



**HAL**  
open science

## **Non-integrability of the minimum-time kepler problem**

Michaël Orioux, Jean-Baptiste Caillau, Thierry Combet, Jacques Fejoz

► **To cite this version:**

Michaël Orioux, Jean-Baptiste Caillau, Thierry Combet, Jacques Fejoz. Non-integrability of the minimum-time kepler problem . 2018. <hal-01679261v1>

**HAL Id: hal-01679261**

**<https://inria.hal.science/hal-01679261v1>**

Preprint submitted on 9 Jan 2018 (v1), last revised 24 Jun 2018 (v2)

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



HAL Authorization

# NON-INTEGRABILITY OF THE MINIMUM-TIME KEPLER PROBLEM

M. ORIEUX, J.-B. CAILLAU, T. COMBOT, AND J. FÉJOZ

ABSTRACT. We prove that the minimum-time controlled Kepler problem is not meromorphically Liouville integrable on the Riemann surface of its Hamiltonian.

**Key words.** Integrability, Hamiltonian systems, optimal control, differential Galois theory.

## 1. INTRODUCTION

The Kepler problem

$$(1) \quad \ddot{q} + \frac{q}{\|q\|^3} = 0, \quad q \in \mathbb{R}^2 \setminus \{0\}.$$

is a classical reduction of the two-body problem [2]. Here, we think of  $q$  as an artificial satellite, and of the attraction as the action of the Earth. We are interested in controlling the transfer of the satellite from one Keplerian orbit towards another, in the plane. Denoting  $v = \dot{q}$  the velocity, and the adjoint variables of  $q$  and  $v$  by  $p_q$  and  $p_v$ , the minimum time dynamics is a Hamiltonian system with Hamiltonian

$$(2) \quad H(q, v, p_q, p_v) = p_q \cdot v - \frac{p_v \cdot q}{\|q\|^3} + \|p_v\|,$$

as is explained in section 2.1. Prior studies of this problem can be found in [4, 6]. The controlled Kepler problem can be embedded in the two parameter family obtained when considering the control of the circular restricted three-body problem:

$$(3) \quad \ddot{q} + \nabla \Omega_\mu(t, q) = \varepsilon u,$$

---

CEREMADE, Univ. Paris Dauphine, Place du Maréchal de Lattre de Tassigny, F-75016 Paris ([orieux@ceremade.dauphine.fr](mailto:orieux@ceremade.dauphine.fr)).

LJAD, Univ. Côte d'Azur & CNRS/Inria, Parc Valrose, F-06108 Nice ([caillau@unice.fr](mailto:caillau@unice.fr)).

Institut math., UBFC & CNRS, 9 avenue Savary, F-21078 Dijon ([thierry.combot@ubfc.fr](mailto:thierry.combot@ubfc.fr)).

CEREMADE, Univ. Paris Dauphine, Place du Maréchal de Lattre de Tassigny, F-75016 Paris and IMCCE, Observatoire de Paris, 77 avenue Denfert Rochereau, F-75014 Paris ([jacques.fejoz@dauphine.fr](mailto:jacques.fejoz@dauphine.fr)).

where

$$\Omega_\mu(t, q) = -\frac{1 - \mu}{\sqrt{(q_1 + \mu \cos t)^2 + (q_2 + \mu \sin t)^2}} - \frac{\mu}{\sqrt{(q_1 - (1 - \mu) \cos t)^2 + (q_2 - (1 - \mu) \sin t)^2}}$$

is the potential parameterized by the ratio of masses,  $\mu \in [0, 1/2]$ , and where  $u \in \mathbb{R}^2$  is the control, whose amplitude is modulated by the second parameter,  $\varepsilon \geq 0$ . Alternatively to time minimization, minimization of the  $\mathbb{L}^2$  norm of the control can be considered,

$$\int_0^{t_f} u^2(t) dt \rightarrow \min.$$

This is the so-called *energy* cost. In the uncontrolled case ( $\varepsilon = 0$ ), it is well known that the Kepler case ( $\mu = 0$ ) is integrable (and geodesic) while there are obstructions to integrability for positive  $\mu$ . In the controlled case ( $\varepsilon > 0$ ), the Kepler problem for the energy cost has been shown to be integrable (and geodesic) when suitably averaged (see [5] for a survey). The aim of this paper is to study the integrability properties of the Kepler problem for time minimization.

