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PARISIAN TIMES FOR LINEAR DIFFUSIONS

CHRISTOPHE PROFETA

ABSTRACT. We compute the joint distribution of the first times a linear diffusion makes an excursion longer than some given duration above (resp. below) some fixed level. In the literature, such stopping times have been introduced and studied in the framework of *Parisian* barrier options, mainly in the case of Brownian motion with drift. We also exhibit several independence properties, and provide some formulae for the associated ruin probabilities.

1. STATEMENT OF THE MAIN RESULTS

1.1. Introduction. Let X be a regular diffusion, without killing, taking values in some interval $I \subseteq \mathbb{R}$. In this note, we are interested in the following *Parisian* stopping times:

$$\kappa_r^{(a,\pm)} = \inf\{t > r, t - g_t^{(a)} \geq r \text{ and } \pm(X_t - a) \geq 0\}$$

where $a \in \text{Int}(I)$, the interior of I , $r > 0$ and

$$g_t^{(a)} = \sup\{s \leq t, X_s = a\}.$$

These stopping times are the first times the process X makes an excursion above a (resp. below a) longer than $r > 0$. They were introduced by Chesney, Jeanblanc & Yor [3] for the pricing of some barrier options, in which the owner loses the option if the underlying asset X remains long enough below or above some fixed level. Several papers have been devoted to the valuation of such options, but in the case of diffusion processes, to our knowledge, only the example of Brownian motion with drift has been studied in details so far. In the case of two barriers, we refer the reader to Dassios & Wu [5] who start from a semi-Markov model and use a limiting procedure, and to Anderluh & van der Weide [1] and Labart & Lelong [10] who use mainly martingales methods.

We shall not tackle here the problem of computing Parisian options, but rather study the distribution of the quadruplet $(\kappa_u^{(a,-)}, \kappa_v^{(b,+)}, X_{\kappa_u^{(a,-)}}, X_{\kappa_v^{(b,+)}})$ where $a < b$ and $u, v > 0$. This distribution is always at the core of all the aforementioned papers. Our approach will essentially rely on diffusion and excursion theory.

The paper is divided as follows. In the remainder of Section 1, we introduce some notation and state our main result, as well as several corollaries and consequences. Section 2 is devoted

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to some preliminary lemmas, which will be needed for the proof of the main theorem in Section 3. Finally, several examples are given in Section 4, including Brownian motion with drift and some Bessel processes.

1.2. Notations. We shall work in the framework of Pitman & Yor [12]. Let us denote by \mathbb{P}_x and \mathbb{E}_x the probability and expectation of X when starting from $x \in I$, and $(\mathcal{F}_t, t \geq 0)$ its natural filtration. We assume that the infinitesimal generator of X has the form

$$\mathcal{G} = \frac{1}{2}\sigma^2(x)\frac{d^2}{dx^2} + \mu(x)\frac{d}{dx}$$

on some appropriate domain, with σ a strictly positive and continuous function and μ a locally bounded function. Recall that X may also be defined by its scale function s and speed measure m which are given by :

$$s'(x) = \exp\left(-\int^x \frac{2\mu(y)}{\sigma^2(y)}dy\right) \quad \text{and} \quad m(dx) = \frac{2}{s'(x)\sigma^2(x)}dx. \quad (1.1)$$

For $a \in \text{Int}(I)$, we define the first hitting time

$$T_a = \inf\{t > 0, X_t = a\}.$$

Its Laplace transform is classically given for $\lambda > 0$ by :

$$\mathbb{E}_x[e^{-\lambda T_a}] = \begin{cases} \Phi_{\lambda,-}(x)/\Phi_{\lambda,-}(a) & \text{if } x \leq a, \\ \Phi_{\lambda,+}(x)/\Phi_{\lambda,+}(a) & \text{if } x > a, \end{cases}$$

where $\Phi_{\lambda,-}$ and $\Phi_{\lambda,+}$ are two solutions (respectively increasing and decreasing) of the ordinary differential equation $\mathcal{G}u = \lambda u$, subject to some appropriate boundary conditions. Applying the Markov property and using the continuity of the paths of X , we also have for $a \leq x \leq b$ and $\lambda > 0$:

$$\mathbb{E}_x[e^{-\lambda T_a} 1_{\{T_a < T_b\}}] = \frac{W_\lambda^{(b)}(x)}{W_\lambda^{(b)}(a)} \quad \text{and} \quad \mathbb{E}_x[e^{-\lambda T_b} 1_{\{T_b < T_a\}}] = \frac{W_\lambda^{(a)}(x)}{W_\lambda^{(a)}(b)}$$

where, for $c \in I$, the function $W_\lambda^{(c)}$ is defined by :

$$W_\lambda^{(c)}(x) = \Phi_{\lambda,+}(x)\Phi_{\lambda,-}(c) - \Phi_{\lambda,-}(x)\Phi_{\lambda,+}(c).$$

Recall next from Itô & McKean [8, Section 6.2] that T_a admits a continuous probability density function and that there exist two Lévy measures $\nu_+^{(a)}$ and $\nu_-^{(a)}$ such that

$$\nu_\pm^{(a)}(dt) = \pm \lim_{\varepsilon \downarrow 0} \frac{\mathbb{P}_{a \pm \varepsilon}(T_a \in dt)}{s(a \pm \varepsilon) - s(a)} = \pm \frac{1}{s'(a)} \frac{d}{dx} \mathbb{P}_x(T_a \in dt) \Big|_{x=a^\pm}.$$

Equivalently, $\nu_+^{(a)}$ and $\nu_-^{(a)}$ may be defined by their (modified) Laplace transforms :

$$\int_{(0,+\infty]} (1 - e^{-\lambda t}) \nu_\pm^{(a)}(dt) = \mp \frac{1}{s'(a)} \frac{\Phi'_{\lambda,\pm}(a)}{\Phi_{\lambda,\pm}(a)}.$$

Some care must be taken when dealing with transient diffusions, as in this case, the measure $\nu_{\pm}^{(a)}$ might have an atom at $t = +\infty$ given by :

$$\nu_{\pm}^{(a)}(\{+\infty\}) = \mp \frac{1}{s'(a)} \lim_{\lambda \downarrow 0} \frac{\Phi'_{\lambda, \pm}(a)}{\Phi_{\lambda, \pm}(a)}.$$

We also introduce the Green kernel of X which is defined by the formula

$$G_{\lambda}(x, y) = \int_0^{+\infty} e^{-\lambda t} p(t; x, y) dt$$

where $p(t; x, y)$ denotes the transition density of X with respect to the speed measure m , i.e.

$$\mathbb{P}_x(X_t \in dy) = p(t; x, y) m(dy).$$

The Green kernel admits the classic representation

$$G_{\lambda}(x, y) = \omega_{\lambda}^{-1} \Phi_{\lambda, -}(x) \Phi_{\lambda, +}(y), \quad x \leq y,$$

where the Wronskian ω_{λ} is defined by

$$\omega_{\lambda} = \frac{\Phi'_{\lambda, -}(x) \Phi_{\lambda, +}(x) - \Phi_{\lambda, -}(x) \Phi'_{\lambda, +}(x)}{s'(x)}$$

and is independent from x .

We finally define the positive and negative meanders of lengths u by :

$$\forall \Lambda_u \in \mathcal{F}_u, \quad \begin{cases} \mathbb{P}_a^{\uparrow, u}(\Lambda_u) = \lim_{\varepsilon \downarrow 0} \mathbb{P}_{a+\varepsilon}(\Lambda_u | T_a > u) \\ \mathbb{P}_a^{\downarrow, u}(\Lambda_u) = \lim_{\varepsilon \downarrow 0} \mathbb{P}_{a-\varepsilon}(\Lambda_u | T_a > u). \end{cases}$$

The occurrence of meanders in the forthcoming formulae is not surprising as the random variables $X_{\kappa_u^{(a, -)}}$ and $X_{\kappa_v^{(b, +)}}$ are respectively a.s. smaller than a and greater than b .

