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# Kinetic entropy for layer-averaged hydrostatic Navier-Stokes equations

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## Abstract

We are interested in the numerical approximation of the hydrostatic free surface incompressible Navier-Stokes equations. By using a layer-averaged version of the equations, we are able to extend previous results obtained for shallow water system and to derive a vertically implicit horizontally explicit finite volume kinetic scheme that ensures the positivity of the approximated water depth and a fully discrete energy inequality.

*Keywords* : Incompressible Euler and Navier-Stokes systems, free surface flows, layer-averaged model, finite volumes, kinetic solver, hydrostatic reconstruction, discrete entropy inequality, IMEX scheme.

*2000 Mathematics Subject Classification*: 65M12, 74S10, 76M12, 35L65.

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# 1 Introduction

Shallow water equations [7, 1] have been widely used to model free surface geophysical fluid flows. This hyperbolic system can be derived from free surface incompressible Navier-Stokes equations by integration along the vertical direction and under a long wave approximation, that implies in particular an hydrostatic distribution of the pressure at the leading order [12]. Due to this reduction of dimension, it is much more easy to deal with in a numerical point of view. Nevertheless it is not suitable for certain situations - stratified flows, wind-driven vertical circulation... In these cases, it is necessary to come back to three dimensional models but the hydrostatic assumption remains mostly valid, leading to consider the hydrostatic Navier-Stokes equations, also known as primitive equations [10, 13, 5, 14, 17], see also [20, 19] for a general introduction to ocean models. In a previous work [4], we introduced a layer-averaged approach to deal with this hydrostatic Navier-Stokes system in a framework that shares (forgetting for a while the viscous part to concentrate on Euler equations and advective processes) some hyperbolic properties with the shallow water equations. Our main result in this work is to propose a vertically implicit and horizontally explicit colocated finite volume scheme to compute approximate solutions of this layer-averaged model for which we are able to prove positivity of the water depth and a fully discrete energy inequality.

It is well known that incompressible hydrostatic Euler and shallow water equations satisfy some invariant domain properties since the water depth of the flow remains non-negative. Moreover regular solutions satisfy an energy equality. Extending the theory developed for scalar hyperbolic equations, one assumes the energy associated to discontinuous solution, that acts as a mathematical entropy, has to decrease through the shocks. To extend these properties to the numerical approximation is not straightforward. A huge literature was devoted to derive explicit colocated finite volume scheme for the shallow water equations including topographic source terms, but, up to our knowledge,

only very few schemes are endowed with these properties [23, 7]. In a recent work [2], a kinetic framework was used to prove that the hydrostatic reconstruction technique [1] associated to a kinetic scheme [21, 22, 3] is able to provide, under a classical CFL condition, approximate solutions of the shallow water system with topographic source terms that are positive and satisfy a fully discrete entropy inequality with a right hand side, or let say an error term, that is proved to be proportional to  $\Delta x^2$  with a constant that is independent of the regularity of the solution - that may develop discontinuities since one deals with hyperbolic system. This inequality is the key step to prove the convergence of the scheme [8]. Thanks to the similarity between the shallow water equations and the layer-averaged model for hydrostatic Euler equations, it is possible to extend the kinetic framework and the hydrostatic reconstruction technique, and then parts of the previous proof, to the latter. But the new terms related to the vertical direction makes the things more intricate, and the fully explicit finite volume scheme originally proposed in [4] may suffer from a very restrictive CFL condition in some situations. In this work, we then propose a new implicit discretization of the exchange terms in the vertical direction that allows us to obtain entropy satisfying approximate solutions under the same CFL condition as for the shallow water system. Up to our knowledge, it is the first time this kind of result is obtained in this three dimensional context. From a practical point of view, the linear problem to solve in the implicit part is restricted to the vertical direction and its size is then proportionnal to the number of layers that are introduced in the model, independently of the horizontal discretization. The added CPU cost is then strictly limited.

