# Uniform approximation of $2 d$ Navier-Stokes equation by stochastic interacting particle systems <br> Franco Flandoli, Christian Olivera, Marielle Simon 

## To cite this version:

Franco Flandoli, Christian Olivera, Marielle Simon. Uniform approximation of $2 d$ Navier-Stokes equation by stochastic interacting particle systems. 2020. hal-02529632v1

HAL Id: hal-02529632

## https://inria.hal.science/hal-02529632v1

Preprint submitted on 6 Apr 2020 (v1), last revised 16 Oct 2020 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Uniform approximation of $2 d$ Navier-Stokes equation by stochastic interacting particle systems 

Franco Flandoli*, Christian Olivera ${ }^{\dagger}$, Marielle Simon ${ }^{\ddagger}$


#### Abstract

We consider an interacting particle system modeled as a system of $N$ stochastic differential equations driven by Brownian motions. We prove that the (mollified) empirical process converges, uniformly in time and space variables, to the solution of the two-dimensional Navier-Stokes equation written in vorticity form. The proofs follow a semigroup approach.


Key words and phrases: Moderatly interacting particle system, Stochastic differential equations, $2 d$ Navier-Stokes equation, Vorticity equation, Analytic semigroup. MSC2010 subject classification: 60H20, 60H10, 60F99.

## 1 Introduction

The main goal of this paper is to provide a stochastic approximation of the twodimensional Navier-Stokes equation. Precisely, we consider the following classical Cauchy problem which describes the evolution of the velocity field $u: \mathbb{R}_{+} \times \mathbb{R}^{2} \rightarrow \mathbb{R}^{2}$ of an incompressible fluid with kinematic viscosity coefficient $\nu>0$ : for any $(t, x) \in$ $\mathbb{R}_{+} \times \mathbb{R}^{2}$,

$$
\left\{\begin{align*}
\partial_{t} u(t, x) & =\nu \Delta u(t, x)-[u(t, x) \cdot \nabla] u(t, x)-\nabla p(t, x)  \tag{1}\\
\operatorname{div} u(t, x) & =0 \\
u(0, x) & =u^{\text {ini }}(x),
\end{align*}\right.
$$

[^0]where: - denotes the standard Euclidean product in $\mathbb{R}^{2}$; the unknown quantities are the velocity $u(t, x)=\left(u_{1}(t, x), u_{2}(t, x)\right) \in \mathbb{R}^{2}$ of the fluid element at time $t$ and position $x$ and the pressure $p(t, x) \in \mathbb{R}$. Such equations are attracting the attention of a large scientific community, with a large amount of publications in the literature. Since this system is very famous, we do not comment here on its derivation and rather refer to the monographs [25] and [26]. For recent developments, see also [17].

The associated (scalar) vorticity field $\xi=\partial_{1} u_{2}-\partial_{2} u_{1}: \mathbb{R}^{2} \rightarrow \mathbb{R}$ satisfies a remarkably simple equation of convection-diffusion propagation, namely

$$
\begin{equation*}
\partial_{t} \xi+u \cdot \nabla \xi=\nu \Delta \xi, \quad x \in \mathbb{R}^{2}, t>0 \tag{2}
\end{equation*}
$$

The velocity field $u(t, x)$ can be reconstructed from the vorticity distribution $\xi(t, x)$ by the convolution with the Biot-Savart kernel $K$ as:

$$
\begin{equation*}
u(t, x)=(K * \xi(t, \cdot))(x)=\frac{1}{\pi} \int_{\mathbb{R}^{2}} \frac{(x-y)^{\perp}}{|x-y|^{2}} \xi(t, y) d y \tag{3}
\end{equation*}
$$

where $\left(x_{1}, x_{2}\right)^{\perp}:=\left(-x_{2}, x_{1}\right)$. It is well-known [19] that there is a constant $c_{K}>0$ such that: for any $\xi \in \mathbb{L}^{1}\left(\mathbb{R}^{2}\right) \cap \mathbb{L}^{\infty}\left(\mathbb{R}^{2}\right)$

$$
\begin{equation*}
\|K * \xi\|_{\mathbb{L}^{\infty}} \leqslant c_{K}\left(\|\xi\|_{\mathbb{L}^{1}}+\|\xi\|_{\mathbb{L}^{\infty}}\right), \tag{4}
\end{equation*}
$$

where $\|\cdot\|_{\mathbb{L}^{p}}$ denotes the usual $\mathbb{L}^{p}\left(\mathbb{R}^{2}\right)$ norm. The existence and uniqueness of a global solution to (2) for an $\mathbb{L}^{p}$ - or measure-valued - initial condition has been proved ; on that subject we refer the reader to $[1,10,11,16]$.

The question of a particle approximation to the $2 d$ Navier-Stokes equation has been already considered in the literature, as recalled in more detail in Section 1.1 below. The aim of this paper is to provide a new rigorous approximation of the vorticity field $\xi$ by stochastic particle systems, stronger than others: contrary to the previous works where only the empirical measure of the density of particles is shown to converge, here we also prove that a mollified empirical measure converges uniformly. More precisely, we consider the $N$-particle dynamics described, for each $N \in \mathbb{N}$, by the following system of coupled stochastic differential equations in $\mathbb{R}^{2}$ : for any $i=1, \ldots, N$,

$$
\begin{equation*}
d X_{t}^{i, N}=F\left(\frac{1}{N} \sum_{k=1}^{N}\left(K * V^{N}\right)\left(X_{t}^{i, N}-X_{t}^{k, N}\right)\right) d t+\sqrt{2 \nu} d W_{t}^{i} \tag{5}
\end{equation*}
$$

where:

- for a given $M>0$ chosen ahead (see Theorem 3 below), the function $F$ is given by

$$
\begin{equation*}
F:\binom{x_{1}}{x_{2}} \mapsto\binom{\left(x_{1} \wedge M\right) \vee(-M)}{\left(x_{2} \wedge M\right) \vee(-M)} ; \tag{6}
\end{equation*}
$$

- $\left\{W_{t}^{i}, \quad i \in \mathbb{N}\right\}$ is a family of independent standard Brownian motions on $\mathbb{R}^{2}$ defined on a filtered probability space $\left(\Omega, \mathcal{F}, \mathcal{F}_{t}, \mathbb{P}\right)$;
- the interaction potential $V^{N}: \mathbb{R}^{2} \rightarrow \mathbb{R}_{+}$is continuous and will be specified later on.

Finally, * stands for the standard convolution product, and $\wedge$ (resp. $\vee$ ) is the usual notation for the minimum (resp. maximum) of two real numbers.

Now let us define the empirical process of this particle system as

$$
S_{t}^{N}:=\frac{1}{N} \sum_{i=1}^{N} \delta_{X_{t}^{i, N}}
$$

which is a (scalar) measure-valued process associated to the $\mathbb{R}^{2}$-valued processes $\left\{t \mapsto X_{t}^{i, N}\right\}_{i=1, \ldots, N}$. Above, $\delta_{a}$ is the delta Dirac measure concentrated at $a \in \mathbb{R}^{2}$. For any test function $\phi: \mathbb{R}^{2} \rightarrow \mathbb{R}$, we use the standard notation

$$
\left\langle S_{t}^{N}, \phi\right\rangle:=\frac{1}{N} \sum_{i=1}^{N} \phi\left(X_{t}^{i, N}\right) .
$$

Our interest lies in the investigation of the dynamical process $t \mapsto S_{t}^{N}$ in the large particle limit $N \rightarrow \infty$.

The dynamics of the empirical measure is determined by the Itô formula, which reads as follows: for any test function $\phi: \mathbb{R}^{2} \rightarrow \mathbb{R}$ of class $C^{2}$, the empirical measure $S_{t}^{N}$ satisfies

$$
\begin{align*}
\left\langle S_{t}^{N}, \phi\right\rangle=\left\langle S_{0}^{N}, \phi\right\rangle & +\int_{0}^{t}\left\langle S_{s}^{N}, F\left(K * V^{N} * S_{s}^{N}\right) \cdot \nabla \phi\right\rangle d s \\
& +\nu \int_{0}^{t}\left\langle S_{s}^{N}, \Delta \phi\right\rangle d s+\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} \nabla \phi\left(X_{s}^{i, N}\right) \cdot d W_{s}^{i} \tag{7}
\end{align*}
$$

Our main result is the uniform convergence (in the space and time variables) of the mollified empirical measure

$$
g_{t}^{N}:=V^{N} * S_{t}^{N}: x \in \mathbb{R}^{2} \mapsto \int_{\mathbb{R}^{2}} V^{N}(x-y) d S_{t}^{N}(y)
$$

to the solution of the Navier-Stokes equation written in vorticity form, given below in Theorem 3. To prove that result, we follow the new approach presented in [5] and then in $[6,7,23]$, based on semigroup theory. Our source of inspiration has been the works of Oelschläger [22] and Jourdain and Méléard [15], where stochastic approximations of PDE's are investigated. We assume that the initial vorticity satisfies $\xi^{\text {ini }} \in \mathbb{L}^{1}\left(\mathbb{R}^{2}\right) \cap \mathbb{L}^{\infty}\left(\mathbb{R}^{2}\right)$, but we believe that our approach can be adapted for more irregular initial data, for instance when $\xi^{\text {ini }}$ belongs to $\mathbb{L}^{1}\left(\mathbb{R}^{2}\right) \cap \mathbb{L}^{p}\left(\mathbb{R}^{2}\right)$ with $p \in(2, \infty)$.

### 1.1 Related works

Thirty years ago, Chorin [4] proposed a heuristic probabilistic algorithm to numerically simulate the solution of the Navier-Stokes equation in two dimensions, by approximating the (scalar) vorticity function, involving cut-off kernels, by random interacting "point vortices". The convergence of Chorin's vortex method was first mathematically proved in 1982 by Marchioro and Pulvirenti [18], who interpreted the vortex equation in two dimensions with bounded and integrable initial condition as a generalized McKean-Vlasov equation. Later Méléard [19, 20] improved the results and showed the convergence in the path space of the empirical measures of the interacting particle system. Fontbona [8] then generalized that result in dimension $d=3$. In 1987, Osada [21] proved a propagation of chaos result for an interacting particle system which approximates the solution of the McKean-Vlasov equation, without cut-off, by an analytical method based on generators of generalized divergence form, but only for large viscosities and bounded density initial data. The convergence of empirical measure and propagation of chaos have then been considered under more general assumptions and with innovative techniques of entropy and Fisher information by Fournier, Hauray and Mischler [9]. Finally we mention that recently, Jabin and Wang [14] showed that a mean field approximation converges to the solution of the Navier-Stokes equation written in vorticity form.

Besides, let us note that Marchioro and Pulvirenti [18] wished to describe a unified approach for both Navier-Stokes and Euler equations, and for that reason they did not fully exploit the stochastic nature of the Navier-Stokes equation, which is exactly what we are doing here. In fact, we strongly exploit the Brownian perturbation of the system, and therefore we cannot cover the results obtained for the Euler equation as in [18].

### 1.2 Notations and results

Before concluding the introduction, late us state the main results of this work. We first need to introduce some of our notations, which are listed below:

- For any measure space $(S, \Sigma, \mu)$, the standard $\mathbb{L}^{p}(S)$-spaces of real-valued functions, with $p \in[1, \infty]$, are provided with their usual norm denoted by $\|\cdot\|_{\mathbb{L}^{p}(S)}$, or $\|\cdot\|_{\mathbb{L}^{p}}$ whenever the space $S$ will be clear to the reader. With a little abuse of notation, and as soon as no confusion regarding the space $S$ arises, we denote by $\langle f, g\rangle$ the inner product on $\mathbb{L}^{2}(S)$ between two functions $f$ and $g$. In more general cases, if the functions take values in some space $X$, the notation will become $\mathbb{L}^{p}(S ; X)$. Finally, the norm $\|\cdot\|_{\mathbb{L}^{p}(S) \rightarrow \mathbb{L}^{p}(S)}$ is the usual operator norm.
- For any $\varepsilon \in \mathbb{R}$, and $d \in \mathbb{N}$, we denote by $\mathbb{H}_{2}^{\varepsilon}\left(\mathbb{R}^{d}\right)$ the Bessel potential space

$$
\mathbb{H}_{2}^{\varepsilon}\left(\mathbb{R}^{d}\right):=\left\{u \in \mathbb{L}^{2}\left(\mathbb{R}^{d}\right) ; \mathcal{F}^{-1}\left(\left(1+|\cdot|^{2}\right)^{\frac{\varepsilon}{2}} \mathcal{F} u(\cdot)\right) \in \mathbb{L}^{2}\left(\mathbb{R}^{d}\right)\right\},
$$

where $\mathcal{F} u$ denotes the Fourier transform of $u$. These spaces are endowed with their norm

$$
\|u\|_{\varepsilon, 2}^{2}:=\left\|\mathcal{F}^{-1}\left(\left(1+|\cdot|^{2}\right)^{\frac{\varepsilon}{2}} \mathcal{F} u(\cdot)\right)\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{d}\right)}^{2}<\infty .
$$

Note that

$$
\|u\|_{0,2}=\|u\|_{\mathbb{L}^{2}\left(\mathbb{R}^{d}\right)} \quad \text { and, for any } \varepsilon \leqslant 0,\|u\|_{\varepsilon, 2} \leqslant\|u\|_{0,2} .
$$

- The space of smooth real-valued functions with compact support in $\mathbb{R}^{d}$ is denoted by $C_{0}^{\infty}\left(\mathbb{R}^{d}\right)$. The space of functions of class $C^{k}$ with $k \in \mathbb{N}$ is denoted by $C^{k}\left(\mathbb{R}^{d}\right)$. Finally, the space of bounded functions is denoted by $C_{b}\left(\mathbb{R}^{d}\right)$.

