





<sup>1</sup> Magique-3D, Inria | <sup>2</sup> TOTAL S.A.

# POLYNOMIAL SPACE-TIME TREFFTZ-DG METHOD FOR ELASTO-ACOUSTICS

MATHIAS'18 | Computational Science Engineering & Data Science by TOTAL

Serris - France | 2018

# CONTEXT

## TREFFTZ-DG METHODS

#### **BIBLIOGRAPHY**

## Frequency domain:

Farhat, Tezaur, Harari, Hetmaniuk (2003 - 2006), Gabard (2007), Badics (2014), Hiptmair, Moiola, Pirugia (2011 - 2016), Barucq, Bendali, Diaz, Tordeux, Fares, Mattesi (2017)

#### Time domain:

Tsukerman, Egger, Kretzschmar, Schnepp, Weiland (2014 - 2015), Wang, Tezaur, Farhat (2014), Moiola, Pirugia (2016 - 2018), Banjai, Lijoka (2017)

## DISCONTINUOUS GALERKIN METHODS

#### PROS AND CONS

Adapted to the complex geometries High-order accuracy and hp-adaptivity Explicit semi-dicrete form Conservation laws

Higher number of degrees of freedom, compared to the methods with continuous approximation

## TREFFTZ METHOD

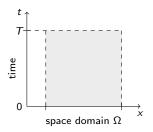
E. TREFFTZ, 1926

Basis functions are local solutions of the initial PDEs

# MATHEMATICAL FORMULATION

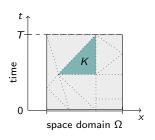
ACOUSTIC SYSTEM

#### DOMAIN DESCRIPTION



$$\begin{cases} \frac{1}{c^2 \rho} \frac{\partial \mathbf{p}}{\partial t} + \operatorname{div} \mathbf{v} = f & \text{in } \Omega \times [0, T], \\ \rho \frac{\partial \mathbf{v}}{\partial t} + \nabla \mathbf{p} = 0 & \text{in } \Omega \times [0, T], \\ \mathbf{v}(\cdot, 0) = \mathbf{v}_0, \ \mathbf{p}(\cdot, 0) = \mathbf{p}_0 & \text{in } \Omega, \\ \mathbf{v} \cdot \mathbf{n}^{\times} = \mathbf{g}_D & \text{in } \partial \Omega \times [0, T]. \end{cases}$$

#### SPACE-TIME DISCRETZATION



$$\mathcal{F}_h = \bigcup_{K \in \mathcal{T}_h} \partial K$$
 - mesh skeleton

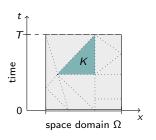
 $\mathcal{F}^I$  - internal

-  $\mathcal{F}^D$  - boundary

 $\longrightarrow$   $\mathcal{F}^0$  - initial time

--  $\mathcal{F}^{T}$  - final time

#### SPACE-TIME DISCRETZATION



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 - mesh skeleton

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-  $\mathcal{F}^D$  - boundary

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--  $\mathcal{F}^T$  - final time

$$\begin{aligned} & (\mathbf{v}_h, \mathbf{p}_h) \in V^h(\mathcal{T}_h)^d \times V^h(\mathcal{T}_h) \equiv \mathbf{V}^h(\mathcal{T}_h) \\ & (\omega, q) \in V^h(\mathcal{T}_h)^d \times V^h(\mathcal{T}_h) \equiv \mathbf{V}^h(\mathcal{T}_h) \end{aligned}$$

#### SPACE-TIME DISCRETIZATION

$$\begin{split} &-\sum_{K}\int_{K} \mathbf{p_h} \Big(\frac{1}{c^2\rho}\frac{\partial \mathbf{q}}{\partial t} + \mathrm{div}\omega\Big) + \mathbf{v_h} \cdot \Big(\rho\frac{\partial\omega}{\partial t} + \nabla q\Big) \\ &+\sum_{K}\int_{\partial K} \Big(\frac{1}{c^2\rho} \check{\mathbf{p}_h} \mathbf{q} + \check{\mathbf{v}_h} \cdot \omega\Big) n_K^t + \Big(q\hat{\mathbf{v}_h} + \hat{\mathbf{v}_h}\omega\Big) \cdot n_K^{\mathsf{x}} = \sum_{K}\int_{K} fq \end{split}$$

