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Hélène Barucq, Henri Calandra, Aurélien Citrain, Julien Diaz, Christian Gout. High-order DG-SE approximations of elastodynamic wave problems with stable Perfectly Matched Layers. MATHIAS 2019, Oct 2019, Serris, France. hal-02425960

HAL Id: hal-02425960 https://inria.hal.science/hal-02425960

Submitted on 7 Jan 2020

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High-order DG-SE approximations of elsatodynamic wave problems with stable Perfectly Matched Layers[†]

Hélène Barucq¹, Henri Calandra², Aurélien Citrain^{3,1}, <u>Julien Diaz</u>¹ and Christian Gout³

Mathias 2019







The authors thank the M2NUM project which is co-financed by the European Union with the European regional development fund (ERDF, HN0002137) and by the Normandie Regional Council.

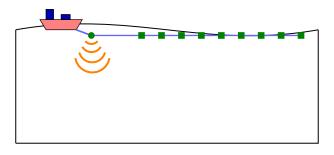
¹ Team project Magique.3D, INRIA, E2S UPPA, CNRS, Pau, France.

² TOTAL SA, CSTJF, Pau, France,

³ LMI. Normandie Université, INSA ROUEN, 76000 Rouen, France.

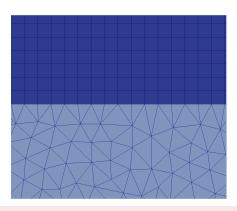
[†] This work is dedicated to the memory of Dimitri Komatitsch.

Seismic imaging



Why using hybrid meshes?

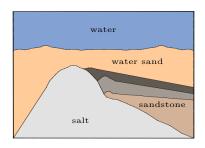




- Useful when the use of unstructured grid is non-sense (e.g. medium with a layer of water).
- Well suited for the coupling of numerical methods in order to reduce the computational cost and improve the accuracy.

Why using hybrid meshes?





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- Well suited for the coupling of numerical methods in order to reduce the computational cost and improve the accuracy.

Elastodynamic system

$$x \in \Omega \subset \mathbb{R}^d$$
, $t \in [0, T]$, $T > 0$:

$$\begin{cases} \rho(x) \frac{\partial v}{\partial t}(x,t) &= \nabla \cdot \underline{\underline{\sigma}}(x,t), \\ \frac{\partial \underline{\underline{\sigma}}}{\partial t}(x,t) &= \underline{\underline{C}}(x)\underline{\underline{\epsilon}}(v(x,t)). \end{cases}$$

With:

- \bullet $\rho(x)$ the density,
- $\underline{C}(x)$ the elasticity tensor,
- $\underline{\epsilon}(x,t)$ the deformation tensor,
- v(x, t), the wavespeed,
- $\underline{\sigma}(x, t)$ the strain tensor.

Elasticus software

Software written in Fortran for wave propagation simulation in the time domain

Features

Simulation:

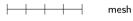
- on various types of meshes (unstructured triangles and tetrahedra),
- on heterogeneous media (acoustic, elastic and elasto-acoustic).
- Discontinuous Galerkin (DG) based on unstructured triangles and unstructured tetrahedra,
- with various time-schemes: Runge-Kutta (2 or 4), Leap-Frog,
- with multi-order computation(p-adaptivity)...

Table of contents

- DGm and SEm
- DG/SEM coupling
- Perfectly Matched Layer(PML)

Discontinuous Galerkin Method

Use discontinuous functions :



continuous

discontinuous

h adaptivity:

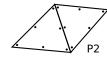


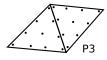
p adaptivity :



Degrees of freedom necessary on each cell:



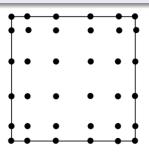




Spectral Element Method

General principle

- Finite Element Method (FEM) discretization + Gauss-Lobatto quadrature,
- ullet Gauss-Lobatto points as degrees of freedom (gives us exponential convergence on L^2 -norm).



$$\bullet \ \varphi_i(\xi_j) = \delta_{ij}.$$

Advantages of each method

DG

- Element per element computation (hp-adaptivity).
- Time discretization quasi explicit (block diagonal mass matrix).
- Simple to parallelize.
- Robust to brutal changes of physics and geometry

SEM

- Couples the flexibility of FEM with the accuracy of the pseudo-spectral method.
- Simplifies the mass and stiffness matrices (mass matrix diagonal).

