

Hybridizable Discontinuous Galerkin Methods for modelling 3D seismic wave propagation in harmonic domain

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Hybridizable Discontinuous Galerkin Methods for modelling 3D seismic wave propagation in harmonic domain

 $M.$ Bonnasse-Gahot 1,2 , H. Calandra³, J. Diaz¹ and S. Lanteri²

 $^{\rm 1}$ INRIA Bordeaux-Sud-Ouest, team-project Magique 3D

³ TOTAL Exploration-Production

Principles of seismic imaging

Examples of seismic imaging campaigns

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Imaging methods

Full Wave Inversion (FWI) : inversion process requiring to solve many forward problems

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Imaging methods

Full Wave Inversion (FWI) : inversion process requiring to solve many forward problems

Seismic imaging : time-domain or harmonic-domain ?

- **Time-domain** : imaging condition complicated but quite low computational cost
- **Harmonic-domain** : imaging condition simple but huge computational cost

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Imaging methods

Full Wave Inversion (FWI) : inversion process requiring to solve many forward problems

Seismic imaging : time-domain or harmonic-domain ?

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Resolution of the forward problem of the inversion process

Elastic wave propagation in the frequency domain : $Helmholtz$ equation

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Resolution of the forward problem of the inversion process

Elastic wave propagation in the frequency domain : Helmholtz equation

First order formulation of Helmholtz wave equations

$$
\mathbf{x} = (x, y, z) \in \Omega \subset \mathbb{R}^3,
$$

$$
\begin{cases} i\omega \rho(\mathbf{x})\mathbf{v}(\mathbf{x}) = \nabla \cdot \underline{\sigma}(\mathbf{x}) \\ i\omega \underline{\sigma}(\mathbf{x}) = \underline{\underline{C}}(\mathbf{x}) \underline{\varepsilon}(\mathbf{v}(\mathbf{x})) + f_s(\mathbf{x})) \end{cases}
$$

- \triangleright **v** : velocity vector
- \triangleright σ : stress tensor
- \blacktriangleright $\underline{\varepsilon}$: strain tensor

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Discontinuous Galerkin Methods

- ✓ unstructured tetrahedral meshes
- ✓ combination between FEM and finite volume method (FVM)
- $\sqrt{h}p$ -adaptivity
- \checkmark easily parallelizable method

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- ✗ ✗ large number of DOF as compared to classical FEM

Discontinuous Galerkin Methods

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Hybridizable Discontinuous Galerkin Methods

✓ same advantages as DG methods : unstructured tetrahedral meshes, hp-adaptivity, easily parallelizable method, discontinuous basis functions

✓ introduction of a new variable defined only on the interfaces ✓ lower number of coupled DOF than classical DG methods

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Hybridizable Discontinuous Galerkin Methods

- 螶 B. Cockburn, J. Gopalakrishnan and R. Lazarov. Unified hybridization of discontinuous Galerkin, mixed and continuous Galerkin methods for second order elliptic problems. SIAM Journal on Numerical Analysis, Vol. 47 :1319-1365, 2009.
- **S.** Lanteri, L. Li and R. Perrussel. Numerical investigation of a high order hybridizable discontinuous Galerkin method for 2d time-harmonic Maxwell's equations. COMPEL, 32(3)1112-1138, 2013.
- 譶 N.C. Nguyen, J. Peraire and B. Cockburn. High-order implicit hybridizable discontinuous Galerkin methods for acoustics and elastodynamics. Journal of Computational Physics, 230 :7151-7175, 2011
	- N.C. Nguyen and B. Cockburn. Hybridizable discontinuous Galerkin methods for partial differential equations in continuum mechanics. Journal of Computational Physics 231 :5955–5988, 2012

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1. Introduction of a Lagrange multiplier λ

1. Introduction of a Lagrange multiplier λ

 \triangleright Weak formulation of waves equations

$$
\left\{\begin{array}{l} \displaystyle\int_{K} i\omega \rho^K \mathbf{v}^K \cdot \mathbf{w} + \int_{K} \underline{\underline{\sigma}}^K : \nabla \mathbf{w} - \int_{\partial K} \underline{\underline{\sigma}} \cdot \mathbf{n} \cdot \mathbf{w} = 0 \\ \displaystyle\int_{K} i\omega \underline{\underline{\sigma}}^K : \underline{\xi} + \int_{K} \mathbf{v}^K \cdot \nabla \cdot \left(\underline{\underline{C}}^K \underline{\xi}\right) - \int_{\partial K} \mathbf{v} \cdot \underline{\underline{C}}^K \underline{\xi} \cdot \mathbf{n} = \int_{K} f_S^K \cdot \underline{\xi} \end{array}\right.
$$