The outline of the paper is as follows : in Section 2, we recall the free surface incompressible hydrostatic Navier-Stokes equations ; in Section 3, we consider the inviscid Euler version of the equations and introduce the layer-averaged model and its kinetic description and we derive the related IMEX kinetic scheme ; in Section 4, we demonstrate the discrete energy inequality, first, on flat bottom and, second, including a bottom topography ; finally, in Section 5 we extend the results to the Navier-Stokes case.

## 2 The Navier-Stokes system

We begin by considering the two-dimensional hydrostatic incompressible Navier–Stokes system [15] describing a free surface gravitational flow moving over a bottom topography  $z_b(x)$ . For free surface flows, the hydrostatic assumption consists in neglecting the vertical acceleration, see [10, 13, 5, 14, 17] for justifications and mathematical analysis of the obtained models.

## 2.1 The hydrostatic incompressible Navier-Stokes system

We denote with  $x$  and  $z$  the horizontal and vertical directions, respectively. The system reads

$$\frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0, \quad (1)$$

$$\frac{\partial u}{\partial t} + \frac{\partial u^2}{\partial x} + \frac{\partial uw}{\partial z} + \frac{\partial p}{\partial x} = \frac{\partial \Sigma_{xx}}{\partial x} + \frac{\partial \Sigma_{xz}}{\partial z}, \quad (2)$$

$$\frac{\partial p}{\partial z} = -g + \frac{\partial \Sigma_{zx}}{\partial x} + \frac{\partial \Sigma_{zz}}{\partial z}, \quad (3)$$

and we consider solutions of the equations for

$$t > t_0, \quad x \in \mathbb{R}, \quad z_b(x) \leq z \leq \eta(t, x) := h(t, x) + z_b(x),$$

where  $\eta(t, x)$  represents the free surface elevation,  $h(t, x)$  the water depth,  $\mathbf{u} = (u, w)^T$  the velocity vector and  $g$  the gravity acceleration.

The chosen form of the viscosity tensor is

$$\begin{aligned} \Sigma_{xx} &= 2\mu \frac{\partial u}{\partial x}, & \Sigma_{xz} &= \mu \left( \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right), \\ \Sigma_{zz} &= 2\mu \frac{\partial w}{\partial z}, & \Sigma_{zx} &= \mu \left( \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right), \end{aligned}$$

where  $\mu$  is a dynamic viscosity.

## 2.2 Boundary conditions

The system (1)-(3) is completed with boundary conditions at the bottom and at the free surface. The outward unit normal vector to the free surface  $\mathbf{n}_s$  and the upward unit normal vector to the bottom  $\mathbf{n}_b$  are given by

$$\mathbf{n}_s = \frac{1}{\sqrt{1 + \left(\frac{\partial \eta}{\partial x}\right)^2}} \begin{pmatrix} -\frac{\partial \eta}{\partial x} \\ 1 \end{pmatrix}, \quad \mathbf{n}_b = \frac{1}{\sqrt{1 + \left(\frac{\partial z_b}{\partial x}\right)^2}} \begin{pmatrix} -\frac{\partial z_b}{\partial x} \\ 1 \end{pmatrix},$$

respectively. We then denote with  $\Sigma_T$  the total stress tensor, which has the form:

$$\Sigma_T = -pI_d + \begin{pmatrix} \Sigma_{xx} & \Sigma_{xz} \\ \Sigma_{zx} & \Sigma_{zz} \end{pmatrix}.$$

### 2.2.1 Free surface conditions

At the free surface we have the kinematic boundary condition

$$\frac{\partial \eta}{\partial t} + u_s \frac{\partial \eta}{\partial x} - w_s = 0, \quad (4)$$

where the subscript  $s$  indicates the value of the considered quantity at the free surface.

Assuming negligible the air viscosity, the continuity of stresses at the free boundary imposes

$$\Sigma_T \mathbf{n}_s = -p^a \mathbf{n}_s, \quad (5)$$

where  $p^a = p^a(t, x)$  is a given function corresponding to the atmospheric pressure.