1.3. Main results. Let α and β be two continuous and bounded functions from I to $[0, +\infty)$. For $a \leq x \leq b$ and $\gamma, \lambda, u, v > 0$ let us define the modified Laplace transform of the quadruplet $(\kappa_u^{(a, -)}, \kappa_v^{(b, +)}, X_{\kappa_u^{(a, -)}}, X_{\kappa_v^{(b, +)}})$ by :

$$F_{\alpha, \beta, \gamma, \lambda}(x) = \mathbb{E}_x \left[e^{-\gamma \kappa_u^{(a, -)} - \lambda \kappa_v^{(b, +)}} \alpha(X_{\kappa_u^{(a, -)}}) \beta(X_{\kappa_v^{(b, +)}}) 1_{\{\kappa_v^{(b, +)} \leq \kappa_u^{(a, -)} < +\infty\}} \right].$$

The main result of the paper is the computation of this function. To simplify the statement, let us set :

$$\Psi_{\lambda}^{(\pm)}(a, b, r) = \frac{\Phi_{\lambda, \pm}(a)}{\Phi_{\lambda, \pm}(b)} \pm \frac{W_{\lambda}^{(b)}(a)}{\omega_{\lambda}} \int_r^{+\infty} e^{-\lambda t} \nu_{\pm}^{(b)}(dt).$$

Theorem 1. For $a \leq x \leq b$, the function $F_{\alpha,\beta,\gamma,\lambda}$ admits the expression :

$$F_{\alpha,\beta,\gamma,\lambda}(x) = \left(\frac{W_{\gamma+\lambda}^{(b)}(x)}{W_{\gamma+\lambda}^{(b)}(a)} \frac{1}{\Psi_{\gamma+\lambda}^{(-)}(b, a, u)} + \frac{W_{\gamma+\lambda}^{(a)}(x)}{W_{\gamma+\lambda}^{(a)}(b)} \right) F_{\alpha,\beta,\gamma,\lambda}(b) \quad (1.2)$$

where

$$\begin{aligned} F_{\alpha,\beta,\gamma,\lambda}(b) \\ = e^{-(\lambda+\gamma)v} \frac{W_{\gamma+\lambda}^{(b)}(a)}{\omega_{\gamma+\lambda}} \frac{\nu_+^{(b)}[v, +\infty] \mathbb{E}_b \left[e^{-\gamma \kappa_u^{(a,-)}} \right]}{\Psi_{\gamma+\lambda}^{(+)}(a, b, v) - \frac{1}{\Psi_{\gamma+\lambda}^{(-)}(b, a, u)}} \mathbb{E}_a^{\downarrow, u} [\alpha(X_u)] \mathbb{E}_b^{\uparrow, v} \left[\frac{\Phi_{\gamma,+}(X_v)}{\Phi_{\gamma,+}(b)} \beta(X_v) \right] \end{aligned}$$

and

$$\mathbb{E}_b \left[e^{-\gamma \kappa_u^{(a,-)}} \right] = e^{-\gamma u} \frac{\Phi_{\gamma,+}(b)}{\Phi_{\gamma,+}(a)} \frac{\nu_-^{(a)}[u, +\infty]}{\frac{1}{G_{\gamma}(a, a)} + \int_u^{+\infty} e^{-\gamma t} \nu_-^{(a)}(dt)}.$$

One may observe several independence properties in this formula.

- (1) Letting $\gamma \downarrow 0$, we deduce that on the set $\{\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)} < +\infty\}$, the random variables $\kappa_v^{(b,+)}$, $X_{\kappa_u^{(a,-)}}$ and $X_{\kappa_v^{(b,+)}}$ are mutually independent.
- (2) Taking $\alpha \equiv 1$, we observe that one may remove the condition $\{\kappa_u^{(a,-)} < +\infty\}$ in the definition of $F_{1,\beta,\gamma,\lambda}$. Letting $\gamma \downarrow 0$ and then $u \rightarrow +\infty$, we deduce that on the set $\{\kappa_v^{(b,+)} < +\infty\}$, the random variables $\kappa_v^{(b,+)}$ and $X_{\kappa_v^{(b,+)}}$ are independent. This independence property, when dealing with a single boundary, was already observed in the Brownian case by Chesney, Jeanblanc & Yor [3].

We mention below two situations where Theorem 1 greatly simplifies. First, when $\alpha \equiv \beta \equiv 1$, one obtains the following corollary :

Corollary 2. For $a < b$ and $\gamma, \lambda > 0$, the Laplace transform of the pair $(\kappa_u^{(a,-)}, \kappa_v^{(b,+)})$ is given by :

$$\mathbb{E}_b \left[e^{-\gamma \kappa_u^{(a,-)} - \lambda \kappa_v^{(b,+)}} 1_{\{\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)}\}} \right] = e^{-\lambda v} \frac{W_{\gamma+\lambda}^{(b)}(a)}{\omega_{\gamma+\lambda}} \frac{\int_v^{+\infty} e^{-\gamma t} \nu_+^{(b)}(dt)}{\Psi_{\gamma+\lambda}^{(+)}(a, b, v) - \frac{1}{\Psi_{\gamma+\lambda}^{(-)}(b, a, u)}} \mathbb{E}_b \left[e^{-\gamma \kappa_u^{(a,-)}} \right].$$

This result follows from the observation that

$$\begin{aligned} \mathbb{E}_b^{\uparrow, v} \left[\frac{\Phi_{\gamma,+}(X_v)}{\Phi_{\gamma,+}(b)} \right] &= \lim_{\varepsilon \downarrow 0} \frac{\mathbb{E}_{b+\varepsilon} [\mathbb{E}_{X_v} [e^{-\gamma T_b}] 1_{\{T_b \geq v\}}]]}{\mathbb{P}_{b+\varepsilon}(T_b \geq v)} \\ &= \lim_{\varepsilon \downarrow 0} \frac{\mathbb{E}_{b+\varepsilon} [e^{-\gamma(T_b-v)} 1_{\{T_b \geq v\}}]]}{\mathbb{P}_{b+\varepsilon}(T_b \geq v)} = \frac{e^{\gamma v} \int_v^{+\infty} e^{-\gamma t} \nu_+^{(b)}(dt)}{\nu_+^{(b)}[v, +\infty]} \end{aligned}$$

where we have applied the Markov property in the second equality. Another simplification is obtained by letting $a \uparrow b$. Using that $W_\lambda^{(a)}(b) \xrightarrow{a \uparrow b} 0$ and computing the limit yields the general formula for the one barrier case :

Corollary 3. *For $\gamma, \lambda > 0$, the Laplace transform of the pair $(\kappa_u^{(b,-)}, \kappa_v^{(b,+)})$ is given by :*

$$\begin{aligned} \mathbb{E}_b \left[e^{-\gamma \kappa_u^{(b,-)} - \lambda \kappa_v^{(b,+)}} 1_{\{\kappa_v^{(b,+)} \leq \kappa_u^{(b,-)}\}} \right] \\ = \frac{e^{-\lambda v} \int_v^{+\infty} e^{-\gamma t} \nu_+^{(b)}(dt)}{\int_v^{+\infty} e^{-(\gamma+\lambda)t} \nu_+^{(b)}(dt) + \int_u^{+\infty} e^{-(\gamma+\lambda)t} \nu_-^{(b)}(dt) + \frac{1}{G_{\gamma+\lambda}(b,b)}} \mathbb{E}_b \left[e^{-\gamma \kappa_u^{(b,-)}} \right]. \end{aligned}$$

In particular, when X is a Brownian motion with drift, by letting $\gamma \downarrow 0$, we recover the formula obtained by Dassios & Wu [4].

1.4. Parisian ruin probabilities. Letting $\gamma \downarrow 0$ in the Laplace transform of $\kappa_u^{(a,-)}$ given in Theorem 1, we obtain the following ruin probabilities :

- (1) if X is recurrent, then $\lim_{\gamma \downarrow 0} G_\gamma(a, a) = +\infty$ and $\nu_-^{(a)}(\{+\infty\}) = 0$, hence

$$\mathbb{P}_a(\kappa_u^{(a,-)} < +\infty) = \frac{\nu_-^{(a)}[u, +\infty]}{\nu_-^{(a)}[u, +\infty)} = 1.$$

- (2) if X is transient, then $G_0(a, a) = \lim_{\gamma \downarrow 0} G_\gamma(a, a) < +\infty$ and

$$\mathbb{P}_a(\kappa_u^{(a,-)} < +\infty) = \frac{\nu_-^{(a)}[u, +\infty]}{\frac{1}{G_0(a,a)} + \nu_-^{(a)}[u, +\infty)}.$$

In particular, if X goes to $\inf(I)$ a.s. as $t \rightarrow +\infty$, then $\frac{1}{G_0(a,a)} = \nu_-^{(a)}(\{+\infty\})$ and this probability also equals 1.