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- 1. Introduction of a Lagrange multiplier λ
	- \triangleright Weak formulation of waves equations

$$
\begin{cases}\n\int_{K} i\omega \rho^{K} \mathbf{v}^{K} \cdot \mathbf{w} + \int_{K} \underline{\underline{\sigma}}^{K} : \nabla \mathbf{w} - \int_{\partial K} \underline{\widehat{\underline{\sigma}}}^{\partial K} \cdot \mathbf{n} \cdot \mathbf{w} = 0 \\
\int_{K} i\omega \underline{\underline{\sigma}}^{K} : \underline{\xi} + \int_{K} \mathbf{v}^{K} \cdot \nabla \cdot \left(\underline{\underline{C}}^{K} \underline{\xi} \right) - \int_{\partial K} \widehat{\mathbf{v}}^{\partial K} \cdot \underline{\underline{C}}^{K} \underline{\xi} \cdot \mathbf{n} = \int_{K} f_{S}^{K} \cdot \underline{\xi} \\
\frac{\widehat{\underline{\sigma}}^{K}}{\partial K} \text{ and } \widehat{\mathbf{v}}^{K} \text{ are numerical traces of } \underline{\underline{\sigma}}^{K} \text{ and } \mathbf{v}^{K} \text{ respectively on} \\
\end{cases}
$$

- 1. Introduction of a Lagrange multiplier λ
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$$
\left\{\begin{array}{l} \displaystyle\int_K i\omega \rho^K \mathbf{v}^K\cdot \mathbf{w} + \int_K \underline{\underline{\underline{\sigma}}}^K : \nabla \mathbf{w} - \int_{\partial K} \underline{\underline{\widehat{\underline{\sigma}}}^{\partial K}} \cdot \mathbf{n} \cdot \mathbf{w} = 0 \\ \displaystyle\int_K i\omega \underline{\underline{\underline{\sigma}}}^K : \underline{\underline{\xi}} + \int_K \mathbf{v}^K \cdot \nabla \cdot \left(\underline{\underline{\underline{C}}}^K \underline{\underline{\xi}}\right) - \int_{\partial K} \widehat{\mathbf{v}}^{\partial K} \cdot \underline{\underline{\underline{C}}}^K \underline{\underline{\xi}} \cdot \mathbf{n} = \int_K f_S^K \cdot \underline{\underline{\xi}} \end{array}\right.
$$

$$
\widehat{\mathbf{v}}^{\partial K} = \lambda^F, \qquad \forall F \in \mathcal{F}_h,
$$

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- 1. Introduction of a Lagrange multiplier λ
	- \blacktriangleright Weak formulation of waves equations

$$
\begin{cases}\n\int_{K} i\omega \rho^{K} \mathbf{v}^{K} \cdot \mathbf{w} + \int_{K} \underline{\sigma}^{K} : \nabla \mathbf{w} - \int_{\partial K} \underline{\widehat{\sigma}}^{\partial K} \cdot \mathbf{n} \cdot \mathbf{w} = 0 \\
\int_{K} i\omega \underline{\sigma}^{K} : \underline{\xi} + \int_{K} \mathbf{v}^{K} \cdot \nabla \cdot \left(\underline{\underline{C}}^{K} \underline{\xi} \right) - \int_{\partial K} \widehat{\mathbf{v}}^{\partial K} \cdot \underline{\underline{C}}^{K} \underline{\xi} \cdot \mathbf{n} = \int_{K} f_{S}^{K} \cdot \underline{\xi} \\
\stackrel{\sim}{\equiv} \\
\stackrel{\sim}{\sim} \\
\frac{\widehat{\mathbf{v}}^{\partial K}}{\underline{\hat{\Xi}}} \cdot \mathbf{n} = \frac{\lambda^{F}}{\underline{\hat{\Xi}}} , \quad \mathbf{n} - \tau \mathbf{I} \left(\mathbf{v}^{K} - \lambda^{F} \right), \quad \text{on } \partial K \\
\text{where } \tau \text{ is the stabilization parameter } (\tau > 0)\n\end{cases}
$$