### 2.2.2 Bottom conditions

The kinematic boundary condition at the bottom consists in a classical no-penetration condition:

$$\frac{\partial z_b}{\partial t} + u_b \frac{\partial z_b}{\partial x} - w_b = 0, \quad (6)$$

that reduces to  $\mathbf{u}_b \cdot \mathbf{n}_b = 0$  when  $z_b$  does not depend on time  $t$ .

For the stresses at the bottom we consider a wall law under the form

$$\mathbf{t}_b \cdot \Sigma_T \mathbf{n}_b = \kappa \mathbf{u}_b \cdot \mathbf{t}_b,$$

where  $\mathbf{t}_b$  is a unit vector satisfying  $\mathbf{t}_b \cdot \mathbf{n}_b = 0$ . If  $\kappa(\mathbf{u}_b, h)$  is constant then we recover a Navier friction condition as in [12]. Introducing a laminar friction  $k_l$  and a turbulent friction  $k_t$ , we use the expression

$$\kappa(\mathbf{u}_b, h) = k_l + k_t h |\mathbf{u}_b|,$$

corresponding to the boundary condition used in [16]. Another form of  $\kappa(\mathbf{u}_b, h)$  is used in [9], and for other wall laws the reader can also refer to [18]. Due to thermo-mechanical considerations, in the sequel we will suppose  $\kappa(\mathbf{u}_b, h) \geq 0$ , and  $\kappa(\mathbf{u}_b, h)$  will be often simply denoted by  $\kappa$ .

## 2.3 Energy balance

We recall the fundamental stability property related to the fact that the hydrostatic Navier-Stokes system admits a mechanical energy

$$E(t, x, z) = \frac{u^2}{2} + gz, \quad (7)$$

leading to the following relation for smooth solutions

$$\begin{aligned} & \frac{\partial}{\partial t} \int_{z_b}^{\eta} (E + p^a) dz + \frac{\partial}{\partial x} \int_{z_b}^{\eta} \left[ u(E + p) - \mu \left( 2u \frac{\partial u}{\partial x} + w \left( \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) \right) \right] dz \\ &= -2\mu \int \left[ \left( \frac{\partial u}{\partial x} \right)^2 + \frac{1}{2} \left( \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right)^2 + \left( \frac{\partial w}{\partial z} \right)^2 \right] dz + h \frac{\partial p^a}{\partial t} + (p|_b - p^a) \frac{\partial z_b}{\partial t} - \kappa u_b^2. \end{aligned} \quad (8)$$

For the sake of simplicity, in the following we neglect the variations of the atmospheric pressure  $p^a$  i.e.  $p^a = p_0^a$  with  $p_0^a = 0$  and we also consider

$$\frac{\partial z_b}{\partial t} = 0. \quad (9)$$

It follows the right hand side of (8) is nonpositive and the mean energy

$$\bar{E}(t, x) = \int_{z_b}^{\eta} E(t, x, z) dz$$

is not increasing in time.

### 3 The layer-averaged Euler system

In this section, we present an simplified derivation of the layer-averaged system that was introduced in [4]. As the main result of this work is related to a discrete entropy inequality, we focus on the energy property of the model.

Neglecting the viscous effects (we come back to the Navier-Stokes system in Section 5), we consider the hydrostatic Euler equations in a conservative form

$$\frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0, \quad (10)$$

$$\frac{\partial u}{\partial t} + \frac{\partial u^2}{\partial x} + \frac{\partial uw}{\partial z} + \frac{\partial p}{\partial x} = 0, \quad (11)$$

$$\frac{\partial p}{\partial z} = -g. \quad (12)$$

Kinematic boundary conditions (4) and (6) remain unchanged while the equality of stresses at the free surface (5) reduces to

$$p_s = 0. \quad (13)$$

The energy balance (8) reduces to the equality

$$\frac{\partial}{\partial t} \int_{z_b}^{\eta} E dz + \frac{\partial}{\partial x} \int_{z_b}^{\eta} u(E + p) dz = 0. \quad (14)$$

Note that using (12) and (13) the pressure  $p(t, x, z)$  can be computed has a function depending only of the free surface and the vertical coordinate. Moreover, using the divergence free condition (10) and the boundary condition (6), the vertical velocity can be computed as a function of the horizontal velocity and the vertical coordinate. It follows the unknowns of the system reduce to the water depth  $h(t, x) = \eta(t, x) - z_b(x)$  and the horizontal velocity  $u(t, x, z)$ , that will appear even more explicitly in the layer averaged version presented below.