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### TREFFTZ SPACE

$$\mathbf{T}(\mathcal{T}_h) \equiv \left\{ (\omega, q) \in \mathbf{V}^h(\mathcal{T}_h) \text{ s.t. } \frac{1}{c^2 \rho} \frac{\partial q}{\partial t} + \mathrm{div}\omega \, = \, \rho \frac{\partial \omega}{\partial t} + \nabla q = 0, \ \, \forall K \in \mathcal{T}_h \right\}$$

#### SPACE-TIME DISCRETIZATION

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#### SPACE-TIME DISCRETIZATION IN T

$$\sum_{K}\int\limits_{\partial K} \left(\frac{1}{c^{2}\rho} \breve{\mathbf{p}}_{\pmb{h}}^{} q + \breve{\mathbf{v}}_{\pmb{h}} \cdot \omega\right) n_{K}^{t} + \left(q \hat{\mathbf{v}}_{\pmb{h}} + \hat{\mathbf{v}}_{\pmb{h}} \omega\right) \cdot n_{K}^{\mathsf{x}} = \sum_{K}\int\limits_{K} \mathit{f} q$$

#### TREFFTZ-DG FORMULATION

$$\sum_{K}\int\limits_{\partial K} \left(\frac{1}{c^{2}\rho} \check{\mathbf{p}}_{\pmb{h}}^{} q + \check{\mathbf{v}}_{\pmb{h}} \cdot \omega\right) n_{K}^{t} + \left(q \hat{\mathbf{v}}_{\pmb{h}}^{} + \hat{\mathbf{v}}_{\pmb{h}}^{} \omega\right) \cdot n_{K}^{\mathsf{x}} = \sum_{K}\int\limits_{K} fq$$

#### NUMERICAL FLUXES THROUGH THE FLEMENT FACES

## TREFFTZ-DG FORMULATION

Seek  $(v_h, p_h) \in \mathbf{V}^h(\mathcal{T}_h)$  s.t.  $\forall (\omega, q) \in \mathbf{T}(\mathcal{T}_h)$  it holds true:

$$\mathcal{A}_{TDG}ig((\mathbf{v_h},\mathbf{p_h});(\omega,q)ig)=\ell_{TDG}ig(\omega,qig)$$

#### WELL POSEDNESS

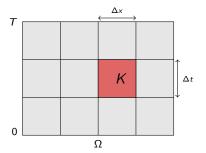
Trefftz-DG formulations for the first order Acoustic System <sup>1,2</sup>, Elastodynamic System <sup>2</sup>, and Elasto-Acoustic System <sup>2</sup> are well-posed.

A space-time Trefftz- dG method for the acoustic wave equation in first-order formulation. (2018) Numerische Mathematik, 138(2):389-435

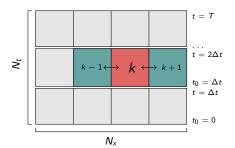
<sup>&</sup>lt;sup>1</sup> A. Moiola & I. Perugia.

<sup>&</sup>lt;sup>2</sup> H. Barucq, H. Calandra, J. Diaz & Elvira Shishenina.
Space-Time Trefftz-Discontinuous Galerkin Approximation for Elasto-Acoustics. (2017) https://hal.inria.fr/hal-01614126/document

### MESH AND ELEMENT NUMBERING



### MESH AND ELEMENT NUMBERING



#### **BILINEAR FORM**

$$\mathcal{A}_{TDG} = \underbrace{\int_{\mathcal{F}^T} + \int_{\mathcal{F}^0} + \int_{\mathcal{F}^I} + \int_{\mathcal{F}^D}}_{\mathcal{A}_{TDG}^I}$$

GLOBAL SPACE-TIME MATRIX

$$\mathbf{M} = \bar{\Delta}_{x} \mathbf{M}_{\Omega} + \bar{\Delta}_{t} M_{I}$$

#### **BILINEAR FORM**

$$\mathcal{A}_{TDG} = \underbrace{\int_{\mathcal{F}^{T}} + \int_{\mathcal{F}^{0}} + \int_{\mathcal{F}^{I}} + \int_{\mathcal{F}^{D}}}_{\mathcal{A}_{TDG}^{I}}$$

#### GLOBAL SPACE-TIME MATRIX

$$\mathbf{M} = \bar{\Delta}_{\mathsf{X}} \mathbf{M}_{\mathsf{\Omega}} + \bar{\Delta}_{\mathsf{t}} M_{\mathsf{I}}$$