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1. Introduction of a Lagrange multiplier λ

 \triangleright Weak formulation of waves equations

$$
\begin{array}{rcl}\n\widehat{\mathbf{v}}^{\partial K} & = & \lambda^F, \\
\widehat{\underline{\underline{\sigma}}}^{\partial K} \cdot \mathbf{n} & = & \underline{\underline{\sigma}}^K \cdot \mathbf{n} - \tau \mathbf{I} \left(\mathbf{v}^K - \lambda^F \right), \quad \text{on } \partial K\n\end{array}
$$

 \blacktriangleright Local HDG formulation

$$
\left\{\begin{array}{l} \displaystyle\int_{K} i\omega \rho^{K} \mathbf{v}^{K}\cdot \mathbf{w}-\int_{K}\left(\nabla \cdot \underline{\underline{\sigma}}^{K}\right) \cdot \mathbf{w}+\int_{\partial K} \tau \mathbf{I}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \mathbf{w}=0\\ \displaystyle\int_{K} i\omega \underline{\underline{\sigma}}^{K}:\underline{\xi}+\int_{K} \mathbf{v}^{K}\cdot \nabla \cdot \left(\underline{\underline{C}}^{K}\underline{\xi}\right)-\int_{\partial K} \lambda^{F}\cdot \underline{\underline{C}}^{K}\underline{\xi} \cdot \mathbf{n}=0 \end{array}\right.
$$

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1. Introduction of a Lagrange multiplier λ

 \triangleright Weak formulation of waves equations

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\begin{array}{rcl}\n\widehat{\mathbf{v}}^{\partial K} & = & \lambda^F, \\
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$$

 \blacktriangleright Local HDG formulation

$$
\mathbb{A}^K \mathbf{W}^K + \mathbb{C}^K \lambda = \mathbf{S}^K
$$

with $\mathbf{W}^K = (\mathbf{v}^K, \underline{\sigma}^K)^T$ and $\lambda = (\lambda^{F_1}, \lambda^{F_2}, ..., \lambda^{F_{n_f}})^T$, where $n_f = \text{card}(\mathcal{F}_h)$

- 1. Introduction of a Lagrange multiplier λ
- 2. Expressing the initial unknowns W^{K} as a function of λ

$$
\mathbb{A}^K \mathbf{W}^K + \mathbb{C}^K \lambda = \mathbf{S}^K \longrightarrow \mathbf{W}^K = \left(\mathbb{A}^K\right)^{-1} \mathbf{S}^K - \left(\mathbb{A}^K\right)^{-1} \mathbb{C}^K \lambda
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\mathbb{A}^K \mathbf{W}^K + \mathbb{C}^K \lambda = \mathbf{S}^K \longrightarrow \mathbf{W}^K = \left(\mathbb{A}^K\right)^{-1} \mathbf{S}^K - \left(\mathbb{A}^K\right)^{-1} \mathbb{C}^K \lambda
$$

3. Transmission condition :

$$
\int_{\mathbf{F}} \underline{\mathbb{I}}\underline{\widehat{\underline{\sigma}}^{\partial K}}\cdot \mathbf{n}]\!]\cdot \boldsymbol{\eta} = 0
$$

- 1. Introduction of a Lagrange multiplier λ
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$$
\mathbb{A}^{K}W^{K} + \mathbb{C}^{K}\lambda = \mathbf{S}^{K} \longrightarrow W^{K} = (\mathbb{A}^{K})^{-1}\mathbf{S}^{K} - (\mathbb{A}^{K})^{-1}\mathbb{C}^{K}\lambda
$$

3. Transmission condition :

$$
\sum_{\mathsf{K}} \mathbb{B}^{\mathsf{K}} \mathsf{W}^{\mathsf{K}} + \mathbb{L}^{\mathsf{K}} \lambda = 0
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$$

4. $M\lambda = S$

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- 1. Introduction of a Lagrange multiplier λ
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3. Transmission condition :

$$
\sum_{K} \mathbb{B}^{K} \mathbf{W}^{K} + \mathbb{L}^{K} \lambda = 0
$$

4. $M\lambda = S$

5. Computation of the solutions of the initial problem **W**^K , element by element

Contents

2D Numerical results : performances comparison of the HDG method

Anisotropic test case Marmousi test-case

3D numerical results

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Main steps of the algorithms