Now, let us give our main assumptions: first, we need to be more precise about the interaction potential $V^{N}$; second, recall that we are interested in the large $N$ limit of the process $t \mapsto S_{t}^{N}$, and we therefore need to specify its initial condition, which is supposed to be random. The expectation with respect to $\mathbb{P}$ is denoted by $\mathbb{E}$. We say that a function $f: \mathbb{R}^{2} \rightarrow \mathbb{R}_{+}$is a probability density if $\int_{\mathbb{R}^{2}} f(x) d x=1$.

Assumption 1 We assume that there exists a probability density $V: \mathbb{R}^{2} \rightarrow \mathbb{R}_{+}$, and a parameter $\beta \in[0,1]$, such that

1. for any $x \in \mathbb{R}^{2}, V^{N}(x)=N^{2 \beta} V\left(N^{\beta} x\right)$;
2. $V \in C_{0}^{\infty}\left(\mathbb{R}^{2}\right)$;
3. there exists $\alpha>1$ such that, for any $p>0$,

$$
\begin{equation*}
\sup _{N \in \mathbb{N}} \mathbb{E}\left[\left\|V^{N} * S_{0}^{N}\right\|_{\alpha, 2}^{p}\right]<\infty \tag{8}
\end{equation*}
$$

4. there exists $\xi^{\mathrm{ini}} \in \mathbb{L}^{1}\left(\mathbb{R}^{2}\right) \cap \mathbb{L}^{\infty}\left(\mathbb{R}^{2}\right)$ such that the sequence of measures $\left\{S_{0}^{N}\right\}_{N}$ weakly converges to the initial measure $\xi^{\mathrm{ini}}(\cdot) d x$, as $N \rightarrow \infty$, in probability.
5. finally, the parameters $(\beta, \alpha)$ satisfy:

$$
\begin{equation*}
0<\beta<\frac{1}{2+2 \alpha}<\frac{1}{4} \tag{9}
\end{equation*}
$$

Remark 2 In [7] the authors provide sufficient conditions for the validity of (8).
In all what follows we fix a time horizon $T \geqslant 0$.
Theorem 3 We assume Assumption 1 and we consider the particle system (5) with the parameter $M$ which satisfies

$$
\begin{equation*}
M \geqslant c_{K}\left(1+\left\|\xi^{\mathrm{ini}}\right\|_{\mathbb{L}^{\infty}}\right) \tag{10}
\end{equation*}
$$

where $c_{K}$ has been defined in (4).
Then, for every $\eta \in(1, \alpha)$, the sequence of processes $\left\{t \mapsto g_{t}^{N}=V^{N} * S_{t}^{N}\right\}_{N \in \mathbb{N}}$ converges in probability with respect to the

- weak topology of $\mathbb{L}^{2}\left([0, T] ; \mathbb{H}_{2}^{\alpha}\left(\mathbb{R}^{2}\right)\right)$,
- strong topology of $C\left([0, T] ; \mathbb{H}_{2, \text { loc }}^{\eta}\left(\mathbb{R}^{2}\right)\right)$,
as $N \rightarrow \infty$, to the unique weak solution of the partial differential equation (PDE)

$$
\left\{\begin{array}{l}
\partial_{t} \xi+\operatorname{div}(\xi(K * \xi))=\nu \Delta \xi  \tag{11}\\
\xi(0, x)=\xi^{\operatorname{ini}}(x), \quad x \in \mathbb{R}^{2}, t>0
\end{array}\right.
$$

Namely, for any real-valued test function $\phi \in C_{0}^{\infty}\left(\mathbb{R}^{2}\right)$ and any $t \geqslant 0$, it holds

$$
\begin{equation*}
\langle\xi(t, \cdot), \phi\rangle=\left\langle\xi^{\mathrm{ini}}, \phi\right\rangle+\int_{0}^{t}\langle\xi(s, \cdot),(K * \xi)(s, \cdot) \cdot \nabla \phi\rangle d s+\nu \int_{0}^{t}\langle\xi(s, \cdot), \Delta \phi\rangle d s \tag{12}
\end{equation*}
$$

Remark 4 Note that the limiting PDE (11) does not depend on the value of the parameter $\beta \in\left(0, \frac{1}{2+2 \alpha}\right)$.

Remark 5 The previous result implies, by the Sobolev embeddings' Theorem (see [28, Section 2] for instance), the strong convergence in $C([0, T] \times K)$ for every compact set $K \subset \mathbb{R}^{2}$.

Here it follows an outline of the paper: we start in Section 2 with an exposition of the strategy to prove Theorem 3. We will prove the technical estimates in Section 3. We chose to investigate in details the case where $\xi^{\text {ini }}$ is a probability density, in particular is non-negative and then, in Appendix A, we show that the same result holds in the general case $\xi^{\text {ini }} \in \mathbb{L}^{1}\left(\mathbb{R}^{2}\right) \cap \mathbb{L}^{\infty}\left(\mathbb{R}^{2}\right)$, without assumptions on the sign and the value of $\int \xi^{\text {ini }}(x) d x$, by a simple decoupling argument. Finally, concerning the importance and geometrical interpretation of the uniform convergence of the mollified empirical measure, we have included a short discussion in Appendix B. In particular, the uniform convergence does not follow from the weak convergence of the empirical measure.

## 2 Strategy of the proof of Theorem 3

There are two main steps to derive the convergence result stated in Theorem 3:

1. First, we write the mild formulation of the identity satisfied by

$$
\begin{equation*}
g_{t}^{N}(x)=\left(V^{N} * S_{t}^{N}\right)(x)=\int_{\mathbb{R}^{2}} V^{N}(x-y) d S_{t}^{N}(y) \tag{13}
\end{equation*}
$$

see Section 2.2. We will obtain a "closed" inequality (note that $g^{N}$ already appears in the right hand side of (7)), and then prove two uniform bounds, see Proposition 6 and Proposition 7. To that aim we will use two main properties of the function $F$ : first $F$ is Lipschitz continuous, and second it is bounded, precisely $|F(x)-F(y)| \leqslant|x-y|$ and $\|F\|_{\mathbb{L}^{\infty}\left(\mathbb{R}^{2}\right)} \leqslant M$.
2. Then we apply compacteness arguments and Sobolev embeddings to have subsequences which converge so as to pass to the limit, see Section 2.3, Section 2.4 and Section 2.6.

The support of a function $f$ is denoted by $\operatorname{Supp} f$. When a constant $C$ will depend on some parameter, say $\alpha$, this will be highlighted in index by $C_{\alpha}$, but the constant may change from line to line.

### 2.1 Analytic semigroup

Let us first introduce the operator

$$
A: \mathcal{D}(A) \subset \mathbb{L}^{2}\left(\mathbb{R}^{d}\right) \rightarrow \mathbb{L}^{2}\left(\mathbb{R}^{d}\right)
$$

defined as $A f=\nu \Delta f$. It is the infinitesimal generator of an analytic semigroup (the heat semigroup) in $\mathbb{L}^{2}\left(\mathbb{R}^{d}\right)$ (see for instance [24]). We denote this semigroup by $\left\{e^{t A}, t \geqslant 0\right\}$, which is simply given by

$$
\left(e^{t A} f\right)(x)=\int_{\mathbb{R}^{d}} \frac{1}{(4 \nu \pi t)^{d / 2}} e^{-|x-y|^{2} /(4 \nu t)} f(y) d y, \quad f \in \mathbb{L}^{2}\left(\mathbb{R}^{d}\right)
$$

Moreover, denoting by I the identity operator, we know that, for any $\varepsilon \in \mathbb{R}$, the domain of the operator $(\mathrm{I}-A)^{\varepsilon / 2}$ is given by

$$
\mathcal{D}\left((\mathrm{I}-A)^{\varepsilon / 2}\right)=\mathbb{H}_{2}^{\varepsilon}\left(\mathbb{R}^{d}\right)
$$

with equivalent norms. Recall also from [24] that, for every $\varepsilon>0$, and $T>0$, there is a constant $C_{\varepsilon, T, \nu}>0$ such that, for any $t \in(0, T]$,

$$
\begin{equation*}
\left\|(\mathrm{I}-A)^{\varepsilon} e^{t A}\right\|_{\mathbb{L}^{2} \rightarrow \mathbb{L}^{2}} \leqslant \frac{C_{\varepsilon, T, \nu}}{t^{\varepsilon}} \tag{14}
\end{equation*}
$$

We are now ready to prove Theorem 3, therefore the dimension is from now on $d=2$.

### 2.2 The equation for $V^{N} * S_{t}^{N}$ in mild form

We want to deduce an identity for $g_{t}^{N}(x)$ from (7). For every $x \in \mathbb{R}^{2}$ take, in identity (7), the test function $\phi_{x}(y)=V^{N}(x-y)$. We get (recall the definition (13) of $g_{t}^{N}$ )

$$
\begin{align*}
g_{t}^{N}(x)=g_{0}^{N}(x) & +\int_{0}^{t}\left\langle S_{s}^{N}, F\left(K * g_{s}^{N}\right) \cdot \nabla V^{N}(x-\cdot)\right\rangle d s \\
& +\nu \int_{0}^{t} \Delta g_{s}^{N}(x) d s+\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} \nabla V^{N}\left(x-X_{s}^{i, N}\right) \cdot d W_{s}^{i} \tag{15}
\end{align*}
$$

In the sequel, let us write for the sake of clarity,

$$
\left\langle S_{s}^{N}, F\left(K * g_{s}^{N}\right) \cdot \nabla V^{N}(x-\cdot)\right\rangle=:\left(\nabla V^{N} *\left(F\left(K * g_{s}^{N}\right) S_{s}^{N}\right)\right)(x)
$$

and similarly for analogous expressions. Hence we can write (using the same idea as in [5]):

$$
\begin{aligned}
g_{t}^{N}=e^{t A} g_{0}^{N} & +\int_{0}^{t} e^{(t-s) A}\left(\nabla V^{N} *\left(F\left(K * g_{s}^{N}\right) S_{s}^{N}\right)\right) d s \\
& +\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} e^{(t-s) A}\left(\nabla V^{N}\left(\cdot-X_{s}^{i, N}\right)\right) \cdot d W_{s}^{i} .
\end{aligned}
$$

By inspection in the convolution explicit formula for $e^{(t-s) A}$, one can see that

$$
e^{(t-s) A} \nabla f=\nabla e^{(t-s) A} f
$$

and then one can use the semigroup property, so as to write

$$
\begin{align*}
g_{t}^{N}=e^{t A} g_{0}^{N} & +\int_{0}^{t} \nabla e^{(t-s) A}\left(V^{N} *\left(F\left(K * g_{s}^{N}\right) S_{s}^{N}\right)\right) d s \\
& +\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} e^{(t-s) A}\left(\nabla V^{N}\left(\cdot-X_{s}^{i, N}\right)\right) \cdot d W_{s}^{i} \tag{16}
\end{align*}
$$

In the following we will prove two important bounds:
Proposition 6 We assume Assumption 1. Let $\varepsilon>0$ and $p \geqslant 2$. Then there exists a positive constant $C_{\varepsilon, T, M, \nu, p}$ such that, for all $t \in(0, T]$ and $N \in \mathbb{N}$, it holds:

$$
\begin{equation*}
\mathbb{E}\left[\left\|(\mathrm{I}-A)^{\varepsilon / 2} g_{t}^{N}\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)}^{p}\right] \leqslant C_{\varepsilon, T, M, \nu, p} \tag{17}
\end{equation*}
$$

Proposition 7 We assume Assumption 1. Let $\gamma \in\left(0, \frac{1}{2}\right)$ and $q \geqslant 2$. There exists a positive constant $C_{T, M, \nu, q}$ such that, for any $N \in \mathbb{N}$, it holds:

$$
\begin{equation*}
\mathbb{E}\left[\int_{0}^{T} \int_{0}^{T} \frac{\left\|g_{t}^{N}-g_{s}^{N}\right\|_{-2,2}^{q}}{|t-s|^{1+q \gamma}} d s d t\right] \leqslant C_{T, M, \nu, q} \tag{18}
\end{equation*}
$$

The proofs of Proposition 6 and Proposition 7 are postponed to Section 3.1 and Section 3.2, respectively.

