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# A direct proof of the equivalence of side conditions for strictly positive real matrix transfer functions

Augusto Ferrante Alexander Lanzon Bernard Brogliato

Abstract—This brief note proves in a direct way that two different side conditions, which have been used in the literature to characterize strictly positive real matrix transfer functions in the frequency domain, are equivalent.

#### I. INTRODUCTION

The frequency domain conditions characterizing the fact that a matrix transfer function F is strictly positive real involve a positivity constraint at infinite frequency. This constraint — usually referred to as *side condition* — has been a source of confusion and controversy in the literature for more than a decade. As pointed out in [3], the side conditions used in [6], [7], [8] were incorrect as they had some inconsistencies. To fix the problem, [3] proposed a new condition at infinite frequency, i.e.

$$\lim_{\omega \to \infty} \omega^{2\rho} \det \left( F(j\omega) + F(-j\omega)^{\top} \right) > 0 \tag{1}$$

where  $\rho$  is the dimension of  $\ker (F(\infty) + F(\infty)^{\top})$ .

On the other hand, a different, but equally valid, condition at infinite frequency was proposed the second edition of the book by Khalil published in 1996 (see [5, Lemma 10.1]); such a condition, that reads as follows, was recently used in [4] to establish a counterpart result for negative imaginary systems (see [4, Remark 1]):  $\exists \delta > 0, \sigma_0 > 0$  such that

$$\underline{\sigma} \left[ \omega^2 \left( F(j\omega) + F(-j\omega)^\top \right) \right] \ge \sigma_0 \quad \forall |\omega| \ge \delta.$$
 (2)

This note is devoted to the analysis of the two side-conditions (1) and (2). We will prove that while they are in general not equivalent at infinite frequency, they are indeed equivalent under the other conditions guaranteeing that F is strictly positive real. Hence both conditions at infinite frequency are equally valid. While this could be deduced from [3] and [5], our results provide a direct proof of such equivalency.

**Notation:** Let the set of real (resp. complex) numbers be denoted by  $\mathbb{R}$  (resp.  $\mathbb{C}$ ) and the corresponding sets of matrices

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of dimension  $m \times n$  be denoted by  $\mathbb{R}^{m \times n}$  (resp.  $\mathbb{C}^{m \times n}$ ). Given  $M \in \mathbb{C}^{m \times m}$ ,  $M^*$  denotes the complex conjugate transpose of matrix M (i.e., if M = A + jB for real matrices A and B, then  $M^* = A^\top - jB^\top$ ). A matrix M is said to be Hermitian if  $M = M^*$  and M > 0 denotes that the matrix M is Hermitian and positive definite. The smallest singular value of M is denoted by  $\underline{\sigma}(M)$ . We recall that the singular values of a positive semi-definite Hermitian matrix are its nonzero eigenvalues [2, p.649].

#### II. MAIN RESULT

The following definition, adapted from [4, Definitions 1 and 2], is the standard definition for strictly positive real systems. It essentially states that a transfer function matrix F(s) is strictly positive real if for some  $\epsilon > 0$ , the transfer function matrix  $F(s-\epsilon)$  is positive real and  $F(s) + F(-s)^{\top}$  has full normal rank. See also [4, Lemma 2] for an equivalent re-characterization.

Definition 2.1: Let  $F:\mathbb{C} \longrightarrow \mathbb{C}^{m \times m}$  be a real transfer function. Then F(s) is said to be Strictly Positive Real (SPR) if there exists a real scalar  $\epsilon>0$  such that F(s) is analytic in  $\{s\in\mathbb{C}:\mathfrak{Re}\{s\}>-\epsilon\}$ ,  $F(s)+F(s)^{\star}\geq0$  for all  $s\in\{s\in\mathbb{C}:\mathfrak{Re}\{s\}>-\epsilon\}$  and  $F(s)+F(-s)^{\top}$  has full normal rank. SPR matrix transfer functions can be characterized in the frequency domain by three conditions: the first two are conditions 1 and 2 in the next proposition, the third is the *side condition* and it has been stated in two different manners: side condition (3a) in Proposition 2.1 can be found in [5], [4], while side condition (3b) can be found in [3]. These side conditions can be interpreted as apparently different conditions on how  $(F(j\omega)+F(-j\omega)^{\top})$  approaches zero for sufficiently large  $|\omega|$  in directions where it looses rank.

