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► **To cite this version:**

Bruno Cessac. Retinal processing: insights from mathematical modelling. J. Imaging, inPress. hal-03454859v1

HAL Id: hal-03454859

<https://inria.hal.science/hal-03454859v1>

Submitted on 29 Nov 2021 (v1), last revised 17 Jan 2022 (v2)

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Retinal processing: some mathematical insights

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Abstract: The retina is the entrance of the visual system. Although based on common biophysical principles the dynamics of retinal neurons is quite different from their cortical counterparts, raising interesting problems for modellers. In this paper I address some mathematically stated questions in this spirit, discussing, in particular: (1) How could lateral (amacrine cells) connectivity shape the spatio-temporal spike response of retinal ganglion cells? (2) How could spatio-temporal stimuli correlations and retinal network dynamics shape the spike train correlations at the output of the retina. I also briefly discuss some potential consequences of these results at the cortical level.

Keywords: Retinal network; visual system; spatio-temporal spike correlations; linear response; non stationarity

0. Introduction

Let us start with a very simple experiment. Look around you... That's it, the experiment is over. A very ordinary experience, isn't it? Is it really though? Let us first point out that looking around you to see, that is, having the sense of sight, is indeed ordinary — except for those who have partially or totally lost their ability to see. We will come back to this point at the end of the paper. Now, excluding visual impairments, vision is everything but ordinary.

Think of it. A flux of photons, with frequencies in the visible spectrum range, emitted by the external world around us enters into our eyes, then "something" happens, and we see. Thanks to constant progress in experimental and theoretical neuroscience, we understand better and better this "something", the mechanisms of vision, although our view of it is far from being complete. Especially, in these times of artificial intelligence, bio-inspired computing, computer vision and ... global warming, it might be helpful to understand how our brain is able to handle the complex visual information coming from the external world so rapidly and efficiently ... with an energy consumption of the order of a few Watts.

Certainly, the retina plays a central role in this process. It is known for long that this is definitely not a camera. The retina is smart [1] and it has to be. Think especially of the difference of scale between the retina and the visual cortex, in terms of size but also numbers of neurons and synapses. As everything that goes to the visual cortex comes from the retina, this little membrane, at the back of the eye, half a millimetre thick, with an area of order a cm^2 (for humans), has somewhat to filter the visual information, leaving out "irrelevant detail" and stick at crucial events, and then, signal them in the proper way to the brain via spike trains. As a matter of fact, the question(s) of "efficiently" encoding information by spikes has been the subject of many fascinating papers [2–4], especially in the seminal paper from Barlow [5] with concepts such as reducing redundancy, information compression and efficient coding. These concepts are regularly updated with recent experimental and theoretical investigations [6–15]. We come back to this point, at the end of the paper too.

The retina has, roughly, the following structure (Fig. 1, left). For more detail see e.g. [16] or <https://webvision.med.utah.edu/book/part-i-foundations/simple-anatomy->

Citation: Cessac B. Retinal processing: mathematical insights. *J. Imaging* 2021, 1, 0. <https://doi.org/>

Received:

Accepted:

Published:

Publisher's Note: MDPI stays neutral with regard to jurisdictional claims in published maps and institutional affiliations.

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of-the-retina/)). It is organized in five neuronal types : Photo-receptors, rod and cones (P), horizontal cells (HCells), bipolar cells (BCells), amacrine cells (ACells), ganglion cells (GCells), to which are added glial cells (Mueller's cells). These neurons types are connected by chemical and electric synapses, in specific functional circuits or "pathways" (like the rod-cone pathway [17,18]) which are a key in the retinal capacity to convert the light coming from a visual scene into spike patterns sent to the visual cortex, through the Lateral Geniculate Nucleus (LGN), via the optic nerve made of GCells axons. In particular, there are in the retina very specific synapses like the ribbon synapse enabling neurons to transmit light signals from photoreceptors to BCells over a dynamic range of several orders of magnitude in intensity [19]. Roughly, two main connectivity structure can be distinguished: feed-forward, the P-B-G path which leads from the photo transduction to the spike trains emitted by the GCells towards the cortex. There is also a lateral connectivity due to Hcells, at the origin of the Center-Surround structure of the receptive fields, and the Acells whose role is still poorly understood and which are one of the main object of study of this paper.

57

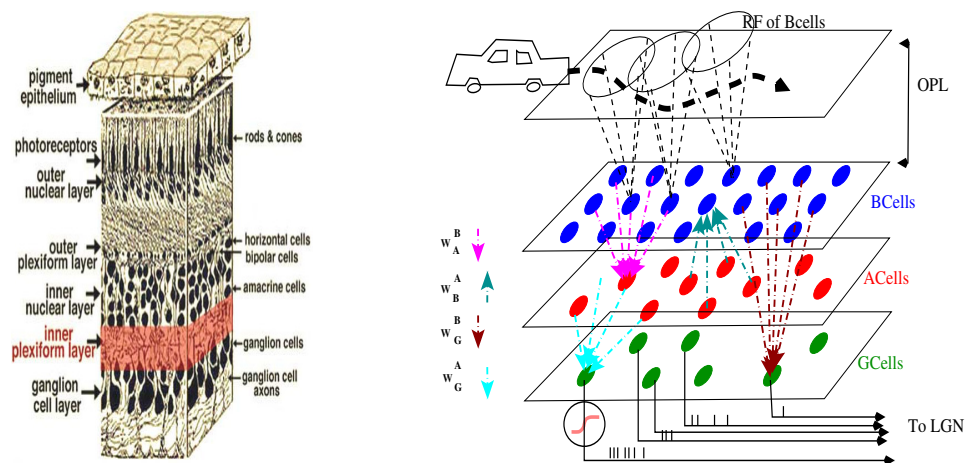


Figure 1. Left. Structure of the biological retina. From <http://webvision.org/es/gross-anatomy-of-the-eye/1-2-simple-anatomy-of-the-retina/> **Right. Structure of the model introduced in section 1.1.** A moving object (here, presumably, a car) moves along a trajectory (dashed black line). Its image is projected by the eye optics to the upper retina layers (Photoreceptors and HCells) and stimulates them. In the model, this corresponds to the convolution of the stimulus with the Receptive Field (RF) of BCells. This provides to BCells what we call the "OPL" input to BCells. BCells (blue points) are connected to ACells (red points) via excitatory synapses (pink arrows, denoted W_A^B) and to GCells (green points) via excitatory synapses (brown arrows, denoted W_G^B). ACells are connected to BCells via inhibitory synapses (green arrows, denoted W_B^A) and to GCells via inhibitory synapses (cyan arrows, denoted W_C^A). The voltage of GCells is sent through a non linearity (pink curve in the black circle) so as to produce spike trains conveyed to the LGN.

58 The structure of the retina and its behaviour are thus well studied on the experi-
 59 mental side. There are comparatively less modelling studies and quite less mathematical
 60 results on the way how retinal structure, especially, lateral ACells connectivity shapes
 61 the spike response to spatio-temporal stimuli. One of the goals of this paper is to elicit
 62 reflections in this direction, grounded on recent mathematical developments fed by the
 63 recent progresses in the knowledge of retina physiology and structure. This is a humble
 64 and partial point of view, issued from my collaboration with neuro-biologists experts
 65 in the retina. More precisely, I am addressing the following problems on mathematical
 66 grounds. In the main text I focus on neuroscience modeling perspective, whereas, in the
 67 conclusion section, I discuss potential consequences of these results out of the field of

68 neuroscience.

69

70 **Problem 1. How does the structure of the retina, in particular, amacrine lateral**
 71 **connectivity condition the retinal response to dynamic stimuli?**

72 The problem can be addressed at several levels.

73

Level 1. Single cell response to stimuli. The individual response of ganglion cells is usually expressed in terms of their receptive field. This notion is on the one hand phenomenological: it is observed that each ganglion cell responds preferentially to stimuli, localized in space, with a characteristic spatio-temporal structure. For example, a ON-Center cell preferentially responds to an increase in luminance in a circular area corresponding to the central part of the receiving field. This notion is also expressed mathematically by a kernel \mathcal{K}_G , i.e. a function of space and time, so that the response of a GCell to a spatio-temporal stimulus $S(x, y, t)$, takes the form:

$$\left[\mathcal{K}_G \begin{smallmatrix} x,y,t \\ * \end{smallmatrix} S \right] (t) = \int_{x=-\infty}^{+\infty} \int_{y=-\infty}^{+\infty} \int_{s=-\infty}^t \mathcal{K}_G(x - x_C, y - y_C, t - s) S(x, y, s) dx dy ds, \quad (1)$$

74 where $\begin{smallmatrix} x,y,t \\ * \end{smallmatrix}$ means space (x, y) -time (t) convolution. x_C, y_C are the coordinates of the RF
 75 center. The integrals are well defined since the kernel decreases fast enough to infinity,
 76 in space and time, to guarantee convergence. The upper bound in time, t , expresses
 77 causality, whereas the lower bound, $-\infty$, implicitly assumes that the stimulus has
 78 been applied in a distant past compared to t , quite longer than the characteristic times
 79 involved in GCell response.

80 Equation (1) corresponds to a linear response. It is therefore only valid for stimuli
 81 of low amplitude in voltage. More generally, the voltage response to the stimulus is
 82 a functional of the stimulus that one can, under well posed mathematical conditions,
 83 write as a Volterra expansion [20], (1) being the lowest order (linear) term. Unfortunately,
 84 higher-order terms are essentially inaccessible experimentally and one usually constrains
 85 instead the non-linearity of the response under other modalities. Especially, taking into
 86 account that the response of a ganglion cell to a stimulus is, ultimately, a sequence of
 87 spikes, one writes the probability density of emitting a spike between t and $t + dt$ in
 88 the form $f\left(\left[\mathcal{K}_G \begin{smallmatrix} x,y,t \\ * \end{smallmatrix} S\right](t) + b\right)$ where f is a nonlinear positive increasing function
 89 (typically, a sigmoid), b is a threshold constraining the level of activity of the GCell in
 90 the absence of stimulation. This procedure defines an inhomogeneous Poisson process
 91 called the linear-non-linear Poisson (LNP) model [21–23]. Experimentally the kernel
 92 \mathcal{K}_G is determined by Spike-Triggered Average or Spike-Triggered Correlation technique,
 93 studying the response at a white noise [21]. Nonlinearity is then determined, typically
 94 by the Levenberg-Marquart method [24]. This modelling asks however the following
 95 questions:

- 96 (i) How is the kernel \mathcal{K}_G of the GCells constrained by the structure/dynamics of
 97 the upper layers of retinal cells ?
- 98 (ii) The forms (1) implicitly assumes that \mathcal{K}_G does not depend on the stimulus. Can
 99 one write mathematical conditions that guarantee such an independence?
- 100 (iii) To which extent is the notion of Ganglion Cells Receptive Field compatible with
 101 non linear effects reported in retinal neurons and synapses, such as voltage
 102 rectification or gain control ?

103 **Level 2. Collective response to stimuli and spike statistics.** GCells do not inter-
 104 act directly, but amacrine connectivity induces an effective, causal interaction between
 105 them. What is therefore the structure of the spatio-temporal correlations induced by the
 106 conjunction of the spatio-temporal stimulus and the response of the retinal network, in
 107 particular, the amacrine lateral connectivity ? A classical paradigm in neural coding is
 108 to assume that the retina decorrelates GCells outputs to maximize information transfer

109 [6–11,13–15]. It is in particular believed that ACells play a central role in this decorrela-
110 tion process (see [15] and references therein). What can be, at the mathematical level, the
111 conditions, on the stimulus and dynamics, that allow a network of neurons interacting
112 with each other to produce vanishing, or at least, *weak* correlations ? When does weak
113 mean negligible ? These questions are actually closely related to the second problem.

114

115 **Problem 2. Spike statistics.**

116 More generally, considering the retina as a dynamical system forced by non-
117 stationary, spatially inhomogeneous stimuli, what could be a general form for the
118 (non-stationary) statistic of spike trains emitted by ganglion cells, taking into account
119 that spike trains emitted by the retina are all what the LGN and cortex see ? One can
120 attempt to construct a canonical form of probability distributions of the retinal spike
121 trains taking into account that:

- 122 (i) Stimuli, thus statistics, are not stationary;
- 123 (ii) The cortex (and before, the LGN) only receive spikes, thus have no information
124 about the biophysical processes which have generated those spikes and no
125 information on the underlying dynamics of the retina (voltages, activation
126 variables, conductances). All the information is contained in the spatio-temporal
127 structure of spikes;
- 128 (iii) Spike trains distribution may have a long memory.

129 In this paper I address these problems with the help of two models. The first,
130 presented in section 1.1, grounded on biology and e.g. the papers [25–28] mimics the
131 Bipolar-Amacrine-Ganglion cells network and is used, in sections 2.1, 2.2, to make pro-
132 gresses in elucidating problem 1. I first show how one can obtain an explicit form for
133 the kernel (1) featuring the ACells lateral connectivity. This RF explicitly depends on
134 the BCs-ACs network through the eigenvalues and eigenvectors of a operator I call
135 "transport operator". I discuss some consequences of this result, especially in terms of
136 response to propagating stimulus. This result is valid when cells act as linear integrators.
137 However, cells are in general rectified by non linearities. I propose piecewise-linear
138 rectifications (as used in several retina model) and I discuss how rectification acts on
139 the RF of eq. (1). A striking conclusion is that, if the convolution form (1) is preserved,
140 this is to the price of having a RF depending on the stimulus. A consequence of this
141 analysis is that spike correlations may depend on the stimulus and are expected to be
142 quite different when considering e.g. objects moving along trajectories in comparison to
143 static images.

144

145 The second model, introduced in section 1.2 and analysed in section 2.3 attempts to
146 propose a canonical form of probability distributions of the retinal spike trains based
147 on the constraints (i), (ii), (iii) above. These sections essentially presents the conclusions
148 of works published elsewhere [29–39]. As I argue, these constraints leads to a natural
149 notion of spike probabilities, somewhat extending the statistical physics notion of Gibbs
150 distribution to the non stationary case. In this setting, one establishes a linear response for
151 a network of interacting spiking cells, that can mimic a set of GCells coupled via effective
152 interactions corresponding to the ACells network influence. This linear response theory
153 not only gives the effect of a non stationary stimulus to first order spike statistics (firing
154 rates) but also its effect on higher order correlations. Indeed, spike correlations are
155 modified by a spatio-temporal stimulus and can be computed thanks to the knowledge
156 of spontaneous correlations. The linear response formula is expressed as a convolution
157 where the kernel can be explicitly computed for an Integrate and Fire conductance
158 based model. Moreover, as I argue, these spike trains distributions have close links with
159 information geometry. Especially, they induce a natural metric in an abstract space of
160 probabilities, with close potential links with the neuro-geometry introduced by Sarti,
161 Citti, Petitot et al [40–43]. This is discussed in the conclusion section.