### 3.1 Discretization of the fluid domain

The interval  $[z_b, \eta]$  is divided into  $N$  layers  $\{L_\alpha\}_{\alpha \in \{1, \dots, N\}}$  of thickness  $l_\alpha h(t, x)$  where each layer  $L_\alpha$  corresponds to the points satisfying  $z \in L_\alpha(t, x) = [z_{\alpha-1/2}, z_{\alpha+1/2}]$  with

$$\begin{cases} z_{\alpha+1/2}(t, x) = z_b(x) + \sum_{j=1}^{\alpha} l_j h(t, x), & z_{1/2}(x) = z_b(x) \\ h_\alpha(t, x) = z_{\alpha+1/2}(t, x) - z_{\alpha-1/2}(t, x) = l_\alpha h(t, x), & \alpha \in [1, \dots, N], \end{cases} \quad (15)$$

with  $l_j > 0$ ,  $\sum_{j=1}^N l_j = 1$ , see Fig. 1.

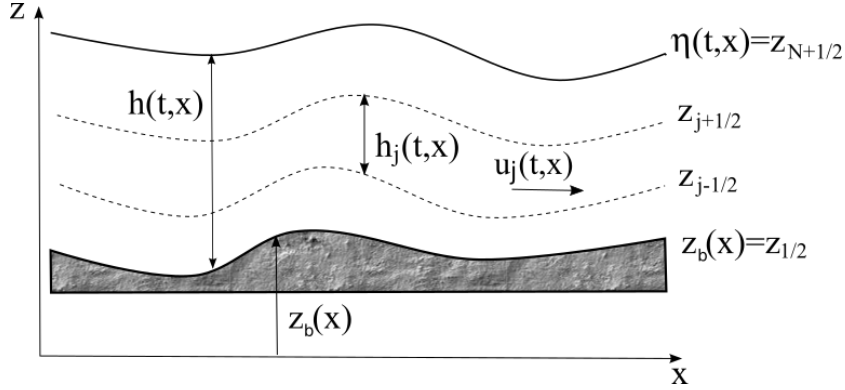


Figure 1: Notations for the multilayer approach.

### 3.2 Layer-averaged model

Let us consider the space  $\mathbb{P}_{0,h}^{N,t}$  of piecewise constant functions defined by

$$\mathbb{P}_{0,h}^{N,t} = \{\mathbf{1}_{z \in L_\alpha(t,x)}(z), \quad \alpha \in \{1, \dots, N\}\},$$

where  $\mathbf{1}_{z \in L_\alpha(t,x)}(z)$  is the characteristic function of the interval  $L_\alpha(t, x)$ . Using this formalism, the projection of  $u, w$  on  $\mathbb{P}_{0,h}^{N,t}$  is a piecewise constant function defined by

$$X^N(x, z, \{z_\alpha\}, t) = \sum_{\alpha=1}^N \mathbf{1}_{[z_{\alpha-1/2}, z_{\alpha+1/2}]}(z) X_\alpha(t, x), \quad (16)$$

for  $X \in (u, w)$ .