Hybrid meshes structures

- Aim at coupling P_k and Q_k structures.
- Need to extend or split some structures (e.g. neighbour indices).
- Define new face matrices:

$$M^{K,L}_{ij} = \int_{K\cap L} \phi^K_i \phi^L_j, \quad M^{K,L}_{ij} = \int_{K\cap L} \psi^K_i \psi^L_j, \quad M^{K,L}_{ij} = \int_{K\cap L} \phi^K_i \psi^L_j.$$

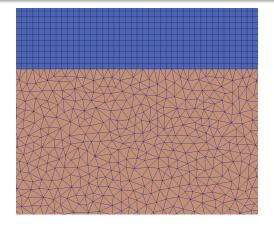


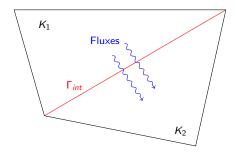




Global context

• Domain in two parts : $\Omega_{h,1}$ (structured quadrangles + SEM), $\Omega_{h,2}$ (unstructured triangles + DG).





Definitions

Jump and average.

$$[[u]] = (u_{K_1} \mathbf{n}_{K_1} + u_{K_2} \mathbf{n}_{K_2})$$

$$\{\{u\}\}=\frac{1}{2}(u_{K_2}+u_{K_1})$$

SEM variational formulation :

$$\begin{cases} \int_{\Omega_{h,1}} \rho \partial_t \mathbf{v}_1 \cdot \mathbf{w}_1 = -\int_{\Omega_{h,1}} \sigma_1 : \nabla \mathbf{w}_1 + \int_{\Gamma_{out,1}} (\sigma_1 \mathbf{n}_1) \cdot \mathbf{w}_1, \\ \\ \int_{\Omega_{h,1}} \partial_t \sigma_1 : \xi_1 = \int_{\Omega_{h,1}} (C\xi_1) : \nabla \mathbf{v}_1. \end{cases}$$

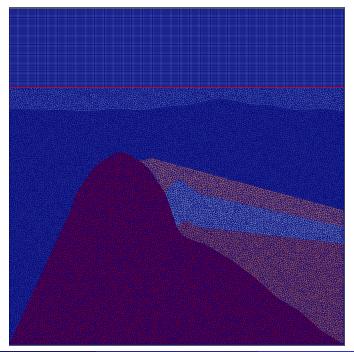
DG variational formulation :

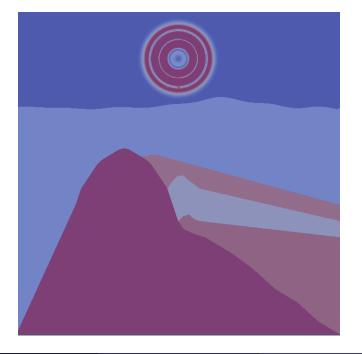
$$\begin{cases} \int_{\Omega_{h,2}} \rho \partial_t v_2 \cdot w_2 = -\int_{\Omega_{h,2}} \sigma_2 : \nabla w_2 + \int_{\Gamma_{out,2}} (\sigma_2 \mathbf{n}_2) \cdot w_2 + \int_{\Gamma_{int}} \{ \{\sigma_2\} \} : [[w_2]], \\ \int_{\Omega_{h,2}} \partial_t \sigma_2 : \xi_2 = -\int_{\Omega_{h,2}} (\boldsymbol{\nabla} \cdot (\boldsymbol{C}\xi_2)) \cdot v_2 + \int_{\Gamma_{out,2}} (\boldsymbol{C}\xi_2 \mathbf{n}_2) \cdot v_2 + \int_{\Gamma_{int}} \{ \{v_2\} \} \cdot [[\boldsymbol{C}\xi_2]]. \end{cases}$$