162 More generally, the application and discussion sections shortly proposes possible
 163 extension of this work to several domains: Retinal prostheses, section 3.1; Convolutional
 164 networks, section 3.2; Implications for cortical models, section 4.1; Neuro-geometry,
 165 section 4.2.

166 1. Materials and Methods

167 1.1. Modelling the retinal network

168 1.1.1. Specificities of the retina

169 Neurons in the retina have the same biophysics as their cortical counterparts.
 170 However, they operate under different modalities. Remarkably, with the exception of the
 171 GCells, the retinal neurons do not emit action potentials. Their activity and interactions
 172 therefore take place through graded (continuous) membrane potentials as opposed to the
 173 sharp peak of an action potential. Furthermore, there is no long-term synaptic plasticity
 174 in the retina. Finally, the main "computational" elements in the retina are functional
 175 circuits [18] made of a few neurons and synapses, in large contrast with "computational"
 176 units in the visual cortex, such as cortical columns, involving thousands of neurons.
 177 A modelling consequence is that mean-field or neural masses description used in the
 178 cortex might not be relevant to study the retina.

179 The goal of this paper is to address mathematical questions about the dynamics and
 180 behaviour of the retina embedded in the visual system. To instantiate these questions on
 181 a firm mathematical ground we are going to consider a model of the retinal network,
 182 based on a few fundamental facts briefly exposed in the previous section:

- 183 1. The retina is a high dimensional, non autonomous and noisy dynamical system,
 184 layered and structured, with non stationary and spatially inhomogeneous entries
 185 (visual scenes).
- 186 2. Most retinal neurons are not spiking, except GCells. Thus, retina performs analogic
 187 computing.
- 188 3. Local retinal circuits efficiently process the local visual information. These local
 189 circuits are connected together, spanning the whole retina in a regular tiling. In
 190 this perspective, it is important to consider individual neurons and synapses, in
 191 contrast, e.g., to cortical modelling where it is relevant to consider mean-field
 192 approaches averaging over populations.

193 Thus, the model presented below and in Fig. 1, right, is non stationary, with a
 194 layered retina's like structure, where dynamics ruling BCells, ACells, and GCells voltage
 195 is piecewise linear. As we discuss, the model affords additional non linearities like gain
 196 control although we will not delve into the maths of the gain controlled case (see [28] for
 197 a study of this case). For GCells, the spiking process is mimic by a non linear firing rate
 198 so that our model enters in the class of LNP models.

199 1.1.2. Structure of the retina model

200 We assimilate the retina to a superimposition of 3 layers, each one being a flat, two
 201 dimensional square of edge length L mm where spatial coordinates are noted x, y (Fig.
 202 1, right). Each layer corresponds to a cell population (BCells, ACells, GCells) where
 203 the density of cells is taken uniform. We note δ_p the lattice spacing in mm, and N_p the
 204 total number of cells in the layer p . Without loss of generality we assume that L , the
 205 retina's edge size, is a multiple of δ_p . We note $L_p = \frac{L}{\delta_p}$, the number of cells p per row or
 206 column so that $N_p = L_p^2$. Each cell in the population p has thus Cartesian coordinates
 207 $(x, y) = (i_x \delta_p, i_y \delta_p)$, $(i_x, i_y) \in \{1, \dots, L_p\}^2$. To avoid multiples indices, we associate to
 208 each pair (i_x, i_y) a unique index $i = i_x + (i_y - 1) L_p$. The cell of population p , located at
 209 coordinates $(i_x \delta_p, i_y \delta_p)$ is then denoted by p_i .

210 One can roughly subdivide the real retina into two blocks (Fig. 1). The first, that we
 211 name in short, for modelling purposes¹ OPL, (Outer Plexiform Layer), includes the P,
 212 HCells, BCells and the related synapses. As an "input" of this block is the flow of photons
 213 emitted by the outside world and picked up by the photo-receptors. In our model, this
 214 corresponds to a "stimulus", i.e. a function $\mathcal{S}(x, y, t)$ where t is the time coordinate. As
 215 we don't consider color sensitivity here \mathcal{S} characterizes a black and white scene, with
 216 a control on the level of contrast $\in [0, 1]$. The "output" of the OPL is sent to Bcells in
 217 the form of a "drive" voltage, defined in eq. (2) below. In the real retina, the voltage
 218 of each BCell integrates, spatially and temporally, the local visual information of the
 219 photo-receptors which are connected to it, with a lateral modulation due to the HCells.
 220 Each BCells is thus sensitive to specific local characteristics of the visual scene, defining
 221 its Receptive Field (RF). Thus, BCells, like GCells, have a receptive field. But, as they are
 222 earlier in the vertical pathway they integrate less features.

We label BCells (layer 1) with the index $i = 1, \dots, N_B$ and we model the RF of BCells
 by a convolution kernel, \mathcal{K}_{B_i} , such that the voltage of BCell i is stimulus-driven by the
 term:

$$V_{i_{drive}}(t) = \left[\mathcal{K}_{B_i} \overset{x,y,t}{*} \mathcal{S} \right](t). \quad (2)$$

223 The center of the RF, located at x_i, y_i , also corresponds to the coordinates of the BCell i .
 224 A typical shape for the RF of BCells is illustrated in Fig. 2, although the explicit form
 225 does not play a role in the subsequent developments.

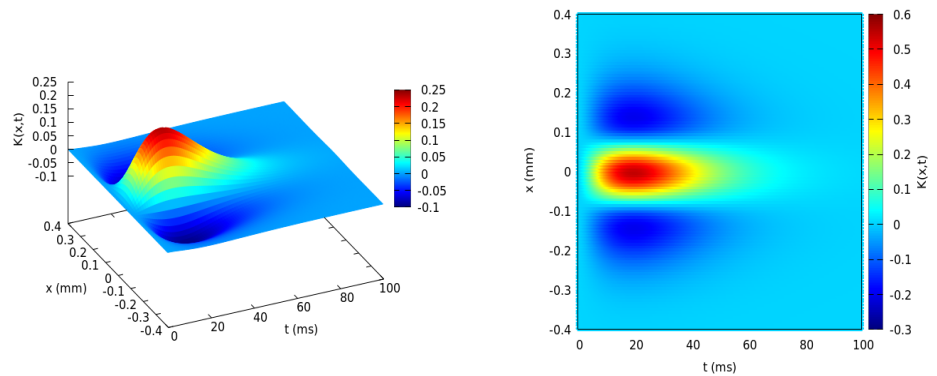


Figure 2. Receptive Field of a ON BCell. **Left.** Example of a spatio-temporal RF of BCells (ON center cell) represented in 3D (one dimension of space, x and time t). There is inhibition at the surround, physiologically due to HCells. **Right.** Spatio-temporal RF representation with a color map.

226 The second block, that we name in short IPL (Inner Plexiform Layer), comprises the
 227 ACells and GCells and the afferent synapses. Its "input" is the output of the OPL, and its
 228 output, the trains of action potentials emitted by the GCells. ACells are difficult to study
 229 experimentally because hardly accessible from electrophysiology measurements. There
 230 are also a large number of cell subtypes in the ACells class (around 40), of which only a
 231 small number have duly identified functions. It is however recognized that they play
 232 an essential role in the treatment of motion [17,44,45]. Here we address mathematically
 233 the question of the RGCell receptive field form, resulting from the pooling of BCells, as
 234 illustrated in Fig. 1, each with a specific RF as exemplified in Fig. 2, and modulated by
 235 ACells lateral connectivity.

¹ Note that the terminology OPL and IPL refers actually to synaptic layers. "The outer plexiform layer has a wide external band composed of inner fibers of rods and cones and a narrower inner band consisting of synapses between photoreceptor cells and cells from the inner nuclear layer." "The inner plexiform layer consists of synaptic connections between the axons of bipolar cells and dendrites of ganglion cells" (ref <https://www.sciencedirect.com/topics/medicine-and-dentistry/>). In our model, these naming are short cuts to distinguish the network input (OPL) and the network processing (IPL).

236 1.1.3. BCells-ACells interactions

237 We label ACells (second layer) with the index $j = 1 \dots N_A$. We note $W_{B_i}^{A_j}$ the
 238 synaptic weight from ACell j to BCell i and $W_{A_j}^{B_i}$ the synaptic weight from BCell i to
 239 ACell j . We set $W_{B_i}^{A_j} \leq 0$, (ACells are in general inhibitory although some excitatory
 240 ACells exist, not considered here) whereas $W_{A_j}^{B_i} \geq 0$. The synaptic weight matrices BCells
 241 to ACells and ACells to BCells are noted $W_{A_j}^B, W_B^A$. They are not squared in general.
 242 There also exist electric synapses (gap junctions) between BCells and ACells (e.g. in the
 243 Rod Cone pathway [17]) but we will not consider them, for simplicity. Note however
 244 that adding gap junctions would simply result in adding linear terms to equations (3),
 245 (6), (9) (when considering passive gap junctions) and modify characteristic time scales,
 246 without changing the global analysis.

The voltage of BCell i , V_{B_i} , evolves according to:

$$\frac{dV_{B_i}}{dt} = -\frac{1}{\tau_{B_i}} V_{B_i} + \sum_{j=1}^{N_A} W_{B_i}^{A_j} \mathcal{N}_A(V_{A_j}) + F_{B_i}(t). \quad (3)$$

Here, τ_{B_i} is the characteristic time scale of BCell i response (in ms). The function:

$$\mathcal{N}_A(V) = \begin{cases} V - \theta_A, & \text{if } V_A > \theta_A; \\ 0, & \text{otherwise} \end{cases}, \quad (4)$$

247 is a linear rectifier ensuring that the synapse $j \rightarrow i$ becomes silent when the voltage
 248 of the pre-synaptic ACell j , V_{A_j} , is lower than a threshold θ_A . This corresponds to a
 249 biophysical fact : a synapse cannot change its sign. For simplicity we consider θ_A to
 250 be the same for all ACells, although the present formalism can be extended, e.g., to
 251 several families of ACells having different thresholds. Note that linear rectifiers of type
 252 (4) rectify cell's voltage "from below". Rectification "from above" also exist, ensuring that
 253 the cell's voltage does not increase without bounds. A typical mechanism is gain control,
 254 where an additional variable, called the activity, increasing as voltage increases, triggers
 255 a gain function non linearly dropping down the voltage when it exceeds an upper
 256 threshold [25,27]. Under some mild assumptions gain control can also be implemented
 257 as a piecewise linear function of the activity (see [28]).

Finally, $F_{B_i}(t)$ is the OPL input term. To match classical retina models as developed
 e.g. in [25,27] it reads:

$$F_{B_i}(t) = \frac{V_{i_{drive}}}{\tau_B} + \frac{dV_{i_{drive}}}{dt} = \left[\mathcal{K}_{B_i}^{x,y,t} * \left(\frac{\mathcal{S}}{\tau_B} + \frac{d\mathcal{S}}{dt} \right) \right](t), \quad (5)$$

258 (where $\mathcal{K}_{B_i}(x, y, 0) = 0$). In short, $F_{B_i}(t)$ is chosen so that, in the absence of ACells
 259 interaction, $V_{B_i}(t) = V_{i_{drive}}(t)$. Note that $F_{B_i}(t)$ implements therefore a time derivative of
 260 the drive, which makes, e.g. BCells response to moving objects sensitive to changes in
 261 directions or speed.

262

ACells are connected to BCells with chemical synapses. The differential equation
 obeyed by the voltage of ACell j is:

$$\frac{dV_{A_j}}{dt} = -\frac{1}{\tau_{A_j}} V_{A_j} + \sum_{i=1}^{N_B} W_{A_j}^{B_i} \mathcal{N}_B(V_{B_i}), \quad (6)$$

263 where τ_{A_j} is the characteristic time scale of ACell j response, and \mathcal{N}_B has the same form
 264 as (4), with a threshold θ_B . Note that, in contrast to BCells, ACells do not receive an OPL
 265 input.

266 1.1.4. GCells.

We label GCells (third layer) with the index $k = 1 \dots N_G$. They are connected to BCells with excitatory synaptic weights, $W_{G_k}^{B_i} \geq 0$ (e.g. glutamatergic synapses) and to ACells with inhibitory synaptic weights, $W_{G_k}^{A_j} \leq 0$ (e.g. glycinergic or GABA-ergic synapses). Their voltage, V_{G_k} , evolves according to:

$$\frac{dV_{G_k}}{dt} = -\frac{1}{\tau_G} V_{G_k} + \sum_{i=1}^{N_B} W_{G_k}^{B_i} \mathcal{N}_B(V_{B_i}) + \sum_{j=1}^{N_A} W_{G_k}^{A_j} \mathcal{N}_A(V_{A_j}). \quad (7)$$

GCells are spiking. In the model their spiking activity (firing rate) is defined by a LNP model [21–23]. It depends on the voltage via a non linear function $\mathcal{N}_G(V_G) \equiv f\left(\frac{V_G(t) - \theta_G}{\sigma_G}\right)$, where f is typically a sigmoid. Although the detailed form of f does not matter here, it will be convenient, in the sequel, to consider:

$$\mathcal{N}_G(V_G) = \frac{A}{\sqrt{2\pi}} \int_{-\infty}^{\frac{V_G - \theta_G}{\sigma_G}} e^{-\frac{x^2}{2}} dx. \quad (8)$$

267 The parameters θ_G (spiking threshold) and σ_G (controlling the slope of the sigmoid at
268 $V_G = \theta_G$) corresponds, in the case where \mathcal{N}_G has the form (8), to the probability that a
269 Gaussian centered Ornstein-Uhlenbeck processes with mean-square deviation σ_G crosses
270 the threshold θ_G .

271 1.1.5. Joint dynamics

The joint dynamics of all cells voltage is given by the dynamical system:

$$\begin{cases} \frac{dV_{B_i}}{dt} = -\frac{1}{\tau_B} V_{B_i} + \sum_{j=1}^{N_A} W_{B_i}^{A_j} \mathcal{N}_A(V_{A_j}) + F_{B_i}(t), & i = 1 \dots N_B; \\ \frac{dV_{A_j}}{dt} = -\frac{1}{\tau_A} V_{A_j} + \sum_{i=1}^{N_B} W_{A_j}^{B_i} \mathcal{N}_B(V_{B_i}), & j = 1 \dots N_A; \\ \frac{dV_{G_k}}{dt} = -\frac{1}{\tau_G} V_{G_k} + \sum_{i=1}^{N_B} W_{G_k}^{B_i} \mathcal{N}_B(V_{B_i}) + \sum_{j=1}^{N_A} W_{G_k}^{A_j} \mathcal{N}_A(V_{A_j}), & k = 1 \dots N_G; \end{cases} \quad (9)$$

272 whereas GCells spikes are produced by the LNP mechanism described above.