### 2.3 Criterion of compactness

In this subsection we follow the arguments of [7, Section 3.1]. We start by constructing a space on which the sequence of the probability laws of $g^{N}$ is tight. For that purpose, we need to consider, for $\varepsilon \in \mathbb{R}$, the Sobolev spaces $\mathbb{W}^{\varepsilon, 2}\left(\mathbb{R}^{d}\right)$ defined for instance in the monograph [28], which contains the material needed here. Whenever $\varepsilon>0$, the restriction of $f \in \mathbb{W}^{\varepsilon, 2}\left(\mathbb{R}^{d}\right)$ to a ball $\mathcal{B}_{R} \subset \mathbb{R}^{d}$ of radius $R>0$ is in $\mathbb{W}^{\varepsilon, 2}\left(\mathcal{B}_{R}\right)$, and the spaces $\mathbb{H}_{2}^{\varepsilon}\left(\mathbb{R}^{d}\right)$ and $\mathbb{W}^{\varepsilon, 2}\left(\mathbb{R}^{d}\right)$ have equivalent norms.

We exploit Corollary 9 of Simon [27], using as far as possible the notations of that paper. Given a ball $\mathcal{B}_{R} \subset \mathbb{R}^{2}$, taken $\alpha>1$ (as in Theorem 3), and $1<\eta<\alpha$, we consider the space

$$
X:=\mathbb{W}^{\alpha, 2}\left(\mathcal{B}_{R}\right), \quad B:=\mathbb{W}^{\eta, 2}\left(\mathcal{B}_{R}\right), \quad Y:=\mathbb{W}^{-2,2}\left(\mathcal{B}_{R}\right)
$$

One can check that $X \subset B \subset Y$, with compact dense embedding. Moreover, we have the interpolation inequality: for any $f \in X$

$$
\|f\|_{B} \leqslant C_{R}\|f\|_{X}^{1-\theta}\|f\|_{Y}^{\theta}
$$

with

$$
\begin{equation*}
\theta:=\frac{\alpha-\eta}{2+\alpha} . \tag{19}
\end{equation*}
$$

Now taking, in the notations of [27, Corollary 9]: $s_{0}:=0, s_{1}:=\gamma \in\left(0, \frac{1}{2}\right)$, and choosing $p, q \geqslant 2$ such that :

$$
\begin{equation*}
s_{1} q=\gamma q>1, \quad \text { and } \quad s_{\theta}:=\theta s_{1}=\theta \gamma>\frac{1-\theta}{p}+\frac{\theta}{q} \tag{20}
\end{equation*}
$$

then the corollary tells us that the space $\mathbb{L}^{p}([0, T] ; X) \cap \mathbb{W}^{\gamma, q}([0, T] ; Y)$ is relatively compact in $C([0, T] ; B)$.

Therefore, for any $\gamma \in\left(0, \frac{1}{2}\right)$, for the parameter $\alpha>1$ given by point 3. of Assumption 1, and for $p, q \geqslant 2$ which satisfy (20), we now consider the space

$$
\mathfrak{Y}_{0}:=\mathbb{L}^{p}\left([0, T] ; \mathbb{H}_{2}^{\alpha}\right) \cap \mathbb{W}^{\gamma, q}\left([0, T] ; \mathbb{H}_{2}^{-2}\right)
$$

We use the Fréchet topology on $C\left([0, T] ; \mathbb{W}_{\text {loc }}^{\eta, 2}\left(\mathbb{R}^{2}\right)\right)$ defined as

$$
d(f, g)=\sum_{n=1}^{\infty} 2^{-n}\left(1 \wedge \sup _{t \in[0, T]}\|(f-g)(t, \cdot)\|_{\mathbb{W}^{\eta, 2}\left(\mathcal{B}_{n}\right)}^{2}\right)
$$

where $\wedge$ denotes the infimum. From the above, we conclude that $\mathfrak{Y}_{0}$ is compactly embedded into $C\left([0, T] ; \mathbb{W}_{\text {loc }}^{\eta, 2}\right)$ for any $1<\eta<\alpha$. Finally, let us denote by $\mathbb{L}_{w}^{2}$ the spaces $\mathbb{L}^{2}$ endowed with the weak topology. We obtain that $\mathfrak{Y}_{0}$ is compactly embedded into

$$
\begin{equation*}
\mathfrak{Y}:=\mathbb{L}_{w}^{2}\left([0, T] ; \mathbb{H}_{2}^{\alpha}\right) \cap C\left([0, T] ; \mathbb{W}_{\mathrm{loc}}^{\eta, 2}\right) \tag{21}
\end{equation*}
$$

Note that

$$
C\left([0, T] ; \mathbb{W}_{\mathrm{loc}}^{\eta, 2}\right) \subset C([0, T] ; C(D))
$$

for every regular bounded domain $D \subset \mathbb{R}^{2}$.
Let us now go back to the sequence of processes $\left\{g^{N}\right\}_{N}$, for which we have proved several estimates. The Chebyshev inequality ensures that

$$
\mathbb{P}\left(\left\|g^{N}\right\|_{\mathfrak{Z}_{0}}^{2}>R\right) \leqslant \frac{\mathbb{E}\left[\left\|g^{N}\right\|_{\mathfrak{Z}_{0}}^{2}\right]}{R}, \quad \text { for any } R>0
$$

Thus by Proposition 6 and Proposition 7 (since $p, q \geqslant 2$ ), we obtain

$$
\mathbb{P}\left(\left\|g^{N}\right\|_{\mathfrak{Y}_{0}}^{2}>R\right) \leqslant \frac{C}{R}, \quad \text { for any } R>0, N \in \mathbb{N}
$$

The process $t \in[0, T] \mapsto g_{t}^{N}$ defines a probability $\mathbf{P}_{N}$ on $\mathfrak{Y}$. Fix $\varepsilon>0$. Last inequality implies that there exists a bounded set $B_{\varepsilon} \in \mathfrak{Y}_{0}$ such that $\mathbf{P}_{N}\left(B_{\varepsilon}\right)<1-\varepsilon$ for all $N$, and therefore, from the previous argument, there exists a compact set $K_{\varepsilon} \subset \mathfrak{Y}$ such that $\mathbf{P}_{N}\left(K_{\varepsilon}\right)<1-\varepsilon$.

Finally, denote by $\left\{\mathrm{L}^{N}\right\}_{N \in \mathbb{N}}$ the laws of the processes $\left\{g^{N}\right\}_{N \in \mathbb{N}}$ on $\mathfrak{Y}_{0}$, we have proved that $\left\{\mathrm{L}^{N}\right\}_{N \in \mathbb{N}}$ is tight in $\mathfrak{Y}$, hence relatively compact, by Prohorov's Theorem. From every subsequence of $\left\{\mathrm{L}^{N}\right\}_{N \in \mathbb{N}}$ it is possible to extract a further subsequence which converges to a probability measure L on $\mathfrak{Y}$. Moreover by a theorem of Skorokhod (see [13, Theorem 2.7]), we are allowed, eventually after choosing a suitable probability space where all our random variables can be defined, to assume

$$
\begin{equation*}
g^{N} \rightarrow \xi \quad \text { in } \mathfrak{Y}, \quad \text { a.s. } \tag{22}
\end{equation*}
$$

where the law of $\xi$ is L .

### 2.4 Passing to the limit

Next step is to characterize the limit. First, recall formula (15), which reads

$$
\begin{aligned}
g_{t}^{N}(x)=g_{0}^{N}(x) & +\int_{0}^{t}\left\langle S_{s}^{N}, F\left(K * g_{s}^{N}\right) \cdot \nabla V^{N}(x-\cdot)\right\rangle d s \\
& +\nu \int_{0}^{t} \Delta g_{s}^{N}(x) d s+\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} \nabla V^{N}\left(x-X_{s}^{i, N}\right) \cdot d W_{s}^{i}
\end{aligned}
$$

Taking a test function $\phi: \mathbb{R}^{2} \rightarrow \mathbb{R}$ we have

$$
\begin{align*}
\left\langle g_{t}^{N}, \phi\right\rangle=\left\langle g_{0}^{N}, \phi\right\rangle & +\int_{0}^{t}\left\langle S_{s}^{N}, F\left(K * g_{s}^{N}\right) \cdot \nabla\left(V^{N} * \phi\right)\right\rangle d s \\
& +\nu \int_{0}^{t}\left\langle g_{s}^{N}, \Delta \phi\right\rangle(x) d s+\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} \nabla\left(V^{N} * \phi\right)\left(X_{s}^{i, N}\right) \cdot d W_{s}^{i} \tag{23}
\end{align*}
$$

It is clear from (22) that

$$
\begin{equation*}
\left\langle g_{t}^{N}, \phi\right\rangle \underset{N \rightarrow \infty}{\longrightarrow}\langle\xi, \phi\rangle, \quad\left\langle g_{0}^{N}, \phi\right\rangle \underset{N \rightarrow \infty}{\longrightarrow}\left\langle\xi^{\mathrm{ini}}, \phi\right\rangle \tag{24}
\end{equation*}
$$

and also that

$$
\begin{equation*}
\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} \nabla\left(V^{N} * \phi\right)\left(X_{s}^{i, N}\right) \cdot d W_{s}^{i} \underset{N \rightarrow \infty}{\longrightarrow} 0 \tag{25}
\end{equation*}
$$

We now claim that

$$
\begin{equation*}
\lim _{N \rightarrow \infty} \int_{0}^{t}\left\langle S_{s}^{N}, F\left(K * g_{s}^{N}\right) \cdot \nabla\left(V^{N} * \phi\right)\right\rangle d s=\int_{0}^{t} \int_{\mathbb{R}^{2}} \xi(s, x) F(K * u) \cdot \nabla \phi(x) d x d s \tag{26}
\end{equation*}
$$

Proof of (26). First, we observe that $\left\|g_{t}^{N}\right\|_{\mathbb{L}^{1}}=1$ and $g^{N}$ is uniformly bounded in $\mathbb{L}^{2}\left([0, T] ; \mathbb{H}_{2}^{\eta}\right)$ for any $\eta>1$. Then by the Sobolev embeddings' Theorem (see [28, Section 2.8.1]), we have that $g^{N}$ is uniformly bounded in $\mathbb{L}^{2}\left([0, T] ; C_{b}(\mathbb{R})\right)$.

By interpolation we also know that $g^{N} \in \mathbb{L}^{2}\left([0, T] ; \mathbb{L}^{p}\left(\mathbb{R}^{2}\right)\right)$ for any $p \in[1,+\infty]$, and there is a constant $C_{K}>0$ such that

$$
\begin{equation*}
\left\|K * g_{t}^{N}\right\|_{\mathbb{L}^{p}} \leqslant C_{K}\left\|g_{t}^{N}\right\|_{\mathbb{L}^{q}} \leqslant\left\|g_{t}^{N}\right\|_{\mathbb{L}^{\infty}}^{1-\frac{1}{q}} \tag{27}
\end{equation*}
$$

with $\frac{1}{p}=\frac{1}{q}-\frac{1}{2}$ and $1<q<2$. The kernel $K$ is singular, of Calderon-Zygmund type. Hence, there exists a constant $C>0$ such that, for any $p \in(1,+\infty)$,

$$
\begin{equation*}
\left\|\nabla\left(K * g_{t}^{N}\right)\right\|_{\mathbb{L}^{p}} \leqslant C\left\|g_{t}^{N}\right\|_{\mathbb{L}^{p}} \tag{28}
\end{equation*}
$$

Let us introduce $f^{N}=K * g^{N}$. By the Sobolev embeddings' Theorem, estimates (27) and (28) imply that: for any $\widetilde{\eta}>0\left(\right.$ take $\widetilde{\eta}=1-\frac{2}{p}$ with $p>2$ ),

$$
\left\|f^{N}\right\|_{\mathbb{L}^{2}\left([0, T] ; C^{\tilde{\eta}}\left(\mathbb{R}^{2}\right)\right)} \leqslant C .
$$

Now, let $\chi: \mathbb{R}^{2} \rightarrow[0,1] \in C_{0}^{\infty}$ be a cut-off function, i.e. such that:

$$
0 \leqslant \chi(x) \leqslant 1, \quad \text { and } \quad \chi(x)=1, \text { if }|x| \leqslant 1
$$

Let $\tilde{\chi}=1-\chi$. We can decompose:

$$
\begin{aligned}
& \left(K *\left(g_{t}^{N}-\xi\right)\right)(x) \\
& \quad=\int_{\mathbb{R}^{2}} \chi(y) K(y)\left(g_{t}^{N}-\xi\right)(x-y) d y+\int_{\mathbb{R}^{2}} \tilde{\chi}(x-y) K(x-y)\left(g_{t}^{N}-\xi\right)(y) d y .
\end{aligned}
$$

We observe that $\chi(\cdot) K(\cdot) \in \mathbb{L}^{a}\left(\mathbb{R}^{2}\right)$ with $a<2$, and $\widetilde{\chi}(x-\cdot) K(x-\cdot) \in \mathbb{L}^{b}\left(\mathbb{R}^{2}\right)$ with $b>2$. Since $g^{N}$ is uniformly bounded in $\mathbb{L}^{2}\left([0, T] ; \mathbb{L}^{p}\left(\mathbb{R}^{2}\right)\right)$ for all $p \in[1,+\infty]$ we obtain that $K * g_{t}^{N}$ converges to $K * \xi$.