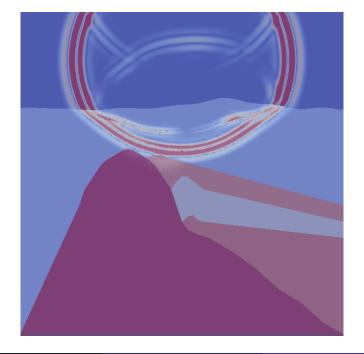
$$\begin{cases} \int_{\Omega_{h,1}} \rho \partial_t \mathbf{v}_1 \cdot \mathbf{w}_1 + \int_{\Omega_{h,2}} \rho \partial_t \mathbf{v}_2 \cdot \mathbf{w}_2 = -\int_{\Omega_{h,1}} \sigma_1 : \nabla \mathbf{w}_1 - \int_{\Omega_{h,2}} \sigma_2 : \nabla \mathbf{w}_2 \\ + \int_{\Gamma_{out,1}} (\sigma_1 \mathbf{n}_1) \cdot \mathbf{w}_1 + \int_{\Gamma_{out,2}} (\sigma_2 \mathbf{n}_2) \cdot \mathbf{w}_2 + \int_{\Gamma_{int}} \{ \{\sigma_2\} \} : [[\mathbf{w}_2]] \\ + \int_{\Gamma_{1/2}} [[\sigma \mathbf{w}]], \\ \int_{\Omega_{h,1}} \partial_t \sigma_1 : \xi_1 + \int_{\Omega_{h,2}} \partial_t \sigma_2 : \xi_2 = \int_{\Omega_{h,1}} (C\xi_1) : \nabla \mathbf{v}_1 - \int_{\Omega_{h,2}} (\nabla \cdot (C\xi_2)) \cdot \mathbf{v}_2 \\ + \int_{\Gamma_{out,2}} (C\xi_2 \mathbf{n}_2) \cdot \mathbf{v}_2 + \int_{\Gamma_{int}} \{ \{\mathbf{v}_2\} \} \cdot [[C\xi_2]] \\ + \int_{\Gamma_{1/2}} [[(C\xi) \mathbf{v}]]. \end{cases}$$

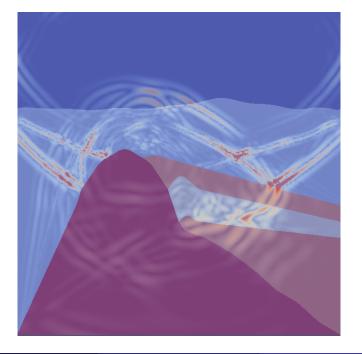
$$\begin{cases} \int_{\Omega_{h,1}} \rho \partial_{t} v_{1} \cdot w_{1} + \int_{\Omega_{h,2}} \rho \partial_{t} v_{2} \cdot w_{2} = -\int_{\Omega_{h,1}} \sigma_{1} : \nabla w_{1} - \int_{\Omega_{h,2}} \sigma_{2} : \nabla w_{2} \\ + \int_{\Gamma_{out,1}} (\sigma_{1} \mathbf{n}_{1}) \cdot w_{1} + \int_{\Gamma_{out,2}} (\sigma_{2} \mathbf{n}_{2}) \cdot w_{2} + \int_{\Gamma_{int}} \{\{\sigma_{2}\}\} : [[w_{2}]] \\ + \int_{\Gamma_{1/2}} \{\{\sigma\}\} : [[w]] + [[\sigma]] \cdot \{\{w\}\}, \\ \int_{\Omega_{h,1}} \partial_{t} \sigma_{1} : \xi_{1} + \int_{\Omega_{h,2}} \partial_{t} \sigma_{2} : \xi_{2} = \int_{\Omega_{h,1}} (C\xi_{1}) : \nabla v_{1} - \int_{\Omega_{h,2}} (\nabla \cdot (C\xi_{2})) \cdot v_{2} \\ + \int_{\Gamma_{out,2}} (C\xi_{2} \mathbf{n}_{2}) \cdot v_{2} + \int_{\Gamma_{int}} \{\{v_{2}\}\} \cdot [[C\xi_{2}]] \\ + \int_{\Gamma_{1/2}} \{\{v\}\} \cdot [[C\xi]] + \{\{C\xi\}\} : [[v]]. \end{cases}$$