273 The system of eq. (9) can be summarized as follows (Fig. 1, right). BCells re-
274 ceive the visual input via the term $F_{B_i}(t)$ which depends on the stimulus and on the
275 Bcell's receptive field. They are inhibited by ACells via the synaptic weights $W_{B_i}^{A_j} < 0$.
276 ACells are excited by BCells via the synaptic weights $W_{A_j}^{B_i} > 0$. BCells are connected to
277 GCells via the synaptic weights $W_{G_k}^{B_i} > 0$. ACells are connected to GCells via the synap-
278 tic weights $W_{G_k}^{A_j} < 0$. Note that we do not impose any constraint on the connectivity here.
279

To study mathematically the dynamical system (9) we write it in a more convenient form. We use Greek indices $\alpha, \beta, \gamma = 1 \dots N \equiv N_A + N_B + N_G$, and define the state vector $\vec{\mathcal{X}}$, with entries:

$$\mathcal{X}_\alpha = \begin{cases} V_{B_i}, & \alpha = i, & i = 1 \dots N_B; \\ V_{A_j}, & \alpha = N_B + j, & j = 1 \dots N_A; \\ V_{G_k}, & \alpha = N_B + N_A + k, & k = 1 \dots N_G. \end{cases} \quad (10)$$

We introduce $\vec{\mathcal{F}}(t)$, the non stationary input, with entries:

$$\mathcal{F}_\alpha(t) = \begin{cases} F_{B_i}(t), & \alpha = i, & i = 1 \dots N_B; \\ 0, & \alpha > N_B; \end{cases}$$

and $\vec{\mathcal{R}}(\vec{\mathcal{X}})$, the rectification term, with entries:

$$\mathcal{R}_\alpha(\vec{\mathcal{X}}) = \begin{cases} \mathcal{N}_B(V_{B_i}), & \alpha = i, & i = 1 \dots N_B; \\ \mathcal{N}_A(V_{A_j}), & \alpha = N_B + j, & j = 1 \dots N_A; \\ 0, & \alpha = N_B + N_A + k, & k = 1 \dots N_G. \end{cases}$$

280 \mathcal{R}_α can be extended to include gain control; it becomes then a function of voltage and
281 cell's activity [25,28].

We use the notation $0_{n_1 n_2}$ for the $n_1 \times n_2$ matrix with zero entries. We introduce the $N \times N$ matrices:

$$\mathcal{T} = \begin{pmatrix} -\text{diag}[\tau_{B_i}]_{i=1 \dots N_B} & 0_{N_B N_A} & 0_{N_B N_G} \\ 0_{N_A N_B} & -\text{diag}[\tau_{A_j}]_{j=1 \dots N_A} & 0_{N_A N_G} \\ 0_{N_G N_B} & 0_{N_G N_A} & -\text{diag}[\tau_{G_k}]_{k=1 \dots N_G} \end{pmatrix}, \quad (11)$$

characterizing the characteristic integration times of cells,

$$\mathcal{W} = \begin{pmatrix} 0_{N_B N_B} & W_B^A & 0_{N_B N_G} \\ W_A^B & 0_{N_A N_A} & 0_{N_A N_G} \\ W_G^B & W_G^A & 0_{N_G N_G} \end{pmatrix}, \quad (12)$$

282 summarizing chemical synapses interactions. Note that, to our best knowledge, there
283 are no synapse from GCells to GCells, but they could be added in this formalism.

Then, the dynamical system (9) reads, in vector form:

$$\frac{d\vec{\mathcal{X}}}{dt} = \mathcal{T}^{-1} \cdot \vec{\mathcal{X}} + \mathcal{W} \cdot \vec{\mathcal{R}}(\vec{\mathcal{X}}) + \vec{\mathcal{F}}(t). \quad (13)$$

284 We remark that (13) has a specific skew-product structure: the dynamics of GCells is
285 driven by BCells and ACells with no feedback. This means that one can study first the
286 coupled dynamics of BCells and ACells and then the effect on GCells. This corresponds
287 to a biological reality as, to our best knowledge, there is no feedback from GCells to
288 BCells or to ACells.

289 1.1.6. Piecewise linear evolution

290 We assume here that $\mathcal{F}_\alpha(t)$ is bounded, as well as synaptic weights. Thus, the phase
291 space Ω of (13) can be taken compact. Indeed, trajectories cannot escape to infinity thanks
292 to the rectification terms $\mathcal{N}_B, \mathcal{N}_A$, (eq. (4)) and thanks to the sign of synaptic weights
293 $W_{A_j}^{B_i}, W_{B_i}^{A_j}$. More precisely, V_{B_i} cannot become arbitrary large and positive because the
294 input term $\mathcal{F}_\alpha(t) \equiv F_{B_i}(t)$ is bounded and because $\sum_{j=1}^{N_A} W_{B_i}^{A_j} \mathcal{N}_A(V_{A_j}) \leq 0$. Assume
295 indeed that V_{B_i} increases (due to a large enough $\mathcal{F}_\alpha > 0$ making the r.h.s. of eq. (3)
296 positive). This leads to an increase of connected ACells voltages V_{A_j} (eq. (6)), thus to
297 a decrease of the term $\sum_{j=1}^{N_A} W_{B_i}^{A_j} \mathcal{N}_A(V_{A_j}) \leq 0$ until the point where the r.h.s. of (3)
298 becomes negative, thereby decreasing V_{B_i} and preventing it from becoming arbitrary
299 large. This, implies as well that V_{A_j} s cannot become arbitrary large. On the opposite, if
300 V_{B_i} (resp. V_{A_j}) becomes smaller than θ_B (resp. θ_A) it does not play any more role in the
301 dynamics because of rectification.

302 Due to the specific form (4) of the rectification terms, the dynamical system (13) is
303 piecewise linear. More precisely, we can partition the phase space Ω into sub domains
304 $\Omega^{(n)}$, $n = 1 \dots 2^{N_B + N_A}$ defined as follow. To each cell $\alpha = 1 \dots N_B + N_A$ (BCell or ACell)
305 we associate a "rectification label" $\eta_\alpha = 1$ if the cell α is rectified and $\eta_\alpha = 0$ otherwise.
306 Because of the form (4) of the rectification, the label η_α corresponds to a partition of
307 the voltage X_α 's domain of variation into two sub domains (e.g., for a BCell, $\eta_\alpha = 1$ if

308 $V_{B_i} < \theta_B$ and $\eta_\alpha = 0$ if $V_{B_i} \geq \theta_B$). One can actually generalize this definition to gain
 309 control by extending the phase space (adding activity variables) and approximating the
 310 gain control function by a step function [28]. Now, the set $\{0, 1\}^{N_B+N_A}$ is made of chains
 311 $\eta = (\eta_1 \dots \eta_{N_B+N_A})$ composed of the rectification labels η_α of all BCells and ACells. To
 312 each such sequence is therefore associated a convex domain $\Gamma^{(n)}$ of $\mathbb{R}^{N_B+N_A}$ where all
 313 cells α such that $\eta_\alpha = 0$ have their voltage X_α larger than the rectification threshold, thus,
 314 are not rectified, and all cells such that $\eta_\alpha = 1$ are rectified. To each such η is associated
 315 a unique integer (e.g. $n = \sum_{\alpha=1}^{N_B+N_A} \eta_\alpha 2^{\alpha-1}$, η is then the binary coding of n). Finally,
 316 we set $\Omega^{(n)} = \Gamma^{(n)} \times \mathbb{R}^{N_G}$, where the product with the subspace \mathbb{R}^{N_G} integrates the
 317 states space of GCells dynamics. They are slaved by BCells and ACells dynamics, but
 318 they are not rectified. In this setting, $\Omega^{(0)}$ is the subset of Ω such that neither BCells nor
 319 ACells are rectified; $\Omega^{(1)}$ the subset of the phase space where only BCell 1 is rectified;
 320 $\Omega^{(3)}$ the subset where only BCells 1, 2 are rectified; $\Omega^{(2^{N_B})}$ the subset where only ACell 1
 321 is rectified and so on.

It is easy to check that the sets $\Omega^{(n)}$ are disjoint and cover \mathbb{R}^N , thus, make a partition
 of the phase space. The vector $\vec{\mathcal{R}}(\vec{\mathcal{X}})$ has now the form:

$$\mathcal{R}_\alpha(\vec{\mathcal{X}}) = \begin{cases} (1 - \eta_\alpha) (X_\alpha - \theta_B), & \alpha = 1 \dots N_B; \\ (1 - \eta_\alpha) (X_\alpha - \theta_A), & \alpha = N_B \dots N_B + N_A; \\ 0, & \alpha = N_B + N_A + k, \quad k = 1 \dots N_G, \end{cases}$$

and is piecewise-linear in $\vec{\mathcal{X}}$. For $\vec{\mathcal{X}} \in \Omega^{(n)}$, the transformation $\mathcal{T}^{-1} \cdot \vec{\mathcal{X}} + \mathcal{W} \cdot \vec{\mathcal{R}}(\vec{\mathcal{X}})$ can
 therefore be written $\mathcal{L}^{(n)} \cdot \vec{\mathcal{X}} + \vec{\mathcal{C}}^{(n)}$, where $\vec{\mathcal{C}}^{(n)}$ is the vector with entries:

$$\vec{\mathcal{C}}^{(n)} = \begin{cases} -\theta_B (1 - \eta_\alpha), & \alpha = 1 \dots N_B; \\ -\theta_A (1 - \eta_\alpha), & \alpha = N_B \dots N_B + N_A; \\ 0, & \alpha = N_B + N_A + 1, \dots, N. \end{cases} \quad (14)$$

The matrix:

$$\mathcal{L}^{(n)} = \begin{pmatrix} -\text{diag} \left[\frac{1}{\tau_{B_i}} \right]_{i=1 \dots N_B} & W_B^A \cdot D_A^{(n)} & & 0_{N_B N_G} \\ & W_A^B \cdot D_B^{(n)} & -\text{diag} \left[\frac{1}{\tau_{A_j}} \right]_{j=1 \dots N_A} & 0_{N_A N_G} \\ & W_G^B \cdot D_B^{(n)} & W_G^A \cdot D_A^{(n)} & -\text{diag} \left[\frac{1}{\tau_{G_k}} \right]_{k=1 \dots N_G} \end{pmatrix}, \quad (15)$$

322 is called the *transport operator in the domain* $\Omega^{(n)}$. This terminology is further explained
 323 in section 1.1.8, but, in short, $\mathcal{L}^{(n)}$ acts as a flow (or a propagator) characterizing the evo-
 324 lution of a trajectory within $\Omega^{(n)}$. In eq. (15), the matrices $D_B^{(n)} = \text{diag}[1 - \eta_\alpha]_{\alpha=1 \dots N_B}$,
 325 $D_A^{(n)} = \text{diag}[1 - \eta_\alpha]_{\alpha=N_B+1 \dots N_B+N_A}$ are projecting onto the subspace of non rectified
 326 cells in the domain $\Omega^{(n)}$. In other words, when the state $\vec{\mathcal{X}}$ is $\Omega^{(n)}$, a rectified cell α gives
 327 a zero contribution to the dynamics of other cells, which corresponds to have a row and
 328 column α made of zeros in $D_A^{(n)}, D_B^{(n)}$.

The dynamical system (13) reads now:

$$\frac{d\vec{\mathcal{X}}}{dt} = \mathcal{L}^{(n)} \cdot \vec{\mathcal{X}} + \vec{\mathcal{F}}^{(n)}(t), \quad \vec{\mathcal{X}} \in \Omega^{(n)}, \quad (16)$$

329 where we wrote $\vec{\mathcal{F}}^{(n)}(t) = \vec{\mathcal{C}}^{(n)} + \vec{\mathcal{F}}(t)$. Thanks to the decomposition of the phase
 330 space into convex sub-domains $\Omega^{(n)}$, (16) is now linear. This technique of phase space
 331 decomposition is classical and has been used in domains such as ergodic theory and
 332 billiards, self-organized criticality [46,47] or neurosciences [29,30,32,48]. See especially
 333 the recent paper by A. Rajakumar et al [49] very much in the spirit of the present model.

334 1.1.7. Spectra and fixed points

335 It is important to consider in detail the spectrum² of $\mathcal{L}^{(n)}$ for further studies. We
 336 note $\lambda_\beta^{(n)}, \beta = 1 \dots N$, the eigenvalues of $\mathcal{L}^{(n)}$ and its right eigenvectors are noted, $\mathcal{P}_\beta^{(n)}$.
 337 These vectors are the columns of the matrix $\mathcal{P}^{(n)}$ transforming $\mathcal{L}^{(n)}$ in diagonal form
 338 (assuming it is diagonalizable). $(\mathcal{P}^{(n)})^{-1}$ is the inverse matrix. Its rows are the left
 339 eigenvectors of $\mathcal{L}^{(n)}$.

340 As $D_A^{(n)}, D_B^{(n)}$ are projections matrices, it is easy to see, from the form (15), that a
 341 rectified cell generates an eigenvalue $-\frac{1}{\tau_\alpha}$ and an eigenvector \vec{e}_α , the canonical basis
 342 vector of \mathbb{R}^N in the direction α . The non rectified cells span a subspace of \mathbb{R}^N and
 343 the projection of $\mathcal{L}^{(n)}$ on this subspace has a spectrum depending on the connectivity
 344 matrices W_A^B, W_B^A and on other parameters like characteristic times.

345 The corresponding eigenvalues $\lambda_\beta^{(n)}, \beta = 1 \dots N$ can be real or complex, with a
 346 positive or a negative real part. In the case where W_A^B and W_B^A commute it is actually
 347 possible to explicitly compute the eigenvalues and the eigenvectors and obtain conditions
 348 for stability (all eigenvalues with real negative part) and real/complex eigenvalues [28].
 349 If we further assume that W_A^B and W_B^A have no zero eigenvalues, the sign constraints
 350 on these matrices implies that $\mathcal{L}^{(n)}$ is invertible for all n . This is what we are going to
 351 assume from now.

352 It follows that, in the absence of external stimulus ($\vec{\mathcal{F}}(t) = \vec{0}$), equation (16) has,
 353 for each n , a unique fixed point $\vec{\mathcal{X}}^{(n)} = -(\mathcal{L}^{(n)})^{-1} \cdot \vec{\mathcal{C}}^{(n)}$. Note, however, that this point
 354 *may not be* in $\Omega^{(n)}$. This is a typical situation for piecewise linear dynamical systems
 355 (like Iterated Function Systems [52–54]) where dynamics can have complex attractors
 356 even if maps are linear into sub-domains of the phase space. The simplest non trivial
 357 case is when dynamics generates a periodic orbit, but more complex attractors (fractal
 358 sets) can be obtained. Here, it is reasonable to assume, at least, that cells at rest are not
 359 rectified. Mathematically, this means that the fixed point of $\mathcal{L}^{(0)}$, $\vec{\mathcal{X}}^* = -(\mathcal{L}^{(0)})^{-1} \cdot \vec{\mathcal{C}}^{(0)}$,
 360 belongs to $\Omega^{(0)}$ and this is what we are going to assume for now. This imposes a set of
 361 constraints linking synaptic weights and thresholds. A simple assumption consists of
 362 having vanishing thresholds $\theta_A = \theta_B = 0$, in which case the rest state is $\vec{0}$. We will also
 363 assume that $\vec{\mathcal{X}}^*$ is stable (eigenvalues of $\mathcal{L}^{(0)}$ have a negative real part), which imposes
 364 additional assumptions on synaptic weights and cells integration times. On biophysical
 365 grounds it means that the rest state is stable to small perturbations, like noise. Because
 366 rectified cells produce stable eigenvalues the following holds. Taking an initial condition
 367 in any domain $\Omega^{(n)}$ spontaneous dynamics (without stimulus) eventually drives the
 368 trajectory back to $\Omega^{(0)}$ and, then, to the rest state. This is further commented below
 369 (section 1.1.8, remark 2).