 $M_{\rm O}$  - block-diagonal

#### **BILINEAR FORM**

$$A_{TDG} = \int_{\mathcal{F}^T} + \int_{\mathcal{F}^0} + \int_{\mathcal{F}^I} + \int_{\mathcal{F}^D}$$

$$A_{TDG}^{\Omega} = \int_{\mathcal{A}_{TDG}^{I}} + \int_{\mathcal{A}_{TDG}^{$$

#### GLOBAL SPACE-TIME MATRIX

$$\mathbf{M} = \bar{\Delta}_{\mathsf{X}} \mathbf{M}_{\mathsf{\Omega}} + \bar{\Delta}_{\mathsf{t}} M_{\mathsf{I}}$$

 $M_{\Omega}$  - block-diagonal  $M_{I}$  - sparse

#### GLOBAL MATRIX INVERSION

$$\mathbf{M}^{-1} \equiv \left[ \bar{\Delta}_{x} \mathbf{M}_{\Omega} + \bar{\Delta}_{t} \mathbf{M}_{I} \right]^{-1}$$

$$=$$

$$\left[ I + \frac{\bar{\Delta}_{t}}{\bar{\Delta}_{x}} \left( \mathbf{M}_{\Omega}^{-1} \mathbf{M}_{I} \right) \right]^{-1} \left[ \bar{\Delta}_{x} \mathbf{M}_{\Omega} \right]^{-1}$$

TAYLOR EXPANSION (SMALL ENOUGH  $\frac{\bar{\Delta}_t}{\bar{\Delta}_r}|M_{\Omega}^{-1}M_I|\propto \frac{\Delta_t}{\Delta_x}|M_{\Omega}^{-1}M_I|$ )

$$\mathbf{M}^{-1} \approx \Big[\sum_{i=0}^{n} (-1)^{i} \Big(\frac{\bar{\Delta}_{t}}{\bar{\Delta}_{x}}\Big)^{i} (\mathbf{M}_{\Omega}^{-1} \mathbf{M}_{I})^{i} \Big] \Big[\bar{\Delta}_{x} \mathbf{M}_{\Omega}\Big]^{-1}$$

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TAYLOR EXPANSION (SMALL ENOUGH  $\frac{\bar{\Delta}_t}{\bar{\Lambda}}|M_\Omega^{-1}M_I|\propto \frac{\Delta_t}{\bar{\Lambda}_X}|M_\Omega^{-1}M_I|)$ 

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block-diagonal

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$$=$$

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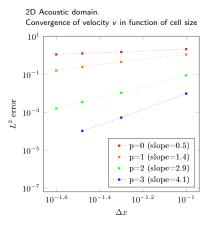
TAYLOR EXPANSION (SMALL ENOUGH  $\frac{\bar{\Delta}_t}{\bar{\Delta}_r}|M_{\Omega}^{-1}M_I|\propto \frac{\Delta_t}{\Delta_x}|M_{\Omega}^{-1}M_I|$ )

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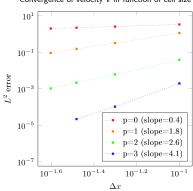
block-diagonal × sparse

# NUMERICAL RESULTS

## 2D ACOUSTIC AND ELASTODYNAMIC SYSTEMS



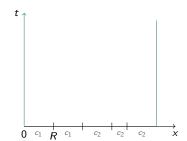
# 2D Elastodynamic domain. Convergence of velocity *v* in function of cell size



APPLICATION TO TREFFTZ-DG METHOD

### 1D WAVE EQUATION

$$\begin{cases} \frac{1}{c^2} \frac{\partial^2 \mathbf{p}}{\partial t^2} - \Delta \mathbf{p} = 0, \\ \\ \mathbf{p}(\cdot, 0) = \mathbf{p}_0, \\ \frac{\partial \mathbf{p}}{\partial t}(\cdot, 0) = \mathbf{p}_1. \end{cases}$$

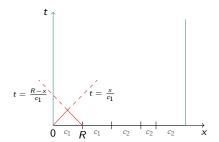


#### SOLUTION

$$p(x,t) = \frac{1}{2} \left( p_0(x - ct) + p_0(x + ct) \right) + \frac{1}{2c} \left( p_1(x - ct) + p_1(x + ct) \right)$$

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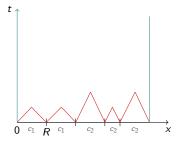


#### **SOLUTION**

$$p(x,t) = \frac{1}{2} \left( p_0(x - ct) + p_0(x + ct) \right) + \frac{1}{2c} \left( p_1(x - ct) + p_1(x + ct) \right)$$

