## Optimality of linear optical Bell measurements. How much can ancillae help? <br> Andrea Olivo, Frédéric Grosshans

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# Optimality of linear optical Bell measurements 

How much can ancillæ help?

Andrea Olivo ${ }^{1,2}$, Frédéric Grosshans ${ }^{1}$

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## The task: unambiguous Bell measurement

$$
\begin{array}{ll}
\left|\phi^{+}\right\rangle=\frac{1}{\sqrt{2}}(|00\rangle+|11\rangle) & \left|\psi^{+}\right\rangle=\frac{1}{\sqrt{2}}(|01\rangle+|10\rangle) \\
\left|\phi^{-}\right\rangle=\frac{1}{\sqrt{2}}(|00\rangle-|11\rangle) & \left|\psi^{-}\right\rangle=\frac{1}{\sqrt{2}}(|01\rangle-|10\rangle)
\end{array}
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## Bell Measurement

Projective measurement on the Bell basis

## Unambiguous

Outcome never wrong, but can fail with probability $\mathcal{P}_{\text {fail }}$

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## Bell Measurement

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Quantum teleportation, dense coding, entanglement swapping...

## The framework: static linear optics

Why? Because...

- ...experimentally easier: no feedforward



## The framework: static linear optics

Why? Because...

- ...experimentally easier: no feedforward
- ...simple mathematical framework to work with



## State of the art: Bell measurement without ancilla

[Braunstein and Mann, 1995]

- Known $\mathcal{P}_{\text {succ }}=\frac{1}{2}$ scheme [Braunstein and Mann, 1995]
- Optimal, even with feedforward [Calsamiglia and Lütkenhaus, 2001]


$$
\begin{aligned}
&\left|\psi^{+}\right\rangle \longrightarrow \frac{1}{\sqrt{2}}(|1100\rangle+|0011\rangle) \\
&\left|\psi^{-}\right\rangle \longrightarrow \frac{1}{\sqrt{2}}(|1001\rangle-|0110\rangle) \\
&\left|\phi^{ \pm}\right\rangle \longrightarrow \frac{i}{2}(|2000\rangle \pm|0200\rangle+ \\
&\quad|0020\rangle \pm|0002\rangle)
\end{aligned}
$$

## Beyond 1/2

## [Grice, 2011]

- $\mathcal{P}_{\text {succ }}=\frac{3}{4}$ with an extra $\left|\phi^{+}\right\rangle$ as ancilla
- $\mathcal{P}_{\text {fail }}=2^{-N}$ with GHZ-like states of $2^{N}-2$ photons



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[Ewert and van Loock, 2014]
- $\mathcal{P}_{\text {succ }}=\frac{3}{4}$ with 4 single photons
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- Cons: limited efficiency Pros: unentangled ancilla (for the two above)


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## Our work

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We provide:

- An analytical bound for polarization-preserving interferometers
- A thorough numerical search for generic (small) interferometers


## Analytical upper bound

## Overview of the proof



Input: $\left|\Psi_{\text {in }}\right\rangle=\left|\beta_{i}\right\rangle|\Upsilon\rangle$, with $|\Upsilon\rangle=\sum_{\lambda=0}^{k} v_{\lambda}|\Upsilon, \lambda\rangle$ a $k$-photon ancilla Each $|\Upsilon, \lambda\rangle$ is a state with $\lambda$ horizontally polarized photons

## Overview of the proof



Output statistics unchanged under a projective measurement of the horizontally-polarized photon number of $\left|\Psi_{\text {out }}\right\rangle$

## Overview of the proof



Measurement commutes $\Longleftrightarrow$ polarization-preserving network

## Overview of the proof


$\left|\psi^{+}\right\rangle,\left|\psi^{-}\right\rangle$and $\left|\phi^{ \pm}\right\rangle$are mapped in three orthogonal subspaces $\forall \lambda$

## Overview of the proof



Non-orthogonal post-measurement states $\left|\Lambda^{ \pm}\right\rangle$corresponding to $\left|\phi^{+}\right\rangle$ and $\left|\phi^{-}\right\rangle$can be discriminated with optimal $\mathcal{P}_{\text {succ }} \leq 1-\left|\left\langle\Lambda^{+} \mid \Lambda^{-}\right\rangle\right|$.

## Upper bounds

Ancilla polarization-based upper bound:

$$
\begin{equation*}
\mathcal{P}_{\text {fail }} \geq \frac{1}{2}\left(\max _{\lambda \text { even }}\left|v_{\lambda}\right|^{2}+\max _{\lambda \text { odd }}\left|v_{\lambda}\right|^{2}\right) \tag{1}
\end{equation*}
$$

(looser) Photon number-based upper bound:

$$
\begin{equation*}
\mathcal{P}_{\text {fail }} \geq \frac{1}{\lceil k+1\rceil_{\text {even }}} \tag{2}
\end{equation*}
$$

N.B. The Grice's schemes saturate both bounds!