370 1.1.8. Solutions

We now consider the general situation where dynamics is in the rest state at times
 $t < 0$, and, from time $t = 0$ on, the stimulus $\mathcal{S}(t)$ is applied, resulting in a non stationary
 drive $\vec{\mathcal{F}}(t)$. In general, the stimulus is applied over a finite time. After this the system
 eventually returns to rest. Under this stimulation the trajectory $\left\{ \vec{\mathcal{X}}(t) \right\}_{t \geq 0}$ is going to
 cross a sequence of domains $\Omega^{(n_k)}, k = 1, \dots$, with $n_1 = 0$, entirely determined by the
 stimulus and the network characteristics. Call $t_-^{(n_{k+1})}$ the time where the trajectory enters

² Another approach consists of considering the Schur decomposition instead of the diagonalisation [49–51].

the domain $\Omega^{(n_{k+1})}$ and $t_+^{(n_{k+1})}$ the time where it gets out. Note that $t_-^{(n_{k+1})} = t_+^{(n_k)}$. By direct integration of eq. (16), we have:

$$\vec{\mathcal{X}}(t) = e^{\mathcal{L}^{(n_{k+1})}(t-t_-^{(n_{k+1})})} \cdot \vec{\mathcal{X}}(t_-^{(n_{k+1})}) + \int_{t_-^{(n_{k+1})}}^t e^{\mathcal{L}^{(n_{k+1})}(t-s)} \cdot \vec{\mathcal{F}}^{(n_{k+1})}(s) ds, \quad t \in [t_-^{(n_{k+1})}, t_+^{(n_{k+1})}], \quad (17)$$

where $\vec{\mathcal{X}}(t_-^{(n_{k+1})})$, corresponding to the state of $\vec{\mathcal{X}}$ when entering $\Omega^{(n_{k+1})}$, is given by the integration of the past trajectory and can be computed explicitly. This is:

$$\vec{\mathcal{X}}(t_-^{(n_{k+1})}) = \vec{\mathcal{X}}(t_+^{(n_k)}) = \sum_{m=0}^k \mathcal{H}_m^k \vec{\Phi}_m, \quad (18)$$

where \mathcal{H}_m^k is a sequence of matrices satisfying:

$$\mathcal{H}_k^k = \mathcal{I}_N; \quad \mathcal{H}_m^k = \mathcal{H}_{k-1}^k \mathcal{H}_m^{k-1}; \quad \mathcal{H}_{k-1}^k = e^{\mathcal{L}^{(n_k)}(t_+^{(n_k)} - t_+^{(n_{k-1})})}, \quad (19)$$

where \mathcal{I}_N is the identity matrix of dimension N . The matrix \mathcal{H}_m^k transports the flow from the exit point of $\Omega^{(n_m)}$ to the exit point of $\Omega^{(n_k)}$. The vectors $\vec{\Phi}_m$ are defined by:

$$\vec{\Phi}_0 = \vec{\mathcal{X}}(0); \quad \vec{\Phi}_m = \int_{t_-^{(n_m)}}^{t_+^{(n_m)}} e^{\mathcal{L}^{(m)}(t_+^{(n_m)} - s)} \cdot \vec{\mathcal{F}}^{(m)}(s) ds. \quad (20)$$

371 The proof of (18) is easily done by recurrence.

372 1.1.9. Remarks

373 Let us now make some remarks on the structure of these solutions.

- 374 1. In the definition of \mathcal{H}_m^k the operators \mathcal{L}_{n_k} do not commute in general.
- 375 2. Eigenvalues of some \mathcal{H}_m^k can have a positive real part leading to exponential
376 increase along the corresponding eigen-direction. This means that some cells
377 voltage increases exponentially in absolute value. However, when voltages become
378 too large, voltage rectification takes place, corresponding to the trajectory entering
379 a new continuity domain. Here, unstable cells do not contribute any more to
380 dynamics which is projected on the subspace of non rectified cells. This has
381 the effect of transforming unstable eigenvalues into stable ones preventing the
382 trajectories $\vec{\mathcal{X}}(t)$ to diverge. Actually, the spectrum of \mathcal{H}_m^k , controlling stability,
383 resembles the Lyapunov spectrum in ergodic theory [55], with two main differences.
384 First, we are simply considering product of matrices without multiplying by the
385 adjoint so that eigenvalues can be complex. Second, we are not assuming stationarity
386 and the existence of an invariant measure. Instead, the product \mathcal{H}_m^k is constrained
387 by the non stationary stimulus and dynamical system parameters which fixes the
388 sequence of times n_k s.
- 389 3. Rectification induces a weak form of non linearity where e.g. the contraction/expansion
390 in the phase space depends on the domain $\Omega^{(n_k)}$ (whereas in a differentiable non
391 linear system it would depend on the point in the phase space). This has deep
392 consequences on cells response, as commented in the results sections.

393 1.2. Spike statistics

394 As pointed out in the introduction, it might be helpful to propose a mathematical
395 setting taking into account non stationarity and potentially long memory in spike trains
396 probabilities. Such a setting exists since long but has not been applied to spike train
397 statistics until recently. It is inherited from statistical physics on one hand [56] and
398 on extensions of Markov chains to unbounded memory on the other hand [57]. The
399 material briefly sketched here has been published in [29–39].

400 1.2.1. Mathematical setting for spike trains

401 Neurons variables such as membrane potential or ionic currents are described
 402 by continuous-time equations. In contrast, spikes resulting from the experimental
 403 observation are discrete events, binned with a certain time resolution δ , say of order
 404 a millisecond. We consider a network of N spiking neurons, labelled with an index
 405 $k = 1 \dots N$. We define a spike variable $\omega_k(n) = 1$ if neuron k has emitted a spike in the
 406 time interval $[n\delta, (n+1)\delta[$, and $\omega_k(n) = 0$ otherwise. We denote by $\omega(n) = [\omega_k(n)]_{k=1}^N$
 407 the spike-state of the entire network at time n , which we call a *spiking pattern*. A *spike*
 408 *block* denoted by ω_m^n , $n \geq m$, is the sequence of spike patterns $\omega(m), \omega(m+1) \dots \omega(n)$.
 409 The range of a block ω_m^n is $n - m + 1$, the number of time steps from m to n . We call
 410 a spike train an infinite sequence of spikes both in the past and in the future, and, to
 411 simplify notations we note a spike train ω (instead of $\omega_{-\infty}^{+\infty}$). Of course, on operational
 412 grounds spike trains are finite, but it is mathematically more convenient to work on a
 413 space of bi-infinite spike sequences.

414 1.2.2. Mathematical setting for spiking probabilities

415 We now consider a family of transition probabilities of the form $\mathbb{P}_n \left[\omega(n) \mid \omega_{-\infty}^{n-1} \right]$,
 416 which represent the probability, that at time n , one observes the spiking pattern $\omega(n)$
 417 given the network spike history, extending to an infinite past. This is an extension of
 418 Markov chains where probabilities have the form $\mathbb{P}_n \left[\omega(n) \mid \omega_{n-D}^n \right]$, where D is the
 419 memory depth of the Markov chain. Letting the memory be possibly infinite corresponds
 420 to situation where one cannot precisely fix the memory depth necessary to characterize
 421 the probability of a spike pattern given the past spike history. An example of a model
 422 requiring this context is presented in section 1.2.3 below. Having infinite memory
 423 imposes mathematical constraints on the memory decay that has to be sufficiently fast
 424 (typically, exponential) so that the situation is close to Markov chains. In addition to the
 425 model presented below, neural models with infinite memories have been considered by
 426 several authors such as E. Loecherbach and A. Galves [58]. A few remarks about this
 427 form of probability:

- 428 1. We do not assume stationarity. \mathbb{P}_n may depend explicitly on time. This is actually
 429 the reason why we have an index n . A time translation invariant probability will
 430 simply be written \mathbb{P} .
- 431 2. For such probabilities to be well defined and useful, one need to make assumptions
 432 on their structure. Beyond technical assumptions such as measurability, summabil-
 433 ity, non nullness and continuity [59,60], the most important assumption here is that
 434 the dependence in the past (memory) decays fast enough, typically, exponentially,
 435 so that, even if this chain has infinite memory it is very close to Markov.
- 436 3. As one can associate to Markov chains an equilibrium probability (under condi-
 437 tions actually quite more general than detailed-balance) the system of transition
 438 probabilities $\{\mathbb{P}_n\}_{n \in \mathbb{Z}}$ also admits, under the mathematical conditions sketched in
 439 the item 2 above, an equivalent notion called “chains with complete connections”
 440 or a “chain swith unbounded memory” [57].
- 441 4. These distributions are formally (left-sided) Gibbs distributions where the Gibbs po-
 442 tential is $\Phi(n, \omega) \stackrel{\text{def}}{=} \log \mathbb{P}_n \left[\omega(n) \mid \omega_{-\infty}^{n-1} \right]$ (the non-nullness assumption imposes
 443 that $\mathbb{P}_n \left[\omega(n) \mid \omega_{-\infty}^{n-1} \right] > 0$). This establishes a formal link to statistical physics.
 444 In particular, when the chain is stationary, expanding the potential in product
 445 of spikes events up to second order one recovers the maximum entropy models
 446 used in the literature of spike trains analysis, including the so-called Ising model
 447 [35,61–63]. However, the chains we consider are not necessarily stationary.

448 1.2.3. A model of effective interactions between GCells

449 The visual cortex has no clue on which biophysical processes are taking place in the
 450 retina. All the visual information it receives is encoded in spike trains. This leads to the
 451 idea of proposing models of spiking GCells network where dynamics of GCells voltage
 452 is only constrained by GCells spikes history. Here, one assumes that GCells dynamics is
 453 controlled by the interactions with hidden layers, for example, the BCells-ACells layers
 454 in the model (13), in a situation where an observer is just recording the spikes emitted by
 455 GCells, while having no clue of the dynamics in the upper layers. These hidden layers
 456 result in providing, *causal, effective interactions* between GCells that one can interpolate
 457 by fitting the statistics. The idea is then to construct a dynamical model where the
 458 spiking of a GCell depends on the spike history emitted by the network, with virtual
 459 interactions that mimic hidden causal effects [64]. This strategy lead us to propose the
 460 model presented in the next paragraph. The advantage of this approach is that one can
 461 explicitly write the transition probabilities $\mathbb{P}_n \left[\omega(n) \mid \omega_{-\infty}^{n-1} \right] > 0$ and infer, from this,
 462 a linear response formula telling us how statistical quantities such as firing rates, but
 463 also spike correlations are modified by a time dependent stimulus. These results are
 464 presented in the "Results" subsection 1.2.3.

465 The model is inspired from the generalized Integrate and Fire model (gIF) proposed
 466 by Rudolph and Destexhe [65] and generalizes the Leaky-Integrate and Fire model
 467 [66,67]. We have N neurons (say GCells) characterized by their voltage $V_k, k = 1 \dots N$.
 468 One fixes a voltage threshold θ such that, whenever $V_k(t) = \theta$ a spike is emitted by
 469 neuron k at time t , and is reset to a reset value (typically, $V_{reset} = 0$). Below θ , the
 470 dynamics of voltage (sub-threshold dynamics) is governed by eq. (23) below.

In the LIF model, synaptic conductances are constant. In the gIF model, in contrast,
 the synaptic conductance g_{kj} between the pre-synaptic neuron j and the post-synaptic
 neuron k depends on spike history as:

$$g_{kj}(t, \omega) = G_{kj} \alpha_{kj}(t, \omega), \quad (21)$$

where:

$$\alpha_{kj}(t, \omega) = \sum_{n=-\infty}^t \alpha_{kj}(t-n) \omega_j(n). \quad (22)$$

471 The notation $g_{kj}(t, \omega)$ means that function g_{kj} depends on spikes occurring before time
 472 t . $G_{kj} \geq 0$ is the maximal conductance between j and k . It is zero when there is
 473 no synaptic connection between neurons j and k . In (22), the function $\alpha_{kj}(t)$, called
 474 α -kernel, summarizes the complex dynamical process underlying the generation of
 475 a post-synaptic potential after the emission of a pre-synaptic spike [68]. It has the
 476 typical form $\alpha_{kj}(t) = P(t) e^{-\frac{t}{\tau_{kj}}} H(t)$ where $P(t)$ is a polynomial in time and $H(t)$ is the
 477 Heaviside function. What matters on mathematical grounds is the exponential tail of
 478 $\alpha_{kj}(t)$ [33]. The function $\alpha_{kj}(t, \omega)$ depends on the spike history preceding t . It records
 479 the spikes emitted by the pre-synaptic neuron j before t , corresponding to $\omega_j(n) = 1$ and
 480 adds up a contribution $\alpha_{kj}(t-n)$ to the post synaptic conductance from pre-synaptic
 481 neuron j to post-synaptic neuron k .

Now, the gIF dynamics reads [30,33,34]:

$$C_k \frac{dV_k}{dt} + g_L (V_k - E_L) + \sum_j g_{kj}(t, \omega) (V_k - E_j) = S_k(t) + \sigma_B \tilde{\zeta}_k(t), \quad \text{if } V_k(t) < \theta,$$

482 where g_L, E_L are respectively the leak conductance and the leak reversal potential, E_j the
 483 reversal potential characterizing the synaptic transmission between j and k . Finally, $\tilde{\zeta}_k(t)$
 484 is a white noise, introducing stochasticity in dynamics. Its intensity is σ_B .

Setting: $W_{kj} = G_{kj}E_j$, $i_k(t, \omega) = g_L E_L + \sum_j W_{kj} \alpha_{kj}(t, \omega) + S_k(t) + \sigma_B \zeta_k(t)$, $g_k(t, \omega) = g_L + \sum_{j=1}^N g_{kj}(t, \omega)$, one can finally write the gIF dynamics in the form:

$$C_k \frac{dV_k}{dt} + g_k(t, \omega) V_k = i_k(t, \omega), \quad \text{if } V_k(t) < \theta, \quad (23)$$

485 where $i_k(t, \omega)$ depends, on the network spike history via $\alpha_{kj}(t, \omega)$, on the stimulus, and
 486 contains a stochastic term. As the reversal potential E_j can be positive or negative,
 487 the synaptic weights W_{kj} define an oriented and signed graph, whose vertices are the
 488 neurons. These weights are what we call effective interactions.

