \cdot \vec{S}(t_0) + \vec{S}(t_0+1) \right] + \varepsilon^2 \vec{\eta}(t_0+1) \\ \hline \vec{\delta} V(t) &= \varepsilon \sum_{\tau=t_0}^{t-1} DG_{\vec{V}(\tau+1)}^{t-\tau-1} \cdot \vec{S}(\tau) + \varepsilon^2 \vec{R}(t) \end{split}$$

Geometric representation g=3.5, k,=0.158

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Time dependent perturbation



Time dependent perturbation



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Time dependent perturbation



Linear response vs chaotic

Butterfly effect Van Kampen objection

Time dependent perturbation





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Bruno Cessac Linear Response in Neuronal Network Models

Time averaging is robust to perturbation.

$$\lim_{T\to\infty}\frac{1}{T}\sum_{t=1}^{T}\Phi(\vec{G}^{t}(\vec{V}))^{"}=^{"}\lim_{T\to\infty}\frac{1}{T}\sum_{t=1}^{T}\Phi(\vec{G}^{t}(\vec{V}+\vec{\delta}V))$$

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 μ_L Lebesgue measure on the phase-space $\mathscr{M}.$

$$\mu \stackrel{\scriptscriptstyle{W}}{=} \lim_{t \to +\infty} \vec{G}^t \mu_L, \quad \text{SRB measure}$$

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Natural notion of averaging "on" the attractor.

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Gibbs distributions

- Equilibrium stat. phys. (Max. Entropy Principle).
- Non equ. stat. phys. Onsager theory.
- Ergodic theory.

 $\mu \stackrel{\scriptscriptstyle W}{=} \lim_{t \to +\infty} \vec{G}^t \mu_L, \quad {\rm SRB \ measure}$

The Sinai-Ruelle-Bowen measure is a Gibbs measure

$$H(\vec{V}) = -\log \det \pi^u_{\vec{V}} D \vec{G}_{\vec{V}}$$

where $\pi^{u}_{\vec{V}}$ is the local projection on the unstable manifold.



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Out of equilibrium SRB state

D. Ruelle, J. Stat. Phys., 1998

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$$\lim_{T \to \infty} \frac{1}{T} \sum_{t=1}^{T} \Phi\left[\vec{G}^{t}(\vec{V})\right]^{\mu_{L}} \stackrel{a.s.}{=} \int_{\mathscr{M}} \Phi(\vec{V}) \mu(d\vec{V})$$



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Non eq. SRB

$$\mu_t = \lim_{n \to +\infty} \vec{G}'_t \dots \vec{G}'_{t-n} \, \mu_L$$

$$\langle \Phi \rangle_t = \int_{\mathscr{M}} \Phi(\vec{V}) \mu_t(d\vec{V})$$



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Non eq. SRB

$$\mu_t \equiv \mu + \delta_t \mu = \lim_{n \to +\infty} \vec{G}'_t \dots \vec{G}'_{t-n} \mu_L$$

$$\delta_{t}\mu[\Phi] = \varepsilon \sum_{\tau=-\infty}^{t-1} \int \mu(d\vec{V}) D\vec{G}_{\vec{V}}^{t-\tau-1}\vec{S}(\tau) \cdot \vec{\nabla}_{\vec{V}(t-\tau-1)} \Phi + O(\varepsilon^{2})$$
$$\delta_{t}\mu[\Phi] = \varepsilon \sum_{\tau=-\infty}^{t-1} \left\langle \kappa_{\tau}\vec{S}_{t-\tau-1} \circ \vec{G}^{-1} | \Phi \right\rangle_{eq}$$



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Brief sketch

The out of equilibrium SRB state is well defined for uniformly hyperbolic dynamical systems (Ruelle, 1998).

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- The out of equilibrium SRB state is well defined for uniformly hyperbolic dynamical systems (Ruelle, 1998).
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Is our model unif. hyperbolic?