Besides Brownian motion with drift, such Parisian ruin probabilities have also been computed for spectrally negative Lévy processes, see for instance Loeffen, Czarna & Palmowski [11].

Another probability of interest is the probability that $\kappa_v^{(b,+)}$ takes place before $\kappa_u^{(a,-)}$, which may be obtained by letting λ and γ go to 0 in Corollary 2 and Formula (1.2).

Corollary 4. *Assume that either X is recurrent, or X is transient and goes a.s to $\sup(I)$ or to $\inf(I)$ as $t \rightarrow +\infty$. Then, for $a \leq x \leq b$:*

$$\begin{aligned} \mathbb{P}_x(\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)} < +\infty) \\ = \left(\frac{s(b) - s(x)}{\frac{1}{\mathbb{P}_a(T_b < +\infty)} + (s(b) - s(a))\nu_-^{(a)}[u, +\infty)} + (s(x) - s(a)) \right) \frac{\mathbb{P}_b(\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)} < +\infty)}{s(b) - s(a)} \end{aligned}$$

with

$$\begin{aligned} \mathbb{P}_b(\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)} < +\infty) \\ = \frac{(s(b) - s(a))\nu_+^{(b)}[v, +\infty)\mathbb{P}_b(T_a < +\infty)\mathbb{P}_a(\kappa_u^{(a,-)} < +\infty)}{\frac{1}{\mathbb{P}_b(T_a < +\infty)} + (s(b) - s(a))\nu_+^{(b)}[v, +\infty) - \frac{1}{\frac{1}{\mathbb{P}_a(T_b < +\infty)} + (s(b) - s(a))\nu_-^{(a)}[u, +\infty)}}. \end{aligned}$$

Remark 5. When X is recurrent, since $\mathbb{P}_b(T_a < +\infty) = \mathbb{P}_a(T_b < +\infty) = 1$, Corollary 4 greatly simplifies and one obtains :

$$\mathbb{P}_x(\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)}) = \frac{\left(1 + (s(x) - s(a))\nu_-^{(a)}[u, +\infty)\right)\nu_+^{(b)}[v, +\infty)}{\nu_-^{(a)}[u, +\infty) + \nu_+^{(b)}[v, +\infty) + (s(b) - s(a))\nu_-^{(a)}[u, +\infty)\nu_+^{(b)}[v, +\infty)}.$$

To prove Corollary 4, observe that by definition of the Wronskian ω_λ , we have

$$\frac{W_\lambda^{(b)}(a)}{\omega_\lambda} = \Phi_{\lambda,+}(a)\Phi_{\lambda,+}(b) \int_a^b \frac{s'(z)}{\Phi_{\lambda,+}^2(z)} dz = \int_a^b \frac{\mathbb{E}_b[e^{-\lambda T_z}]}{\mathbb{E}_z[e^{-\lambda T_a}]} s'(z) dz.$$

Then, on the one hand, when X is recurrent or X is transient and goes a.s. to $\inf(I)$, we obtain

$$\frac{W_\lambda^{(b)}(a)}{\omega_\lambda} \xrightarrow[\lambda \downarrow 0]{} \int_a^b s'(z) dz = s(b) - s(a).$$

On the other hand, when X is transient and goes a.s. to $\sup(I)$, we may assume without loss of generality that $s(\sup(I)) = 0$ and again

$$\frac{W_\lambda^{(b)}(a)}{\omega_\lambda} \xrightarrow[\lambda \downarrow 0]{} \int_a^b \frac{s(b)/s(z)}{s(z)/s(a)} s'(z) dz = s(b) - s(a).$$

Corollary 4 now follows by letting γ and λ go to 0 in Corollary 2 and Formula (1.2) and using the above limit.

1.5. Remarks on the meanders. We shall give in Section 4 a few examples where all the terms appearing in Theorem 1 may be computed explicitly. To this end, it is worthwhile to note that the terms involving the meanders in Theorem 1 only require the knowledge of the transition density p and of the Lévy measures $\nu_\pm^{(a)}$. Indeed, by applying the Markov property

$$\mathbb{E}_x[\alpha(X_u)1_{\{T_a \geq u\}}] = \mathbb{E}_x[\alpha(X_u)] - \int_0^u \mathbb{E}_a[\alpha(X_{u-t})] \mathbb{P}_x(T_a \in dt).$$

Inverting this transform, we obtain after some simple manipulations :

$$\begin{aligned} \mathbb{P}_x(X_u \in dz | T_a \geq u) / m(dz) \\ = \frac{p(u; x, z) - p(u; a, z)}{\mathbb{P}_x(T_a \geq u)} + p(u; a, z) + \int_0^u \frac{p(u; a, z) - p(u-t; a, z)}{\mathbb{P}_x(T_a \geq u)} \mathbb{P}_x(T_a \in dt). \end{aligned}$$

Letting then $x \uparrow a$, we deduce that for $z < a$:

$$\begin{aligned} & \mathbb{P}_a^{(\downarrow, u)}(X_u \in dz) / m(dz) \\ &= -\frac{1}{s'(a)\nu_-^{(a)}[u, +\infty]} \frac{\partial p}{\partial x}(u; x, z)|_{x=a} + p(u; a, z) + \int_0^u \frac{p(u; a, z) - p(u-t; a, z)}{\nu_-^{(a)}[u, +\infty]} \nu_-^{(a)}(dt). \end{aligned}$$

One may also give an alternative definition of the meanders as derivatives. Indeed, for example, by definition of the left derivative :

$$\begin{aligned} \mathbb{P}_a^{(\downarrow, u)}(\Lambda_u) &= \lim_{\varepsilon \downarrow 0} \frac{-\varepsilon}{\mathbb{P}_{a-\varepsilon}(T_a \geq u)} \frac{\mathbb{P}_{a-\varepsilon}(\Lambda_u 1_{\{T_a \geq u\}})}{-\varepsilon} \\ &= -\frac{1}{s'(a)\nu_-^{(a)}[u, +\infty]} \frac{d}{dx} \mathbb{P}_x(\Lambda_u 1_{\{T_a \geq u\}}) \Big|_{x=a^-}. \end{aligned} \quad (1.3)$$

Formula (1.3) will be the one appearing during the proofs.

2. PRELIMINARY COMPUTATIONS

Our proof will rely on some local estimates for $t \rightarrow p(t; a, a)$. Let us introduce the semimartingale local time of X :

$$L_t^a = \sigma^2(a) \lim_{\varepsilon \downarrow 0} \frac{1}{2\varepsilon} \int_0^t 1_{\{|X_s - a| \leq \varepsilon\}} ds, \quad t \geq 0.$$

From the occupation time formula (see [14, Chapter VI]) and the expression (1.1) of the speed measure m , we have

$$\mathbb{E}_a[L_t^a] = \frac{2}{s'(a)} \int_0^t p(t; a, a) dt.$$

In particular, using Fatou's lemma, and the fact that $\sigma > 0$, we deduce that

$$\frac{2}{s'(a)} \liminf_{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \int_0^\varepsilon p(t; a, a) dt = \liminf_{\varepsilon \downarrow 0} \mathbb{E}_a \left[\frac{L_\varepsilon^a - L_0^a}{\varepsilon} \right] \geq \mathbb{E}_a \left[\liminf_{\varepsilon \downarrow 0} \frac{L_\varepsilon^a - L_0^a}{\varepsilon} \right] = +\infty$$

which yields a first estimate :

$$\lim_{\varepsilon \downarrow 0} \frac{\varepsilon}{\int_0^\varepsilon p(t; a, a) dt} = 0. \quad (2.1)$$

We now prove a lemma which will be repeatedly used in the sequel.