489

490 What makes the gIF model very rich is that it proposes a biophysically grounded
 491 way to construct a dynamical system where the variables (here, voltages) are constrained
 492 by the only information of spike train history. The price to pay is that dynamics actually
 493 depends on the *whole spike history*, which is potentially infinite. Actually, the gIF model
 494 has an infinite memory. This is essentially because the conductance depends on the
 495 whole history, and, contrarily to voltages is not reset when the neuron fires. Nevertheless,
 496 the exponential decay in the alpha profile actually ensures the existence (and uniqueness)
 497 of transition probabilities of the form $\mathbb{P}_n \left[\omega(n) \mid \omega_{-\infty}^{n-1} \right]$ [34–38].

498 Note that the integration of (23) does not only requires the knowledge of voltages
 499 V_k , stimulus and noise at time t . It requires, in addition, the knowledge of the spike train
 500 ω emitted by the network before t . In this sense, this is not a classical dynamical system.
 501 Nevertheless, eq. (23) can be explicitly integrated [34,39].

502 2. Results

503 2.1. How could lateral ACells connectivity shape the receptive field of a ganglion cell ?

504 The response of a GCell to visual stimuli is shaped by the retina structure depicted
 505 in Fig. 1. Here, with the model introduced in section 1.1, we would like to characterize
 506 the respective effects of the stimulus and of the network connectivity, especially ACells,
 507 and understand under which condition can the conjugated effect of network dynamics
 508 and stimulus be represented by a convolution of the form (1) where the kernel \mathcal{K}_{G_α} is
 509 *intrinsic* to the cell, i.e. does not depend on the stimulus ?

510 2.1.1. Non rectified case

511 The answer is relatively easy when no rectification takes place, i.e. when the
 512 trajectory of (13) stays in the domain $\Omega^{(0)}$ (see section 1.1.6 for the definition). Indeed,
 513 in this case evolution is ruled by equation (17) which holds from the initial time $t = t_0$
 514 where the stimulus starts to be applied, to the current time t . Actually, we can consider
 515 that t_0 starts far in the past and let it tend to $-\infty$. This corresponds to considering that
 516 the stimulus is applied on a time scale quite longer than the characteristic times in the
 517 problem (i.e. the inverse of the real part of eigenvalues). Then eq. (17) reads $\vec{\mathcal{X}}(t) =$
 518 $\int_{-\infty}^t e^{\mathcal{L}^{(0)}(t-s)} \cdot \vec{\mathcal{F}}^{(0)}(s) ds$, which is $\vec{\mathcal{X}}(t) = \left[e^{\mathcal{L}^{(0)} \cdot t} * \vec{\mathcal{F}}_0 \right](t)$. This equation actually makes
 519 sense only if all eigenvalues of $\mathcal{L}^{(0)}$ are stable, as we assumed above. Note also that
 520 $\vec{\mathcal{F}}^{(0)} = \vec{\mathcal{C}}^{(0)} + \vec{\mathcal{F}}$ where $\vec{\mathcal{C}}^{(0)}$ is a constant, depending on thresholds (eq. (14)) and
 521 whose integration in the convolution product gives $-\left(\mathcal{L}^{(0)}\right)^{-1} \cdot \vec{\mathcal{C}}^{(0)} = \vec{\mathcal{X}}^*$, the base line
 522 activity of $\vec{\mathcal{X}}(t)$ without stimulus. We may ignore this constant in the sequel and focus
 523 on the time varying part of the response, $\left[e^{\mathcal{L}^{(0)} \cdot t} * \vec{\mathcal{F}} \right](t)$. As $\vec{\mathcal{F}}$ is itself defined in terms
 524 of a convolution (eq. (5)) with the stimulus and its derivative, $\vec{\mathcal{X}}(t)$ is a convolution with
 525 the stimulus and its derivative. Here, it is useful to express $\vec{\mathcal{X}}(t)$ in components.

One can then show that [28]:

$$\mathcal{X}_\alpha(t) = V_{\alpha_{drive}}(t) + \mathcal{E}_{\alpha_{net}}^{(0)}(t), \quad \alpha = 1 \dots N, \quad (24)$$

where:

$$\mathcal{E}_{\alpha_{net}}^{(0)}(t) = \sum_{\beta=1}^N \sum_{\gamma=1}^{N_B} \mathcal{P}_{\alpha\beta}^{(0)} \left(\mathcal{P}_{\beta\gamma}^{(0)} \right)^{-1} \omega_{\beta\gamma}^{(0)} \int_{-\infty}^t e^{\lambda_{\beta}^{(0)}(t-s)} V_{\gamma_{drive}}(s) ds, \quad (25)$$

532 where $\omega_{\beta\gamma}^{(0)} = \lambda_{\beta}^{(0)} + \frac{1}{\tau_{B\gamma}}$. The term $V_{\alpha_{drive}}(t)$ in eq. (24) is the stimulus drive and acts
 537 only on BCells (it vanishes for $\alpha > N_B$). The term (25) contains the network effects.
 538 The drive imposed on BCells impacts ACells via the connectivity and, thereby, have a
 539 feedback effect on BCells. In addition, the join activity of BCells and ACells drive the
 540 GCells response ($\alpha > N_B + N_A$). In particular, this equation allows to compute explicitly
 541 the RF of a GCell.

For this, we introduce the function $e_{\beta}^{(0)}(t) \equiv e^{\lambda_{\beta}^{(0)} t} H(t)$ so that $\int_{-\infty}^t e^{\lambda_{\beta}^{(0)}(t-s)} V_{\gamma_{drive}}(s) ds \equiv$
 $\left[e_{\beta}^{(0)} \overset{t}{*} V_{\gamma_{drive}} \right](t)$, which according to (2) is $\left[e_{\beta}^{(0)} \overset{t}{*} \mathcal{K}_{B\gamma} \overset{x,y,t}{*} \mathcal{S} \right](t)$. Thus, by identifica-
 tion with (1), the kernel of GCell $\alpha = N_B + N_A + 1 \dots N_G$ is:

$$\mathcal{K}_{G_\alpha}(x, y, t) = \sum_{\beta=1}^N \sum_{\gamma=1}^{N_B} \mathcal{P}_{\alpha\beta}^{(0)} \left(\mathcal{P}_{\beta\gamma}^{(0)} \right)^{-1} \omega_{\beta\gamma}^{(0)} \left[e_{\beta}^{(0)} \overset{t}{*} \mathcal{K}_{B\gamma} \right]. \quad (26)$$

532 This provides an explicit equation for the kernel of a GCell, embedded in a network of
 533 BCells, ACells, GCells with dynamics (13), when no rectification take place.

534 2.1.2. Consequences

535 **Interpretation.** The kernel obtained in (26) is the response of the GCell to a Dirac
 536 pulse corresponding, in experiments, to a brief light (or dark) full-field flash. It can also
 537 obtained from a white noise stimulus, corresponding, in experiments, to the so-called
 538 Spike Triggered Average (STA) [21–23]. It corresponds therefore to the functional def-
 539 inition of the receptive field of GCells used in experiments. In addition, eq. (24), (25)
 540 give us the voltage of *all* cells in the network at time t under the influence of a stimulus.
 541 Interestingly, thus, these equations allow us to visualize the join evolution of BCells
 542 and ACells as well as their action of GCells. Note that BCells and ACells are difficult to
 543 access experimentally. Given a prescribed connectivity (matrices $W_A^B, W_B^A, W_G^B, W_G^A$), eq.
 544 (24) provides us, therefore, a mathematical insight on the potential, hidden, dynamics
 545 of BCells and ACells leading to the experimentally observed response of GCells. Thus,
 546 this give us possible scenarios characterizing the potential effects of ACells networks on
 547 GCells response. In addition, eq. (26) also provides the RF for BCells ($\alpha = 1 \dots N_B$) and
 548 ACells ($\alpha = N_B + 1 \dots N_B + N_A$). We observe in particular that, in a network, the RF of
 549 a BCell is therefore not only what comes from the OPL - the term $V_{\alpha_{drive}}(t)$ - it integrates
 550 as well lateral ACells connectivity. This is similar to the center-surround shaping of
 551 OPL output due to HCells, but here, we might have different effects, due to the different
 552 physiology of ACells.

553 **Space-time separability.** The GCell kernel, in general, does not factorise into a
 product of a function of space and a function of time (separability). Even in the case
 where the BCells RF is separable, i.e. $\mathcal{K}_{B\gamma}(x, y, t) = \mathcal{K}_{B_{S\gamma}}(x, y) \mathcal{K}_{B_{T\gamma}}(t)$ where $\mathcal{K}_{B_{S\gamma}}$ is the
 spatial part, centred at x_γ, y_γ and $\mathcal{K}_{B_{T\gamma}}$ the temporal part, the GCell kernel reads:

$$\mathcal{K}_{G_\alpha}(x, y, t) = \sum_{\beta=1}^N \mathcal{P}_{\alpha\beta}^{(0)} \left(\sum_{\gamma=1}^{N_B} \left(\mathcal{P}_{\beta\gamma}^{(0)} \right)^{-1} \omega_{\beta\gamma}^{(0)} \left[e_{\beta}^{(0)} \overset{t}{*} \mathcal{K}_{B_{T\gamma}} \right] \times \mathcal{K}_{B_{S\gamma}}(x, y) \right), \quad (27)$$

and is not separable either. Now, if BCells have the same temporal kernel \mathcal{K}_{B_T} , independent of γ and the same characteristic time τ_B , such that $\omega_{\beta\gamma}^{(0)} = \lambda_{\beta}^{(0)} + \frac{1}{\tau_B}$ is independent of γ , we can write :

$$\mathcal{K}_{G_{\alpha}}(x, y, t) = \sum_{\beta=1}^N \mathcal{P}_{\alpha\beta}^{(0)} \omega_{\beta}^{(0)} \left[e_{\beta}^{(0)} \overset{t}{*} \mathcal{K}_{B_T} \right] \left(\sum_{\gamma=1}^{N_B} \left(\mathcal{P}_{\beta\gamma}^{(0)} \right)^{-1} \mathcal{K}_{B_{S_{\gamma}}}(x, y) \right). \quad (28)$$

554 This kernel is not yet strictly separable as the term $\sum_{\gamma=1}^{N_B} \left(\mathcal{P}_{\beta\gamma}^{(0)} \right)^{-1} \mathcal{K}_{B_{S_{\gamma}}}(x, y)$ still de-
 555 pends on β , the eigenmode index, via $\left(\mathcal{P}_{\beta\gamma}^{(0)} \right)^{-1}$. Now, the eigenmodes depends
 556 on connectivity. Especially, the BCell to GCell connectivity corresponds to a pool-
 557 ing of BCells located in the vicinity of GCell α . The simplest case is when there is
 558 no lateral connectivity and where each GCell α is contacted by only BCell with in-
 559 dex γ_{α} (this implies $N_B = N_G$). In this case: $\mathcal{P}_{\alpha\beta}^{(0)} = \delta_{\alpha\beta}$, $\left(\mathcal{P}_{\beta\gamma}^{(0)} \right)^{-1} = \delta_{\beta\gamma}$ so that

560 $\mathcal{K}_{G_{\alpha}}(x, y, t) = \omega_{\alpha}^{(0)} \left[e_{\alpha}^{(0)} \overset{t}{*} \mathcal{K}_{B_T} \right] \mathcal{K}_{B_{S_{\alpha}}}(x, y)$ is separable. More generally, pooling im-

561 plies that $\mathcal{P}_{\alpha\beta}^{(0)}$ and $\left(\mathcal{P}_{\beta\gamma}^{(0)} \right)^{-1}$ are locally spread around α resulting in a spatial part

562 $\sum_{\gamma=1}^{N_B} \left(\mathcal{P}_{\beta\gamma}^{(0)} \right)^{-1} \mathcal{K}_{B_{S_{\gamma}}}(x, y)$ depending only on α .

563

564 **Resonances.** The eigenvalues of $\mathcal{L}^{(0)}$ can be complex, going by conjugated pairs. It
 565 is actually quite easy to obtain such a situation mathematically, even considering nearest
 566 neighbours interactions [28]. A straightforward consequence is the existence of preferred
 567 time frequencies (resonance) for a GCell. In other words, applying periodic sequences
 568 of brief flashes with a varying frequency, one might observe a peak in the amplitude of
 569 the GCell response, for specific frequencies. This remark could, e.g., explain the "bump"
 570 observed in experiments when the retina is submitted to the so-called "Chirp" stimulus
 571 [69], a stimulus composed of different phases of flashes stimulation where one varies
 572 duration, frequency, and amplitude. In the phase where the amplitude is constant but
 573 frequency is varying, some GCells exhibit a resonance like peak (see e.g. Fig. 1 b in [69]).
 574 Of course, such resonances could also be explained by intrinsic cells properties, like
 575 ion channels response. The potential effect of lateral ACells connectivity would have
 576 to be tested experimentally by, e.g. inhibiting ACells synaptic transmission for GCells
 577 exhibiting resonance peaks.

578

579 **Stimulus induced waves.** This is a general fact that networks of coupled units can
 580 produce waves. Spontaneous waves are actually reported in the developmental retina,
 581 induced, in the so-called stage II and stage III by ACells [70]. They are generated by non
 582 linear mechanisms and closeness to bifurcations [71]. This is not the type of wave we
 583 are dealing with here, though. Instead, we are referring to waves triggered by a moving
 584 stimulus, say a moving bar. The idea is that such a stimulus can induce, via ACells con-
 585 nectivity, a wave of connectivity which can be *ahead* of the stimulus, for a certain range
 586 of parameters (e.g. synaptic coupling intensity) compatible with physiology. Stimulus
 587 induced waves, in advance with respect to the stimulus, have been reported in the visual
 588 cortex [72]. They are due to lateral cortical activity and induce cortical anticipation. The
 589 mathematical analysis made in [28] suggests that such anticipatory waves could as well
 590 exist in the retina thanks to ACells lateral connectivity, conjugated with non linear gain
 591 control already known to induce a form of retinal anticipation [25,27].

592

593 **Stimulus adaptation.** Short term plasticity has been reported in the retina at the
 594 synapses between BCells-ACells and ACells-GCells[26,73]. Note actually that, although
 595 most models of plasticity, referring to cortical neurons, are considering spiking neurons
 596 [74], the physiology of short term synapse adaptation does not necessarily require spikes

and is compatible with inner retinal networks dynamics. The effect of synaptic plasticity can be integrated in the model (13). It will result in a variations of eigenvalues and eigenvectors of the transport operator \mathcal{L} with potential changes in dynamics. Although, potential and highly relevant phenomena such as bifurcations induced by plasticity would require considering a non linear version of (13) (at least, rectification to avoid exponential instability), we can ask about simple linear effect of plasticity on the GCells response. A straightforward potential effect could be frequency adaptation to periodic flashes.