In the following we focus on the following layer-averaged model approximating the incompressible hydrostatic Euler system (10)-(12)

$$\frac{\partial h}{\partial t} + \sum_{\alpha=1}^N \frac{\partial h_\alpha u_\alpha}{\partial x} = 0. \quad (17)$$

$$\frac{\partial h_\alpha u_\alpha}{\partial t} + \frac{\partial}{\partial x} \left( h_\alpha u_\alpha^2 + \frac{g}{2} h_\alpha h \right) = -g h_\alpha \frac{\partial z_b}{\partial x} + u_{\alpha+1/2} G_{\alpha+1/2} - u_{\alpha-1/2} G_{\alpha-1/2}, \quad (18)$$



where the mass exchange terms  $G_{\alpha+1/2}$  satisfy

$$G_{\alpha+1/2} = \sum_{j=1}^{\alpha} \left( \frac{\partial h_j}{\partial t} + \frac{\partial h_j u_j}{\partial x} \right), \quad (19)$$

$$G_{N+1/2} = G_{1/2} = 0. \quad (20)$$

and the interface velocities  $u_{\alpha+1/2}$  are defined using an upwinding strategy

$$u_{\alpha+1/2} = \begin{cases} u_{\alpha} & \text{if } G_{\alpha+1/2} \leq 0 \\ u_{\alpha+1} & \text{if } G_{\alpha+1/2} > 0 \end{cases} \quad (21)$$

**Remark 3.1** *In the monolayer case  $\alpha = 1$ , system (17)-(18) with boundary conditions (20) reduces to the classical shallow water equations, see [7, 1].*

Relation (18) is obtained by integrating the momentum equation (11) on the layer  $L_{\alpha}$  while the integration of the divergence equation (10) on the same domain leads to

$$\frac{\partial h_{\alpha}}{\partial t} + \frac{\partial h_{\alpha} u_{\alpha}}{\partial x} = G_{\alpha+1/2} - G_{\alpha-1/2}. \quad (22)$$

Relations (20) directly follow from the kinematic boundary conditions (4) and (6). Then global mass equation (17) and definition of the mass exchange terms (19) are deduced by summing relations (22) over layers  $L_j$  for  $j$  varying from 1 to  $N$  or from 1 to  $\alpha$  respectively.

**Remark 3.2** *The mass exchange terms  $G_{\alpha+1/2}$  can be defined by using only derivatives in space that correspond to partial mass fluxes. It follows from the global mass equation (17) and the definition of the layer depth (15) that relation (19) may be written*

$$G_{\alpha+1/2} = \sum_{j=1}^{\alpha} \left( \frac{\partial h_j u_j}{\partial x} - l_j \sum_{k=1}^N \frac{\partial h_k u_k}{\partial x} \right). \quad (23)$$

**Proposition 3.3** *The layer-averaged system (17)-(18) admits, for smooth solutions, the layer energy balance*

$$\begin{aligned} \frac{\partial}{\partial t} E_{\alpha} + \frac{\partial}{\partial x} \left( u_{\alpha} \left( E_{\alpha} + \frac{gh}{2} h_{\alpha} \right) \right) &= \left( u_{\alpha+1/2} u_{\alpha} - \frac{u_{\alpha}^2}{2} + g\eta \right) G_{\alpha+1/2} \\ &- \left( u_{\alpha-1/2} u_{\alpha} - \frac{u_{\alpha}^2}{2} + g\eta \right) G_{\alpha-1/2}, \end{aligned} \quad (24)$$

where  $E_{\alpha}$  is defined by

$$E_{\alpha} = h_{\alpha} \frac{u_{\alpha}^2}{2} + h_{\alpha} g \left( \frac{h}{2} + z_b \right).$$

It follows the global energy inequality

$$\frac{\partial}{\partial t} \left( \sum_{\alpha=1}^N E_{\alpha} \right) + \frac{\partial}{\partial x} \left( \sum_{\alpha=1}^N u_{\alpha} \left( E_{\alpha} + \frac{gh}{2} h_{\alpha} \right) \right) = - \sum_{\alpha=1}^{N-1} \frac{1}{2} (u_{\alpha+1} - u_{\alpha})^2 |G_{\alpha+1/2}|. \quad (25)$$

where the right hand side is obviously nonpositive.