## Linear network optimizer

## Polynomial representation

Second quantization, $n$-modes interferometer.
Input/output state $\longrightarrow$ polynomial in the input/output mode operators:

$$
\left|\psi_{\text {in }}\right\rangle=P_{\text {in }}\left(a_{1}^{\dagger}, \ldots, a_{n}^{\dagger}\right)|0\rangle \quad\left|\psi_{\text {out }}\right\rangle=P_{\text {out }}\left(c_{1}^{\dagger}, \ldots, c_{n}^{\dagger}\right)|0\rangle
$$

Action of the interferometer on $\left|\psi_{\text {in }}\right\rangle \longrightarrow$ unitary transformation $U$ acting on the mode operators:

$$
a_{i}^{\dagger}=\sum_{j=1}^{n} u_{i j} c_{j}^{\dagger}
$$

Measurement of $\left|\psi_{\text {out }}\right\rangle$ by an array of PNRD.
Detection event $\longrightarrow$ a configuration of clicks at the output.

## The method

Task: maximize $\mathcal{P}_{\text {succ }}$ over all possible $n$-modes interferometers, using $k$-photon ancilla $|\Upsilon\rangle$

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Above steps are automatic symbolic computation. Then:
3. Functions $p_{i}^{e}(U)$ are hardcoded in $C$
4. A f.o.m. $f(U)$ is numerically optimized over $U(n)$

## An example: no ancilla on $\left|\phi^{+}\right\rangle$

$$
n=4, k=0 \quad P_{\mathrm{in}}=\frac{1}{\sqrt{2}}\left(a_{1}^{\dagger} a_{3}^{\dagger}+a_{2}^{\dagger} a_{4}^{\dagger}\right)
$$

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\begin{gathered}
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\downarrow \\
P_{\text {out }}=\frac{1}{\sqrt{2}}\left(\sum_{j_{1}} u_{1 j_{1}} c_{j_{1}}^{\dagger}\right)\left(\sum_{j_{2}} u_{3 j_{2}} j_{j_{2}}^{\dagger}\right)+\frac{1}{\sqrt{2}}\left(\sum_{j_{3}} u_{2 j_{3}} c_{j_{3}}^{\dagger}\right)\left(\sum_{j_{4}} u_{4 j_{4}} c_{j_{4}}^{\dagger}\right)
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\downarrow \\
2000 \\
0200 \\
0020 \\
0002 \\
1100 \\
1010 \\
1001 \\
0110 \\
0101 \\
0011
\end{gathered} \begin{gathered}
\left(u_{11} u_{31}+u_{21} u_{41}\right. \\
\left(u_{11} u_{32}+u_{12} u_{31}+u_{13} u_{31}+u_{21} u_{42} u_{43}+u_{23} u_{41}\right) / \sqrt{2} \\
\left(u_{11} u_{34}+u_{14} u_{31}+u_{21} u_{44}+u_{24} u_{41}\right) / \sqrt{2} \\
\left(u_{12} u_{33}+u_{13} u_{32}+u_{22} u_{43}+u_{23} u_{42}\right) / \sqrt{2} \\
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## Optimization: a case for symbolic computation

Why not using pen and paper for $p_{i}^{e}(U)$ ?
For (small!) $\frac{3}{4}$ Grice's scheme:
Total number of functions 1320
Independent just 5

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For (small!) $\frac{3}{4}$ Grice's scheme:

| Total number of functions | 1320 |
| :---: | :---: |
| Independent | just 5 |

But 2 ${ }^{\text {nd }}$ iteration, $\frac{7}{8}$ :

| Total number of functions | 490314 |
| :---: | :---: |
| Independent | 22 |
| Total number of terms | $1,8 \cdot 10^{6}$ |

## Results

## Ancilla: a single photon

Simplest type of ancilla, no known scheme

- Polarization-preserving bound predicts $\mathcal{P}_{\text {succ }} \leq \frac{1}{2}$
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We find no advantage using just one extra photon.

## More single, unentangled photons

| $k$ photons | Pol. Pres. bound | Num. search | Explicit scheme |
| :---: | :---: | :---: | :---: |
| 2 | $5 / 8$ | $5 / 8$ | $5 / 8$ |
| 4 | $3 / 4$ | $3 / 4$ | $3 / 4$ |
| 6 | $13 / 16$ | $3 / 4$ | - |
| 8 | $25 / 32$ | - | $49 / 64$ |
| 12 | $13 / 16$ | - | $25 / 32$ |

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just 2 single photons beat 50\% limit!

## The "half"-Ewert \& van Loock scheme



Also independently found by Ewert \& van Loock

## Entangled ancillæ

Do entangled ancillæ help?

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- Pol. pres. bound predicts $\mathcal{P}_{\text {fail }} \gtrsim \frac{1}{\sqrt{\pi k}} \geq \frac{1}{k}$


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- Pol. pres. bound predicts $\mathcal{P}_{\text {fail }} \gtrsim \frac{1}{\sqrt{\pi k}} \geq \frac{1}{k}$
- Grice's $\frac{7}{8}$ scheme needs $n=16, k=6$ : barely out of reach

Conclusion

## Achievements

What can we say on the optimal $\mathcal{P}_{\text {succ }}$ of unambiguous Bell measurement?

- Upper bound for polarization-preserving interferometers, saturated by known schemes
- Hybrid numerical/symbolical search
- Confirms optimality of (some) known schemes
- New 2-photon scheme with $\mathcal{P}_{\text {succ }}=\frac{5}{8}$
- Explored several ancillæ $\longrightarrow \frac{3}{4}$ stays the best (for small networks)

Thank you