Finally, we can bound as follows:

$$
\begin{aligned}
& \left|\left\langle S_{s}^{N}, F\left(K * g_{s}^{N}\right) \cdot \nabla\left(V^{N} * \phi\right)\right\rangle-\left\langle g_{s}^{N}, F\left(K * g_{s}^{N}\right) \cdot \nabla\left(V^{N} * \phi\right)\right\rangle\right| \\
& \left.\quad \leqslant \sup _{x \in \mathbb{R}^{2}} \mid F\left(K * g_{s}^{N}\right) \cdot \nabla\left(V^{N} * \phi\right)(x)-\left(F\left(K * g_{s}^{N}\right) \cdot \nabla\left(V^{N} * \phi\right)\right) * V^{N}\right)(x) \mid .
\end{aligned}
$$

Let us control the last term, using the facts that

- $V$ is a density (denoted below by $\left.\left(\int V=1\right)\right)$,
- $F$ is Lipschitz and bounded (denoted below by $\left(F \in \operatorname{Lip} \cap L^{\infty}\right)$ ),
- $V$ is compactly supported (denoted below by ( $V$ is c.s.)),
- and $\phi$ is compactly supported and smooth,
as follows (the norm $\|\cdot\|$ below is the Euclidean norm on $\mathbb{R}^{2}$ ): for any $x \in \mathbb{R}^{2}$,

$$
\begin{aligned}
\mid F\left(K * g_{s}^{N}\right) & \left.(x) \cdot \nabla\left(V^{N} * \phi\right)(x)-\left(F\left(K * g_{s}^{N}\right) \cdot \nabla\left(V^{N} * \phi\right)\right) * V^{N}\right)(x) \mid \\
\stackrel{(S V=1)}{\leqslant} & \int_{\mathbb{R}^{2}} V(y)\left\|\nabla\left(V^{N} * \phi\right)(x)\right\|\left\|F\left(K * g_{s}^{N}\right)(x)-F\left(K * g_{s}^{N}\right)\left(x-\frac{y}{N^{\beta}}\right)\right\| d y \\
& +\int_{\mathbb{R}^{2}} V(y)\left\|\nabla\left(V^{N} * \phi\right)(x)-\nabla\left(V^{N} * \phi\right)\left(x-\frac{y}{N^{\beta}}\right)\right\|\left\|\left(K * g_{s}^{N}\right)(x)\right\| d y \\
\begin{array}{c}
\left(F \in \operatorname{Lip} \cap L^{\infty}\right) \\
\leqslant
\end{array} & C \int_{\mathbb{R}^{2}} V(y)\left\|\nabla\left(V^{N} * \phi\right)(x)\right\|\left\|f_{s}^{N}(x)-f_{s}^{N}\left(x-\frac{y}{N^{\beta}}\right)\right\| d y \\
& +\frac{C}{N^{\beta}} \int_{\mathbb{R}^{2}} V(y)\|y\| d y \\
(V \text { is c.s. }) & \frac{C}{N_{\tilde{\eta} \beta}^{\beta}} \sup _{x, y \in \mathbf{K}}^{\leqslant} \frac{\left\|f_{s}^{N}(x)-f_{s}^{N}(y)\right\|}{\|x-y\|^{\tilde{\eta}}} \int_{\mathbb{R}^{2}} V(y)\|y\|^{\tilde{\eta}} d y \\
& +\frac{C}{N^{\beta}} \int_{\mathbb{R}^{2}} V(y)\|y\| d y,
\end{aligned}
$$

where $\mathbf{K} \subset \mathbb{R}^{2}$ is a compact set. Therefore we have obtained

$$
\left.\mid F\left(K * g_{s}^{N}\right)(x) \cdot \nabla\left(V^{N} * \phi\right)(x)-\left(F\left(K * g_{s}^{N}\right) \cdot \nabla\left(V^{N} * \phi\right)\right) * V^{N}\right)(x) \left\lvert\, \leqslant \frac{C}{N \tilde{\eta} \beta}\right.,
$$

where the constant $C$ depend on $\left\|f^{N}\right\|_{C_{\tilde{n}}\left(\mathbb{R}^{2}\right)}$. Thus,

$$
\begin{aligned}
\lim _{N \rightarrow \infty} \int_{0}^{t}\left\langle S_{s}^{N}, F\left(K * g_{s}^{N}\right)\right. & \left.\bullet \nabla\left(V^{N} * \phi\right)\right\rangle d s \\
& =\lim _{N \rightarrow \infty} \int_{0}^{t}\left\langle g_{s}^{N}, F\left(K * g_{s}^{N}\right) \bullet \nabla\left(V^{N} * \phi\right)\right\rangle d s \\
& =\lim _{N \rightarrow \infty} \int_{0}^{t} \int_{\mathbb{R}^{2}} g_{s}^{N}(x) F\left(K * g_{s}^{N}\right)(x) \cdot \nabla\left(V^{N} * \phi\right)(x) d x d s \\
& =\int_{0}^{t} \int_{\mathbb{R}^{2}} \xi(s, x) F(K * \xi(s, \cdot))(x) \bullet \nabla \phi(x) d x d s
\end{aligned}
$$

where in the last equality we used that $g_{s}^{N} \rightarrow \xi$ strongly in $\mathbb{L}^{2}\left([0, T] ; C\left(\mathbb{R}^{2}\right)\right)$. We have proved (26).

Therefore from (23), (24), (25) and (26), we conclude that the limit point $\xi$ is then a weak solution to the PDE

$$
\begin{equation*}
\partial_{t} \xi+\operatorname{div}(\xi F(K * \xi))=\nu \Delta \xi,\left.\quad \xi\right|_{t=0}=\xi^{\mathrm{ini}} \tag{29}
\end{equation*}
$$

Note that there is one more step to recover (46) in Theorem 3, which will be achieved in Section 2.7 below. Before that, we need to prove that the solution to (11) is unique.

### 2.5 Uniqueness of the solution

Theorem 8 We assume that $\xi^{\text {ini }} \in \mathbb{L}^{1} \cap \mathbb{L}^{\infty}\left(\mathbb{R}^{2}\right)$. Then there is at most one weak solution of the equation (11) which belongs to $\mathbb{L}^{2}\left([0, T] ; \mathbb{L}^{1} \cap \mathbb{L}^{\infty}\left(\mathbb{R}^{2}\right)\right)$.

Proof. For any function $u:[0, T] \times \mathbb{R}^{2} \rightarrow \mathbb{R}$ we introduce the notation

$$
\|u\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}:=\|u(t, \cdot)\|_{\mathbb{L}^{1}}+\|u(t, \cdot)\|_{\mathbb{L}^{\infty}} .
$$

Let $\xi^{1}, \xi^{2}$ be two weak solutions of (11), with the same initial condition $\xi^{\text {ini }}$ which satisfies $\xi \in \mathbb{L}^{1} \cap \mathbb{L}^{\infty}\left(\mathbb{R}^{2}\right)$. By hypothesis, from $[1], \xi^{1}, \xi^{2} \in \mathbb{L}^{2}\left([0, T] ; \mathbb{L}^{\infty}\right)$.

Let $\left\{h_{\varepsilon}\right\}_{\varepsilon}$ be a family of standard symmetric mollifiers on $\mathbb{R}^{2}$. For any $\varepsilon>0$ and $x \in \mathbb{R}^{2}$ we can use $h_{\varepsilon}(x-\cdot)$ as a test function in the equation (12). Therefore we set $\xi_{\varepsilon}^{i}(t, x)=\left(\xi^{i}(t, \cdot) * h_{\varepsilon}\right)(x)$ for $i=1,2$. Then we have, for any $(t, x) \in[0, T] \times \mathbb{R}^{2}$,

$$
\xi_{\varepsilon}^{i}(t, x)=\left(\xi^{\mathrm{ini}} * h_{\varepsilon}\right)(x)+\nu \int_{0}^{t} \Delta \xi_{\varepsilon}^{i}(s, x) d s+\int_{0}^{t}\left(\left(\nabla h_{\varepsilon} \cdot F\left(K * \xi^{i}\right)\right) * \xi^{i}\right)(s, x) d s
$$

Writing this identity in mild form we obtain (writing with a little abuse of notation $\xi^{i}(t)$ for the function $\xi^{i}(t, \cdot)$ and $\mathfrak{S}(t)$ for $\left.e^{t A}\right)$ :

$$
\xi_{\varepsilon}^{i}(t)=\mathfrak{S}(t) \cdot\left(\xi^{\mathrm{ini}} * h_{\varepsilon}\right)+\int_{0}^{t} \mathfrak{S}(t-s) \cdot\left(\left(\nabla h_{\varepsilon} \cdot F\left(K * \xi^{i}\right)\right) * \xi^{i}\right)(s) d s
$$

The function $X=\xi^{1}-\xi^{2}$ satisfies

$$
h_{\varepsilon} * X(t)=\int_{0}^{t} \nabla \mathfrak{S}(t-s) \cdot\left(h_{\varepsilon} *\left(F\left(K * \xi^{1}\right) \xi^{1}-F\left(K * \xi^{2}\right) \xi^{2}\right)\right) d s
$$

Thus we obtain

$$
\left\|h_{\varepsilon} * X(t)\right\|_{\mathbb{L}^{\infty}} \leqslant \int_{0}^{t}\left\|\nabla \mathfrak{S}(t-s) \cdot\left(h_{\varepsilon} *\left(F\left(K * \xi^{1}\right) \xi^{1}-F\left(K * \xi^{2}\right) \xi^{2}\right)\right)\right\|_{\mathbb{L}^{\infty}} d s
$$

Now, we have

$$
\left\|h_{\varepsilon} * X(t)\right\|_{\mathbb{L}^{\infty}} \leqslant \int_{0}^{t} \frac{1}{(t-s)^{\frac{1}{2}}}\left\|h_{\varepsilon} *\left(F\left(K * \xi^{1}\right) \xi^{1}-F\left(K * \xi^{2}\right) \xi^{2}\right)\right\|_{\mathbb{L}^{\infty}} d s
$$

Taking the limit as $\varepsilon \rightarrow 0$ we arrive at

$$
\|X(t)\|_{\mathbb{L}^{\infty}} \leqslant \int_{0}^{t} \frac{1}{(t-s)^{\frac{1}{2}}}\left\|F\left(K * \xi^{1}\right) \xi^{1}-F\left(K * \xi^{2}\right) \xi^{2}\right\|_{\mathbb{L}^{\infty}} d s .
$$

With similar arguments we have the same estimate in $\mathbb{L}^{1}$-norm as follows:

$$
\|X(t)\|_{\mathbb{L}^{1}} \leqslant \int_{0}^{t} \frac{1}{(t-s)^{\frac{1}{2}}}\left\|F\left(K * \xi^{1}\right) \xi^{1}-F\left(K * \xi^{2}\right) \xi^{2}\right\|_{\mathbb{L}^{1}} d s
$$