**Remark 3.4** *The layer energy  $E_\alpha$  is not the meanvalue on a layer of the pointwise energy  $E$  associated to Euler equations and defined by (7). The reason is we are interested in a local kinetic energy associated to momentum equation per layer (18) but in a global potential energy associated to the global mass conservation (17). According to that, it is easy to check that*

$$\sum_{\alpha=1}^N E_\alpha = \sum_{\alpha=1}^N \frac{h_\alpha u_\alpha^2}{2} + \frac{gh^2}{2}$$

*Relation (25) has then to be compared to relation (14) for Euler equations. It appears that the vertical layer-averaging introduces numerical diffusion, as it is usually the case for spatial discretization associated to an upwinding strategy, see (21).*

**Proof of prop. 3.3** *Numerical computations to obtain relation (24) from the layer mass (22) and momentum (18) equations are a straightforward generalization of what is usually done for the classical shallow water model. More precisely, multiplying momentum equation (18) by  $u_\alpha$  and using (22) leads to*

$$\begin{aligned} \frac{\partial}{\partial t} \left( \frac{h_\alpha u_\alpha^2}{2} \right) + \frac{\partial}{\partial x} \left( u_\alpha \frac{h_\alpha u_\alpha^2}{2} \right) + \frac{\partial}{\partial x} \left( \frac{gh}{2} h_\alpha \right) u_\alpha + gh_\alpha u_\alpha \frac{\partial z_b}{\partial x} \\ = \left( u_{\alpha+1/2} u_\alpha - \frac{u_\alpha^2}{2} \right) G_{\alpha+1/2} - \left( u_{\alpha-1/2} u_\alpha - \frac{u_\alpha^2}{2} \right) G_{\alpha-1/2}, \end{aligned}$$

*Now multiplying mass equation (22) by  $g\eta = g(h + z_b)$  and using relations (9) and (15) leads to*

$$\frac{\partial}{\partial t} (gz_b h_\alpha) + \frac{\partial}{\partial t} \left( \frac{gh}{2} h_\alpha \right) + gz_b \frac{\partial h_\alpha u_\alpha}{\partial x} + gh \frac{\partial h_\alpha u_\alpha}{\partial x} = (G_{\alpha+1/2} - G_{\alpha-1/2}) g\eta.$$

*Adding both relations and using (15) for the pressure term, we obtain the energy relation per layer (24). The global energy inequality (3.3) follows by adding these relations for all layers and using the upwind definition of the interface velocities (21). ■*

**Remark 3.5** *Note that multiplying relation (10) and integrating over the layer  $L_\alpha$  leads to the following equality*

$$\begin{aligned} \frac{\partial}{\partial t} \left( \frac{z_{\alpha+1/2}^2 - z_{\alpha-1/2}^2}{2} \right) + \frac{\partial}{\partial x} \left( \frac{z_{\alpha+1/2}^2 - z_{\alpha-1/2}^2}{2} u_\alpha \right) dz = h_\alpha w_\alpha \\ + z_{\alpha+1/2} G_{\alpha+1/2} - z_{\alpha-1/2} G_{\alpha-1/2}, \end{aligned}$$

*that may be used as a postprocessing to compute the layer vertical velocity  $w_\alpha$ .*

### 3.3 Kinetic description

In this paragraph we first give a kinetic interpretation of the system (17)-(24) and then we establish some properties of the proposed discrete scheme. It is a generalization to the layer-averaged framework of kinetic interpretations proposed for other fluid models, see [3, 23, 22]. Note that a first kinetic interpretation of the layer-averaged model was introduced in [4]. Here we propose an improved version and we derive an energy balance at the kinetic level.