*Proposition 2.1:* Let  $F: \mathbb{C} \longrightarrow \mathbb{C}^{m \times m}$  be a real, rational, proper transfer function such that the following two conditions hold:

- 1) F(s) has no poles in  $\{s \in \mathbb{C} : \Re \{s\} \ge 0\}$ ;
- 2)  $F(j\omega) + F(-j\omega)^{\top} > 0$  for all  $\omega \in \mathbb{R}$ .

Then, the following two side conditions are equivalent:

3a)  $\exists \delta > 0, \sigma_0 > 0$ :

$$\underline{\sigma} \left[ \omega^2 \left( F(j\omega) + F(-j\omega)^\top \right) \right] \ge \sigma_0 \quad \forall |\omega| \ge \delta \quad (3)$$

3b) 
$$\lim_{\omega \to \infty} \omega^{2\rho} \det[F(j\omega) + F(-j\omega)^{\top}] \neq 0$$
 (4)

where  $\rho = \dim \ker(F(\infty) + F(\infty)^{\top})$ .

*Proof.* Since F(s) is proper, let  $F(s) = C(sI - A)^{-1}B + D$  for some state-space realization (A, B, C, D). Near  $s = \infty$ , the expansion

$$F(s) = F_0 + F_1/s + F_2/s^2 + \dots$$
 (5)

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holds (with  $F_0 = F(\infty) = D$ ), so that

$$F(j\omega) + F(-j\omega)^{\top} = Q + \frac{H}{\omega} + \frac{K(\omega)}{\omega^2}$$
 (6)

with  $Q = D + D^{\top}$ ,  $H = (F_1 - F_1^{\top})/j$  and  $K(\omega) = \sum_{i=1}^{\infty} \left(\frac{F_{i+1}}{j^{i+1}\omega^{i-1}} + \frac{F_{i+1}^{T}}{(-j)^{i+1}\omega^{i-1}}\right)$ . Then it is easy to see that

$$\lim_{\omega \to \infty} K(\omega) = K_0 \in \mathbb{R}^{m \times m},$$

with  $K_0 := -F_2 - F_2^T$ . If Q is positive definite, in view of the Hermitian symmetry of both H and  $K(\omega)$ , both conditions 3a) and 3b) are clearly satisfied and the result is obvious.

Assume than that Q is singular with rank  $m-\rho$ . Since Q is symmetric, there exists an orthogonal matrix U such that  $UQU^\top=\begin{bmatrix}Q_1&0\\0&0\end{bmatrix}$  with  $Q_1\in\mathbb{R}^{(m-\rho)\times(m-\rho)}$  nonsingular. Hence, without changing the essence of the problem, we assume that Q has the form

$$Q = \left[ \begin{array}{cc} Q_1 & 0 \\ 0 & 0 \end{array} \right] \tag{7}$$

with  $Q_1 \in \mathbb{R}^{(m-\rho)\times (m-\rho)}$  nonsingular, and hence  $Q_1 > 0$ . We partition H and  $K(\omega)$  accordingly, as follows:

$$H = \begin{bmatrix} H_1 & H_{12} \\ H_{21} & H_2 \end{bmatrix}, \qquad K(\omega) = \begin{bmatrix} K_1(\omega) & K_{12}(\omega) \\ K_{21}(\omega) & K_2(\omega) \end{bmatrix}. \tag{8}$$

Thus

$$[F(j\omega) + F(-j\omega)^{\top}]_{22} = \frac{1}{\omega} (H_2 + \frac{K_2(\omega)}{\omega}).$$
 (9)