By easy calculation we have

$$
\begin{aligned}
\|X(t)\|_{\mathbb{L}^{\infty}} & \leqslant \int_{0}^{t} \frac{1}{(t-s)^{\frac{1}{2}}}\left(\left\|X F\left(K * \xi^{1}\right)\right\|_{\mathbb{L}^{\infty}}+\left\|\xi^{2}\left(F\left(K * \xi^{1}\right)-F\left(K * \xi^{2}\right)\right)\right\|_{\mathbb{L}^{\infty}}\right) d s . \\
& \leqslant \int_{0}^{t} \frac{\left\|\xi^{1}\right\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}}{(t-s)^{\frac{1}{2}}}\|X\|_{\mathbb{L}^{\infty}} d s+\int_{0} \frac{\left\|\xi^{2}\right\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}}{(t-s)^{\frac{1}{2}}}\left(\|X\|_{\mathbb{L}^{\infty}}+\|X\|_{\mathbb{L}^{1}}\right) d s \\
& \leqslant \int_{0}^{t} \frac{\left\|\xi^{2}\right\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}+\left\|\xi^{2}\right\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}}{(t-s)^{\frac{1}{2}}}\|X\|_{\mathbb{L}^{\infty}}+\int_{0}^{t} \frac{\left\|\xi^{2}\right\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}+\left\|\xi^{2}\right\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}}{(t-s)^{\frac{1}{2}}}\|X\|_{\mathbb{L}^{1}} d s .
\end{aligned}
$$

On the other hand, in a similar way we have

$$
\begin{aligned}
\|X(t)\|_{\mathbb{L}^{1}} & \leqslant \int_{0}^{t} \frac{1}{(t-s)^{\frac{1}{2}}}\left(\left\|X F\left(K * \xi^{1}\right)\right\|_{\mathbb{L}^{1}}+\left\|\xi^{2}\left(F\left(K * \xi^{1}\right)-F\left(K * \xi^{2}\right)\right)\right\|_{\mathbb{L}^{1}}\right) d s \\
& \leqslant \int_{0}^{t} \frac{\left\|\xi^{2}\right\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}+\left\|\xi^{2}\right\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}}{(t-s)^{\frac{1}{2}}}\|X\|_{\mathbb{L}^{\infty}}+\int_{0}^{t} \frac{\left\|\xi^{2}\right\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}+\left\|\xi^{2}\right\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}}{(t-s)^{\frac{1}{2}}}\|X\|_{\mathbb{L}^{1}} d s
\end{aligned}
$$

Therefore we have that

$$
\|X(t)\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}} \leqslant C \int_{0}^{t} \frac{\left\|\xi^{2}\right\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}+\left\|\xi^{2}\right\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}}}{(t-s)^{\frac{1}{2}}}\|X(s)\|_{\mathbb{L}^{1} \cap \mathbb{L}^{\infty}} d s
$$

By Gronwall's Lemma we conclude $X=0$.

### 2.6 Convergence in probability

Corollary 9 The sequence $\left\{g^{N}\right\}_{N \in \mathbb{N}}$ converges in probability to $\xi$.
Proof. We denote the joint law of $\left(g^{N}, g^{M}\right)$ by $\nu^{N, M}$. Similarly to the proof of tightness for $g^{N}$ (Section 2.3) we have that the family $\left\{\nu^{N, M}\right\}$ is tight in $\mathfrak{Y} \times \mathfrak{Y}$, where $\mathfrak{Y}$ has been defined in (21).

Let us take any subsequence $\nu^{N_{k}, M_{k}}$. By Prohorov's Theorem, it is relatively weakly compact hence it contains a weakly convergent subsequence. Without loss of generality we may assume that the original sequence $\left\{\nu^{N, M}\right\}$ itself converges weakly to a measure $\nu$. According to Skorokhod immersion's Theorem, we infer the existence of a probability space $(\bar{\Omega}, \overline{\mathcal{F}}, \overline{\mathbb{P}})$ with a sequence of random variables $\left(\bar{g}^{N}, \bar{g}^{M}\right)$ converging almost surely in $\mathfrak{Y} \times \mathfrak{Y}$ to random variable ( $\bar{u}, \check{u}$ ) and the laws of ( $\bar{g}^{N}, \bar{g}^{M}$ ) and $(\bar{u}, \breve{u})$ under $\overline{\mathbb{P}}$ coincide with $\nu^{N, M}$ and $\nu$, respectively.

Analogously, it can be applied to both $\bar{g}^{N}$ and $\bar{g}^{M}$ in order to show that $\bar{u}$ and $\check{u}$ are two solutions of the PDE (11). By Theorem 8 which gives the uniqueness of the solution we have $\bar{u}=\check{u}$. Therefore

$$
\nu((x, y) \in \mathfrak{Y} \times \mathfrak{Y} ; x=y)=\overline{\mathbb{P}}(\bar{u}=\check{u})=1 .
$$

Now, we have all in hands to apply Gyongy-Krylov's characterization of convergence in probability, which reads as follows:

Lemma 10 (Gyongy-Krylov [12]) Let $\left\{X_{n}\right\}$ be a sequence of random elements in a Polish space $\Psi$ equipped with the Borel $\sigma$-algebra. Then $X_{n}$ converges in probability to an $\Psi$-valued random element if, and only if, for each pair $\left(X_{\ell}, X_{m}\right)$ of subsequences, there exists a subsequence $\left\{v_{k}\right\}$ given by

$$
v_{k}=\left(X_{\ell(k)}, X_{m(k)}\right),
$$

converging weakly to a random element $v(x, y)$ supported on the diagonal set

$$
\{(x, y) \in \Psi \times \Psi: x=y\} .
$$

This lemma implies that the original sequence is defined on the initial probability space converges in probability in the topology of $\mathfrak{Y}$ to a random variable $\mu$.

### 2.7 Conclusion

Let $\xi(t, x)$ be the unique solution of the vorticity equation (11) with initial condition $\xi^{\text {ini }} \in \mathbb{L}^{\infty} \cap \mathbb{L}^{1}\left(\mathbb{R}^{2}\right)$. From [1] we have $\|\xi\|_{\infty} \leqslant\left\|\xi^{\text {ini }}\right\|_{\infty}$. Then, by definition of the Biot-Savart kernel $K$, there is a positive constant $c_{K}$ (given by (4)) such that

$$
\|K * \xi\|_{\infty} \leqslant c_{K}\left(1+\left\|\xi^{\mathrm{ini}}\right\|_{\infty}\right)
$$

Therefore if we take $M \geqslant c_{K}\left(1+\left\|\xi^{\text {ini }}\right\|_{\infty}\right)$, we conclude that $\xi(t, x)$ coincides with the unique solution of (29), which is satisfied by the limit point of the sequence $\left\{g^{N}\right\}$.

## 3 Technical proofs

In this last section we prove Proposition 6 and Proposition 7.

### 3.1 Proof of Proposition 6

Let us prove the first estimate on $g^{N}$ given in Proposition 6, namely (17).
Step 1. From (16) after a multiplication by $(\mathrm{I}-A)^{\varepsilon / 2}$ and by triangular inequality we have

$$
\begin{align*}
&\left\|(\mathrm{I}-A)^{\varepsilon / 2} g_{t}^{N}\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)} \leqslant\left\|(\mathrm{I}-A)^{\varepsilon / 2} e^{t A} g_{0}^{N}\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)}  \tag{30}\\
&+\int_{0}^{t}\left\|(\mathrm{I}-A)^{\varepsilon / 2} \nabla e^{(t-s) A}\left(V^{N} *\left(F\left(K * g_{s}^{N}\right) S_{s}^{N}\right)\right)\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)} d s  \tag{31}\\
&+\left\|\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t}(\mathrm{I}-A)^{\varepsilon / 2} \nabla e^{(t-s) A}\left(V^{N}\left(\cdot-X_{s}^{i, N}\right)\right) d W_{s}^{i}\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)} \tag{32}
\end{align*}
$$

We denote $H:=\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)$. Then

$$
\begin{align*}
&\left\|(\mathrm{I}-A)^{\varepsilon / 2} g_{t}^{N}\right\|_{\mathbb{L}^{p}(\Omega, H)} \leqslant\left\|(\mathrm{I}-A)^{\varepsilon / 2} e^{t A} g_{0}^{N}\right\|_{\mathbb{L}^{p}(\Omega, H)}  \tag{33}\\
&+\int_{0}^{t}\left\|(\mathrm{I}-A)^{\varepsilon / 2} \nabla e^{(t-s) A}\left(V^{N} *\left(F\left(K * g_{s}^{N}\right) S_{s}^{N}\right)\right)\right\|_{\mathbb{L}^{p}(\Omega, H)} d s  \tag{34}\\
&+\left\|\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t}(\mathrm{I}-A)^{\varepsilon / 2} \nabla e^{(t-s) A}\left(V^{N}\left(\cdot-X_{s}^{i, N}\right)\right) d W_{s}^{i}\right\|_{\mathbb{L}^{p}(\Omega, H)} \tag{35}
\end{align*}
$$

Step 2. The first term (30) can be estimated by

$$
\left\|(\mathrm{I}-A)^{\varepsilon / 2} e^{t A} g_{0}^{N}\right\|_{\mathbb{L}^{p}(\Omega, H)} \leqslant\left\|(\mathrm{I}-A)^{\varepsilon / 2} g_{0}^{N}\right\|_{\mathbb{L}^{p}(\Omega, H)} \leqslant C_{\varepsilon} .
$$

The boundedness of $g_{0}^{N}$ we follow from Assumption 1, item 3.

Step 3. Let us come to the second term (31):

$$
\begin{aligned}
& \int_{0}^{t}\left\|(\mathrm{I}-A)^{\varepsilon / 2} \nabla e^{(t-s) A}\left(V^{N} *\left(F\left(K * g_{s}^{N}\right) S_{s}^{N}\right)\right)\right\|_{\mathbb{L}^{p}(\Omega, H)} d s \\
& \leqslant C \int_{0}^{t}\left\{\left\|(\mathrm{I}-A)^{1 / 2} e^{((t-s) / 2) A}\right\|_{\mathbb{L}^{2} \rightarrow \mathbb{L}^{2}}\right. \\
& \left.\quad \times\left\|(\mathrm{I}-A)^{\varepsilon / 2} e^{((t-s) / 2) A}\left(V^{N} *\left(F\left(K * g_{s}^{N}\right) S_{s}^{N}\right)\right)\right\|_{\mathbb{L}^{p}(\Omega, H)}\right\} d s .
\end{aligned}
$$

We have

$$
\left\|(\mathrm{I}-A)^{1 / 2} e^{((t-s) / 2) A}\right\|_{\mathbb{L}^{2} \rightarrow \mathbb{L}^{2}} \leqslant \frac{C_{\nu}}{(t-s)^{1 / 2}},
$$

On the other hand, for any $x \in \mathbb{R}^{2}$,

$$
\left|\left(V^{N} *\left(F\left(K * g_{s}^{N}\right) S_{s}^{N}\right)\right)(x)\right| \leqslant\left\|F\left(K * g_{s}^{N}\right)\right\|_{\infty}\left|V^{N} * S_{s}^{N}(x)\right| \leqslant M\left|g_{s}^{N}(x)\right| .
$$

Using Lemma 16 in [5] we obtain

$$
\begin{aligned}
& \left\|(\mathrm{I}-A)^{\varepsilon / 2} e^{((t-s) / 2) A}\left(V^{N} *\left(F\left(K * g_{s}^{N}\right) S_{s}^{N}\right)\right)\right\|_{\mathbb{L}^{p}(\Omega, H)} \\
& \quad \leqslant C_{M}\left\|e^{((t-s) / 2) A}(\mathrm{I}-A)^{\varepsilon / 2} g_{s}^{N}\right\|_{\mathbb{L}^{p}(\Omega, H)} \leqslant C_{M}\left\|(\mathrm{I}-A)^{\varepsilon / 2} g_{s}^{N}\right\|_{\mathbb{L}^{p}(\Omega, H)} .
\end{aligned}
$$

To summarize, we have proved

$$
\begin{aligned}
\int_{0}^{t} \|(\mathrm{I}-A)^{\varepsilon / 2} \nabla e^{(t-s) A}\left(V^{N} *(F\right. & \left.\left.\left(K * g_{s}^{N}\right) S_{s}^{N}\right)\right) \|_{\mathbb{L}^{p}(\Omega, H)} d s \\
& \leqslant C_{\varepsilon, \nu, M} \int_{0}^{t}(t-s)^{\frac{1}{2}}\left\|(\mathrm{I}-A)^{\varepsilon / 2} g_{s}^{N}\right\|_{\mathbb{L}^{p}(\Omega, H)} d s .
\end{aligned}
$$

This bounds the second term.
Step 4. The estimate of the third term (32) is quite tricky and we postpone it to Lemma 11 below, see (36). Collecting the three bounds together, we get

$$
\left\|(\mathrm{I}-A)^{\varepsilon / 2} g_{t}^{N}\right\|_{\mathbb{L}^{p}(\Omega, H)} \leqslant C_{\varepsilon}+C_{\varepsilon, T, M} \int_{0}^{t}(t-s)^{\frac{1}{2}}\left\|\left(\mathrm{I}-A_{2}\right)^{\varepsilon / 2} g_{s}^{N}\right\|_{\mathbb{L}^{p}(\Omega, H)} d s
$$

We may apply Gronwall's Lemma we deduce

$$
\left\|(\mathrm{I}-A)^{\varepsilon / 2} g_{t}^{N}\right\|_{\mathbb{L}^{p}(\Omega, H)} \leqslant C_{\varepsilon, T, M, \nu},
$$

and Proposition 6 follows.