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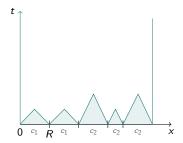


### **SOLUTION**

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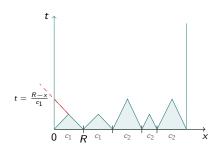


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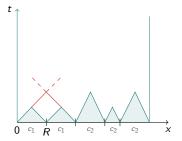


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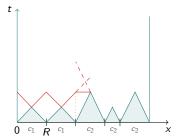


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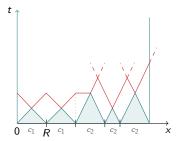


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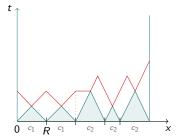


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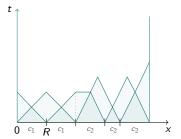


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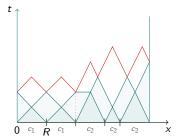


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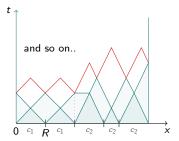


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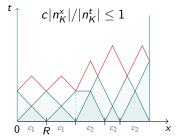


### **SOLUTION**

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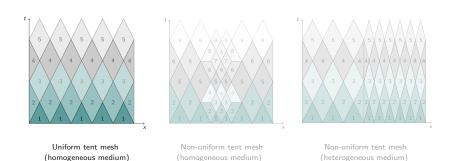
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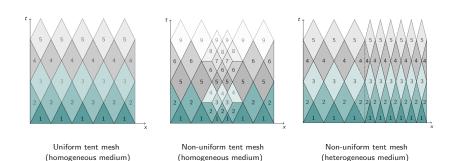
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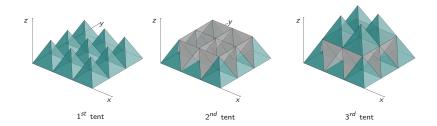
#### 1D TENT PITCHING MESHES



#### 1D TENT PITCHING MESHES



#### 2D TENT PITCHING MESHES



ANALYTICAL STABILITY CONDITIONS 1

$$c|n_K^{\mathsf{x}}|/|n_K^t| \leq 1$$

- TDG acoustic formulation

$$\sqrt{|\mathbf{C}|} |n_K^{\mathsf{x}}| / \sqrt{\rho} |n_K^t| \leq 1$$

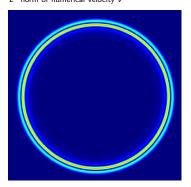
- TDG elastodynamic formulation

<sup>&</sup>lt;sup>1</sup> H. Barucq, H. Calandra, J. Diaz & Elvira Shishenina.
Space-Time Discretization of Elasto-Acoustic Wave Equation in Polynomial Trefftz-DG Bases. (2018)
PhD thesis [under review]

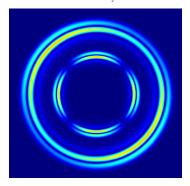
#### PROPAGATION OF NUMERICAL VELOCITY

2D Acoustic domain.

L<sup>2</sup> norm of numerical velocity v

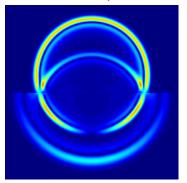


2D Elastodynamic domain.  $L^2$  norm of numerical velocity v



### PROPAGATION OF NUMERICAL VELOCITY

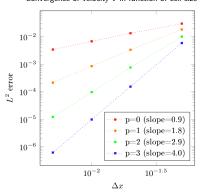
2D Elasto-Acoustic domain.  $L^2$  norm of numerical velocity v



### CONVERGENCE CURVES

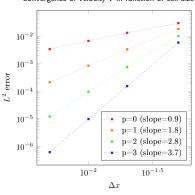
2D Acoustic domain.

Convergence of velocity v in function of cell size

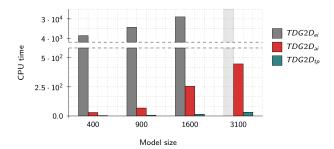


2D Elastodynamic domain.