The pioneering work of Ziglin in the 80's [17], followed by the modern formulation of differential Galois theory in the late 90's by Morales, Ramis and Simó [11, 12], have led to a very diverse literature on the integrability of Hamiltonian systems. According to Pontrjagin's Maximum principle, one can turn general optimization problems with dynamical constraints into Hamiltonian systems, which are generally not everywhere differentiable. Optimal control theory thus provides an abundant class of dynamical systems, for which integrability naturally is a central question. Yet, differential Galois theory has not so often been applied in this context, in part because of the difficulty brought by the singularities. Notwithstanding these singularities (vanishing of the adjoint variable  $p_v$ , here), we show how to apply these ideas to the system (2).

## 2. SETTING

**2.1. The minimum-time controlled Kepler problem.** We first recall some classical facts on optimal control. We refer for example to the book of Agrachev [1] for more details. Let  $M$  be an  $n$ -dimensional smooth manifold and  $U$  an arbitrary subset of  $\mathbb{R}^m$ . A *controlled dynamical system* is a smooth family of vector fields

$$f : M \times U \rightarrow TM$$

(weaker smoothness assumptions may also be made). Let

$$\mathcal{U} = \{u : t \mapsto u(t) \in L_{\text{loc}}^\infty(\mathbb{R}, U)\}$$

be the set of *admissible controls*. A preliminary question is the following: Is some final state  $x_f$  accessible from some initial state  $x_0$ , *i.e.* does the system

$$(4) \quad \begin{cases} \dot{x}(t) = f(x(t), u(t)), \\ x(0) = x_0, \\ x(t_f) = x_f \end{cases}$$

have a solution for some control  $u$  in  $\mathcal{U}$ ? The system is said *controllable* if the answer is positive for all possible initial and final states  $x_0, x_f \in M$ . The controlled Kepler problem, associated with (1), is

$$\begin{cases} \ddot{q} + \frac{q}{\|q\|^3} = u, & q, u \in \mathbb{R}^2, & u_1^2 + u_2^2 \leq 1, \\ (q(0), \dot{q}(0)) = (q_0, v_0), \\ (q(t_f), \dot{q}(t_f)) = (q_f, v_f), \end{cases}$$

where  $q$  is always the position vector of a satellite, and, in this mechanical context, the control  $u$  is the thrust of the engine. The thrust is naturally bounded; here we assume it takes values in the euclidean unit ball. (Note that, with respect to (3), we have chosen  $\varepsilon = 1$ ; as will be clear from Section 3, this does not restrict the generality of the analysis.)

**Proposition 1** ([6]). *The Kepler problem is controllable.*

This is a consequence of two facts : the Lie algebra generated by the drift and the vector field supporting the control generate the whole tangent space in each point (which entails some local controllability), and the uncontrolled flow (or *drift*) of the Kepler problem is recurrent. Under some additional compactness assumptions, one is then able to retrieve existence of optimal controls.

We now deal with such optimal controls. Here we will restrict ourselves to integral cost functions, that is to problems of the form

$$(5) \quad \begin{cases} \dot{x}(t) = f(x(t), u(t)), \\ x(0) = x_0, \\ x(t_f) = x_f, \\ \int_0^{t_f} \varphi(x(t), u(t)) dt \rightarrow \min. \end{cases}$$

In the early 60's, Pontrjagin and his coauthors realized that necessary conditions for optimality could be stated in Hamiltonian terms.

**Definition 1.** *The associated pseudo-Hamiltonian is*

$$H : T^*M \times \mathbb{R} \times U \rightarrow \mathbb{R}, \quad (x, p, p_0, u) \mapsto \langle p, f(x, u) \rangle + p_0 \varphi(x, u).$$

The fundamental result is the following Pontrjagin Maximum Principle (PMP), for which we refer to [14] or, [1] for a more recent presentation.