2.1. Limiting lemma.

Lemma 6. *Let $a \in \text{Int}(I)$ and let φ be a positive and continuous function such that $\mathbb{E}_x[\varphi(X_t)] < +\infty$ is finite for t small enough and x in the vicinity of a . We assume that $\varphi(a) = 0$ and that φ admits a left and a right derivative at the point a . Then the following asymptotics hold:*

$$\mathbb{E}_a[\varphi(X_\varepsilon) 1_{\{X_\varepsilon \geq a\}}] \underset{\varepsilon \downarrow 0}{\sim} \frac{1}{s'(a)} \varphi'(a^+) \int_0^\varepsilon p(t; a, a) dt$$

and

$$\mathbb{E}_a [\varphi(X_\varepsilon) 1_{\{X_\varepsilon \leq a\}}] \underset{\varepsilon \downarrow 0}{\sim} -\frac{1}{s'(a)} \varphi'(a^-) \int_0^\varepsilon p(t; a, a) dt.$$

Proof. We only prove the first asymptotics, as the second one will follow from a similar argument. Take δ small enough and let us decompose the first expectation in Lemma 6 according as whether X_ε is greater or smaller than $a + \delta$:

$$\mathbb{E}_a [\varphi(X_\varepsilon) 1_{\{X_\varepsilon \geq a\}}] = \mathbb{E}_a [\varphi(X_\varepsilon) 1_{\{X_\varepsilon \geq a+\delta\}}] + \mathbb{E}_a [\varphi(X_\varepsilon) 1_{\{a+\delta > X_\varepsilon \geq a\}}].$$

Applying the Markov property, the first term on the right-hand side is smaller than

$$\begin{aligned} \mathbb{E}_a [\varphi(X_\varepsilon) 1_{\{X_\varepsilon \geq a+\delta\}}] &= \mathbb{E}_a \left[\mathbb{E}_{a+\delta} [\varphi(X_{\varepsilon-s}) 1_{\{X_{\varepsilon-s} \geq a+\delta\}}] \Big|_{s=T_{a+\delta}} 1_{\{T_{a+\delta} \leq \varepsilon\}} \right] \\ &\leq \sup_{s \in [0, \varepsilon]} \mathbb{E}_{a+\delta} [\varphi(X_s)] \mathbb{P}_a (T_{a+\delta} \leq \varepsilon) \end{aligned}$$

From Kotani & Watanabe [9], there is the asymptotics for $z \in \text{Int}(I)$:

$$\lim_{\varepsilon \downarrow 0} -2\varepsilon \ln (\mathbb{P}_a (T_z \leq \varepsilon)) = \left(\int_a^z \frac{1}{\sigma(u) \sqrt{s'(u)}} du \right)^2. \quad (2.2)$$

As a consequence, we deduce from (2.1) that

$$\lim_{\varepsilon \rightarrow 0} \frac{\mathbb{E}_a [\varphi(X_\varepsilon) 1_{\{X_\varepsilon \geq a+\delta\}}]}{\int_0^\varepsilon p(t; a, a) dt} = 0$$

which proves that this term is negligible. Then, since $\varphi(a) = 0$, the second term may be bounded by :

$$\inf_{a \leq z \leq a+\delta} \frac{\varphi(z)}{z-a} \leq \frac{\mathbb{E}_a [\varphi(X_\varepsilon) 1_{\{a+\delta > X_\varepsilon \geq a\}}]}{\mathbb{E}_a [(X_\varepsilon - a) 1_{\{a+\delta > X_\varepsilon \geq a\}}]} \leq \sup_{a \leq z \leq a+\delta} \frac{\varphi(z)}{z-a}. \quad (2.3)$$

Recall now that for any $x \in \text{Int}(I)$, the function $(t, y) \rightarrow p(t; x, y)$ is a solution of the partial differential equation $\frac{\partial}{\partial t} p = \mathcal{G}p$ on $(0, +\infty) \times \text{Int}(I)$, and that for $z \neq a$, from Kotani & Watanabe [9] :

$$\lim_{t \downarrow 0} p(t; a, z) = 0. \quad (2.4)$$

Since the scale function s and the speed measure m are assumed to be absolutely continuous, an adaptation of their proof shows also that

$$\lim_{t \downarrow 0} \frac{\partial p}{\partial z}(t; a, z) = 0. \quad (2.5)$$

Now, integrating the partial differential equation for p on $[0, \varepsilon]$ and using (2.4), we deduce that for $z \neq a$,

$$p(\varepsilon; a, z) = \int_0^\varepsilon \left(\frac{1}{2} \sigma^2(z) \frac{\partial^2 p}{\partial z^2}(t; a, z) + \mu(z) \frac{\partial p}{\partial z}(t; a, z) \right) dt.$$

Then applying Fubini theorem, and recalling that $m(dx) = 2/(s'(x)\sigma^2(x))dx$, we may write

$$\begin{aligned} \mathbb{E}_a [(X_\varepsilon - a)1_{\{a+\delta > X_\varepsilon \geq a\}}] \\ = \int_0^\varepsilon \left(\int_a^{a+\delta} \frac{z-a}{s'(z)} \frac{\partial^2 p}{\partial z^2}(t; a, z) dz + \int_a^{a+\delta} (z-a)\mu(z) \frac{\partial p}{\partial z}(t; a, z) m(dz) \right) dt. \end{aligned}$$

Integrating by parts the first term, and using the expression (1.1) of s , we observe that the terms in μ cancel and we obtain :

$$\mathbb{E}_a [(X_\varepsilon - a)1_{\{a+\delta > X_\varepsilon \geq a\}}] = \int_0^\varepsilon \left(\frac{\delta}{s'(a+\delta)} \frac{\partial p}{\partial z}(t; a, a+\delta) - \int_a^{a+\delta} \frac{1}{s'(z)} \frac{\partial p}{\partial z}(t; a, z) dz \right) dt.$$

Using yet another integration by parts, we finally obtain :

$$\begin{aligned} \mathbb{E}_a [(X_\varepsilon - a)1_{\{a+\delta > X_\varepsilon \geq a\}}] &= \frac{1}{s'(a)} \int_0^\varepsilon p(t; a, a) dt \\ &+ \int_0^\varepsilon \left(\frac{1}{s'(a+\delta)} \left(\delta \frac{\partial p}{\partial z}(t; a, a+\delta) - p(t; a, a+\delta) \right) + \mathbb{E}_a [\mu(X_t)1_{\{a+\delta \geq X_t \geq a\}}] \right) dt. \end{aligned}$$

As a consequence, going back to (2.3), we deduce from (2.4), (2.5), and (2.1) that

$$\begin{aligned} \inf_{a \leq z \leq a+\delta} \frac{\varphi(z)}{z-a} &\leq \liminf_{\varepsilon \downarrow 0} \frac{s'(a)\mathbb{E}_a [\varphi(X_\varepsilon)1_{\{a+\delta > X_\varepsilon \geq a\}}]}{\int_0^\varepsilon p(t; a, a) dt} \\ &\leq \limsup_{\varepsilon \downarrow 0} \frac{s'(a)\mathbb{E}_a [\varphi(X_\varepsilon)1_{\{a+\delta > X_\varepsilon \geq a\}}]}{\int_0^\varepsilon p(t; a, a) dt} \leq \sup_{a \leq z \leq a+\delta} \frac{\varphi(z)}{z-a} \end{aligned}$$

and the result follows by letting $\delta \downarrow 0$. □

2.2. The distribution of $(\kappa_u^{(a,-)}, X_{\kappa_u^{(a,-)}})$. We now prove a simpler version of Theorem 1, involving only $\kappa_u^{(a,-)}$ and $X_{\kappa_u^{(a,-)}}$.

Proposition 7. *On the set $\{\kappa_u^{(a,-)} < +\infty\}$ the random variables $\kappa_u^{(a,-)}$ and $X_{\kappa_u^{(a,-)}}$ are independent, and their distributions are respectively given by*

$$\mathbb{E}_a \left[e^{-\gamma \kappa_u^{(a,-)}} \right] = e^{-\gamma u} \frac{\nu_-^{(a)}[u, +\infty]}{\frac{1}{G_\gamma(a, a)} + \int_u^{+\infty} e^{-\gamma t} \nu_-^{(a)}(dt)}$$

and

$$\mathbb{E}_a \left[\alpha(X_{\kappa_u^{(a,-)}}) 1_{\{\kappa_u^{(a,-)} < +\infty\}} \right] = \mathbb{E}_a^{\downarrow, u} [\alpha(X_u)].$$

Proof. We set

$$d_t^{(a)} = \inf\{s > t, X_s = a\}$$

and

$$F_{\alpha, \gamma}(a) = \mathbb{E}_a \left[e^{-\gamma \kappa_u^{(a,-)}} \alpha(X_{\kappa_u^{(a,-)}}) 1_{\{\kappa_u^{(a,-)} < +\infty\}} \right].$$