Lemma 11 We assume Assumption 1. Then there exists a constant $C_{\varepsilon, T}>0$ such that for all $t \in[0, T]$,

$$
\begin{equation*}
\left\|\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t}(\mathrm{I}-A)^{\varepsilon / 2} \nabla e^{(t-s) A}\left(V^{N}\left(\cdot-X_{s}^{i, N}\right)\right) d W_{s}^{i}\right\|_{\mathbb{L}^{p}\left(\Omega, \mathbb{L}^{2}\left(\mathbb{R}^{2}\right)\right)}^{p} \leqslant C_{\varepsilon, T} . \tag{36}
\end{equation*}
$$

Proof. Recall the notation $H:=\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)$. From the Burkholder-Davis-Gundy inequality (see [29] for instance) we obtain

$$
\begin{aligned}
& \left\|\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t}(\mathrm{I}-A)^{\varepsilon / 2} \nabla e^{(t-s) A}\left(V^{N}\left(\cdot-X_{s}^{i, N}\right)\right) d W_{s}^{i}\right\|_{\mathbb{L}^{p}(\Omega, H)}^{p} \\
& \quad \leqslant C_{p} \mathbb{E}\left[\frac{1}{N^{2}} \sum_{i=1}^{N} \int_{0}^{t}\left\|(\mathrm{I}-A)^{\varepsilon / 2} \nabla e^{(t-s) A}\left(V^{N}\left(\cdot-X_{s}^{i, N}\right)\right)\right\|_{H}^{2} d s\right]^{p / 2}
\end{aligned}
$$

Moreover, we can estimate

$$
\begin{align*}
\frac{1}{N^{2}} \int_{\mathbb{R}^{2}} \sum_{i=1}^{N} \int_{0}^{t} & \left|\left((\mathrm{I}-A)^{\varepsilon / 2} \nabla e^{(t-s) A}\left(V^{N}\left(\cdot-X_{s}^{i, N}\right)\right)\right)(x)\right|^{2} d s d x \\
& =\frac{1}{N} \int_{0}^{t}\left\|(\mathrm{I}-A)^{\varepsilon / 2} \nabla e^{(t-s) A} V^{N}\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)}^{2} d s \\
& =\frac{1}{N} \int_{0}^{t}\left\|(\mathrm{I}-A)^{-\delta / 2} \nabla e^{(t-s) A}(\mathrm{I}-A)^{(\varepsilon+\delta) / 2} V^{N}\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)}^{2} d s \\
& \leqslant \frac{1}{N} \int_{0}^{t} \frac{1}{(t-s)^{1-\delta}}\left\|V^{N}\right\|_{\varepsilon+\delta, 2}^{2} d s \\
& \leqslant C_{T, \varepsilon, \delta} N^{\beta(2+2 \delta+2 \varepsilon)-1} \tag{37}
\end{align*}
$$

Therefore (37) is bounded by some constant $C_{\varepsilon, T}$ if we take $\beta<\frac{1}{2+2 \varepsilon}$, $\delta$ close enough to zero. This provides the bound of the lemma.

### 3.2 Proof of Proposition 7

Let us now prove the second estimate on $g^{N}$ given in Proposition 7, namely (18). In this proof we use the fact that $\mathbb{L}^{2}\left(\mathbb{R}^{2}\right) \subset \mathbb{H}_{2}^{-2}$ with continuous embedding, and that the linear operator $\Delta$ is bounded from $\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)$ to $\mathbb{H}_{2}^{-2}$.

We first observe that

$$
\begin{aligned}
g_{t}^{N}(x)-g_{s}^{N}(x)= & \int_{s}^{t}\left\langle S_{r}^{N},\left(K * F\left(g_{r}^{N}\right)\right) \nabla V^{N}(x-\cdot)\right\rangle d r+\nu \int_{s}^{t} \Delta g_{r}^{N}(x) d r \\
& +\frac{1}{N} \sum_{i=1}^{N} \int_{s}^{t} \nabla\left(V^{N}\right)\left(x-X_{r}^{i, N}\right) d W_{r}^{i}
\end{aligned}
$$

Therefore we have

$$
\begin{align*}
& \mathbb{E}\left[\left\|g_{t}^{N}(x)-g_{s}^{N}(x)\right\|_{-2,2}^{q}\right] \\
& \leqslant(t-s)^{q-1} \int_{s}^{t} \mathbb{E}\left[\left\|\left\langle S_{r}^{N}, F\left(K * g_{r}^{N}\right) \nabla V^{N}(x-\cdot)\right\rangle\right\|_{-2,2}^{q}\right] d r  \tag{38}\\
&+(t-s)^{q-1} \frac{1}{2} \int_{s}^{t} \mathbb{E}\left[\left\|\Delta g_{r}^{N}(x)\right\|_{-2,2}^{q}\right] d r  \tag{39}\\
&+\mathbb{E}\left[\left\|\frac{1}{N} \sum_{i=1}^{N} \int_{s}^{t} \nabla\left(V^{N}\right)\left(x-X_{r}^{i, N}\right) d W_{r}^{i}\right\|_{-2,2}^{q}\right] . \tag{40}
\end{align*}
$$

To estimate the first term (38) we observe first that

$$
\begin{align*}
\mathbb{E}\left[\left\|\left\langle S_{r}^{N}, F\left(K * g_{r}^{N}\right) \nabla V^{N}(x-\cdot)\right\rangle\right\|_{-2,2}^{q}\right] & =\mathbb{E}\left[\left\|\nabla\left(S_{r}^{N} F\left(K * g_{r}^{N}\right) * V^{N}\right)\right\|_{-2,2}^{q}\right] \\
& \leqslant \mathbb{E}\left[\|\left(S_{r}^{N} F\left(K * g_{r}^{N}\right) * V^{N} \|_{-1,2}^{q}\right]\right. \\
& \leqslant C_{M} \mathbb{E}\left[\left\|g_{t}^{N}\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)}^{q}\right] \leqslant C \tag{41}
\end{align*}
$$

Moreover, for the second term (39) we write

$$
\begin{equation*}
\mathbb{E}\left[\left\|\Delta g_{r}^{N}\right\|_{-2,2}^{q}\right] \leqslant C \mathbb{E}\left[\left\|g_{r}^{N}\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)}^{q}\right] \leqslant C \tag{42}
\end{equation*}
$$

Finally we bound the last term (40):

$$
\begin{aligned}
\mathbb{E}\left[\| \frac{1}{N} \sum_{i=1}^{N} \int_{s}^{t} \nabla\left(V^{N}\right)\left(x-X_{r}^{i, N}\right)\right. & \left.d W_{r}^{i} \|_{-2,2}^{q}\right] \\
& \leqslant C_{q} \mathbb{E}\left[\frac{1}{N^{2}} \sum_{i=1}^{N} \int_{s}^{t}\left\|\nabla\left(V^{N}\right)\left(x-X_{r}^{i, N}\right)\right\|_{-2,2}^{q} d r\right]^{q / 2}
\end{aligned}
$$

and we observe that

$$
\begin{aligned}
& \frac{1}{N^{2}} \int_{\mathbb{R}} \sum_{i=1}^{N} \int_{s}^{t}\left\|(\mathrm{I}-A)^{-1} \nabla\left(V^{N}\right)\left(x-X_{r}^{i, N}\right)\right\|^{2} d r d x \\
& \quad=(t-s) \frac{1}{N}\left\|V^{N}\right\|_{-1,2}^{2} \leqslant(t-s) \frac{1}{N}\left\|V^{N}\right\|_{0,2}^{2} \leqslant C N^{2 \beta-1}(t-s) \leqslant C(t-s) .
\end{aligned}
$$

In order to conclude the lemma, we need to divide (38)-(40) by $|t-s|^{1+q \gamma}$. From the previous estimates, we always get a term of the form $|t-s|^{\varepsilon}$ with $\varepsilon<1$ (using the assumption $\gamma<\frac{1}{2}$ ).

## A More general initial data

Assume that the initial vorticity $\xi^{\text {ini }}$ has variable sign. Define

$$
\xi_{+}^{\mathrm{ini}}:=\xi^{\mathrm{ini}} \vee 0, \quad \xi_{-}^{\mathrm{ini}}:=\left(-\xi^{\mathrm{ini}}\right) \vee 0
$$

and

$$
\Gamma_{ \pm}:=\int \xi_{ \pm}^{\mathrm{ini}}(x) d x>0
$$

(they are finite, since we assume $\xi^{\text {ini }} \in \mathbb{L}^{1}$ ). Let $\left\{X_{0}^{i, \pm}\right\}_{i \in \mathbb{N}}$ be a double sequence of random variables in $\mathbb{R}^{2}$ such that, for the empirical measures

$$
S_{0}^{N, \pm}:=\frac{\Gamma_{ \pm}}{N} \sum_{i=1}^{N} \delta_{X_{0}^{i, \pm}}
$$

one has, for some $\alpha>1$, for any $p>0$,

$$
\begin{equation*}
\sup _{N \in \mathbb{N}} \mathbb{E}\left[\left\|V^{N} * S_{0}^{N, \pm}\right\|_{\alpha, 2}^{p}\right]<\infty \tag{43}
\end{equation*}
$$

and the two sequences of measures $\left\{S_{0}^{N, \pm}\right\}_{N \in \mathbb{N}}$ weakly converge to the initial measures $\xi_{ \pm}^{\text {ini }}(\cdot) d x$, as $N \rightarrow \infty$, in probability: i.e.

$$
\begin{equation*}
S_{0}^{N, \pm} \underset{N \rightarrow \infty}{ } \xi_{ \pm}^{\mathrm{ini}}(\cdot) d x \quad \text { in probability } . \tag{44}
\end{equation*}
$$

Consider the system of PDEs, given for $x \in \mathbb{R}^{2}, t>0$ by

$$
\begin{align*}
\partial_{t} \xi^{+}+u \bullet \nabla \xi^{+} & =\nu \Delta \xi^{+}, \\
\partial_{t} \xi^{-}+u \bullet \nabla \xi^{-} & =\nu \Delta \xi^{-}, \\
u & =K *\left(\xi^{+}-\xi^{-}\right)  \tag{45}\\
\left.\xi^{+}\right|_{t=0} & =\xi_{0}^{+},\left.\quad \xi^{-}\right|_{t=0}=\xi_{0}^{-} .
\end{align*}
$$

It is not difficult to prove the same results of existence and uniqueness as the ones obtained for the usual Navier-Stokes equations

$$
\begin{equation*}
\partial_{t} \xi+u \cdot \nabla \xi=\nu \Delta \xi, \quad x \in \mathbb{R}^{2}, t>0 \tag{46}
\end{equation*}
$$

Moreover, if $\left(\xi^{+}, \xi^{-}\right)$is a solution of the system, then $\xi=\xi^{+}-\xi^{-}$is a solution of (46); if $\xi$ is a solution of (46) and $\xi^{+}$is a solution of the first equation of the system, with $u=K * \xi$, then $\xi^{-}=\xi^{+}-\xi$ is a solution of the second equation of the system, and $u=K *\left(\xi^{+}-\xi^{-}\right)$holds. In this sense the system and (46) are equivalent.