**Theorem 1** (PMP). *If  $(x, u)$  solves (5), there exists a Lipschitz curve  $p(t) \in T_{x(t)}^*M$ ,  $t \in [0, t_f]$ , a constant  $p_0 \leq 0$ , such that  $(p(t), p_0) \neq 0$  and such that, almost everywhere,*

- $(x, p)$  is solution of the Hamiltonian system associated with  $H(\cdot, \cdot, u)$ :

$$\begin{cases} \dot{x} = \frac{\partial H}{\partial p}(x, p, u), \\ \dot{p} = -\frac{\partial H}{\partial x}(x, p, u), \end{cases}$$

- $H(x(t), p(t), u(t)) = \max_{\tilde{u} \in U} H(x(t), p(t), \tilde{u})$ .

Such curves  $(x, p)$  are called extremals. As a consequence of the maximization condition, the pseudo-Hamiltonian evaluated along an extremal is constant. If, moreover, the final time is free, this constant is zero.

This powerful result has some downsides. The Hamiltonian is defined on the cotangent bundle of the original phase space, and thus the dimension is doubled. Moreover, the maximization condition, which is the key point to obtain a Hamiltonian system, generates singularities. The above theorem applies to time minimization with  $\varphi \equiv 1$  (and free final time). In this case, the constant  $p_0$  is only related to the level of the Hamiltonian, and we will not mention it in the sequel.

**2.2. Main result.** The minimum-time Kepler problem can be stated according to

$$(6) \quad \begin{cases} \ddot{q} + \frac{q}{\|q\|^3} = u, \|u\| \leq 1, \\ (q(0), \dot{q}(0)) = (q_0, v_0), \\ (q(t_f), \dot{q}(t_f)) = (q_f, v_f), \\ t_f \rightarrow \min, \end{cases}$$

where as before  $q \in \mathbb{R}^2$  is the position vector and  $u \in \mathbb{R}^2$  the control. It will be convenient to use the same notation as in the general problem (5), and let

$$q = (x_1, x_2), \quad \dot{q} = (x_3, x_4),$$

be the coordinates on the initial phase space  $M = (\mathbb{R}^2 \setminus \{0\}) \times \mathbb{R}^2$ . The pseudo-Hamiltonian is then

$$H(x, p, u) = p_1 x_3 + p_2 x_4 - \frac{p_3 x_1 + p_4 x_2}{(x_1^2 + x_2^2)^{3/2}} + p_3 u_1 + p_4 u_2.$$

By the Maximum Principle, the extremals of (6) must be projections on  $M$  of solutions of the maximized Hamiltonian system

$$H(x, p) = p_1 x_3 + p_2 x_4 - \frac{p_3 x_1 + p_4 x_2}{(x_1^2 + x_2^2)^{3/2}} + \sqrt{p_3^2 + p_4^2}$$

on  $T^*M$ , while the control is given by

$$u = \frac{1}{\sqrt{p_3^2 + p_4^2}}(p_3, p_4).$$

Now, let

$$\mathcal{M} = \{(x, p, r) \in \mathbb{C}^8 \times \mathbb{C}_*^2, r_1^2 = x_1^2 + x_2^2, r_2^2 = p_3^2 + p_4^2\}$$

be the Riemann surface of  $H$ . It is a complex symplectic manifold (with local Darboux coordinates  $(x, p)$  outside the singular hypersurface  $r_1 r_2 = 0$ ), over which  $H$  extends meromorphically, and even rationally, since

$$(7) \quad H(x, p, r) = p_1 x_3 + p_2 x_4 - \frac{p_3 x_1 + p_4 x_2}{r_1^3} + r_2.$$

The Hamiltonian  $H$  has four degrees of freedom, hence (see [2]) the meromorphic Liouville integrability of  $H$  over  $\mathcal{M}$  would mean that there would exist three independent first integrals, in addition to  $H$  itself, almost everywhere in  $\mathcal{M}$ . This does not happen.