Let $0 < \varepsilon < u$. We shall decompose the expectation of $F_{\alpha,\gamma}(a)$ according as whether X_ε is greater or smaller than a . When $X_\varepsilon > a$, applying the Markov property at the time $d_\varepsilon^{(a)}$, we obtain

$$\mathbb{E}_a \left[e^{-\gamma \kappa_u^{(a,-)}} \alpha(X_{\kappa_u^{(a,-)}}) 1_{\{\kappa_u^{(a,-)} < +\infty\}} 1_{\{X_\varepsilon > a\}} \right] = \mathbb{E}_a \left[e^{-\gamma d_\varepsilon^{(a)}} 1_{\{X_\varepsilon > a\}} \right] F_{\alpha,\gamma}(a). \quad (2.6)$$

When $X_\varepsilon \leq a$, we need to take care as whether or not the excursion below a straddling ε is longer than u . On the one hand, if this excursion is shorter than u , then

$$\begin{aligned} \mathbb{E}_a \left[e^{-\gamma \kappa_u^{(a,-)}} \alpha(X_{\kappa_u^{(a,-)}}) 1_{\{\kappa_u^{(a,-)} < +\infty\}} 1_{\{X_\varepsilon \leq a\}} 1_{\{d_\varepsilon^{(a)} < g_\varepsilon^{(a)} + u\}} \right] \\ = \mathbb{E}_a \left[e^{-\gamma d_\varepsilon^{(a)}} 1_{\{X_\varepsilon \leq a\}} 1_{\{d_\varepsilon^{(a)} < g_\varepsilon^{(a)} + u\}} \right] F_{\alpha,\gamma}(a). \end{aligned} \quad (2.7)$$

On the other hand, if this excursion is longer than u , then $\kappa_u^{(a,-)} = g_\varepsilon^{(a)} + u$ a.s. and

$$\begin{aligned} \mathbb{E}_a \left[e^{-\gamma \kappa_u^{(a,-)}} \alpha(X_{\kappa_u^{(a,-)}}) 1_{\{\kappa_u^{(a,-)} < +\infty\}} 1_{\{X_\varepsilon \leq a\}} 1_{\{d_\varepsilon^{(a)} \geq g_\varepsilon^{(a)} + u\}} \right] \\ = \mathbb{E}_a \left[e^{-\gamma (g_\varepsilon^{(a)} + u)} \alpha(X_{g_\varepsilon^{(a)} + u}) 1_{\{X_\varepsilon \leq a\}} 1_{\{d_\varepsilon^{(a)} \geq g_\varepsilon^{(a)} + u\}} \right]. \end{aligned} \quad (2.8)$$

Gathering (2.6), (2.7) and (2.8) together, we deduce that

$$F_{\alpha,\gamma}(a) = \frac{\mathbb{E}_a \left[e^{-\gamma (g_\varepsilon^{(a)} + u)} \alpha(X_{g_\varepsilon^{(a)} + u}) 1_{\{X_\varepsilon \leq a\}} 1_{\{d_\varepsilon^{(a)} \geq g_\varepsilon^{(a)} + u\}} \right]}{1 - \mathbb{E}_a \left[e^{-\gamma d_\varepsilon^{(a)}} 1_{\{X_\varepsilon > a\}} \right] - \mathbb{E}_a \left[e^{-\gamma d_\varepsilon^{(a)}} 1_{\{X_\varepsilon \leq a\}} 1_{\{d_\varepsilon^{(a)} < g_\varepsilon^{(a)} + u\}} \right]}$$

and it remains to let $\varepsilon \downarrow 0$. We start with the numerator. Fix $\delta \in (\varepsilon, u)$. Since $g_\varepsilon^{(a)} \leq \varepsilon$, we have

$$\begin{aligned} \mathbb{E}_a \left[e^{-\gamma g_\varepsilon^{(a)}} \alpha(X_{g_\varepsilon^{(a)} + u}) 1_{\{X_\varepsilon \leq a\}} 1_{\{d_\varepsilon^{(a)} \geq g_\varepsilon^{(a)} + u\}} \right] \\ \leq \mathbb{E}_a \left[\max_{s \in [0, \delta]} \alpha(X_{s+u}) 1_{\{X_\varepsilon \leq a\}} 1_{\{d_\varepsilon^{(a)} \geq u\}} \right] \\ = \mathbb{E}_a \left[1_{\{X_\varepsilon \leq a\}} \mathbb{E}_{X_\varepsilon} \left[\max_{s \in [0, \delta]} \alpha(X_{s+u-\varepsilon}) 1_{\{\varepsilon + T_a \geq u\}} \right] \right] \\ \leq \mathbb{E}_a \left[1_{\{X_\varepsilon \leq a\}} \mathbb{E}_{X_\varepsilon} \left[\mathbb{E}_{X_{u-\delta}} \left[\max_{s \in [0, \delta]} \alpha(X_s) \right] 1_{\{T_a \geq u-\delta\}} \right] \right]. \end{aligned}$$

From Lemma 6 and Formula (1.3) we deduce that :

$$\begin{aligned} \mathbb{E}_a \left[1_{\{X_\varepsilon \leq a\}} \mathbb{E}_{X_\varepsilon} \left[\mathbb{E}_{X_{u-\delta}} \left[\max_{s \in [0, \delta]} \alpha(X_s) \right] 1_{\{T_a \geq u-\delta\}} \right] \right] \\ \underset{\varepsilon \downarrow 0}{\sim} \mathbb{E}_a^{\downarrow, u-\delta} \left[\mathbb{E}_{X_{u-\delta}} \left[\max_{s \in [0, \delta]} \alpha(X_s) \right] \right] \nu_-^{(a)}[u-\delta, +\infty] \int_0^\varepsilon p(t; a, a) dt. \end{aligned}$$

As a consequence, we obtain the upper bound :

$$\limsup_{\varepsilon \downarrow 0} \frac{\mathbb{E}_a \left[e^{-\gamma g_\varepsilon^{(a)}} \alpha(X_{g_\varepsilon^{(a)}+u}) 1_{\{X_\varepsilon \leq a\}} 1_{\{d_\varepsilon^{(a)} \geq g_\varepsilon^{(a)}+u\}} \right]}{\nu_-^{(a)}[u - \delta, +\infty] \int_0^\varepsilon p(t; a, a) dt} \leq \mathbb{E}_a^{\downarrow, u-\delta} \left[\mathbb{E}_{X_{u-\delta}} \left[\max_{s \in [0, \delta]} \alpha(X_s) \right] \right].$$

Similarly, we have the lower bound :

$$\liminf_{\varepsilon \downarrow 0} \frac{\mathbb{E}_a \left[e^{-\gamma g_\varepsilon^{(a)}} \alpha(X_{g_\varepsilon^{(a)}+u}) 1_{\{X_\varepsilon \leq a\}} 1_{\{d_\varepsilon^{(a)} \geq g_\varepsilon^{(a)}+u\}} \right]}{\nu_-^{(a)}[u, +\infty] \int_0^\varepsilon p(t; a, a) dt} \geq e^{-\lambda \delta} \mathbb{E}_a^{\downarrow, u} \left[\mathbb{E}_{X_u} \left[\min_{s \in [0, \delta]} \alpha(X_s) \right] \right].$$

Letting $\delta \downarrow 0$, we deduce by continuity that the numerator is equivalent to :

$$\mathbb{E}_a \left[e^{-\gamma g_\varepsilon^{(a)}} \alpha(X_{g_\varepsilon^{(a)}+u}) 1_{\{X_\varepsilon \leq a\}} 1_{\{d_\varepsilon^{(a)} \geq g_\varepsilon^{(a)}+u\}} \right] \underset{\varepsilon \downarrow 0}{\sim} \mathbb{E}_a^{\downarrow, u} [\alpha(X_u)] \nu_-^{(a)}[u, +\infty] \int_0^\varepsilon p(t; a, a) dt.$$

Looking now at the denominator, we obtain, adding and subtracting indicators functions :

$$\begin{aligned} \frac{1}{s'(a)} \left(\frac{d}{dx} \mathbb{E}_x [e^{-\gamma T_a} 1_{\{T_a \leq u\}}] \Big|_{x=a^-} - \frac{d}{dx} \mathbb{E}_x [e^{-\gamma T_a}] \Big|_{x=a^+} \right) \int_0^\varepsilon p(t; a, a) dt \\ = \left(\frac{1}{G_\gamma(a, a)} + \int_u^{+\infty} e^{-\gamma t} \nu_-^{(a)}(dt) \right) \int_0^\varepsilon p(t; a, a) dt. \end{aligned}$$