Since  $H=(F_1-F_1^\top)/j$ , we have that: (i) H is Hermitian, so that  $H_2$  is Hermitian as well, and (ii) the elements on the diagonal of H are zero, so that also the elements on the diagonal of  $H_2$  are zero and hence  $H_2$  is traceless. Therefore,  $H_2$  has only real eigenvalues and the sum of the eigenvalues of  $H_2$  is zero. Hence, either  $H_2=0$  or  $H_2$  has at least a negative eigenvalue. But for  $\omega$  sufficiently large, and by continuity of the eigenvalues as functions of  $\omega$ , the eigenvalues of  $H_2+\frac{K_2(\omega)}{\omega}$  are arbitrarily close to those of  $H_2$ , so that if  $H_2\neq 0$ , then  $[F(j\omega)+F(-j\omega)^\top]_{22}$  has a negative eigenvalue for a sufficiently large  $\omega$ , and this is against our assumptions because  $[F(j\omega)+F(-j\omega)^\top]>0$  for all  $\omega\in\mathbb{R}$  so, in turn,  $[F(j\omega)+F(-j\omega)^\top]_{22}>0$  for all  $\omega\in\mathbb{R}$ . In conclusion,  $H_2=0$ .

Now let

$$\Phi(\omega) = \begin{bmatrix} \Phi_1(\omega) & \Phi_{12}(\omega) \\ \Phi_{21}(\omega) & \Phi_2(\omega) \end{bmatrix} = F(j\omega) + F(-j\omega)^{\top}.$$

By continuity, as  $\omega \to \infty$ ,  $m-\rho$  of the eigenvalues of  $\Phi$ , *i.e.*, the eigenvalues of  $\Phi_1(j\omega)$ , tend to the eigenvalues of  $Q_1$  (that are strictly positive) and the remaining  $\rho$  eigenvalues tend to zero. Let  $\lambda(\omega)$  be one of the eigenvalues of  $\Phi$  that tends to zero as  $\omega \to \infty$ . We now show that  $\lambda(\omega)$  tends to zero at least as fast as  $1/\omega^2$ . In fact, provided that  $\omega$  is large enough so that  $\lambda(\omega)$  is not an eigenvalue of  $\Phi_1(\omega)$ , then  $\lambda(\omega)$  must be an eigenvalue of

$$R(\omega) = \Phi_2(\omega) - \Phi_{21}(\omega) [\Phi_1(\omega) - \lambda(\omega)I]^{-1} \Phi_{12}(\omega) \quad (10)$$

because through Schur complements (e.g. [1])

$$\det[\Phi(\omega) - \lambda(\omega)I] = \det[\Phi_1(\omega) - \lambda(\omega)I] \det[R(\omega) - \lambda(\omega)I].$$

Using the fact that  $\Phi_1(\omega) - \lambda(\omega)I = Q_1 + \frac{H_1}{\omega} + \frac{K_1(\omega)}{\omega^2} - \lambda(\omega)I$ , and using [1, Fact 9.9.43],  $R(\omega)$  in (10) can be written equivalently as

$$R(\omega) = \Phi_2(\omega) - \Phi_{21}(\omega)[Q_1^{-1} + \Delta(\omega)]\Phi_{12}(\omega),$$

where  $\Delta(\omega) \to 0$  as  $\omega \to \infty$ , since  $\lambda(\omega) \to 0$  as  $\omega \to \infty$ . Finally, using (6) (7) and (8), we have  $R(\omega) = \frac{1}{\omega^2} \left[ K_2(\omega) - P(\omega) \right]$  where

$$P(\omega) := (H_{21} + \frac{K_{21}(\omega)}{\omega})[Q_1^{-1} + \Delta(\omega)](H_{12} + \frac{K_{12}(\omega)}{\omega}).$$

Since for  $\omega$  sufficiently large  $[K_2(\omega) - P(\omega)]$  is bounded, then the eigenvalues of  $R(\omega)$  tend to zero at least as fast as  $1/\omega^2$ .