**Theorem 2.** *The minimum time Kepler problem is not meromorphically Liouville integrable on  $\mathcal{M}$ .*

It is well known that the classical Kepler problem is integrable, and even super integrable (since there are more first integrals than degrees of freedom, which relates to Kepler's first law and the dynamical degeneracy of the Newtonian potential see [8], for example). But the three-body problem is not, as Poincaré, Julliard-Tosel, Tsygvintsev or Combot have proved in various settings [7, 9, 13, 16]. Similarly, the above theorem asserts that lifting the Kepler problem to the cotangent bundle and introducing the singular control term  $r_2$  break integrability.

This result prevents the existence of enough real analytic (and even meromorphic) first integrals to ensure integrability over  $\mathcal{M}$ . Of course, it does not prevent the existence of an additional first integral which would have a natural frontier outside asymptotic to the real domain and which would thus not extend to the complex plane. Future work may be dedicated to extend this result to real first integrals: The Hamiltonian is invariant by two rotations, and homogeneous with respect to several coordinates. If it is real Liouville integrable, one can find real first integrals have the same symmetries. This may allow us to extend them to the complex domain, creating a contradiction with theorem 2.

### 3. PROOF OF THEOREM 2

The rest of the article is devoted to proving the theorem. Our proof consists in studying the variational equation along some integral curve of (7). In order to carry out this computation, we choose a collision orbit, with the drawback that it requires some regularization. The algebraic obstruction to Liouville integrability will come from the following theorem of Moralès and Ramis.

**Theorem 3** (Moralès-Ramis [12]). *Let us consider a Hamiltonian  $H$  analytic on a complex analytic symplectic manifold and a non constant solution  $\Gamma$ . If  $H$  is integrable in the Liouville sense with meromorphic first integrals, then the first*

order variational equation along  $\Gamma$  has a virtually Abelian Galois group over the base field of meromorphic functions on  $\Gamma$ .

**3.1. A collision orbit.** In order to find an explicit solution, let us define the 4-dimensional symplectic submanifold

$$S = \{(x, p, r), \quad x_2 = x_4 = p_2 = p_4 = 0, r_1 = x_1, r_2 = -p_3\} \cap \mathcal{M}.$$

As  $S$  is the phase space of the controlled Kepler problem on the line parameterized by  $q_1$ , it is invariant. On the interior of  $S$ ,  $(x_1, x_3, p_1, p_3)$  is a set of (Darboux) coordinates and, in restriction to  $S$ , the Hamiltonian reduces to

$$H(x, p) = p_1 x_3 - \frac{p_3}{x_1^2} - p_3,$$

so the Hamiltonian vector field on  $S$  is

$$\begin{cases} \dot{x}_1 = x_3 \\ \dot{x}_3 = -1 - \frac{1}{x_1^2} \\ \dot{p}_1 = -\frac{2p_3}{x_1^3} \\ \dot{p}_3 = -p_1. \end{cases}$$

In particular,

$$(8) \quad \begin{cases} \ddot{x}_1 = -1 - \frac{1}{x_1^2} \\ \ddot{p}_3 - 2\frac{p_3}{x_1^3} = 0. \end{cases}$$

As is known since the work of Charlier and Saint Germain on the Kepler problem with a constant force (see [3]), the function

$$C = \frac{1}{2}x_3^2 + x_1 - \frac{1}{x_1}$$

is a first integral on  $S$  and  $H|_S$  is integrable. Let us switch to the time  $s = x_1(t)$ , in which an explicit solution will be written, and denote by  $' = \frac{d}{ds}$  derivatives with respect to the new time. It suffices to find an obstruction in this modified time, as explained at the end of the proof.