Gathering the two previous asymptotics, we finally obtain

$$\mathbb{E}_a \left[e^{-\gamma \kappa_u^{(a,-)}} \alpha(X_{\kappa_u^{(a,-)}}) 1_{\{\kappa_u^{(a,-)} < +\infty\}} \right] = e^{-\gamma u} \frac{\nu_-^{(a)}[u, +\infty]}{\frac{1}{G_\gamma(a, a)} + \int_u^{+\infty} e^{-\gamma t} \nu_-^{(a)}(dt)} \mathbb{E}_a^{\downarrow, u} [\alpha(X_u)]$$

which ends the proof of Proposition 7. \square

Remark 8. We briefly sketch here another proof for the distribution of $\kappa_u^{(a,-)}$ which is inspired from Gettoor [7]. Assume that X is recurrent. The right-continuous inverse $(\tau_\ell^a, \ell \geq 0)$ of L^a is a subordinator with Laplace transform :

$$\mathbb{E}_a [e^{-\gamma \tau_\ell^a}] = \exp \left(-\ell \frac{s'(a)}{G_\gamma(a, a)} \right) = \exp \left(-\ell s'(a) \int_0^{+\infty} (1 - e^{-\gamma t}) \nu^{(a)}(dt) \right).$$

where $\nu^{(a)} = \nu_-^{(a)} + \nu_+^{(a)}$. Furthermore the two processes

$$\left(\int_0^{\tau_\ell} 1_{\{X_u > a\}} du, \ell \geq 0 \right) \quad \text{and} \quad \left(\int_0^{\tau_\ell} 1_{\{X_u < a\}} du, \ell \geq 0 \right) \quad (2.9)$$

are independent subordinators with respective Lévy measures given by $s'(a) \nu_+^{(a)}$ and $s'(a) \nu_-^{(a)}$. As a consequence, we may decompose $\tau^{(a)}$ into two independent subordinators

$$\tau^{(a)} = \eta^{(0)} + \eta^{(1)}$$

with respective Lévy measures $s'(a)(\nu_+^{(a)}(dt) + 1_{\{t < u\}}\nu_-^{(a)}(dt))$ and $s'(a)1_{\{t \geq u\}}\nu_-^{(a)}(dt)$. Now $\kappa_u^{(a,-)}$ is the starting point of the first excursion below a of length greater than u . Define the first jump of $\eta^{(1)}$ by

$$R = \inf\{\ell > 0, \eta_\ell^{(1)} \neq 0\}.$$

The random variable R is exponentially distributed with parameter $s'(a)\nu_-^{(a,-)}[u, +\infty]$. Then

$$\begin{aligned} \mathbb{E}_a \left[e^{-\gamma(\kappa_u^{(a,-)} - u)} \right] &= \mathbb{E} \left[e^{-\gamma\eta_{R-}^{(0)}} \right] \\ &= s'(a)\nu_-^{(a)}[u, +\infty] \int_0^{+\infty} \mathbb{E} \left[e^{-\gamma\eta_{r-}^{(0)}} \right] e^{-s'(a)\nu_-^{(a)}[u, +\infty]r} dr \\ &= \frac{\nu_-^{(a)}[u, +\infty]}{\int_0^{+\infty} (1 - e^{-\gamma t})\nu_+^{(a)}(dt) + \int_0^u (1 - e^{-\gamma t})\nu_-^{(a)}(dt) + \nu_-^{(a)}[u, +\infty]} \\ &= \frac{\nu_-^{(a)}[u, +\infty]}{\frac{1}{G_\gamma(a,a)} + \int_u^{+\infty} e^{-\gamma t}\nu_-^{(a)}(dt)}, \end{aligned}$$

which yields the Laplace transform of $\kappa_u^{(a,-)}$ when X is recurrent. However, when X is transient, the two subordinators given in Formula (2.9) are no longer independent as explained in [12]. They are only conditionally independent given L_∞^a , since when $\ell = L_\infty^a$, one will be stopped at the time L_∞^a while the other one will jump to $+\infty$. As a consequence, in the transient case, one would need to condition by L_∞^a before mimicking the previous argument.

3. PROOF OF THEOREM 1

The proof of Theorem 1 is similar to the proof of Proposition 7, except that we shall work with two boundaries $a < b$, and thus write down a system of equations involving $F_{\alpha,\beta,\gamma,\lambda}(a)$ and $F_{\alpha,\beta,\gamma,\lambda}(b)$. Let $0 < \varepsilon < \min(u, v)$. To compute $F_{\alpha,\beta,\gamma,\lambda}(a)$ and $F_{\alpha,\beta,\gamma,\lambda}(b)$, we shall decompose the expectation in three terms, according to the position of X_ε with respect to a and b .

3.1. Study of $F_{\alpha,\beta,\gamma,\lambda}(a)$. First, if $X_\varepsilon \leq a$, then the excursion below a straddling ε must remain shorter than u . Applying the Markov property at the time $d_\varepsilon^{(a)}$, this yields the formula

$$\begin{aligned} \mathbb{E}_a \left[e^{-\gamma\kappa_u^{(a,-)} - \lambda\kappa_v^{(b,+)}} \alpha(X_{\kappa_u^{(a,-)}}) \beta(X_{\kappa_v^{(b,+)}}) 1_{\{\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)} < +\infty\}} 1_{\{X_\varepsilon \leq a\}} \right] \\ = \mathbb{E}_a \left[e^{-(\gamma+\lambda)d_\varepsilon^{(a)}} 1_{\{d_\varepsilon^{(a)} < g_\varepsilon^{(a)} + u\}} 1_{\{X_\varepsilon \leq a\}} \right] F_{\alpha,\beta,\gamma,\lambda}(a). \end{aligned} \quad (3.1)$$

Next, if $a < X_\varepsilon < b$, we obtain, applying the Markov property at the times $d_\varepsilon^{(a)}$ and $d_\varepsilon^{(b)}$:

$$\begin{aligned} \mathbb{E}_a \left[e^{-\gamma\kappa_u^{(a,-)} - \lambda\kappa_v^{(b,+)}} \alpha(X_{\kappa_u^{(a,-)}}) \beta(X_{\kappa_v^{(b,+)}}) 1_{\{\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)} < +\infty\}} 1_{\{a < X_\varepsilon < b\}} \right] \\ = \mathbb{E}_a \left[e^{-(\gamma+\lambda)d_\varepsilon^{(b)}} 1_{\{a < X_\varepsilon < b\}} 1_{\{d_\varepsilon^{(b)} < d_\varepsilon^{(a)}\}} \right] F_{\alpha,\beta,\gamma,\lambda}(b) \\ + \mathbb{E}_a \left[e^{-(\gamma+\lambda)d_\varepsilon^{(a)}} 1_{\{a < X_\varepsilon < b\}} 1_{\{d_\varepsilon^{(a)} < d_\varepsilon^{(b)}\}} \right] F_{\alpha,\beta,\gamma,\lambda}(a). \end{aligned} \quad (3.2)$$

The third and last term will have no contribution in the result :

$$\Delta_a(\varepsilon) = \mathbb{E}_a \left[e^{-\gamma \kappa_u^{(a,-)} - \lambda \kappa_v^{(b,+)}} \alpha(X_{\kappa_u^{(a,-)}}) \beta(X_{\kappa_v^{(b,+)}}) 1_{\{\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)} < +\infty\}} 1_{\{X_\varepsilon \geq b\}} \right]. \quad (3.3)$$