2.1.3. Rectification.

Let us now investigate the role of rectification. In the general case, a trajectory crosses several domains, and is characterized by eq. (17). Starting from the domain $\Omega^{(0)}$ (rest state) the state of the network submitted to a stimulus, enters a new domain $\Omega^{(n_1)}$ at time $t_+^{(n_1)}$ where some cells are rectified and so on. Can one still define a response formula of type (1) ? This raises several technical difficulties, first because some eigenvalues can be unstable. As we have seen above, this does not lead to an exponential explosion though precisely rectification prevents cells voltage to diverge. Mathematically, this is expressed by the exit of the trajectory from the domain with positive eigenvalue and a projection on the subspace spanned by non-rectified cells. Another difficulty also comes from the constants $\vec{c}^{(n)}$ defined in (14) coming from the threshold in the rectification function. They can be removed by assuming that all thresholds are equal to 0. This is what we are going to do now for the sake of simplicity. One can then define domain-dependent flows $\Phi^{(n)}(\vec{\mathcal{X}}, t) \equiv e^{\mathcal{L}^{(n)} t} \Theta(\vec{\mathcal{X}}(t) \in \Omega^{(n)})$, where Θ is the indicator function so that $\left[\Phi^{(n)}(\vec{\mathcal{X}}, \cdot) \overset{t}{*} \vec{\mathcal{F}} \right](t) = \sum_{n_m=n} \int_{t_-^{(n_m)}}^{t_+^{(n_m)}} e^{\mathcal{L}^{(m)}(t_+^{(n_m)}-s)} \cdot \vec{\mathcal{F}}(s) ds$ where the sum holds on indices n_m in the trajectory such that $n_m = n$. This allows to express the recurrence formula (18) in terms of a convolution and thereby to express the whole trajectory in terms of a convolution with a transport operator.

However, there are several important differences with the non rectified case. First, the kernel defined this way *depends on the trajectory*. As the sequence of domains met by the trajectory, and the time where the trajectory enters in these domains, depend on the stimulus, the RF of rectifiable cells *depends now on the stimulus*. Note that the situation would actually be even worse for non linear cells. Indeed, the question hidden behind these remarks is: "to what extent the *linear response* assumption defining a RF via a convolution equation such as (2) is valid". We will actually come back to a similar question in section 2.3 for a network of spiking neurons. Linear response essentially requires the perturbation to be "weak enough", which in our case, means that cells are not rectified. The formulation in terms of a piecewise linear system allows to extend the notion of RF to rectified cells, but the price to pay is that RF now depends on the stimulus. With respect to biology, this effect would for example mean that cells identified e.g. to be ON with a STA approach, responds differently (e.g. ON-OFF) to a more sophisticated stimulus like the "chirp" stimulus [69].

In the rectified cases, the eigenvalues $\lambda_\beta^{(n)}$, $\beta = 1 \dots N$ and eigenvectors $\mathcal{P}_\beta^{(n)}$ depend on the domain, i.e. on the list of rectified cells and are different from the domain $\Omega^{(0)}$ of the rest state. They actually differ in two ways. First, rectified cells provide eigenvalues $-\frac{1}{\tau_\beta}$ and eigenvectors \vec{e}_β so that $\mathcal{P}_{\alpha\beta}^{(n)} = \delta_{\alpha\beta}$ for these cells so that they do not contribute any more to the network response. The second effect is more intricate. Indeed, the mere fact of rectifying one cell, has, in general, the effect of *modifying the whole spectrum and eigenvectors*, with strong effects on the cells response. This can be easily understood. Consider the (not really retinal-realistic) situation where a cell is a hub in a network. Silencing it have in general dramatic effects on the global dynamics of this network.

646 2.1.4. Conclusion of section 2.1

647 In this section, we have given a mathematical answer to the problem 1, level 1,
 648 posed in the introduction. On the basis of a simplified model of BCells - ACells -
 649 GCells interactions, we have produced a formalism allowing us to compute this network
 650 response to spatio-temporal stimuli. We have been able to write explicitly the RF of
 651 individual GCells appearing in eq. (1) where the kernel depends explicitly on lateral
 652 connectivity. As we showed, however, the linear response formula (1), where the kernel
 653 is independent of the stimulus, holds when the stimuli has a weak enough amplitude so
 654 that cells are not rectified. As soon as rectification takes place the convolution form (1)
 655 implies, in general, that the kernel can change with the stimulus. This effect could be
 656 observed in experiments if the cell type, characterized via STA, provides a different type
 657 of response to other stimuli.

658 2.2. *How could spatio-temporal stimuli correlations and retinal network dynamics shape the*
 659 *spike train correlations at the output of the retina ?*

660 In the section, we extrapolate the previous analysis of the model (13) to analyse
 661 how spike trains emitted by GCells can be correlated via the network and especially
 662 ACells connectivity. We especially want to make mathematical statements on how could
 663 ACells decorrelate GCells, as claimed on the basis of experiments [15] ? We consider first
 664 the non rectified case and then analyse how rectification can modify correlations.

665 2.2.1. Voltages correlations

666 We now compute the GCells voltage correlations induced by a non stationary
 667 spatio-temporal stimulus in the model (13). Note that correlations requires some notion
 668 of probability, thus, of randomness. Moreover, it is more convenient when such a
 669 probability is stationary, while we want here to consider a non-stationary problem. This
 670 is not contradictory though. There are two simple (not incompatible) ways to address
 671 this point. First, one may consider that the dynamical system (13) has random initial
 672 conditions, drawn with respect to a stationary probability measure. Second, one can add
 673 to the dynamics (13) noise, which always present in biological systems. We can make the
 674 assumption that noise is stationary and that it is Brownian (which is a pure mathematical
 675 convenience). In biology, spike correlations are usually obtained by averaging over
 676 repeats of the same experiment where a stimuli is presented to the retinal network. This
 677 corresponds therefore to averaging over initial conditions in the presence of noise.

We consider therefore a noisy version of (13), first in the non rectified case. For
 simplicity we set $\vec{C}^{(0)} = \vec{0}$. Dynamics of cells voltage is now given by:

$$\frac{d\vec{\mathcal{X}}}{dt} = \mathcal{L}^{(0)}.\vec{\mathcal{X}} + \vec{\mathcal{F}}(t) + \sigma_B \vec{\zeta}(t), \quad (29)$$

678 where $\vec{\zeta}(t)$ is a white noise (we wrote the stochastic dynamics the physicists way). The
 679 noise intensity σ_B is here a constant. We could consider more elaborated versions where
 680 noise depends on neurons indices, but this would unnecessarily make the analysis more
 681 complex. Note however that we add noise to GCells too, but this was already implicitly
 682 done in the LNP formulation (8). As a consequence, these two noise will somehow add
 683 up, as made clear in eq. (30). This is harmless however to our analysis on correlations.
 684 As in the previous section, we consider that the stimulus is applied from an initial time
 685 far in the past, so that we integrate the trajectory from $-\infty$ to t .

The solution of (29) is a Gaussian process with probability law \mathbb{P} , mean:

$$\vec{m}^{(0)}(t) = \int_{-\infty}^t e^{\mathcal{L}^{(0)}(t-s)}.\vec{\mathcal{F}}(s) ds \equiv \left[e^{\mathcal{L}^{(0)} \cdot t} * \vec{\mathcal{F}} \right](t), \quad (30)$$

and correlation matrix:

$$\mathcal{C}(t_1, t_2) = \sigma_B^2 e^{\mathcal{L}^{(0)}(t_2-t_1)} \cdot \int_{-\infty}^{t_1} e^{\tilde{\mathcal{L}}^{(0)}(t_1-s)} \cdot e^{\mathcal{L}^{(0)}(t_1-s)} ds, \quad (31)$$

686 for $t_2 \geq t_1$, where $\tilde{\mathcal{L}}^{(0)}$ is the transpose of $\mathcal{L}^{(0)}$. The integral is a function of t_1 and $\mathcal{L}^{(0)}$.
 687 From these equations we see that, in general, the probability law of \vec{X} depends on the
 688 spectrum of $\mathcal{L}^{(0)}$. Let us make this dependence more explicit.

689

In the general case $\mathcal{L}^{(0)}$ is not symmetric and does not commute with $\tilde{\mathcal{L}}^{(0)}$. One can then compute $\mathcal{C}(t_1, t_2)$ in terms of the (common) spectrum of $\mathcal{L}^{(0)}$, $\tilde{\mathcal{L}}^{(0)}$ using the spectral decomposition theorem $\mathcal{L}^{(0)} = \sum_{\alpha=1}^N \lambda_{\alpha}^{(0)} v_{\alpha}^{(0)} \cdot \tilde{w}_{\alpha}^{(0)}$ where $v_{\alpha}^{(0)}$ is the right eigenvector α of $\mathcal{L}^{(0)}$ (the α -th column of $\mathcal{P}^{(0)}$) and $\tilde{w}_{\alpha}^{(0)}$ is the left eigenvector α of $\mathcal{L}^{(0)}$ (the α -th row of $(\mathcal{P}^{(0)})^{-1}$). In general, right (left) eigenvectors are not mutually orthogonal but $\tilde{w}_{\alpha}^{(0)} \cdot v_{\beta}^{(0)} = \delta_{\alpha\beta}$ so that $v_{\alpha}^{(0)} \cdot \tilde{w}_{\alpha}^{(0)}$ is the projector on eigendirection α . From this, one obtains the correlation matrix:

$$\mathcal{C}(t_1, t_2) = -\sigma_B^2 \sum_{\alpha=1}^N e^{\lambda_{\alpha}^{(0)}(t_2-t_1)} v_{\alpha}^{(0)} \cdot \tilde{w}_{\alpha}^{(0)} \sum_{\beta=1}^N \frac{v_{\beta}^{(0)} \cdot \tilde{w}_{\beta}^{(0)}}{\lambda_{\alpha}^{(0)} + \lambda_{\beta}^{(0)}}, \quad (32)$$

690 where eigenvalues are real or complex conjugate and are assumed to be stable (negative
 691 real part). Note that eigenvalues and projectors combine so that, in fine, the correlation
 692 matrix is real. It is stationary and *stimulus independent*.

693

We will keep this general form for further discussions on the rectified case, but here, it is insightful to consider the case where $\mathcal{L}^{(0)}$ is symmetric. Here, it is diagonalizable in a orthogonal basis, with $(\mathcal{P}^{(0)})^{-1} = \tilde{\mathcal{P}}^{(0)}$ and with real eigenvalues $\lambda_{\beta} \equiv -s_{\beta}$, $\beta = 1 \dots N$, where s_{β} is real, positive. Then, (32) reduces, in form of components, to:

$$C_{\alpha_2, \alpha_1}(t_2 - t_1) = \frac{\sigma_B^2}{2} \sum_{\beta=1}^N \frac{P_{\alpha_2\beta} P_{\alpha_1\beta}}{s_{\beta}} e^{-s_{\beta}(t_2-t_1)}. \quad (33)$$

This is the noise-induced correlation between cell α_2 at time t_2 and cell α_1 at time t_1 . It is useful to express, from (33), the variance of cell α_1 's voltage (independent of time due to stationarity):

$$\sigma_{\alpha_1}^2 = \frac{\sigma_B^2}{2} \sum_{\beta=1}^N \frac{P_{\alpha_2\beta} P_{\alpha_1\beta}}{s_{\beta}}. \quad (34)$$

694

These computations provide the so-called "noise correlations" between cells voltage.
 695 Note that this concerns *all* cell types (not only GCells). The wording "noise correlations"
 696 is actually a bit misleading. Indeed, a simple glance at (33) (or even (31)), shows that
 697 cells correlations are essentially constrained by the network structure, so that, a change
 698 of parameters, e.g. plasticity or pharmacology, induces dramatic variations in this corre-
 699 lation.

700

We now compute spike correlations of GCells. We assume a spiking probability of the form (8). The probability that GCell $\alpha_1 (> N_B + N_A)$ spikes at time t_1 is induced

by the voltage probability \mathbb{P} and is given by $v_{\alpha_1}(t_1) \equiv \mathbb{E}\left[f\left(\frac{V_G(t)-\theta_G}{\sigma_G}\right)\right]$ where the expectation is taken with respect to \mathbb{P} . Taking the form (8) for f this is:

$$v_{\alpha_1}(t_1) = f\left(\frac{m_{\alpha_1}(t_1) - \theta_G}{\sqrt{\sigma_G^2 + \sigma_{\alpha_1}^2}}\right). \quad (35)$$

As pointed out above, two sources of noise add up here: the implicit noise, with variance σ_G^2 appearing in the LNP formulation (8), which is intrinsic to the cell, and the network induced noise, explicit in the term $\sigma_{\alpha_1}^2$.

Likewise, the probability that GCell $\alpha_1 (> N_B + N_A)$ spikes at time t_1 and GCell $\alpha_2 (> N_B + N_A)$ spikes at time t_2 is:

$$v_{\alpha_1\alpha_2}(t_1, t_2) = \int f\left(\frac{\sqrt{\mu_1} \cos(\phi) y_1 - \sqrt{\mu_2} \sin(\phi) y_2 + m_{\alpha_1}(t_1) - \theta_G}{\sigma_G}\right) f\left(\frac{\sqrt{\mu_1} \sin(\phi) y_1 + \sqrt{\mu_2} \cos(\phi) y_2 + m_{\alpha_2}(t_2) - \theta_G}{\sigma_G}\right) DY, \quad (36)$$

where the integral holds on \mathbb{R}^2 and where $DY = \frac{1}{2\pi} e^{-\frac{y_1^2 + y_2^2}{2}} dy_1 dy_2$. Here, μ_1, μ_2 are the eigenvalues of the pairwise correlation matrix $\mathcal{C} = \begin{pmatrix} \sigma_{\alpha_1}^2 & \mathcal{C}_{\alpha_1\alpha_2}(t_1 - t_2) \\ \mathcal{C}_{\alpha_2\alpha_1}(t_2 - t_1) & \sigma_{\alpha_2}^2 \end{pmatrix}$ which is diagonalizable in an orthogonal basis with an orthogonal transformation, a rotation with angle ϕ determined by the coefficients of \mathcal{C} .

2.2.2. Consequences

Decorrelation. It is evident that the double integral (36) factorizes only in the case where \mathcal{C} is diagonal ($\phi = 0, \mu_1 = \sigma_{\alpha_1}, \mu_2 = \sigma_{\alpha_2}$), and it reduces to $v_{\alpha_1\alpha_2}(t_1, t_2) = v_{\alpha_1}(t_1)v_{\alpha_2}(t_2)$. Thus, spikes of GCell α_1 at time t_1 and of GCell α_2 at time t_2 are decorrelated if and only if the correlation matrix (33) is diagonal. In particular, as we saw, simple networks with ACells have, in general, the effect to correlate voltages and thereby spikes. This is in contrast with the claim, found in deep experimental papers stating that "the inhibition" (mediated by ACells) "decorrelates visual feature representations in the inner retina [15]". What could be the origin of this discrepancy?