Using (8), we see that the variable  $p_3$ , as a function of the new time, satisfies the linear differential equation

$$2 \left( C + \frac{1}{x_1} - x_1 \right) p_3''(x_1) - \left( 1 + \frac{1}{x_1^2} \right) p_3'(x_1) - \frac{2}{x_1^3} p_3(x_1) = 0,$$

which yields

$$p_3(x_1) = \frac{\sqrt{-Cx_1 + x_1^2 - 1}}{\sqrt{x_1}} \left( c_1 \int \frac{x_1^{3/2}}{(-Cx_1 + x_1^2 - 1)^{3/2}} dx_1 + c_2 \right)$$

for some constants of integration  $c_1$  and  $c_2$ ; here the symbol  $\int f(x_1)dx_1$  denotes some primitive of  $f$  with respect to the variable  $x_1$ . Let us consider the simple—but rich enough—case when  $c_1 = 0$ ,  $c_2 = 1$ :

$$p_3(x_1) = \frac{\sqrt{-Cx_1 + x_1^2 - 1}}{\sqrt{x_1}}.$$

Using the expression of the first integral  $C$  and of the vector field, we deduce

$$x_3(x_1) = \sqrt{2} \frac{\sqrt{-Cx_1 + x_1^2 - 1}}{\sqrt{x_1}}, \quad p_1(x_1) = -\frac{1}{\sqrt{2}} \frac{x_1^2 + 1}{x_1^2}.$$

Choosing  $C = 2i$  and some determination of the squares yields a particularly simple solution  $\Gamma$  drawn on  $S \subset \mathcal{M}$ ,

$$(9) \quad \begin{cases} x_1 = x_1 \\ x_2 = 0 \\ x_3 = \sqrt{2} \frac{x_1 - i}{\sqrt{x_1}} \\ x_4 = 0 \end{cases} \quad \begin{cases} p_1 = -\frac{x_1^2 + 1}{\sqrt{2}x_1^2} \\ p_2 = 0 \\ p_3 = \frac{x_1 - i}{\sqrt{x_1}} \\ p_4 = 0. \end{cases}$$

**3.2. Normal variational equation.** In the initial time, the linearized equation along  $\Gamma$  is the Hamiltonian vector field associated with the Hamiltonian  $DH$  along  $\Gamma$ :

$$\dot{Z}(t) = A(t)Z(t), \quad A(t) = J D^2 H(\Gamma(t)),$$

where  $J$  is the Poisson structure. In the coordinates  $(x_1, \dots, x_4, p_1, \dots, p_4)$ ,

$$J = \begin{pmatrix} 0_4 & I_4 \\ -I_4 & 0_4 \end{pmatrix}.$$

But we will keep using time  $x_1$ , instead of the initial time  $t$ , writing

$$Z'(x_1(t)) = \frac{1}{x_3(t)} A(x_1(t)) Z(x_1(t)).$$

Let us now order coordinates as  $(x_1, x_3, p_1, p_3, x_2, x_4, p_2, p_4)$ . Since  $S$  is an invariant submanifold, the  $8 \times 8$  matrix  $A$  has an upper triangular bloc structure

$$A = \begin{pmatrix} A_1 & A_2 \\ 0 & A_3 \end{pmatrix}$$

with

$$A_3 = \begin{pmatrix} 0 & 0 & 0 & \frac{1}{\sqrt{2}p_3} \\ -\frac{1}{\sqrt{2}p_3^2} & 0 & -\frac{1}{\sqrt{2}x_1^3 p_3} & 0 \\ 0 & \frac{1}{\sqrt{2}p_3} & 0 & 0 \\ -\frac{1}{\sqrt{2}x_1^3 p_3} & 0 & \frac{3}{\sqrt{2}x_1^4} & 0 \end{pmatrix}.$$