Let $\left\{W_{t}^{i, \pm}\right\}_{i \in \mathbb{N}}$ be a family of independent $\mathbb{R}^{2}$-valued Brownian motions, defined on the same probability space as $\left\{X_{0}^{i, \pm}\right\}_{i \in \mathbb{N}}$ and independent of them. Given $N \in \mathbb{N}$ consider particles with positions $\left\{X_{t}^{i, N, \pm}\right\}_{i \in \mathbb{N}}$ satisfying

$$
\begin{aligned}
d X_{t}^{i, N,+}= & F\left(\Gamma_{+} \frac{1}{N} \sum_{k=1}^{N}\left(K * V^{N}\right)\left(X_{t}^{i, N,+}-X_{t}^{k, N,+}\right)\right. \\
& \left.-\Gamma_{-} \frac{1}{N} \sum_{k=1}^{N}\left(K * V^{N}\right)\left(X_{t}^{i, N,+}-X_{t}^{k, N,-}\right)\right) d t \\
+ & \sqrt{2 \nu} d W_{t}^{i,+} \\
d X_{t}^{i, N,-}= & F\left(\Gamma_{+} \frac{1}{N} \sum_{k=1}^{N}\left(K * V^{N}\right)\left(X_{t}^{i, N,-}-X_{t}^{k, N,+}\right)\right. \\
& \left.-\Gamma_{-} \frac{1}{N} \sum_{k=1}^{N}\left(K * V^{N}\right)\left(X_{t}^{i, N,-}-X_{t}^{k, N,-}\right)\right) d t \\
+ & \sqrt{2 \nu} d W_{t}^{i,-}
\end{aligned}
$$

with initial conditions $\left\{X_{0}^{i, \pm}\right\}_{i \in \mathbb{N}}$. Consider the associated empirical measures

$$
S_{t}^{N, \pm}:=\frac{\Gamma_{ \pm}}{N} \sum_{i=1}^{N} \delta_{X_{t}^{i, \pm}}
$$

and empirical densities

$$
g_{t}^{N, \pm}:=V^{N} * S_{t}^{N, \pm}
$$

Theorem 12 Assume on $V, \beta, \alpha$, the same conditions of Assumption 1 and in addition assume (43) and (44). Consider the particle system $\left\{X_{t}^{i, \pm}\right\}_{i \in \mathbb{N}}$ with the parameter $M$ which satisfies

$$
\begin{equation*}
M \geqslant c_{K}\left(1+\left\|\xi^{\mathrm{ini}}\right\|_{\mathbb{L}^{\infty}}\right) \tag{47}
\end{equation*}
$$

Then, for every $\eta \in(1, \alpha)$, the sequence of processes $\left\{g_{t}^{N,+}, g_{t}^{N,-}\right\}$ converges in probability with respect to the

- weak topology of $\left(\mathbb{L}^{2}\left([0, T] ; \mathbb{H}_{2}^{\alpha}\left(\mathbb{R}^{2}\right)\right)\right)^{2}$,
- strong topology of $\left(C\left([0, T] ; \mathbb{H}_{2, \mathrm{loc}}^{\eta}\left(\mathbb{R}^{2}\right)\right)\right)^{2}$,
as $N \rightarrow \infty$, to the unique weak solution of the PDE system (45) and thus $g_{t}^{N,+}-g_{t}^{N,-}$ converges, in the same topologies, to the unique weak solution of the PDE (46).

We do not repeat the full proof in this case but only sketch the main points. The empirical measures satisfy

$$
\begin{aligned}
\left\langle S_{t}^{N,+}, \phi\right\rangle=\left\langle S_{0}^{N,+}, \phi\right\rangle & +\int_{0}^{t}\left\langle S_{s}^{N,+}, F\left(K * V^{N} * S_{s}^{N,+}-K * V^{N} * S_{s}^{N,-}\right) \cdot \nabla \phi\right\rangle d s \\
& +\nu \int_{0}^{t}\left\langle S_{s}^{N,+}, \Delta \phi\right\rangle d s+\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} \nabla \phi\left(X_{s}^{i, N,+}\right) \cdot d W_{s}^{i,+}
\end{aligned}
$$

and

$$
\begin{aligned}
\left\langle S_{t}^{N,-}, \phi\right\rangle=\left\langle S_{0}^{N,-}, \phi\right\rangle & +\int_{0}^{t}\left\langle S_{s}^{N,-}, F\left(K * V^{N} * S_{s}^{N,+}-K * V^{N} * S_{s}^{N,-}\right) \cdot \nabla \phi\right\rangle d s \\
& +\nu \int_{0}^{t}\left\langle S_{s}^{N,-}, \Delta \phi\right\rangle d s+\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} \nabla \phi\left(X_{s}^{i, N,-}\right) \cdot d W_{s}^{i,-}
\end{aligned}
$$

and the empirical densities satisfy

$$
\begin{aligned}
g_{t}^{N,+}(x)=g_{0}^{N,+}(x) & +\int_{0}^{t}\left\langle S_{s}^{N,+}, F\left(K * g_{s}^{N,+}-K * g_{s}^{N,-}\right) \cdot \nabla V^{N}(x-\cdot)\right\rangle d s \\
& +\nu \int_{0}^{t} \Delta g_{s}^{N,+}(x) d s+\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} \nabla V^{N}\left(x-X_{s}^{i, N,+}\right) \cdot d W_{s}^{i,+}
\end{aligned}
$$

and

$$
\begin{aligned}
g_{t}^{N,-}(x)=g_{0}^{N,-}(x) & +\int_{0}^{t}\left\langle S_{s}^{N,-}, F\left(K * g_{s}^{N,+}-K * g_{s}^{N,-}\right) \cdot \nabla V^{N}(x-\cdot)\right\rangle d s \\
& +\nu \int_{0}^{t} \Delta g_{s}^{N,-}(x) d s+\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} \nabla V^{N}\left(x-X_{s}^{i, N,-}\right) \cdot d W_{s}^{i,-}
\end{aligned}
$$

which in mild form are

$$
\begin{aligned}
g_{t}^{N,+}=e^{t A} g_{0}^{N,+} & +\int_{0}^{t} \nabla e^{(t-s) A}\left(V^{N} *\left(F\left(K * g_{s}^{N,+}-K * g_{s}^{N,-}\right) S_{s}^{N,+}\right)\right) d s \\
& +\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} e^{(t-s) A}\left(\nabla V^{N}\left(\cdot-X_{s}^{i, N,+}\right)\right) \cdot d W_{s}^{i,+}
\end{aligned}
$$

and

$$
\begin{aligned}
g_{t}^{N,-}=e^{t A} g_{0}^{N,-} & +\int_{0}^{t} \nabla e^{(t-s) A}\left(V^{N} *\left(F\left(K * g_{s}^{N,+}-K * g_{s}^{N,-}\right) S_{s}^{N,-}\right)\right) d s \\
& +\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} e^{(t-s) A}\left(\nabla V^{N}\left(\cdot-X_{s}^{i, N,-}\right)\right) \cdot d W_{s}^{i,-}
\end{aligned}
$$

The proof of the estimate

$$
\mathbb{E}\left[\left\|\left(\mathrm{I}-A_{2}\right)^{\varepsilon / 2} g_{t}^{N,+}\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)}^{p}\right] \leqslant C_{\varepsilon, T, M, \nu, p}
$$

(and similarly for $g_{t}^{N,-}$ ) is similar to the case of a single sign (Proposition 6): precisely, the estimates on

$$
\left\|(\mathrm{I}-A)^{\varepsilon / 2} e^{t A} g_{0}^{N,+}\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)}
$$

and on

$$
\left\|\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t}(\mathrm{I}-A)^{\varepsilon / 2} \nabla e^{(t-s) A}\left(V^{N}\left(\cdot-X_{s}^{i, N,+}\right)\right) d W_{s}^{i,+}\right\|_{\mathbb{L}^{2}\left(\mathbb{R}^{2}\right)}
$$

are obviously the same. But also the middle term can be studied in the same way: one has

$$
\begin{aligned}
& \left|\left(V^{N} *\left(F\left(K * g_{s}^{N,+}-K * g_{s}^{N,-}\right) S_{s}^{N,+}\right)\right)(x)\right| \\
& \quad \leqslant\left\|F\left(K * g_{s}^{N,+}-K * g_{s}^{N,-}\right)\right\|_{\infty}\left|V^{N} * S_{s}^{N,+}(x)\right| \leqslant M\left|g_{s}^{N,+}(x)\right|
\end{aligned}
$$

which, by Lemma 16 in [5], implies

$$
\begin{aligned}
& \left\|(\mathrm{I}-A)^{\varepsilon / 2} e^{((t-s) / 2) A}\left[V^{N} *\left(F\left(K * g_{s}^{N,+}-K * g_{s}^{N,-}\right) S_{s}^{N,+}\right)\right]\right\|_{\mathbb{L}^{p}(\Omega, H)} \\
& \quad \leqslant C_{M}\left\|e^{((t-s) / 2) A}(\mathrm{I}-A)^{\varepsilon / 2} g_{s}^{N,+}\right\|_{\mathbb{L}^{p}(\Omega, H)} \leqslant C_{M}\left\|(\mathrm{I}-A)^{\varepsilon / 2} g_{s}^{N,+}\right\|_{\mathbb{L}^{p}(\Omega, H)}
\end{aligned}
$$

yielding to the same conclusion.
Moreover, in the proof of

$$
\mathbb{E}\left[\int_{0}^{T} \int_{0}^{T} \frac{\left\|g_{t}^{N,+}-g_{s}^{N,+}\right\|_{-2,2}^{q}}{|t-s|^{1+q \gamma}} d s d t\right] \leqslant C_{T, M, \nu, q}
$$

(and similarly for $g_{t}^{N,-}$ ) the only part which a priori may change is

$$
\begin{aligned}
\mathbb{E} & {\left[\left\|\left\langle S_{r}^{N,+}, F\left(K * g_{s}^{N,+}-K * g_{s}^{N,-}\right) \nabla V^{N}(x-\cdot)\right\rangle\right\|_{-2,2}^{q}\right] } \\
& =\mathbb{E}\left[\left\|\nabla\left(S_{r}^{N,+} F\left(K * g_{s}^{N,+}-K * g_{s}^{N,-}\right) * V^{N}\right)\right\|_{-2,2}^{q}\right] \\
& \leqslant \mathbb{E}\left[\|\left(S_{r}^{N,+} F\left(K * g_{s}^{N,+}-K * g_{s}^{N,-}\right) * V^{N} \|_{-1,2}^{q}\right]\right. \\
& \leqslant C_{M} \mathbb{E}\left[\left\|g_{t}^{N}\right\|_{\mathbb{L}^{2}}^{q}\right] \leqslant C,
\end{aligned}
$$

so in fact this part remains the same. Then one can apply the same arguments of tightness (see Section 2.3).

In the passage to the limit (Section 2.4), the arguments are similar. We use the weak formulation

$$
\begin{aligned}
\left\langle g_{t}^{N,+}, \phi\right\rangle= & \left\langle g_{0}^{N,+}, \phi\right\rangle+\int_{0}^{t}\left\langle S_{s}^{N,+}, F\left(K * g_{s}^{N,+}-K * g_{s}^{N,-}\right) \cdot \nabla\left(V^{N} * \phi\right)\right\rangle d s \\
& +\nu \int_{0}^{t}\left\langle g_{s}^{N,+}, \Delta \phi\right\rangle(x) d s+\frac{1}{N} \sum_{i=1}^{N} \int_{0}^{t} \nabla\left(V^{N} * \phi\right)\left(X_{s}^{i, N,+}\right) \cdot d W_{s}^{i,+}
\end{aligned}
$$

where $\phi$ is a smooth test function with compact support. The only difficult step is proving that

$$
\begin{aligned}
& \lim _{N \rightarrow \infty} \int_{0}^{t}\left\langle S_{s}^{N,+}, F\left(K * g_{s}^{N,+}-K * g_{s}^{N,-}\right) \cdot \nabla\left(V^{N} * \phi\right)\right\rangle d s \\
&=\int_{0}^{t} \int_{\mathbb{R}^{2}} \xi^{+}(s, x) F\left(K *\left(\xi^{+}(s, x)-\xi^{-}(s, x)\right)\right) \cdot \nabla \phi(x) d s
\end{aligned}
$$

The proof is analogous to the case of a single sign; let us recall the main steps. After application of Skorohod's Theorem, due to the a.s. convergence of $g_{s}^{N, \pm}$ to $\xi^{ \pm}$in $\mathbb{L}_{w}^{2}\left([0, T] ; \mathbb{H}_{2}^{\varepsilon}\right)$ and called $f_{s}^{N, \pm}:=K * g_{s}^{N, \pm}$, thanks to the properties of the BiotSavart operator we get that (a.s.) $f_{s}^{N, \pm}$ converge to $K * \xi^{ \pm}$in $\mathbb{L}_{w}^{2}\left([0, T] ; \mathbb{L}^{p}\right)$; and moreover $f_{s}^{N, \pm}$ are (a.s.) bounded in $\mathbb{L}^{2}\left([0, T] ; C^{\tilde{\eta}}\right)$. The last property is used to prove that

$$
\begin{aligned}
\mid F\left(f_{s}^{N,+}-f_{s}^{N,-}\right)(x) \cdot \nabla\left(V^{N}\right. & * \phi)(x) \\
& -\left(F\left(f_{s}^{N,+}-f_{s}^{N,-}\right) \cdot \nabla\left(V^{N} * \phi\right)\right) * V^{N}(x) \mid \leqslant C / N^{\tilde{\eta} \beta} .
\end{aligned}
$$

From this the other steps are easier and equal to the one-sign case.