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Main steps of the algorithms

embarrassingly parallel

Anisotropic test case

- \blacktriangleright Three meshes :
	- \triangleright 600 elements
	- \blacktriangleright 3000 elements
	- \blacktriangleright 28000 elements

Anisotropic case : Memory consumption

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Anisotropic case : CPU time (s)

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Marmousi test-case

Computational domain Ω composed of 235000 triangles

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Cluster configuration

- Hardware specification : 16 nodes, 12 cores by nodes
- Caracteristics of computing nodes :
	- ▶ 2 Hexa-core Westmere Intel $\begin{matrix} (R) & X\tan (R) \end{matrix}$ X5670
	- \blacktriangleright Frequency : 2,93 GHz
	- \triangleright Cache L3 : 12 Mo
	- \blacktriangleright RAM : 96 Go
	- Infiniband DDR : $20Gb/s$
	- \blacktriangleright Ethernet : 1Gb/s

Efficiency of the parallelism of the construction of the global matrix

$$
\text{Efficiency } = \frac{\text{t}_{\text{ref}}}{\text{nb_cores} \times \text{t}_{\text{one core}}}
$$

where :

t $_{ref} = t_{seq} = sequential computational time$ t _{one core} $=$ computational time with one core

Efficiency of the parallelism of the construction of the global matrix

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Efficiency of the parallelism of the whole algorithm

Speed up for the global matrix construction

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Memory required (GB) for the simulation

Comparison with Finite Elements method

Comparison with Finite Elements method

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Contents

2D Numerical results : performances comparison of the HDG method

3D numerical results Performances analysis of the HDG method Impact of the linear solver

Methods - HDG method for 3D Helmholtz wave equations May 15, 2017 - 20/31

Epati test-case

 V_p -velocity model (m.s⁻¹), vertical section at $y = 700$ m Mesh composed of 25 000 tetrahedrons

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Epati test-case : Scalability, efficiency of the HDG global system construction

Efficiency of the parallelism of the whole algorithm (construction and resolution)

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Resolution of the linear system $M\lambda = S$

- \blacktriangleright Direct solvers
- \blacktriangleright Iterative solvers
- \blacktriangleright Hybrid solvers, combination between direct and iterative solvers
- \triangleright Direct solvers combined with H-matrices

^I ...

Resolution of the linear system $M\lambda = S$

Direct solvers

- \blacktriangleright Iterative solvers
- **Hybrid solvers**, combination between direct and iterative solvers
- \triangleright Direct solvers combined with H-matrices

^I ...

Resolution of the linear system $M\lambda = S$

- **In direct solver : MUMPS** (MUltifrontal Massively Parallel sparse direct Solver) :
	- Direct factorization $A = LU$ or $A = LDL^T$
	- \blacktriangleright Multiple RHS

P.R. Amestoy, I.S. Duff and J.-Y. L'Excellent. Multifrontal parallel distributed symmetric and unsymmetric solvers. Computational Methods in Applied Mechanics and Engineering, Vol. 184 :501-520, 2000.

• hybrid solver: MaPhys (Massively Parallel Hybrid Solver):

- \triangleright combination of direct and iterative methods
- \triangleright non-overlapping algebraic domain decomposition method (Schur complement method)
- \triangleright resolution of each local problem thanks to direct solver such as MUMPS

https ://gitlab.inria.fr/solverstack/maphys

Epati test-case : Memory consumption

Average memory for one node (8 MPI by node and 3 threads by MPI)

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Epati test-case : Execution time

Execution time for the resolution of the HDG- \mathbb{P}_3 system

Epati test-case : Cumulative time of the two solvers using 96 cores

Conclusion-Perspectives

Conclusions

- \blacktriangleright HDG method more efficient than classical DG methods for a same accuracy (memory and CPU time)
- \triangleright In 2D, HDG method has equivalent computational performances than FE methods
- \triangleright Maphys good alternative to MUMPS for the resolution of the HDG system with one RHS
	- ✗ Multiple RHS not yet implemented in Maphys

Conclusion-Perspectives

Perspectives

- \triangleright more detailed analysis of the comparison between solvers
	- \blacktriangleright larger meshes
	- \blacktriangleright more powerful clusters
- \blacktriangleright extension to H-matrices
- \triangleright 3D comparison between FE and HDG methods
- \blacktriangleright extension to elasto-acoustic case
- \blacktriangleright application to inverse problems

Conclusions-Perspectives

Thank you !