The Morales-Ramis Theorem gives necessary conditions for Liouville integrability in terms of the Galois group of this linear differential system over the base field of meromorphic functions *on*  $\Gamma$ . Looking at the expression (9) of  $\Gamma$ , we see that

meromorphic functions on  $\Gamma$  are just meromorphic functions in  $\sqrt{x_1} \in \mathbb{C} \setminus \{0, \pm\sqrt{i}\}$ . The block  $A_3$  corresponds to infinitesimal variations in the normal direction to  $S$ , which is the part where interesting phenomena might occur. As the Picard-Vessiot field is generated by all the components of the solutions, the Picard Vessiot field  $K$  generated by the normal variational equation

$$(L) : \quad X' = A_3 X, \quad X = (X_1, X_2, X_3, X_4)$$

is a subfield of the Picard Vessiot field of the whole variational equation, and thus  $\text{Gal}(A) \supset \text{Gal}(A_3)$ . That  $\text{Gal}(A_3)$  is not virtually Abelian will thus imply that  $\text{Gal}(A)$  itself is not virtually Abelian. In order to reduce the system to a one dimensional linear equation, we use the cyclic vector method on  $A_3$  : from  $(L)$  we get  $X_1' = L_1(X_1, X_2, X_3, X_4)$ , where  $L_1$  is a linear form on  $\mathbb{R}^4$ , thus by derivation,

$$\begin{aligned} X_1'' &= L_1(X_1', X_2, X_3, X_4) + L_1(X_1, X_2', X_3, X_4) \\ &\quad + L_1(X_1, X_2, X_3', X_4) + L_1(X_1, X_2, X_3, X_4') \\ &= L_2(X_1, X_2, X_3, X_4). \end{aligned}$$

Iterating, we obtain

$$\begin{cases} X_1 = X_1 \\ X_1' = L_1(X_1, X_2, X_3, X_4) \\ X_1'' = L_2(X_1, X_2, X_3, X_4) \\ X_1^{(3)} = L_3(X_1, X_2, X_3, X_4) \\ X_1^{(4)} = L_4(X_1, X_2, X_3, X_4), \end{cases}$$

where the  $L_i$ 's are linear forms on  $\mathbb{R}^4$ . As five linear forms in  $\mathbb{R}^4$  are linked,  $X_1$  must satisfy some linear differential equation of order 4, which we compute to be

$$(10) \quad X_1^{(4)} + \frac{2(3i - 5x_1)}{x_1(i - x_1)} X_1^{(3)} + \frac{(-3x_1 + i)(-29x_1 + 23i)}{4(x_1 - i)^2 x_1^2} X_1'' - \frac{(i - 3x_1)(7x_1 + i)}{4(x_1 - i)^2 x_1^3} X_1' + \frac{3x_1 + i}{4(x_1 - 1)^3 x_1^4} X_1 = 0.$$

We find a solution of this equation of the form

$$X_1(x_1) = \frac{i - x_1}{\sqrt{x_1}} \left( c_1 + c_2 \int \sqrt{x_1} (1 + ix_1)^{-\frac{3}{2} - i\frac{\sqrt{3}}{2}} \cdot {}_2F_1(\gamma(x_1)) dx_1 \right),$$

where  ${}_2F_1$  is the Gauss hypergeometric function and

$$\gamma(x_1) = \left( \frac{5}{2} - \frac{\sqrt{3}}{2}i, \frac{1}{2} + \frac{\sqrt{3}}{2}i, 1 + \sqrt{3}i, 1 + ix_1 \right).$$

The Picard-Vessiot field  $K$  contains this solution, and as it is a differential field, it also contains

$$\sqrt{x_1} (1 + ix_1)^{-\frac{3}{2} - i\frac{\sqrt{3}}{2}} {}_2F_1(\gamma(x_1)).$$

But the Galois group of  ${}_2F_1(\gamma(x_1))$  over  $\mathbb{C}(x_1)$  is  $SL_2(\mathbb{C})$  (see Kimura's table, [10]). As the hypergeometric equation (10) is Fuchsian (meaning that all its singular points are regular), we have the following.