First, even if the correlations we compute are non vanishing they can nevertheless be weak. The weakness of pairwise correlations in the retina has actually been reported by many authors [61,62]. It is known since Lancaster, 1957 [75] that the passage of two correlated Gaussian variables through a subsequent non linearity always reduces the correlation of the two signals, regardless of the shape of the nonlinearity. Thus, in our case, the non linear function of the LNP model reduces the decorrelation. This fact was reported by Pitkow and Meister in their paper "Decorrelation and efficient coding by retinal ganglion cells" [12] where they say: "The classical theory of retinal decorrelation attributed that phenomenon to filtering by center-surround receptive fields and explained its purpose as serving the efficient transmission of visual information through the optic nerve". Although their experiments conclude "that the spike trains of retinal ganglion cells were indeed decorrelated compared with the visual input" they insist on the prominent role of nonlinearities: "most of the decorrelation was accomplished not by the receptive fields, but by nonlinear processing in the retina" and "in the LN model, the nonlinearity decorrelates if it has a high threshold, ensuring that each neuron spends much of the time silent except for sharp and sparse firing events." From these remarks they conclude about information transmission by the retinal network: "At very high thresholds, the information transmission is poor. Notably, transmission also drops at low thresholds. Thus, the choice of threshold involves a trade-off between rarely using reliable symbols, such as high spike counts, or frequently using unreliable symbols, such as low spike counts".

Thus, non linearities play a role in retinal coding making the spike rate of RGCs as sparse as possible, so that these cells are silent most of the time and fire at high rate

741 only when salient features of the stimulus make it necessary. This effect should be even
 742 more prominent for moving objects which is clearly an example of a stimulus with salient
 743 features and strong spatio-temporal correlations induced by its trajectory, especially if
 744 this trajectory shows sharp changes.

745

746 **Rectification.** Now, the LNP non linearity is not the only source of decorrelation. Rec-
 747 tification also plays a crucial role. What happens, indeed, in the rectified case? Math-
 748 ematically, one can extend equations (17) providing the general solution of (29) in the
 749 presence of noise, but the main, quite intricate problem is now that the entrance and
 750 exit time of domains $t_{-}^{(n_k)}, t_{+}^{(n_k)}$ are *themselves* random. This is again a consequence of
 751 the stimulus dependence of these times. The computation of the voltage correlations in
 752 this case being, for the moment, out of reach, I am going to give some straightforward
 753 although insightful remarks.

754

The non rectified case correspond to a trajectory staying in the domain $\Omega^{(0)}$ (forget-
 755 ting about conditions on noise ensuring that this holds for an infinite time). Now, the
 756 computation of voltage correlation is essentially the same if the trajectory stays in the
 757 domain $\Omega^{(n)}$. The only difference is that eigenvalues and projectors have a superscript
 758 (n) instead of (0). This difference is essential though, because rectification induces a
 759 projection on the space of non rectified cells. The contribution to rectified cells to voltage
 760 correlations with other cells vanishes thereby transforming the voltage correlations
 761 matrix. By permutation of rows and columns, one can convert this matrix in a form
 762 containing a diagonal block (correlations rectified cells \leftrightarrow rectified cells) and a block
 763 characterizing the correlations non rectified cells \leftrightarrow all cells. This reduces the model
 764 dimensionality and the global correlations. This effect, composed with the LN non
 765 linearity can even reduce correlation.

766

Now, the most important remark here is that rectification implies that GCells corre-
 767 lations are *stimulus dependent*. This effect might not be noticeable with full field stimuli
 768 or white noise, which weakly solicit the lateral ACell connectivity, but it could be more
 769 prominent when studying spatio-temporal stimuli, in particular moving trajectories.
 770 From this perspective, the standard decorrelation argument, where centre and surround
 771 components of Bcells receptive fields interact via ACells lateral inhibition to decorrelate
 772 bipolar cell output in the spatial and temporal domains, may not apply for non stationary
 773 stimuli, which constitute most of real visual scenes.

774 2.2.3. Conclusion of section 2.2

775

In this section, we have mathematically investigated the structure of correlations
 776 induced by the model (13), Fig. 1 right. Our conclusion is essentially that the stimulus
 777 generates GCells spike correlations modulated by the BCells-ACells networks, and
 778 more precisely, by the eigenvalues-eigenmodes of the transport operator. While, for
 779 non rectified cells, the spike correlations are independent of the stimulus, introducing
 780 rectification dramatically changes the picture. As well as the RF, correlations are stimulus
 781 dependent. However, non linearities (LNP transfer function and rectification) have the
 782 general effect of reducing pairwise correlations, confirming a result emphasized by
 783 Pitkow and Meister [12]. Let us also remark that rectification makes the stochastic
 784 process of voltages non Gaussian, because the times of entering and exiting domains
 785 are now random variables too. As a consequence, spike statistics involves higher order
 786 correlations. Although it has to be further investigated on experimental grounds, this
 787 would lead to important consequences in terms of coding. As pointed out, again, by
 788 Pitkow and Meister [12], "for highly non-Gaussian signals, such as neural spike trains
 789 and natural images, correlation may be only weakly related to redundancy."

790

Sticking at the model we may ask the following questions. Assume that we submit
 791 the model to different type of stimuli: the "classical" ones such as white noise, "Chirp"
 792 stimulus, natural images; but also more elaborated ones such as moving objects with
 793 different type of trajectories, or "natural movies" including motion and "surprise". For

794 example, a bird crossing the visual scene, with, on the background, a forest of trees in the
795 wind. It is known that the retina is able to filter the "noisy" motion of tree leaves while
796 signalling the bird, thanks to dedicated circuits involving ACells [1,44]. Such circuits
797 can be easily implemented in the model (13) [28]. What will be the structure of its spike
798 trains, depending on the different type of stimuli ? How can one "efficiently" decode the
799 stimulus from the mere knowledge of those spike trains ? How efficient is a decoding
800 scheme based on independent, decorrelated Gcells ? In contrast, would cooperative
801 network effects make the code more precise affording faster responses to motion [14] ?

802 Although we are not going to answer these questions here (there is still a long
803 way to it), we give, in the next sections, several insightful mathematical results in this
804 direction.

805 2.3. Computing the mixed effect of network and stimulus on spike correlations

806 Let us now consider the retina from the point of view of its output. We sit on the
807 optic nerve and measure the spikes sent to the LGN and cortex via the optic nerve. We
808 have no access to the biophysical machinery taking place in the retina and generating
809 those spikes, but we know that the spike trains contains information about the external
810 world stimuli that we want to extract. We can measure as many quantities as we want
811 such as firing rate, or higher correlations. More generally, we are seeking the (time
812 dependent) joint probability of spikes adopting the approach described in section 1.2,
813 Methods.

814 In this context, assume a retina "at rest" i.e. receiving no stimulus or stationary
815 stimuli like noise. We can describe the spike trains emitted by this retina by a stationary
816 transition probability \mathbb{P} , associated with a stationary probability $\mu^{(sp)}$ (for "spontaneous").
817 In general, this probability has spike correlations of order 2 and higher. Assume now
818 that, from time t_0 on a stimulus (say a moving object) is getting through the visual field
819 of this retina. As exposed in section 2.2 one expects the spike correlations (at any order)
820 to be modified by this stimulation. Typically, a moving object carries spatio-temporal
821 correlations in its trajectory which will superimposed upon the network correlations,
822 resulting in a mixed effect where non linearities can also play a role. Can we predict, for
823 a given stimulus, how correlations will be modified ?

824 Let us give an example. Consider a linear chain of neurons, as depicted in Fig. 3 top.
825 Each neuron (black points), is connected to its neighbours with an excitatory connection
826 (red arrows) and to its second nearest neighbours with an inhibitory connection (blue
827 arrows). In spontaneous activity, we assume an asynchronous dynamics, such as Fig.
828 3 bottom, left. In addition, each neuron is able to sense external stimuli. Consider a
829 moving stimulus propagating from left to right (cyan, bell shaped curve) Fig. 3 top. This
830 stimulus is going to modify the spike patterns, as seen in Fig. 3 bottom, left, where one
831 sees clearly the effect of the stimulus, nearest neighbours excitation and second nearest
832 neighbours inhibition. The remarkable fact is that the stimulus not only modifies the
833 firing rates of neurons, but also *their correlations*. The question is: can we compute this
834 effect ?

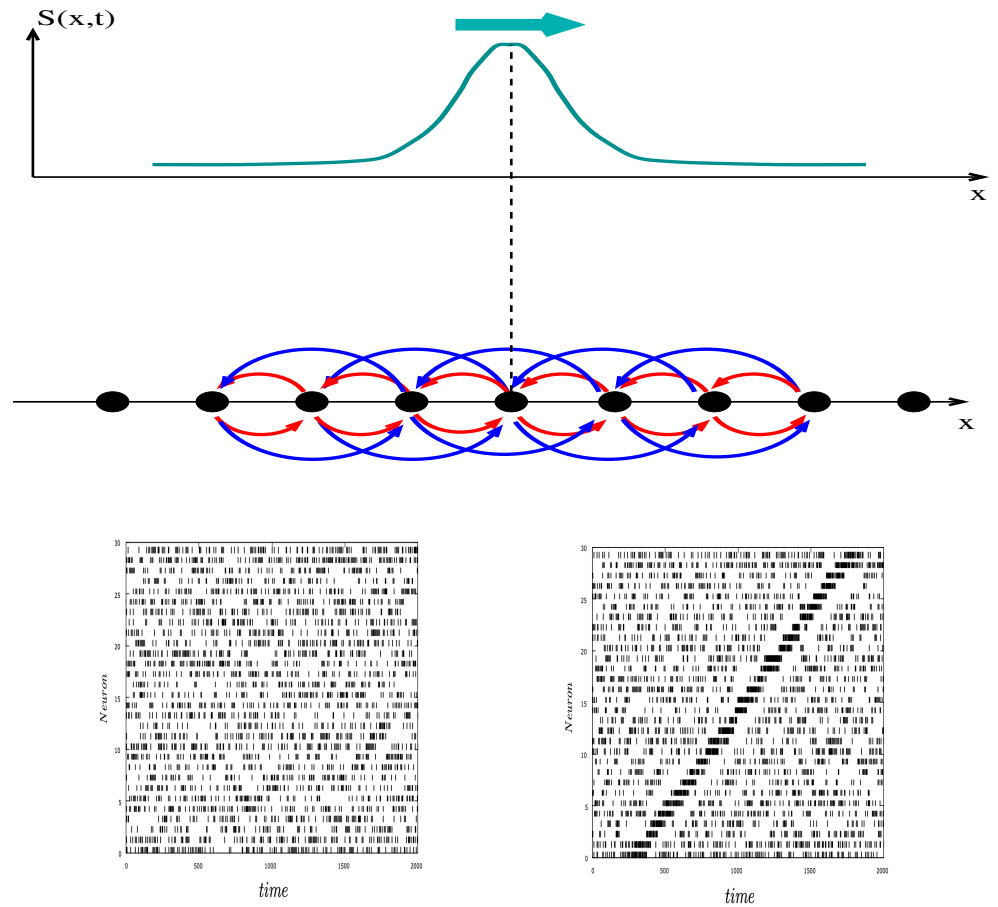


Figure 3. Top. Network of spiking neurons sensing a stimulus. Each neuron, represented as a black point, is connected to its neighbours with an excitatory connection (red arrows) and to its second nearest neighbours with a inhibitory connection (blue arrows). In addition, each neuron is able to sense external stimuli $S(x, t)$ (cyan, bell shaped curve). **Bottom. Left.** Spontaneous spiking activity. **Bottom. Right.** Spiking activity in the presence of the moving stimulus.

835 This question has been solved in the paper [39] for the gIF model (23). Here, we
 836 briefly state the main result (see the paper for technical details). Consider a function
 837 $f(t, \omega)$ (observable) depending on time and spike history up to time t . Let $\mu^{(sp)}$ be
 838 the joint probability distribution of spikes in spontaneous activity (no stimulus), and μ
 839 the joint probability distribution of spikes in the presence of a spatio-temporal stimulus
 840 $S(x, t)$. We note $\delta\mu[f](t) = \mu[f](t) - \mu^{(sp)}[f]$, where $\mu[f](t)$ is the average of f , at
 841 time t , in the presence of the stimulus and $\mu^{(sp)}[f]$ the average of f in spontaneous
 842 activity (which does not depend of time because spontaneous dynamics is stationary).
 843 $\delta\mu[f(t)]$ characterizes how much the time dependent mean of $f(t, \omega)$ under stimulation
 844 departs from the spontaneous mean at time t . In the simplest case $\delta\mu[f(t)]$ characterizes
 845 the variation in the firing rate of neuron k , if $f(t, \omega) = \omega_k(t)$, or the variation in the
 846 correlation between neuron k_1 at time t_1 and neuron k_2 at time $t_1 + t$ if $f(t, \omega) =$
 847 $(\omega_{k_1}(t_1) - \mu^{(sp)}[\omega_{k_1}]) (\omega_{k_2}(t_1 + t) - \mu^{(sp)}[\omega_{k_2}])$, and so on.

One can show that, when the stimulus amplitude is weak enough, $\delta\mu[f(t)]$ is given by a linear response formula of the form:

$$\delta\mu[f(t)] = [\kappa_f * S](t) \quad (37)$$

848 That is, by the convolution of the stimulus with a specific kernel, K_f , depending on the
849 observable f and on the spontaneous distribution $\mu^{(sp)}$. We do not give the expression of this
850 kernel here, for simplicity, but the reader can refer to the paper [39].

851 2.3.1. Consequences

852 Let us comment this result.

853
854 **Convolution.** Similarly to (1) (GCells response to stimuli) or (2) (BCells response
855 to stimuli), we have here again a linear response where the effect of a stimulus on a
856 system is expressed by a convolution. We are however in a completely different per-
857 spective. Indeed, while we were considering formerly voltage response of individuals
858 cells (shaped by network effects), we are now working on a more abstract level, where
859 we attempt to measure the effect of a stimulus on *statistics*. This is of course due to the
860 difference in what is accessible by experiments, what the observer is able to deal with
861 in his observations, here spikes. Thus, the mathematical machinery allowing to extract
862 the response requires to define spike statistics in a non stationary setting, where the
863 influence of the stimulus can be inferred.