Convergence of velocity *v* in function of cell size



### SPEEDUP FOR DIFFERENT MODEL SIZES



# CONCLUSION

### CONCLUSION

#### ON-GOING WORK AND PERSPECTIVES

#### Standard Trefftz-DG method:

- Development of 2D-3D hp propagators
- Analytical study of convergence, comparison with numerical results
- Study of CFL condition, comparison with standard DG method

### Tent pitching Trefftz-DG method:

- Mesh generation study
- Implementation of the boundary conditions
- 3D case

# THANK YOU!

### IMPLEMENTATION OF THE ALGORITHM

### POLYNOMIAL BASIS (ORDER 3)

$$\begin{array}{llll} \mathbf{p=0, \, ndof=2} & \phi_1^{\rm v} = 0 & \phi_1^{\rm p} = -c_F \\ & \phi_2^{\rm v} = 1 & \phi_2^{\rm p} = 0 \\ \end{array}$$
 
$$\begin{array}{lll} \mathbf{p=1, \, ndof=4} & \phi_3^{\rm v} = x & \phi_3^{\rm p} = -c_F^2 t \\ & \phi_4^{\rm v} = c_F t & \phi_4^{\rm p} = -c_F x \\ \end{array}$$
 
$$\mathbf{p=2, \, ndof=6} & \phi_5^{\rm v} = -\frac{x^2}{2} - \frac{c_F^2 t^2}{2} & \phi_5^{\rm p} = c_F^2 x t \\ & \phi_5^{\rm v} = -c_F x t & \phi_6^{\rm p} = c_F (\frac{x^2}{2} + \frac{c_F^2 t^2}{2}) \\ \end{array}$$
 
$$\mathbf{p=3, \, ndof=8} & \phi_7^{\rm v} = -\frac{x^3}{6} - \frac{x c_F^2 t^2}{2} & \phi_7^{\rm p} = c_F (\frac{c_F^3 t^3}{6} + \frac{x^2 c_F t}{2}) \\ & \phi_8^{\rm w} = -\frac{c_F^3 t^3}{6} - \frac{x^2 c_F t}{2} & \phi_8^{\rm p} = c_F (\frac{x^3}{6} + \frac{x^2 c_F t}{2}) \end{array}$$

### ACOUSTIC SYSTEM

#### TREFFTZ-DG FORMULATION

$$\sum_{K}\int\limits_{\partial K}\left(\frac{1}{c^{2}\rho}\breve{\mathbf{p}}_{\pmb{h}}^{\pmb{h}}q+\breve{\mathbf{v}}_{\pmb{h}}\cdot\boldsymbol{\omega}\right)n_{K}^{t}+\left(q\hat{\mathbf{v}}_{\pmb{h}}+\hat{\mathbf{v}}_{\pmb{h}}\boldsymbol{\omega}\right)\cdot n_{K}^{\mathbf{x}}=\sum_{K}\int\limits_{K}fq$$

#### NUMERICAL FLUXES THROUGH THE FLEMENT FACES

$$\begin{pmatrix} \hat{v}_h \\ \hat{\rho}_h \end{pmatrix} \qquad \equiv \qquad \begin{pmatrix} \{v_h\} + \beta[\rho_h]_X \\ \{\rho_h\} + \alpha[v_h]_X \end{pmatrix} \qquad \text{on } \mathcal{F}^I \qquad \cdots$$

$$\begin{pmatrix} \hat{v}_h \\ \hat{\rho}_h \end{pmatrix} \qquad \equiv \qquad \begin{pmatrix} \{v_h\} \\ \{\rho_h\} \end{pmatrix} \qquad \text{on } \mathcal{F}^I \qquad \cdots$$

$$\begin{pmatrix} \hat{v}_h \cdot n_K^x \\ \hat{\rho}_h \end{pmatrix} \qquad \equiv \qquad \begin{pmatrix} g_D \\ \rho_h + \alpha(v_h \cdot n_K^x - g_D) \end{pmatrix} \qquad \text{on } \mathcal{F}^D \qquad \cdots$$

$$\begin{pmatrix} \hat{v}_h \\ \hat{\rho}_h \end{pmatrix} \qquad \equiv \qquad \begin{pmatrix} v_h \\ \rho_h \end{pmatrix} \qquad \text{on } \mathcal{F}^T \qquad - \cdots$$

$$\begin{pmatrix} \hat{v}_h \\ \hat{\rho}_h \end{pmatrix} \qquad \equiv \qquad \begin{pmatrix} v_0 \\ e_D \end{pmatrix} \qquad \text{on } \mathcal{F}^D \qquad \cdots$$

$$\begin{pmatrix} \hat{v}_h \\ \hat{\sigma}_h \end{pmatrix} \qquad \equiv \qquad \begin{pmatrix} v_0 \\ e_D \end{pmatrix} \qquad \text{on } \mathcal{F}^D \qquad \cdots$$

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