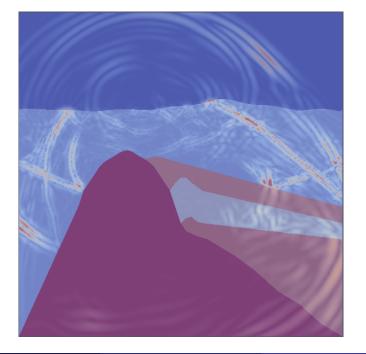
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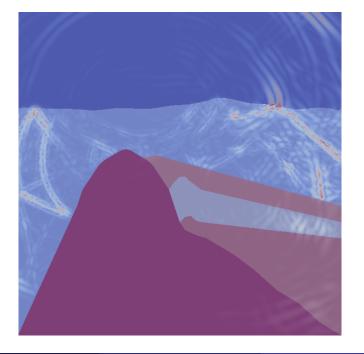


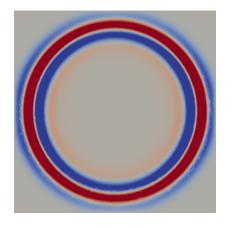


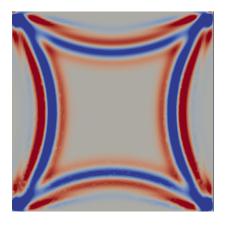


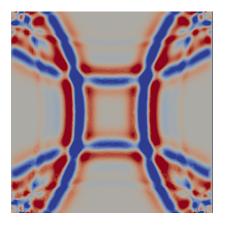


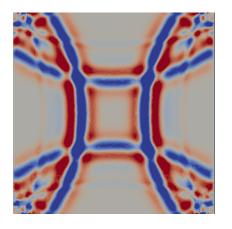




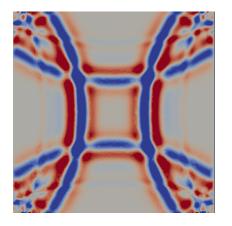




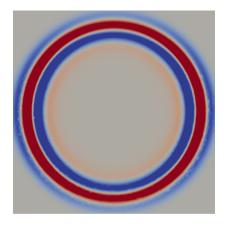


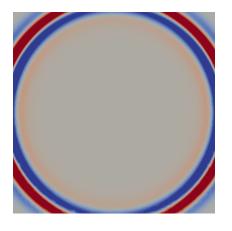


• Easy to implement



- Easy to implement
- Not really adapted to the simulation of "infinite" domain

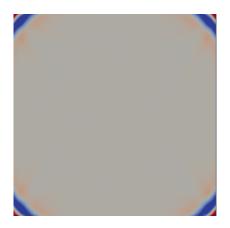




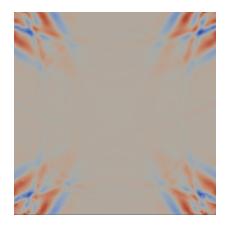




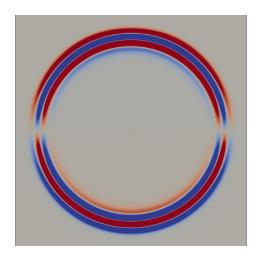
• Possibility to simulate an infinite domain

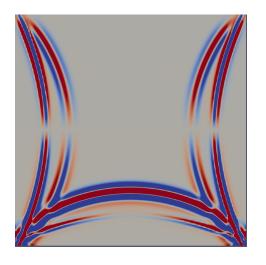


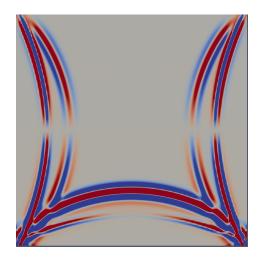
- Possibility to simulate an infinite domain
- Complicated to implement when the order of the condition increase



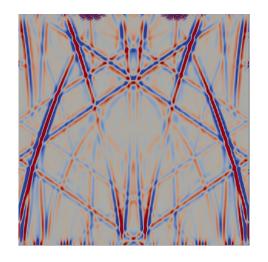
- Possibility to simulate an infinite domain
- Complicated to implement when the order of the condition increase
- Apparition of reflection when it comes to waves with grazing incidence



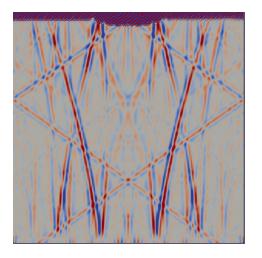




• Possibility to simulate an infinite domain



• Possibility to simulate an infinite domain



- Possibility to simulate an infinite domain
- Some stability problems on elastic case using DGM