864
865 **Kernel.** The kernel K_f can be explicitly computed in the gIF model. It depends on
866 several features. First, on network characteristics (especially the effective interaction
867 W_{kj} , and more generally, the parameters shaping the model dynamics). It also depends
868 on the observable f . However, the main content of this result is that the kernel K_f is
869 actually determined by spike correlations in *spontaneous activity*. In other word, it is
870 possible to anticipate the response to a non stationary stimulus from the knowledge
871 of the spontaneous activity. Although this result is expected from Kubo theory in non
872 equilibrium statistical physics [76,77] or from Volterra-Wiener expansions [20], it has
873 interesting consequences when dealing with neural dynamics, and more specifically here,
874 with retina outputs. First, it provides a consistent treatment of the expected perturbation
875 of higher-order correlations, beyond the known linear perturbation of firing rates and
876 instantaneous pairwise correlations; in particular, it extends to time-dependent corre-
877 lations. In addition it reveals how the stimulus-response and dynamics are entangled
878 in a complex manner. For example, the response of a neuron k to a stimulus applied
879 on neuron i does not only depends on the synaptic weight W_{ki} but, in general, on all
880 synaptic weights, because the dynamics create complex causality loops which build up
881 the response of neuron k [36,78,79]. The linear response formula is written in terms of the
882 parameters of a spiking neuronal network model and the spike history of the network.
883 Although a linear treatment may seem a strong simplification, our results suggest that al-
884 ready, in this case, the connectivity architecture should not be neglected. In the presence
885 of stimuli, the whole architecture of synaptic connectivity, history and the dynamical
886 properties of the networks are playing a role in the spatio-temporal correlations structure.

887
888 **Linear response and higher order corrections.** The derived formula provides
889 a good agreement with simulations in the gIF model under time dependent stimuli
890 (typically, a moving object). It requires however that the stimulus amplitude is weak
891 enough. That is, higher corrections are weaker than the leading order. For larger
892 amplitude stimuli one would to compute higher order correlations. This can be done
893 using the same formalism [80] although it might not be the best approach. Indeed,
894 this method requires to measure spontaneous correlations which are difficult to obtain
895 experimentally for orders higher than 2. This is actually one of the reason why LNP-like
896 models exist. The expected non linearity in the response is handled by a static non linear
897 function. Exploring what could be the best non linear correction to the linear response
898 in such models is definitely an interesting mathematical challenge.

899 2.4. Conclusion

900 **Beyond naive RF description.** This linear response theory actually shows how the
 901 neuronal network substrate and stimulus response are entangled. Indeed, in contrast to
 902 naive RF representation where the convolution kernel is assumed to depend only on the
 903 *cell*, here, mathematics show that it depends as well on the *observable*. The explicit form of
 904 the kernel is also tightly constrained by the neurons connections. Finally, a convolution
 905 implies an integration over histories, requiring thereby to consider spikes probabilities
 906 with memory, instead of "instantaneous" spikes probabilities (not or weakly depending
 907 on the past). Of course, one may always argue that, on experimental grounds, long tail
 908 memory are just impossible to measure so "instantaneous" [61] or first order Markov
 909 [81] models are largely sufficient. But what means "sufficient" ? This is a difficult ques-
 910 tion which requires sophisticated methods to determine the "best performing" memory
 911 depth from data [82–84]. Actually, numerical computations of the kernel use, of course
 912 Markovian approximations [39], although, with a memory depth that can be controlled.

913
 914 **Link with the retina model.** Can we relate the formalism developed here with
 915 the retinal model presented in the section 1.1 ? As GCells voltage are Gaussian it is
 916 in principle possible to compute transition probabilities using the transport operator
 917 formalism. However, even in the non rectified case, the computation promises to be
 918 a formidable task, unless one adds some additional constraints. For example, a big
 919 advantage of Integrate and Fire models is that a spiking neuron loses memory after
 920 spiking, a property which is not implemented in LNP like models.

921
 922 **Information geometry.** There is a close link between Gibbs distributions and in-
 923 formation geometry. This theory, developed by Shun'ichi Amari and his collaborators
 924 (see [85] and references therein) on the basis of early work from Rao [86] establishes a
 925 geometric theory of information where probabilities are considered as points on Riemannian
 926 manifolds. A prominent family of probability measures is called the exponential
 927 family. It contains the Gibbs distributions in the standard statistical physics sense, i.e.
 928 probabilities having the form $\frac{e^{-\beta H}}{Z}$ where the energy H does not depend on time. In this
 929 case, the metric is given by the Hessian of $\log(Z)$, the free energy, and is tightly linked
 930 to Fisher information on one hand and to linear response on the other hand. The linear
 931 response is actually a correlation function, from the fluctuation dissipation theorem.
 932 Thus, correlation functions induce a natural geometry for Gibbs distributions providing
 933 strong insights on how these distribution are modified by smooth, local, transformations
 934 of their parameters (like learning [87]) or under a stimulation of weak amplitude. In this
 935 last case, the stimulus action corresponds to a perturbation in the tangent space of the
 936 manifold [88,89]. Although information geometry has not been extended, to our best
 937 knowledge, to the type of Gibbs distribution we study here (they are non stationary)
 938 the mathematical formalism is similar when considering finite memory and stationary
 939 distributions. This essentially tells us that the structure of spatio-temporal correlations
 940 observed in spike trains reveals an hidden geometrical structure which, somewhat,
 941 shapes the response of the retina, and, henceforth of cortex, to stimuli. We come back to
 942 this point in the conclusion section.

943 3. Applications

944 The OPL-BCells-ACells processing is based on graded potentials departing from
 945 the classical paradigm of binary spike processing. Mathematically, this has strong conse-
 946 quences in terms of response to a spatio-temporal stimulus: existence of eigenmodes,
 947 potentially modulated by non linear effects, inducing properties such as activity waves
 948 ahead of the stimulus (anticipation), resonances, correlations modified by the stimulus.
 949 In this section, and although this paper is essentially theoretical I would like to shortly
 950 propose possible applications of these results, outside the field of neuroscience.

951 3.1. Retinal prostheses

952 Retinal pathologies such as Age Macular Degeneration or Retinitis Pigmentosa, are
953 due to the degeneration of photo-receptors [90]. In addition, they induce morphological
954 and structural changes in the retina with significant pathogenic effects: inflammation,
955 change in connectivity, the appearance of large-scale spontaneous electrical oscillations,
956 and, of course, attenuation of response to visual stimuli [91–94]. In this process of
957 degeneration, however, the Gcells are the last to be deficient, maintaining, therefore, a
958 link between the retina and the brain, provided they are suitably stimulated. The strategy
959 of retinal prostheses is to stimulate the retina electrically by an array of electrodes.
960 Stimulation of an electrode generates, in the visual cortex, a phosphene, the perception of
961 a light spot. By stimulating the electrodes one induces in the cortex an image "pixelised"
962 by the phosphenes, with resolution limited, on the one hand, by the number of electrodes,
963 and, on the other hand, by the size of the phosphenes, which can be enlarged by diffusion
964 and non linear effects [95]. Technological solutions, taking into account the physiological
965 limitation on the electrical power that can be injected in an electrode, improve resolution
966 [96]. However, there are still obstacles which cannot be resolved by purely technological
967 solutions (hardware). In addition, a valid stimulation strategy at a given period of the
968 pathology may not be later because the retina degeneration evolves in time.

969 Stimulation strategies use processor pre-processing to calculate, from a given image
970 (captured by a camera) the pattern of stimulation of the prosthesis, by mimicking the
971 calculation that a healthy retina would make, or by incorporating corrections taking
972 into account the pathology [97]. These algorithms might be improved using what
973 we know about the retinal structure, especially ACells lateral connectivity, where a
974 model like (13) can be easily implemented with a relatively low energy consumption
975 cost. The idea would be to improve electrodes stimulation sequences so as to allow an
976 implanted patient to perceive in real time a moving object. The model (13) with ACells
977 lateral connectivity and gain control is known to produce a wave of activity ahead of a
978 stimulus, performing a form of anticipation [28]. This could be used to compensate for
979 the processing times imposed by the equipment, in the same way that the visual system
980 knows how to compensate for the delays induced by photo-transduction [25]. The ideal
981 would also be to have adaptive algorithms, i.e., depending on parameters adjustable
982 according to the patient and the course of his pathology.

983 3.2. Convolutional networks

984 Several recent studies attempt to understand how retinal response to stimuli is
985 related to circuit processes using convolutional neural network models [98] to under-
986 stand the structure of retinal prediction [99]. Reciprocally, these networks can be used to
987 encode deep-learning models to encode dynamic visual scenes with important potential
988 outcomes in the domain of computer vision. Especially, a recent work by Zheng et
989 al, [100], shows the important role played by recurrence in encoding complex natural
990 scenes. To my best knowledge (which is quite scarce in this field) there is no mathemati-
991 cal analysis of the dynamics of these models, especially the dependence in parameters
992 and robustness of the training schemes. The present study could bring some insights in
993 this perspective. Even if the model (13) is different from what these researchers were
994 using, the techniques of piecewise linear phase decomposition and eigenmodes study
995 could be insightful to better understand the dynamical evolution of these convolutional
996 networks and the role played by rectification.

997 4. Discussion

998 In this paper we have addressed mathematically the potential effect of ACells
999 lateral connectivity on retinal response to spatio-temporal stimuli. We have seen how,
1000 mathematically, the retina structure and the collective dynamics of retinal cells organized
1001 in local circuits spanning the whole retina might constrain this response. In particular,

1002 the structure of correlations is expected to depend on the stimulus, as soon as non linear
1003 effects are involved.

1004 These properties are established on the basis of theoretical results though, based on
1005 incomplete modelling of the retina, and specific assumptions. Their validation would
1006 require experiments, some of which may require time and others are not yet accessible,
1007 for example, simultaneously measuring retina and cortex. As a matter of fact, one may
1008 argue that the models presented here are far too simplistic compared to the real retina(s)
1009 having a large number of BCells, ACells, GCells type, making complex circuits [18]
1010 and whose characteristics depend, in addition, on species, age, pathologies. However,
1011 the idea behind mathematical modelling is precisely to try and infer some generic
1012 mechanism underlying the real object under study, here the retina. This is the simplicity
1013 of the structure which makes it generic. The question is: "Would the addition of more
1014 elaborated retinal features make the response to stimuli simpler?"

1015 In the next section I further discuss some further implications of this work leading
1016 to some new questions.

1017 4.1. Cortical response

1018 If a dynamical stimulus, combined with the retinal network and non linearities
1019 produces non negligible dynamical spatio-temporal correlations, what could be the
1020 consequences at the cortical level³? There is a physiological transformation, called
1021 retinotopy, which maps smoothly the retina topology to the cortical V1 topology. In
1022 models, it is usually considered to be the identity map, although it is not. This is a non
1023 linear transformation, depending, in addition on the species [101–103]. Nevertheless,
1024 what matters here is that this mapping is smooth and invertible. Therefore, retinotopy
1025 transports, in a smooth and invertible way the spatio-temporal retinal correlations to the
1026 visual cortex. This leads to a first question.

1027
1028 **How to define a cortical model taking into account spatio-temporal spike corre-**
1029 **lations.** Cortical models are usually based on mean-field approximations where one
1030 features firing rates evolution, but not spike correlations. This is the case of the Wilson-
1031 Cowan model [104–106] or neural field models [107–109]. I know about 2 mean-field
1032 approaches taking care of spike correlations.

1033 The first approach is the one initiated by S. El Boustani et A. Destexhe [110] using a
1034 Markovian approach to write down mean field equations of second order (i.e. includ-
1035 ing pairwise spatial correlations) and a non static thalamic entry, that can feature the
1036 retinal-LGN input. This model can be used to construct a retino-cortical model [111],
1037 although the mathematical consequences of having correlated retinal entries have not
1038 been explored yet.

1039 The second approach is based on the so-called Ott-Antonsen Ansatz [112] and
1040 has been used by Montbrio, Pazo, Roxin to propose an exact mean field approach
1041 with second order statistics [113]. Since their paper there has been a lot of activity in
1042 developing this model, especially in connection with cortical imaging, with impressive
1043 results [114–117]. It is a promising track.

1044 All these approaches could certainly provide powerful numerical and mathematical
1045 tools to better understand how spatio-temporal retinal correlations could be processed.
1046 In particular, having a retino-(LGN)-cortical model allows to do a task which is currently
1047 impossible experimentally: measure simultaneously retina and cortex.

1048 4.2. Retinal correlations and neurogeometry.

1049 We have also seen that retinal correlations and Gibbs distributions naturally define
1050 a metric on a Riemannian manifold where probabilities are points on this manifold. In
1051 particular, the application of a weak amplitude stimulus corresponds to a perturbation

³ For simplicity, I am going to consider the LGN as a simple relay

1052 along the tangent space of this manifold. What is the image of this metric under the
 1053 retinotopic transformation ? Let us make this question a bit more precise.

1054 The visual system has evolved to map as efficiently as possible retinal output to
 1055 cortical structures. The shaping of the visual system during development is actually a
 1056 highly dynamical process involving retinal waves and synaptic plasticity [70]. These
 1057 processes provides the visual system a structure allowing to respond in a fast and efficient
 1058 way to the stimuli coming from the external world, via the retina. In particular, the
 1059 capacity of the visual cortex to respond to spike trains with spatio-temporal correlations
 1060 induced by natural stimuli should be somewhat imprinted in the cortical connectivity.

1061 The visual perception is actually highly geometrically structured and shaped by the
 1062 structure of the cortical connectivity. This leads researchers to introduce a link between
 1063 the geometry of cortex and the geometry of vision in the concept of neurogeometry
 1064 (or neuromathematics) where the functional architecture of V1 is considered as a Lie
 1065 group of symmetry with a Riemannian geometry (see [40–43] and reference therein). In
 1066 this approach cortical columns are point-like processors detecting visual features where
 1067 functional connectivity is represented in terms of geodesics. To my best knowledge,
 1068 neurogeometry essentially deals with V1 and static percepts, although extensions to
 1069 motor cortex [118] and motion areas [119] have been done. Now, a natural question is:
 1070 "Is there a relation between the cortical metric of neurogeometry and the metric induced
 1071 by spatio-temporal spike correlations observed by the retina ?".

1072 Let us address the problem the other way round: Project the cortical metric back
 1073 to the retina via the inverse retinotopy map, what do we find ? Is there a physiological
 1074 correspondence with the retina structure and especially lateral connectivity ? What
 1075 could be the consequences on spike trains statistics and on the way retina processes
 1076 visual stimuli ? What does cortical metric tell us about retinal spikes correlations ?

1077 Dealing with neural coding of vision, the simplest assumption consists of assuming
 1078 that GCells are independent encoders and that cortex makes the job of restoring the
 1079 spatio-temporal correlations existing in the visual scene (e.g. in the trajectory of a moving
 1080 object). The alternative proposition, where spatio-temporal correlations imprinted in the
 1081 RGCs spike trains are deciphered by the cortex makes the question of stimuli decoding
 1082 by the cortex more challenging, but opens up far more possibilities.

1083 **Acknowledgments:** I would like to warmly acknowledge all the neuroscientist collaborators from
 1084 which I learned about this beautiful object, the retina, and, more generally, about vision: Frédéric
 1085 Chavane, Gerrit Hilgen, Olivier Marre, Adrian Palacios, Serge Picaud and Evelyne Sernagor.

1086 **Conflicts of Interest:** The authors declare no conflict of interest.

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