As a consequence, gathering (3.1), (3.2) and (3.3), we deduce that :

$$F_{\alpha,\beta,\gamma,\lambda}(a) = \frac{\mathbb{E}_a \left[e^{-(\gamma+\lambda)d_\varepsilon^{(b)}} 1_{\{a < X_\varepsilon < b\}} 1_{\{d_\varepsilon^{(b)} < d_\varepsilon^{(a)}\}} \right] F_{\alpha,\beta,\gamma,\lambda}(b) + \Delta_a(\varepsilon)}{1 - \mathbb{E}_a \left[e^{-(\gamma+\lambda)d_\varepsilon^{(a)}} 1_{\{d_\varepsilon^{(a)} < g_\varepsilon^{(a)} + u\}} 1_{\{X_\varepsilon \leq a\}} \right] - \mathbb{E}_a \left[e^{-(\gamma+\lambda)d_\varepsilon^{(a)}} 1_{\{a < X_\varepsilon < b\}} 1_{\{d_\varepsilon^{(a)} < d_\varepsilon^{(b)}\}} \right]}$$

and it remains to let $\varepsilon \downarrow 0$. Applying the Markov property and Lemma 6, we deduce that the first term in the numerator is equivalent to :

$$\begin{aligned} \mathbb{E}_a \left[e^{-(\gamma+\lambda)d_\varepsilon^{(b)}} 1_{\{a < X_\varepsilon < b\}} 1_{\{d_\varepsilon^{(b)} < d_\varepsilon^{(a)}\}} \right] &= \mathbb{E}_a \left[1_{\{a < X_\varepsilon < b\}} \mathbb{E}_{X_\varepsilon} \left[e^{-(\gamma+\lambda)T_b} 1_{\{T_b < T_a\}} \right] \right] \\ &\sim \frac{1}{\varepsilon \downarrow 0} \frac{d}{s'(a)} \mathbb{E}_x \left[e^{-(\gamma+\lambda)T_b} 1_{\{T_b < T_a\}} \right] \Big|_{x=a^+} \times \int_0^\varepsilon p(t; a, a) dt. \end{aligned}$$

Similarly, the denominator is equivalent to :

$$\frac{1}{s'(a)} \left(\frac{d}{dx} \mathbb{E}_x \left[e^{-(\gamma+\lambda)T_a} 1_{\{T_a < u\}} \right] \Big|_{x=a^-} - \frac{d}{dx} \mathbb{E}_x \left[e^{-(\gamma+\lambda)T_a} 1_{\{T_a < T_b\}} \right] \Big|_{x=a^+} \right) \times \int_0^\varepsilon p(t; a, a) dt.$$

Finally, from (2.1) and (2.2), the last term is negligible since

$$\frac{\Delta_a(\varepsilon)}{\int_0^\varepsilon p(t; a, a) dt} \leq \max_{z \in \mathbb{R}} (\alpha(z) \beta(z)) \frac{\mathbb{P}_a(T_b \leq \varepsilon)}{\int_0^\varepsilon p(t; a, a) dt} \xrightarrow{\varepsilon \downarrow 0} 0$$

and we obtain, after some simplifications,

$$\begin{aligned} F_{\alpha,\beta,\gamma,\lambda}(a) &= \frac{\frac{d}{dx} \mathbb{E}_x \left[e^{-(\gamma+\lambda)T_b} 1_{\{T_b < T_a\}} \right] \Big|_{x=a^+}}{\frac{d}{dx} \mathbb{E}_x \left[e^{-(\gamma+\lambda)T_a} 1_{\{T_a < u\}} \right] \Big|_{x=a^-} - \frac{d}{dx} \mathbb{E}_x \left[e^{-(\gamma+\lambda)T_a} 1_{\{T_a < T_b\}} \right] \Big|_{x=a^+}} F_{\alpha,\beta,\gamma,\lambda}(b) \\ &= \frac{\frac{W_{\gamma+\lambda}^{(a)'}(a)}{W_{\gamma+\lambda}^{(a)}(b)}}{\frac{\Phi_{\gamma+\lambda,-}^{(a)}(a)}{\Phi_{\gamma+\lambda,-}^{(a)}(a)} + s'(a) \int_u^{+\infty} e^{-(\gamma+\lambda)t} \nu_-^{(a)}(dt) - \frac{W_{\gamma+\lambda}^{(b)'}(a)}{W_{\gamma+\lambda}^{(b)}(a)}} F_{\alpha,\beta,\gamma,\lambda}(b) \\ &= \frac{1}{\Psi_{\gamma+\lambda}^{(-)}(b, a, u)} F_{\alpha,\beta,\gamma,\lambda}(b). \end{aligned} \quad (3.4)$$

Remark 9. By identification, the Laplace transform of T_b on the event $\{T_b < \kappa_u^{(a,-)}\}$ is given by

$$\mathbb{E}_a \left[e^{-\lambda T_b} 1_{\{T_b < \kappa_u^{(a,-)}\}} \right] = \frac{1}{\Psi_\lambda^{(-)}(b, a, u)}.$$

3.2. Study of $F_{\alpha,\beta,\gamma,\lambda}(b)$. First, if $X_\varepsilon > b$, we need to take care as whether or not the excursion above b straddling ε is longer than v . On the one hand, if this excursion is longer than v , then $\kappa_v^{(b,+)} = g_\varepsilon^{(b)} + v$, hence, applying as before the Markov property at the time $d_\varepsilon^{(b)}$:

$$\begin{aligned} & \mathbb{E}_b \left[e^{-\gamma\kappa_u^{(a,-)} - \lambda\kappa_v^{(b,+)}} \alpha(X_{\kappa_u^{(a,-)}}) \beta(X_{\kappa_v^{(b,+)}}) 1_{\{\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)} < +\infty\}} 1_{\{X_\varepsilon \geq b\}} 1_{\{d_\varepsilon^{(b)} \geq v + g_\varepsilon^{(b)}\}} \right] \\ &= \mathbb{E}_b \left[e^{-\gamma d_\varepsilon^{(b)} - \lambda(g_\varepsilon^{(b)} + v)} \beta(X_{g_\varepsilon^{(b)} + v}) 1_{\{X_\varepsilon \geq b\}} 1_{\{d_\varepsilon^{(b)} \geq v + g_\varepsilon^{(b)}\}} \right] \mathbb{E}_b \left[e^{-\gamma\kappa_u^{(a,-)}} \alpha(X_{\kappa_u^{(a,-)}}) 1_{\{\kappa_u^{(a,-)} < +\infty\}} \right]. \end{aligned}$$

On the other hand, if this excursion is shorter than v , then

$$\begin{aligned} & \mathbb{E}_b \left[e^{-\gamma\kappa_u^{(a,-)} - \lambda\kappa_v^{(b,+)}} \alpha(X_{\kappa_u^{(a,-)}}) \beta(X_{\kappa_v^{(b,+)}}) 1_{\{\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)} < +\infty\}} 1_{\{X_\varepsilon \geq b\}} 1_{\{d_\varepsilon^{(b)} < v + g_\varepsilon^{(b)}\}} \right] \\ &= \mathbb{E}_b \left[e^{-(\gamma+\lambda)d_\varepsilon^{(b)}} 1_{\{X_\varepsilon \geq b\}} 1_{\{d_\varepsilon^{(b)} < v + g_\varepsilon^{(b)}\}} \right] F_{\alpha,\beta,\gamma,\lambda}(b). \end{aligned}$$

Next, if $a < X_\varepsilon < b$, we obtain, applying the Markov property at the times $d_\varepsilon^{(a)}$ and $d_\varepsilon^{(b)}$:

$$\begin{aligned} & \mathbb{E}_b \left[e^{-\lambda\kappa_v^{(b,+)} - \gamma\kappa_u^{(a,-)}} \alpha(X_{\kappa_u^{(a,-)}}) \beta(X_{\kappa_v^{(b,+)}}) 1_{\{\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)} < +\infty\}} 1_{\{a < X_\varepsilon < b\}} \right] \\ &= \mathbb{E}_b \left[e^{-(\gamma+\lambda)d_\varepsilon^{(b)}} 1_{\{a < X_\varepsilon < b\}} 1_{\{d_\varepsilon^{(b)} < d_\varepsilon^{(a)}\}} \right] F_{\alpha,\beta,\gamma,\lambda}(b) \\ &+ \mathbb{E}_b \left[e^{-(\gamma+\lambda)d_\varepsilon^{(a)}} 1_{\{a < X_\varepsilon < b\}} 1_{\{d_\varepsilon^{(a)} < d_\varepsilon^{(b)}\}} \right] F_{\alpha,\beta,\gamma,\lambda}(a). \quad (3.5) \end{aligned}$$

Finally, the last term when $X_\varepsilon \leq a$ is again negligible :

$$\Delta_b(\varepsilon) = \mathbb{E}_b \left[e^{-\gamma\kappa_u^{(a,-)} - \lambda\kappa_v^{(b,+)}} \alpha(X_{\kappa_u^{(a,-)}}) \beta(X_{\kappa_v^{(b,+)}}) 1_{\{\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)} < +\infty\}} 1_{\{X_\varepsilon \leq a\}} \right].$$