Step 1: Rewrite the system in the frequency domain

$$\begin{cases} i\omega\rho\mathbf{v} = \mathbf{\nabla}\cdot\underline{\underline{\sigma}} \\ i\omega\underline{\underline{\sigma}} = \underline{\underline{C}}:\nabla\mathbf{v} \end{cases}$$

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$$\begin{cases} i\omega\rho\mathbf{v} = \nabla \cdot \underline{\underline{\sigma}} \\ i\omega\underline{\underline{\sigma}} = \underline{\underline{C}} : \nabla\mathbf{v} \end{cases}$$

Step 2: Introduce a new system of complex coordinates

Define d_z the damping :

$$\begin{cases} d_z > 0 & \text{if } z \in \Omega_{PML} \\ d_z = 0 & \text{if } z \notin \Omega_{PML} \end{cases}$$

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$$\tilde{z}(z) = z - \frac{i}{\omega} \int_0^z d_z(s) ds$$

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$$\tilde{z}(z) = z - \frac{i}{\omega} \int_0^z d_z(s) ds$$

Define the associate differential operator :

$$\partial_{\tilde{z}} = \frac{i\omega}{i\omega + d_z} \partial_z = 1 - \frac{d_z}{i\omega + d_z} \partial_z$$

Step 3: Rewrite the system

$$\begin{split} i\omega\rho \mathbf{v}_{\mathbf{x}} &= \quad \partial_{\mathbf{x}}\sigma_{\mathbf{x}\mathbf{x}} + \partial_{\bar{\mathbf{z}}}\sigma_{\mathbf{x}\mathbf{z}}, \\ i\omega\rho \mathbf{v}_{\mathbf{z}} &= \quad \partial_{\mathbf{x}}\sigma_{\mathbf{x}\mathbf{z}} + \partial_{\bar{\mathbf{z}}}\sigma_{\mathbf{z}\mathbf{z}}, \\ i\omega\sigma_{\mathbf{x}\mathbf{x}} &= \quad (\lambda + 2\mu)\partial_{\mathbf{x}}\mathbf{v}_{\mathbf{x}} + \lambda\partial_{\bar{\mathbf{z}}}\mathbf{v}_{\mathbf{z}}, \\ i\omega\sigma_{\mathbf{z}\mathbf{z}} &= \quad \lambda\partial_{\mathbf{x}}\mathbf{v}_{\mathbf{x}} + (\lambda + 2\mu)\partial_{\bar{\mathbf{z}}}\mathbf{v}_{\mathbf{z}}, \\ i\omega\sigma_{\mathbf{x}\mathbf{z}} &= \quad \mu(\partial_{\mathbf{x}}\mathbf{v}_{\mathbf{z}} + \partial_{\bar{\mathbf{z}}}\mathbf{v}_{\mathbf{z}}). \end{split}$$



Martin, Roland and Komatitsch, Dimitri and Gedney, Stephen D and Bruthiaux, Emilien and others

A high-order time and space formulation of the unsplit perfectly matched layer for the seismic wave equation using Auxiliary Differential Equations (ADE-PML)

Computer Modeling in Engineering and Sciences (CMES),2010

Introduction of 2^d+1 new variables $\psi_\star\in H^1 o$ Simplify the implementation BUT add new equation to the linear system



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Back to original coordinates



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$$\downarrow \qquad \text{Back to original coordinates}$$

$$i\omega\rho v_z = \partial_x\sigma_{xz} + \partial_z\sigma_{zz} - \frac{d_z}{d_z + i\omega}\partial_z\sigma_{zz}$$



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$$\downarrow \qquad \text{Back to time domain}$$



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Introduction of 2^d+1 new variables $\psi_\star\in H^1\to \mathsf{Simplify}$ the implementation BUT add new equation to the linear system