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Linear response in the firing rate neural network

B. Cessac, J.A. Sepulchre, PRE (2004); Chaos (2006); Physica D (2006); B. Cessac Chaos (2019)

Convolution

$$\delta_{t}\mu\left[\Phi\right] = \varepsilon \sum_{\tau=-\infty}^{t-1} \int \mu(d\vec{V}) D\vec{G}_{\vec{V}}^{t-\tau-1} \vec{S}(\tau) \cdot \vec{\nabla}_{\vec{V}(t-\tau-1)} \Phi$$

$$\Phi(\vec{V}) = \vec{V}$$

$$\delta_{t}\mu\left[\vec{V}\right] = \varepsilon \sum_{\tau=-\infty}^{t-1} \chi(t-\tau-1) \cdot \vec{S}(\tau) + O(\varepsilon^{2}) \equiv \varepsilon [\chi * S]_{i}(t)$$

$$\chi(t-\tau-1) = \left\langle D\vec{G}_{\vec{V}}^{t-\tau-1} \right\rangle_{eq}$$

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Numerical computation

Complex susceptibility

$$\hat{\chi}_{ij}(\omega) = \sum_{t=-\infty}^{+\infty} e^{i\,\omega\,t}\chi_{ij}(t)$$

Consider two perturbations, $\varepsilon \cos(\omega t)\vec{e}_j$ and $-\varepsilon \sin(\omega t)\vec{e}_j$ where \vec{e}_j , is the canonical basis vector in direction j and denote $\vec{V}^{(1)}, \vec{V}^{(2)}$ the corresponding perturbed dynamics.

Ergodic average

$$\hat{\chi}_{ij}(\omega) = \lim_{T \to +\infty} \frac{1}{T \varepsilon} \sum_{t=0}^{T-1} e^{i \, \omega \, (t-1)} \left[V_i^{(1)}(t) + i \, V_i^{(2)}(t) \right]$$

Ruelle resonances

The tangent part of the response is a correlation function \Rightarrow Green-Kubo relation



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- ۰ Ruelle-Pollicott resonances: In the power spectrum. Absolutely continuous part of the SRB measure. Do not depend on the observable (spectrum of the Ruelle-Perron-Frobenius operator).

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- The transverse part (which does not exist in conservative systems) is not a correlation function
- Exotic resonances. Not in the power spectrum. Singular part of the SRB measure.
- Predicted by D. Ruelle (J. Stat. Phys, 1999)
- Exhibited (numerically) in B. Cessac, J.A. Sepulchre, PRE, 2004.

Perturbations effect

Complex susceptibility



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Introduction

From firing rate neurons dynamics to linear response.

From spiking neurons dynamics to linear response. Conclus

Response to a time-dependent stimulus





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Signal propagation by amplitude modulation

B. Cessac, J.A. Sepulchre, Chaos (2006):

Amplitude modulation

- Fix a resonant frequency ω_0 . ۰
- Define a slow amplitude modulation signal

 $S_i(t) = \cos(\omega_M t) \sin(\omega_0 t)$

where ω_M corresponds to a very long period.

$$\hat{\chi}_{ij}(\omega) = \lim_{T \to +\infty} \frac{1}{T\varepsilon} \sum_{t=0}^{T-1} e^{i\omega(t-1)} \left[V_i^{(1)}(t) + i V_i^{(2)}(t) \right]$$

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Conclus

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Can we recover the signal "hidden" in the chaotic background ?

We have 3 different notions of connectivity which do not coincide.



Introduction From firing rate neurons dynamics to linear response. From spiking neurons dynamics to linear response. Conclus 00000000000000000000000

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Linear response theory

Gibbs like distributions

- Equilibrium stat. phys. •
- Non equilibrium stat. phys. Onsager theory. ۰
- Green-Kubo relations.
- ٠ Ergodic theory for "chaotic" systems. SRB measure.
- Non equilibrium SRB state. •

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Firing rate model

- Linear response is possible in a chaotic neural network
- Convolution kernel depending on synaptic graph and dynamics.
- The response graph is different from the synaptic weights graph and depends on the stimulus.
- Further steps in Onsager theory ? Transport of entropy ? Information current ?