Passing to the limit as $\varepsilon \downarrow 0$ and proceeding as in Proposition 7, we deduce that

$$\begin{aligned} & \Psi_{\gamma+\lambda}^{(+)}(a, b, v) F_{\alpha,\beta,\gamma,\lambda}(b) - F_{\alpha,\beta,\gamma,\lambda}(a) \\ &= e^{-\lambda v - \gamma v} \frac{W_{\gamma+\lambda}^{(b)}(a)}{\omega_{\gamma+\lambda}} \mathbb{E}_b \left[e^{-\gamma\kappa_u^{(a,-)}} \alpha(X_{\kappa_u^{(a,-)}}) 1_{\{\kappa_u^{(a,-)} < +\infty\}} \right] \mathbb{E}_b^{\uparrow, v} \left[\frac{\Phi_{\gamma,+}(X_v)}{\Phi_{\gamma,+}(b)} \beta(X_v) \right] \nu_+^{(b)}[v, +\infty] \end{aligned}$$

and the expression for $F_{\alpha,\beta,\gamma,\lambda}(b)$ follows by plugging (3.4) in this last formula. Finally, to get the general expression of Theorem 1 for $a \leq x \leq b$, one simply need to apply the Markov property

$$F_{\alpha,\beta,\gamma,\lambda}(x) = \mathbb{E}_x \left[e^{-(\gamma+\lambda)T_a} 1_{\{T_a < T_b\}} \right] F_{\alpha,\beta,\gamma,\lambda}(a) + \mathbb{E}_x \left[e^{-(\gamma+\lambda)T_b} 1_{\{T_b < T_a\}} \right] F_{\alpha,\beta,\gamma,\lambda}(b)$$

and then use (3.4). \square

4. APPLICATIONS

We briefly give a few examples where the expressions for the Lévy measures $\nu_-^{(a)}$ and $\nu_+^{(b)}$ may be written down explicitly.

4.1. **Brownian motion with drift μ .** Assume that X is a Brownian motion with drift μ . Its scale function may be chosen such that :

$$s'(x) = e^{-2\mu x}.$$

The two eigenfunctions $\Phi_{\lambda,+}$ and $\Phi_{\lambda,-}$ are given by :

$$\Phi_{\lambda,+}(x) = \exp\left(-x\left(\sqrt{2\lambda + \mu^2} + \mu\right)\right) \quad \text{and} \quad \Phi_{\lambda,-}(x) = \exp\left(x\left(\sqrt{2\lambda + \mu^2} - \mu\right)\right)$$

and

$$\mathbb{P}_x(T_a \in dt) = \frac{|a-x|}{\sqrt{2\pi}t^{3/2}} \exp\left(-\frac{(a-x-\mu t)^2}{2t}\right) dt.$$

As a consequence, the Lévy measures $\nu_{\pm}^{(a)}$ only depend on a through the scale function s :

$$\nu_{\pm}^{(a)}(dt) = e^{2\mu a} \frac{1}{\sqrt{2\pi}t^{3/2}} e^{-\frac{\mu^2}{2}t} dt,$$

and

$$\nu_{\pm}^{(a)}(\{+\infty\}) = e^{2\mu a} (|\mu| \pm \mu).$$

When $\mu = 0$, one may choose $s(x) = x$ and applying Remark 5, we recover for instance the formula for a standard Brownian motion, i.e. for $a \leq x \leq b$ (see [5]) :

$$\mathbb{P}_x(\kappa_v^{(b,+)} \leq \kappa_u^{(a,-)}) = \frac{\sqrt{u} + (x-a)\sqrt{\frac{2}{\pi}}}{\sqrt{v} + \sqrt{u} + (b-a)\sqrt{\frac{2}{\pi}}}.$$

4.2. **Three-dimensional Bessel process with drift $\mu > 0$ in the wide sense.** Let us consider the solution of the following SDE :

$$X_t = x + B_t + \mu \int_0^t \frac{\cosh(\mu X_s)}{\sinh(\mu X_s)} ds$$

where $x > 0$ and B is a standard Brownian motion. The law of X is that of a Brownian motion with drift $\mu > 0$ conditioned to stay positive, or alternatively that of a three-dimensional Bessel process with drift μ in the wide sense, see Pitman & Yor [13] or Watanabe [15]. Its scale function may be chosen such that

$$s'(x) = \frac{\mu^2}{\sinh^2(\mu x)}$$

and the two eigenfunctions are given by

$$\Phi_{\lambda,+}(x) = \mu \frac{\exp\left(-x\sqrt{2\lambda + \mu^2}\right)}{\sinh(\mu x)} \quad \text{and} \quad \Phi_{\lambda,-}(x) = \mu \frac{\sinh(x\sqrt{2\lambda + \mu^2})}{\sinh(\mu x)}.$$

As a consequence, we deduce from [6, p.258] that

$$\mathbb{P}_x(T_a \in dt)/dt = \begin{cases} \frac{\sinh(\mu a)}{\sinh(\mu x)} \frac{1}{\sqrt{2\pi t^{3/2}}} \sum_{n=-\infty}^{+\infty} ((2n+1)a - x) e^{-\frac{((2n+1)a-x)^2}{2t} - \frac{\mu^2}{2}t} & \text{if } x < a, \\ \frac{\sinh(\mu a)}{\sinh(\mu x)} \frac{(x-a)}{\sqrt{2\pi t^{3/2}}} \exp\left(-\frac{(x-a)^2}{2t} - \frac{\mu^2}{2}t\right) & \text{if } x > a, \end{cases}$$

hence we obtain

$$\nu_-^{(a)}(dt) = \frac{\sinh^2(\mu a)}{\mu^2} \frac{1}{\sqrt{2\pi t^{3/2}}} e^{-\frac{\mu^2}{2}t} \sum_{n=-\infty}^{+\infty} \left(\frac{\mu \cosh(\mu a)}{\sinh(\mu a)} 2na + 1 - \frac{4n^2 a^2}{t} \right) e^{-\frac{2n^2 a^2}{t}} dt$$

and

$$\nu_+^{(b)}(dt) = \frac{\sinh^2(\mu b)}{\mu^2} \frac{1}{\sqrt{2\pi t^{3/2}}} e^{-\frac{\mu^2}{2}t} dt.$$

Note that the measure $\nu_+^{(b)}$ admits an atom at $+\infty$ which is given by :

$$\nu_+^{(b)}(\{+\infty\}) = \frac{e^{2\mu b} - 1}{2\mu}.$$

As a consequence, we deduce for instance that :

$$\mathbb{P}_a(\kappa_u^{(a,-)} < +\infty) = \frac{\nu_-^{(a)}[u, +\infty)}{\frac{e^{2\mu a} - 1}{2\mu} + \nu_-^{(a)}[u, +\infty)}.$$

4.3. Reflected Brownian motion. Assume that $X = |B|$ with B a standard Brownian motion. The two eigenfunctions $\Phi_{\lambda,+}$ and $\Phi_{\lambda,-}$ are given by :

$$\Phi_{\lambda,+}(x) = \exp(-x\sqrt{2\lambda}) \quad \text{and} \quad \Phi_{\lambda,-}(x) = \cosh(x\sqrt{2\lambda})$$

and from Borodin & Salminen [2], the distribution of T_a admits the density

$$\mathbb{P}_x(T_a \in dt)/dt = \begin{cases} \frac{1}{\sqrt{2\pi t^{3/2}}} \sum_{n=-\infty}^{+\infty} (-1)^n ((2n+1)a + x) e^{-\frac{((2n+1)a+x)^2}{2t}} & \text{if } x < a, \\ \frac{(x-a)}{\sqrt{2\pi t^{3/2}}} \exp\left(-\frac{(x-a)^2}{2t}\right) & \text{if } x > a, \end{cases}$$

Therefore, we deduce that

$$\nu_-^{(a)}(dt) = \frac{1}{\sqrt{2\pi t^{3/2}}} \sum_{n=-\infty}^{+\infty} (-1)^{n+1} \left(\frac{4a^2 n^2}{t} - 1 \right) e^{-\frac{2a^2 n^2}{t}} dt$$

and

$$\nu_+^{(b)}(dt) = \frac{1}{\sqrt{2\pi}t^{3/2}}dt.$$

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