Definition of the memory variable $\psi_{\sigma_{\it ZZ}}$:

$$\psi_{\sigma_{zz}} = -\frac{d_z}{d_z + \partial_t} \partial_z \sigma_{zz}$$

Definition of the memory variable $\psi_{\sigma_{\it zz}}$:

$$\psi_{\sigma_{\sf zz}} = -rac{d_{\sf z}}{d_{\sf z}+\partial_{\sf t}}\partial_{\sf z}\sigma_{\sf zz}$$

$$\rho \partial_t \mathbf{v}_{\mathbf{z}} = \partial_{\mathbf{x}} \sigma_{\mathbf{x}\mathbf{z}} + \partial_{\mathbf{z}} \sigma_{\mathbf{z}\mathbf{z}} + \psi_{\sigma_{\mathbf{z}\mathbf{z}}}$$

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$$\partial_t \psi_{\sigma_{zz}} = - \textit{d}_z \partial_z \sigma_{zz} - \textit{d}_z \psi_{\sigma_{zz}}$$

ADE with C-PML (Convolutional PML)

$$\partial_{\tilde{z}} = \left(1 - \frac{d_z}{\alpha_z + i\omega + d_z}\right) \partial_z$$

Definition of the memory variable $\psi_{\sigma_{zz}}$:

$$\psi_{\sigma_{zz}} = -\frac{d_z}{d_z + \partial_t} \partial_z \sigma_{zz}$$

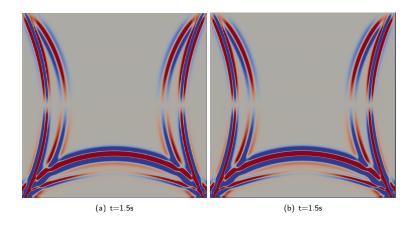
$$\rho \partial_t \mathbf{v}_{\mathbf{z}} = \partial_{\mathbf{x}} \sigma_{\mathbf{x}\mathbf{z}} + \partial_{\mathbf{z}} \sigma_{\mathbf{z}\mathbf{z}} + \psi_{\sigma_{\mathbf{z}\mathbf{z}}}$$

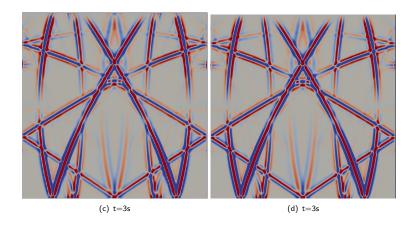
$$\partial_t \psi_{\sigma_{zz}} = - \mathit{d}_z \partial_z \sigma_{zz} - \mathit{d}_z \psi_{\sigma_{zz}}$$

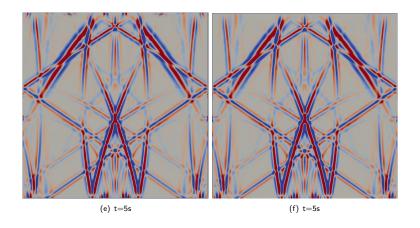
ADE with C-PML (Convolutional PML)

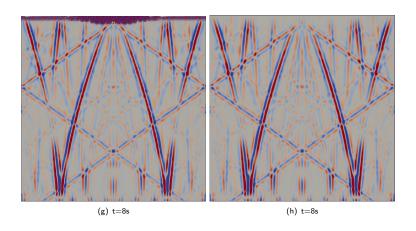
$$\partial_{\check{z}} = \left(1 - rac{d_z}{lpha_z + i\omega + d_z}
ight)\partial_z$$

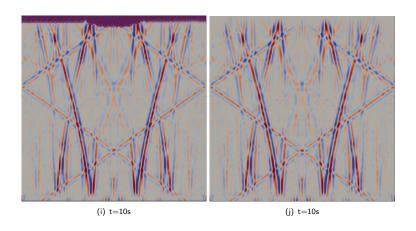
$$\partial_t \psi_{\sigma_{zz}} = -d_z \partial_z \sigma_{zz} - (d_z + \alpha_z) \psi_{\sigma_{zz}}$$





















Conclusion and perspectives

Conclusion

- SEM is more efficient on structured quadrangle mesh than DG
- Build a variational formulation for DG/SEM coupling and find a CFL condition that ensures stability
- § Show the utility of using hybrid meshes and method coupling (reduce computational cost,...)

Achievements

- Implement DG/SEM coupling on the code (2D) √
- Develop DG/SEM coupling in 3D √
- Develop PML in the hexahedral part √

Thank you for your attention !

Questions?