Then  $\Phi(\omega)$  has m strictly positive eigenvalues for all  $\omega \in \mathbb{R}$ ,  $\rho$  of which are going down to zero at least as fast as  $1/\omega^2$  as  $\omega \to \infty$ . The remaining  $(m-\rho)$  eigenvalues tend to the strictly positive eigenvalues of  $Q_1$  as  $\omega \to \infty$ . Now, order the eigenvalues of  $\Phi(\omega)$  in non-decreasing size.

Since  $\det(\cdot)$  is the product of eigenvalues, then  $\omega^{2\rho} \det[\Phi(\omega)]$  is the same as the product of  $(\omega^2 \lambda_1), \ldots, (\omega^2 \lambda_\rho), \lambda_{\rho+1}, \ldots, \lambda_m$  because for each  $i \in \{1, \ldots, \rho\}, \lambda_i$  tends to zero at least as fast as  $1/\omega^2$ . Then, we have that the side condition (3b):

$$\lim_{\omega \to \infty} w^{2\rho} \det[\Phi(\omega)] \neq 0$$

is equivalent to (since  $\lambda_{\rho+1},...,\lambda_m$  do not tend to zero)

$$\lim_{w\to\infty} (w^2\lambda_i) \neq 0 \text{ for all } i \in \{1,...,\rho\},\$$

which is equivalent to

$$\lim_{\omega \to \infty} w^2 \lambda_1 > 0,$$

which, finally, is equivalent to the first side condition (3a). The proposition is proved.  $\hfill\Box$ 

*Examples.* Consider the case of  $F(s) = \frac{1}{s+1} \begin{bmatrix} 1 & 0 \\ 0 & s+2 \end{bmatrix}$ . In this case  $\rho = 1$  and, by direct computation, we see that

$$\lim_{\omega \to \infty} \omega^{2\rho} \det[F(j\omega) + F(-j\omega)^{\top}] = \lim_{\omega \to \infty} \frac{8\omega^2 + 4\omega^4}{(1 + \omega^2)^2} = 4 \neq 0.$$

On the other hand, we have

$$\underline{\sigma} \left[ \omega^2 \left( F(j\omega) + F(-j\omega)^\top \right) \right] = \frac{2\omega^2}{1 + \omega^2}$$

which is clearly greater than 1 for all  $|\omega| > 1$ .

The pathological case, corresponding to the situation in which some of the eigenvalues of the spectrum go to zero faster than  $\frac{1}{\omega^2}$ , as  $\omega$  tends to infinity, is more interesting: let

$$F(s) = \begin{bmatrix} \frac{s+2}{(s+1)^2} & 0\\ 0 & \frac{s+2}{(s+1)} \end{bmatrix}.$$
 In this case  $\rho=1$  and by direct computation, we easily see that

$$\lim_{\omega \to \infty} \omega^{2\rho} \det[F(j\omega) + F(-j\omega)^{\top}] = \lim_{\omega \to \infty} \frac{\omega^2 (16 + 8\omega^2)}{(1 + \omega^2)^3} = 0.$$

The same conclusion is obtained with the other side condition which, in this case, reads:

$$\underline{\sigma} \left[ \omega^2 \left( F(j\omega) + F(-j\omega)^\top \right) \right] = \frac{4\omega^2}{(1+\omega^2)^2}$$

which is clearly not bounded away from zero as  $|\omega|$  diverges. Finally, the trivial example  $F(s)=\begin{bmatrix}1&0\\0&-1\end{bmatrix}$  immediately shows that, if we do not assume the first two conditions of Proposition 2.1, the two side conditions are not necessarily equivalent.

#### III. CONCLUSION

This note shows that two different side conditions used in the control literature to characterize strictly positive real matrix transfer functions are equivalent.

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