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From spiking neurons dynamics to linear response.

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R.Cofré, B. Cessac, Chaos, Solitons and Fractals, 2013



R.Cofré, B. Cessac, Chaos, Solitons and Fractals, 2013



Spikes

- Voltage dynamics is time-continuous.
- Spikes are time-discrete events (time resolution $\Delta > 0$).

$$t_{k}^{(l)} \in [n\Delta, (n+1)\Delta[$$

$$\Rightarrow$$

$$\omega_{k}(n) = 1$$

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Time (ms)

R.Cofré, B. Cessac, Chaos, Solitons and Fractals, 2013

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Spike state ω_k(n) ∈ 0,1.

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R.Cofré, B. Cessac, Chaos, Solitons and Fractals, 2013

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R.Cofré, B. Cessac, Chaos, Solitons and Fractals, 2013

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- Spike block ωⁿ_m.

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M. Rudolph, A. Destexhe, Neural Comput. 2006, (GIF model) R.Cofré, B. Cessac, Chaos, Solitons and Fractals, 2013



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Sub-threshold dynamics:

$$C_{k} \frac{dV_{k}}{dt} + g_{k}(t,\omega) V_{k} = i_{k}(t,\omega)$$

$$W_{kj} \stackrel{\text{def}}{=} G_{kj} E_j$$

 $lpha_{kj}(t, \omega) = \sum_{n \ge 0} \alpha_{kj}(t - n\delta)\omega_j(n)$

$$i_{k}(t, \omega) = \underbrace{g_{L,k}E_{L}}_{Leak} + \underbrace{\sum_{j}W_{kj}\alpha_{kj}(t, \omega)}_{Synapses} + \underbrace{S_{k}(t)}_{Stimulus} + \sigma_{B}\underbrace{\xi_{k}(t)}_{Noise}$$
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- Linear in V.
- Spike history-dependent.

Dynamics integration

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$$\Gamma_k(t_1, t, \omega) = \begin{cases} \\ \end{cases}$$

,

$$e^{-\frac{1}{C_k}\int_{t_1}^t g_k(u,\omega)\,du}$$

if $t \in [t_1, \tau_k(t,\omega)];$
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$$\Gamma_{k}(t_{1},t,\omega) = \begin{cases} e^{-\frac{1}{C_{k}}\int_{t_{1}}^{t}g_{k}(u,\omega) du} \\ \text{if } t \in [t_{1},\tau_{k}(t,\omega)]; \\ 0, \text{otherwise.} \end{cases}$$

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$$V_k(t,\omega) = V_k^{(sp)}(t,\omega) + V_k^{(S)}(t,\omega) + V_k^{(noise)}(t,\omega)$$

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$$V_k^{(syn)}(t,\omega) = \frac{1}{C_k} \sum_{j=1}^N W_{kj} \int_{\tau_k(t,\omega)}^t \Gamma_k(t_1,t,\omega) \alpha_{kj}(t_1,\omega) dt_1$$

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$$\begin{split} \Gamma_{k}(t_{1},t,\omega) &= \begin{cases} e^{-\frac{1}{C_{k}}\int_{t_{1}}^{t}g_{k}(u,\omega)\,du}, \\ &\text{if } t \in [t_{1},\tau_{k}(t,\omega)]; \\ 0, \text{otherwise}. \end{cases}, \\ V_{k}(t,\omega) &= \underbrace{V_{k}^{(sp)}(t,\omega) + V_{k}^{(S)}(t,\omega)}_{V_{k}^{(det)}(t,\omega)} + V_{k}^{(noise)}(t,\omega)}_{V_{k}^{(det)}(t,\omega)} \end{cases}$$

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Variable length Markov chain

$$P_n\left[\omega(n) \left| \omega_{-\infty}^{n-1} \right. \right] \equiv \Pi\left(\omega(n), \frac{V_{th} - V_k^{(det)}(n-1,\omega)}{\sigma_k(n-1,\omega)}\right)$$

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Time (ms)