**Theorem 4** (Schlesinger density theorem [15]). *Let  $(E) : Y' = AY$  be a Fuchsian differential linear equation and let  $\Pi$  be its monodromy group. Then  $\Pi$  is dense for the Zariski topology in the Galois group of the Picard-Vessiot extension of  $(E)$  over the base field of rational functions:  $\bar{\Pi} = \text{Gal}_{\mathbb{Q}}(A)$ .*

Besides, the Galois group over meromorphic functions contains the monodromy group, and of course, is included in the Galois group over rational functions. Eventually, the meromorphic Galois group in  $x_1 \in \mathbb{C} \setminus \{0, i\}$  is  $SL_2(\mathbb{C})$ . Thus, adding the algebraic extension  $\sqrt{x_1}$ , the Galois group can be reduced to at most one subgroup of index 2, and thus there can only be  $SL_2(\mathbb{C})$  again. So the Galois group of  $K$  over the base field of meromorphic functions in  $\sqrt{x_1} \in \mathbb{C} \setminus \{0, \pm\sqrt{i}\}$  contains  $SL_2(\mathbb{C})$  and is not virtually Abelian. According to Morales-Ramis, this concludes the proof.

## REFERENCES

- [1] A. A. Agrachev and Y. L. Sachkov. *Control theory from the geometric viewpoint*, volume 87 of *Encyclopaedia of Mathematical Sciences*. Springer-Verlag, Berlin, 2004. Control Theory and Optimization, II.
- [2] V. I. Arnold. *Mathematical methods of classical mechanics*, volume 60 of *Graduate Texts in Mathematics*. Springer-Verlag, 1989.
- [3] V. V. Beletsky. *Essays on the motion of celestial bodies*. Springer, 2001.
- [4] J.-B. Caillau and B. Daoud. Minimum time control of the restricted three-body problem. *SIAM Journal on Control and Optimization*, 50(6):3178–3202, 2012.
- [5] J.-B. Caillau, B. Daoud, and J. Gergaud. On some Riemannian aspects of the two and three-body controlled problems. In *Recent Advances in Optimization and its Applications in Engineering*, Springer, pages 205–224, 2009.
- [6] J.-B. Caillau and J. Noailles. Coplanar control of a satellite around the earth. *ESAIM: Control, Optimisation and Calculus of Variations*, 6:239–258, 2001.
- [7] T. Combot. Non-integrability of the equal mass  $n$ -body problem with non-zero angular momentum. *Celestial Mech. Dynam. Astronom.*, 114(4):319–340, 2012.
- [8] J. Féjoz and L. Kaczmarek. Sur le théorème de Bertrand (d'après Michael Herman). *Ergodic Theory Dynam. Systems*, 24(5):1583–1589, 2004.
- [9] E. Julliard-Tosel. Bruns' theorem: the proof and some generalizations. *Celestial Mech. Dynam. Astronom.*, 76(4):241–281, 2000.
- [10] Tosihusa Kimura. On riemann's equations which are solvable by quadratures. *Funkcial. Ekvac.*, 12(269-281):1970, 1969.
- [11] J. J. Morales-Ruiz, J.-P. Ramis, and C. Simo. Integrability of Hamiltonian systems and differential Galois groups of higher variational equations. *Ann. Sci. École Norm. Sup. (4)*, 40(6):845–884, 2007.
- [12] Juan J. Morales-Ruiz and Jean Pierre Ramis. Galoisian obstructions to integrability of Hamiltonian systems. *Methods Appl. Anal.*, 8(1):33–95, 97–111, 2001.
- [13] H. Poincaré. *Les méthodes nouvelles de la mécanique céleste*. Gauthier-Villars, 1892.
- [14] L. S. Pontryagin. The mathematical theory of optimal processes and differential games. *Trudy Mat. Inst. Steklov.*, 169:119–158, 254–255, 1985. Topology, ordinary differential equations, dynamical systems.

- [15] J. Sauloy. *Differential Galois Theory through Riemann-Hilbert Correspondence: An Elementary Introduction*, volume 177. American Mathematical Soc., 2016.
- [16] A. V. Tsygvintsev. Non-existence of new meromorphic first integrals in the planar three-body problem. *Celestial Mech. Dynam. Astronom.*, 86(3):237–247, 2003.
- [17] S. L. Ziglin. Branching of solutions and nonexistence of first integrals in hamiltonian mechanics. i. *Functional Analysis and Its Applications*, 16(3):181–189, 1982.