## B Uniform convergence

All past papers dealing with particle approximation of $2 d$ Navier-Stokes equations prove weak convergence of the empirical measures $S_{t}^{N}$ to the probability law $\xi_{t} d x$. The novelty here is that we prove a stronger convergence, namely the convergence in suitable function spaces of the mollified empirical measure $g_{t}^{N}:=V^{N} * S_{t}^{N}$. Consider for instance the property of uniform convergence in space on compact sets $\left(\mathbb{L}^{2}\right.$ in time). It is not possible to deduce this result from the weak convergence $S_{t}^{N} \rightharpoonup \xi_{t} d x$ (see below). If one only knows that $S_{t}^{N} \rightharpoonup \xi_{t} d x$, and one considers a classical kernel smoothing algorithm $\theta_{\varepsilon_{N}} * S_{t}^{N}$ to approximate the profile $\xi_{t}$ by means of $S_{t}^{N}$, it is not clear how to choose $\theta_{\varepsilon_{N}}$ in such a way to have uniform convergence of $\theta_{\varepsilon_{N}} * S_{t}^{N}$ to $\xi_{t}$. The method described in this paper indicates a strategy for a better particle approximation of solutions to $2 d$ Navier-Stokes equations.

Let us understand more closely the strength of the uniform convergence. It is a strong indication that the particles are distributed quite uniformly in space, they do not have too much concentration, aggregation. Let us make this remark more quantitative.

Proposition 13 Assume that the probability density $V$ has the property

$$
h \mathbf{1}_{[-r, r]^{2}} \leqslant V
$$

for some constants $h, r>0$. Assume that for some $(t, \omega, R)$ we have

$$
\sup _{|x| \leqslant R} g_{t}^{N}(x) \leqslant C .
$$

Then

$$
\begin{equation*}
\sup _{|x| \leqslant R} \operatorname{Card}\left\{i=1, \ldots, N ; X_{t}^{i, N} \in\left[x-\frac{r}{N^{\beta}}, x+\frac{r}{N^{\beta}}\right]^{2}\right\} \leqslant \frac{C}{h} N^{1-2 \beta} \tag{48}
\end{equation*}
$$

Proof. The result simply follows from the inequalities

$$
h N^{2 \beta} \mathbf{1}_{\left[-\frac{r}{N^{\beta}}, \frac{r}{N^{\beta}}\right]^{2}} \leqslant V^{N}
$$

and thus

$$
h \frac{\operatorname{Card}\left\{i=1, \ldots, N: X_{t}^{i, N} \in\left[x-\frac{r}{N^{\beta}}, x+\frac{r}{N^{\beta}}\right]^{2}\right\}}{N^{1-2 \beta}} \leqslant g_{t}^{N}(x) .
$$

Let us give a heuristic interpretation of the previous result.
First consider the case of points $X_{t}^{i, N}$ geometrically uniform in $\left[-\frac{1}{2}, \frac{1}{2}\right]^{2}$ : we consider a uniform grid in $\left[-\frac{1}{2}, \frac{1}{2}\right]^{2}$ of side length $\frac{1}{N^{1 / 2}}$, hence with (roughly) $N$ grid points, and we put one particle per grid point.

In a box $\left[x-\frac{r}{N^{\beta}}, x+\frac{r}{N^{\beta}}\right]^{2} \subset\left[-\frac{1}{2}, \frac{1}{2}\right]^{2}$ we have (approximatively)

$$
\left(\frac{\frac{2 r}{N^{\beta}}}{\frac{1}{N^{1 / 2}}}\right)^{2}=4 r^{2} N^{1-2 \beta}
$$

points. This is exactly estimate (48).
Now, let us break this uniformity. Divide $\left[-\frac{1}{2}, \frac{1}{2}\right]^{2}$ in two sets:

$$
Q=\left[0, \frac{1}{N^{1 / 2}}\right]^{2}, \quad \text { and } \quad Q^{c}:=\left[-\frac{1}{2}, \frac{1}{2}\right]^{2} \backslash Q
$$

Call $X_{t}^{i}$ the first $N-N^{1-\beta}$ particles, $\widetilde{X}_{t}^{i}$ the last $N^{1-\beta}$ ones. Put the $N-N^{1-\beta}$ particles $X_{t}^{i}$ in the grid points of $Q^{c}$, no more than one particle per grid point. Put the $N^{1-\beta}$ particles $\widetilde{X}_{t}^{i}$ in $Q_{0}=\left[0, \frac{r}{N^{\beta}}\right]^{2}$, which is contained in $Q$ (for large $N$ ), being $\beta<\frac{1}{4}$.

Then, all cubes $Q_{x}:=\left[x-\frac{r}{N^{\beta}}, x+\frac{r}{N^{\beta}}\right]^{2} \subset Q^{c}$ contain, as above, at most $4 r^{2} N^{1-2 \beta}$ particles. But $Q_{0}$ contains $N^{1-\beta}$ particles, much more than $4 r^{2} N^{1-2 \beta}$ for large $N$. This is an example of configuration which does not fulfill estimate (48). For such configuration, however, we still have

$$
\left\langle S_{t}^{N}, \phi\right\rangle \rightarrow\left\langle\xi_{t}, \phi\right\rangle
$$

for all continuous bounded $\phi$. Let us only check that this is true for $\phi=\mathbf{1}_{\mathcal{O}}$ when $\mathcal{O}$ is an open set (it is not continuous but it is a good heuristic indication). We have

$$
\begin{aligned}
\left\langle S_{t}^{N}, \phi\right\rangle & =\left\langle S_{t}^{N}, \mathbf{1}_{\mathcal{O} \cap Q}\right\rangle+\left\langle S_{t}^{N}, \mathbf{1}_{\mathcal{O} \cap Q^{c}}\right\rangle \leqslant\left\langle S_{t}^{N}, \mathbf{1}_{Q}\right\rangle+\left\langle S_{t}^{N}, \mathbf{1}_{\mathcal{O} \cap Q^{c}}\right\rangle \\
& \sim \frac{1}{N} N^{1-\beta}+|\mathcal{O}| \xrightarrow[N \rightarrow \infty]{ }|\mathcal{O}|=\left\langle\xi_{t}, \phi\right\rangle
\end{aligned}
$$

where $\xi_{t}=1$. The excessive concentration in $Q_{0}$ does not prevent weak convergence but it is not allowed by the stronger convergence proved here.

## Acknowledgements

C. O. is partially supported by FAPESP by the grant 2018/15258-7 and by CNPq by the grant $426747 / 2018-6$. This project has received funding from the CNRSFAPESP cooperation, grant $n^{\circ}$ PRC2726. M.S. also thanks Labex CEMPI (ANR-11-LABX-0007-01), and the ANR grant MICMOV (ANR-19-CE40-0012) of the French National Research Agency (ANR), and finally the European Research Council (ERC) under the European Union's Horizon 2020 research and innovative program (grant agreement $\mathrm{n}^{\circ} 715734$ ).

## References

[1] M. Ben-Artzi, Global solutions of two-dimensional Navier-Stokes and Euler equations, Arch. Rational Mech. Anal. 128, 4, 329-358, 1994.
[2] A. Bertozzi and A. Majda, Vorticity and incompressible flow, Cambridge University Press, Cambridge, 2002.
[3] J.-Y. Chemin, About Navier-Stokes system, Publication du Laboratoire d'Analyse Numérique R 96023, 1996.
[4] A. J. Chorin, Vorticity and Turbulence, Applied Mathematical Sciences, 103, Springer-Verlag, 1994.
[5] F. Flandoli, M. Leimbach, C. Olivera, Uniform approximation of FKPP equation by stochastic particle systems, Journal of Mathematical Analysis and Applications 473, 1, 27-52, 2019.
[6] F. Flandoli, M. Leocata, C. Ricci, The Vlasov-Navier-Stokes equations as a mean field limit, Discrete \& Continuous Dynamical Systems - B. 22. 2018.
[7] F. Flandoli, M. Leocata, A particle system approach to aggregation phenomena, J. Appl. Probab. 56, no. 1, 282-306, 2019.
[8] J. Fontbona A probabilistic interpretation and stochastic particle approximations of the 3-dimensional Navier-Stokes equations, Probab. Theory Relat. Fields 136, 102-156, 2006.
[9] N. Fournier, M. Hauray, S. Mischler, Propagation of chaos for the $2 d$ viscous vortex model, J. Eur. Math. Soc. 16, 1423-1466, 2014.
[10] I. Gallagher, T. Gallay, Uniqueness for the two-dimensional Navier-Stokes equation with a measure as initial vorticity, Math. Ann. 332, 2, 287-327, 2005.
[11] Y. Giga, T. Miyakawa, H. Osada, Two-Dimensional Navier-Stokes Flow with Measures as Initial Vorticity, Arch. Rational Mech. Anal. 104, 223-250, 1988.
[12] I. Gyongy, N. Krylov, Existence of strong solutions for Itô stochastic equations, via approximations, Prob. Th. Relat. Fields 105, 143-158, 1996.
[13] N. Ikeda, S. Watanabe, Stochastic Differential Equations and Diffusion Processes, North Holland Pub. Co. Amsterdam, Oxford, New York, 1981.
[14] P. E. Jabin, Zhenfu Wa, Quantitative estimates of propagation of chaos for stochastic systems with $W^{-1, \infty}$ kernel, Inventiones mathematicae 214, 523-591, 2018.
[15] B. Jourdain, S. Méléard, Propagation of chaos and fluctuations for a moderate model with smooth initial data, Ann. IHP (B) Probabilités et Statistiques 34, 727-766, 1998.
[16] T. Kato, The Navier-Stokes equation for an incompressible fluid in with a measure as the initial vorticity Differential Integral Equations 7, 3-4, 949-966, 1994.
[17] P. G. Lemarie-Rieusset, Recent developments in the Navier-Stokes problem, Chapman Hall/CRC, 2002.
[18] C. Marchioro, M. Pulvirenti, Hydrodynamics in two dimensions and vortex theory, Commun. Maths. Phys 84, 483-503, 1982.
[19] S. Meleard, A trajectorial proof of the vortex method for the $2 d$ Navier-Stokes equation Ann. of Appl. Prob. 10, 1197-1211, 2000.
[20] S. Meleard, Monte-Carlo approximations of the solution of the $2 d$ Navier-Stokes equation with finite measure initial data. Prob. Th. Rel. Fields. 121, 367-388, 2001.
[21] H. Osada, Propagation of Chaos for the Two Dimensional Navier-Stokes Equations, Probabilistic Methods in Math. Phys., (K. Ito and N. Ikeda Eds.), Tokyo: Kinokuniya, 303-334, 1987.
[22] K. Oelschläger, A law of large numbers for moderately interacting diffusion processes, Zeitschrift fur Wahrsch. Verwandte Gebiete 69, 279-322, 1985.
[23] C. Olivera, M. Simon, Non-local Conservation Law from Stochastic Particle Systems, Journal of Dynamics and Differential Equations 30, 4, 1661-1682, 2018.
[24] A. Pazy, Semi-groups of Linear Operators and Applications to Partial Differential Equations, Springer, 1983.
[25] H. Sohr, The Navier-Stokes equations. An elementary functional analytic approach, Birkhäuser Advanced Texts, Birkhäuser Verlag, Basel, 2001.
[26] R. Temam, Navier-Stokes equations and nonlinear functional analysis, Second edition. CBMS-NSF Regional Conference Series in Applied Mathematics, 66. Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 1995.
[27] J. Simon, Compact Sets in the Space $L^{p}(0, T B)$, Annali di Matematica Pura ed Applicata 146, 65-96, 1986.
[28] H. Triebel, Interpolation Theory, Function Spaces, Differential Operators, North-Holland, Amsterdam, 1978.
[29] J. M. A. M. van Neerven, M. C. Veraar, L. Weis, Stochastic integration in UMD Banach spaces, The Annals of Probab. 35, 1438-1478, 2007.


[^0]:    *Scuola Normale Superiore Italy, E-mail: franco.flandoli@sns.it
    ${ }^{\dagger}$ Departamento de Matemática, Universidade Estadual de Campinas, Brazil. E-mail: colivera@ime.unicamp.br
    ${ }^{\ddagger}$ Inria, Univ. Lille, CNRS, UMR 8524 - Laboratoire Paul Painlevé, F-59000 Lille. E-mail: marielle.simon@inria.fr.