Markov chain: $P\left[\omega(n) \middle| \omega_{n-D}^{n-1} \right] > 0$



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$$\mu\left[\omega_{m}^{n}\right] = \prod_{l=m+D}^{n} P\left[\omega(l) \middle| \omega_{l-D}^{l-1} \right] \mu\left[\omega_{m}^{m+D-1}\right], \quad \forall m < n \in \mathbb{Z}$$

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$$\phi\left(\omega_{l-D}^{l}\right) = \log P\left[\omega(l) \middle| \omega_{l-D}^{l-1} \right]$$

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Markov chain:
$$P\left[\omega(n) \middle| \omega_{n-D}^{n-1} \right] > 0$$

$$\mu\left[\omega_{m}^{n}\right] = \exp \sum_{l=m+D}^{n} \phi_{l}\left(\omega_{l-D}^{l}\right) \mu\left[\omega_{m}^{m+D-1}\right], \quad \forall m < n \in \mathbb{Z}$$

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Gibbs distribution

Gibbs like distributions Equilibrium stat. phys. • Non equilibrium stat. phys. Onsager theory. Green-Kubo relations. • Ergodic theory for "chaotic" systems. SRB measure. Non equilibrium SRB state. ۰ Markov chains.

Potential

$$\mu\left[\omega_{m}^{n} \mid \omega_{m}^{m+D-1}\right] = \exp\sum_{l=m+D}^{n} \phi_{l}\left(\omega_{l-D}^{l}\right)$$

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Infinite memory depth $D \rightarrow +\infty$.

$$\phi\left(\left.\omega_{-\infty}^{l}\right) = \log P\left[\left.\omega(l)\right|\omega_{-\infty}^{l-1}\right]$$

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Response to stimuli

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GIF model without stimulus \Rightarrow stationary Gibbs distribution.

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History dependence.

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History dependence, observable

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History dependence, observable, network dynamics

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History dependence, (spontaneous) correlation between observable and network dynamics

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Firing rates

Computations have been done using the INRIA Pranas software, B. Cessac et al, Frontiers in Neuroinformatics, 2017.



Figure: Linear response of $f(\omega, n) = \omega_{k_c}(n)$ for different values of stimulus amplitude A.

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From spiking neurons dynamics to linear response. Conclus

Higher order observable

Computations have been done using the INRIA Pranas software, B. Cessac et al, Frontiers in Neuroinformatics, 2017.



Figure: Linear response of the observable $f(n, \omega) = \omega_{k_c-2}(n-3)\omega_{k_c}(n)$.

Introduction	From firing rate neurons dynamics to linear response.	From spiking neurons dynamics to linear response.	Conclus ○
Range	of validity		

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Figure: L^2 distance between the curves $\delta\mu^{(1)}[f(n)]$, $\delta\mu^{(HC_1)}[f(n)]$ and the empirical curve, as a function of the stimulus amplitude A. Left panel show distance between rate curves and right panel distance between pairwise observable with delay.

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Introduction From firing rate neurons dynamics to linear response. From spiking neurons dynamics to linear response. Conclus

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Main conclusions

- ۰ Linear response in a spiking neural network.
- Convolution kernel depending on synaptic graph and dynamics. •
- Link with receptive fields for sensory neurons ?

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General conclusion

Gibbs like distributions

- Equilibrium stat. phys.
- Non equilibrium stat. phys. Onsager theory.
- Green-Kubo relations.
- Ergodic theory for "chaotic" systems. SRB measure.
- Non equilibrium SRB state.
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Extensions

The framework of Gibbs distributions, in a very wide sense, provides a useful way of studying the response of neuronal networks to non stationary stimuli from the knowledge of spontaneous dynamics.

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Analysis of experimental neural data (see e.g. https://team.inria.fr/biovision/ macular-software/).

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- Analysis of experimental neural data (see e.g. https://team.inria.fr/biovision/ macular-software/).
- Handling non linear effects \Rightarrow computing higher order terms might not be efficient from a practical purpose.
- Extend the Gallavotti-Ruelle idea of currents associated to entropy production to information fluxes in neuronal networks.