### 3.3.1 Kinetic interpretation

Let us define the vector of unknowns

$$U_\alpha = (h_\alpha, h_\alpha u_\alpha)^T, \quad U = (h, h_1 u_1, \dots, h_N u_N)^T. \quad (26)$$

To build Gibbs equilibria, we choose the function

$$\chi_0(z) = \frac{1}{\pi} \sqrt{1 - \frac{z^2}{4}}.$$

This choice corresponds to the classical kinetic maxwellian used in [23] for example and that is defined for  $\xi \in \mathbb{R}$  by

$$M_\alpha = M(U_\alpha, \xi) = \frac{h_\alpha}{c} \chi_0\left(\frac{\xi - u_\alpha}{c}\right) = \frac{l_\alpha}{g\pi} \left(2gh - (\xi - u_\alpha)^2\right)_+^{1/2} = l_\alpha \overline{M}_\alpha, \quad (27)$$

with

$$c = \sqrt{\frac{g}{2}h},$$

It satisfies the following moment relations,

$$\begin{aligned} \int_{\mathbb{R}} M(U_\alpha, \xi) d\xi &= h_\alpha, & \int_{\mathbb{R}} \xi M(U_\alpha, \xi) d\xi &= h_\alpha u_\alpha, \\ \int_{\mathbb{R}} \xi^2 M(U_\alpha, \xi) d\xi &= h_\alpha u_\alpha^2 + h_\alpha \frac{gh}{2}. \end{aligned} \quad (28)$$

Now we introduce a second list of Gibbs equilibria  $N_{\alpha+1/2}$  associated to the mass exchange terms between layers and defined by

$$N_{\alpha+1/2} = N(U_\alpha, U_{\alpha+1}, \xi) = \frac{G_{\alpha+1/2}}{c} \chi_0\left(\frac{\xi - u_{\alpha+1/2}}{c}\right), \quad \alpha = 1, \dots, N-1 \quad (29)$$

completed by the boundary conditions  $N_{1/2} = N_{N+1/2} = 0$ , see (20). Due to the upwind definition of the interface velocity (21), definition (29) is equivalent to

$$N_{\alpha+1/2} = \frac{G_{\alpha+1/2}}{h} \overline{M}_{\alpha+1/2}, \quad \overline{M}_{\alpha+1/2} = \begin{cases} \overline{M}_\alpha & \text{if } G_{\alpha+1/2} \leq 0 \\ \overline{M}_{\alpha+1} & \text{if } G_{\alpha+1/2} \geq 0 \end{cases} \quad (30)$$

It follows from relations (28) that

$$\int_{\mathbb{R}} N_{\alpha+1/2} d\xi = G_{\alpha+1/2}, \quad \int_{\mathbb{R}} \xi N_{\alpha+1/2} d\xi = G_{\alpha+1/2} u_{\alpha+1/2} \quad (31)$$

**Remark 3.6** *The kinetic mass exchange term  $N_{\alpha+1/2}$  defined by relation (30) is well defined even when the water depth  $h$  vanishes. Indeed it is easy to see from definition (19) that*

$$\frac{G_{\alpha+1/2}}{\sqrt{h}} = \mathcal{O}(h^{1/3}).$$

Now definition (27) implies that

$$g\pi M_j \leq \sqrt{h}, \quad j = \alpha, \alpha + 1,$$

It follows that

$$\frac{G_{\alpha+1/2}}{h} M_j = \mathcal{O}(h^{1/3}).$$

We are now equipped to exhibit the kinetic interpretation of the layer-averaged model (17)-(18).

**Proposition 3.7** *The functions  $(h, u^N)$  defined by (16) are strong solutions of the system (17)-(18) if and only if the sets of equilibria  $\{M_\alpha\}_{\alpha=1}^N$ ,  $\{N_{\alpha+1/2}\}_{\alpha=0}^N$  are solutions of the kinetic equations defined by*

$$(\mathcal{B}_\alpha) \quad \frac{\partial M_\alpha}{\partial t} + \xi \frac{\partial M_\alpha}{\partial x} - g \frac{\partial z_b}{\partial x} \frac{\partial M_\alpha}{\partial \xi} - N_{\alpha+1/2} + N_{\alpha-1/2} = Q_\alpha, \quad \alpha = 1, \dots, N \quad (32)$$

The quantities  $Q_\alpha = Q_\alpha(t, x, \xi)$  are “collision terms” equal to zero at the macroscopic level, i.e. they satisfy a.e. for values of  $(t, x)$

$$\int_{\mathbb{R}} Q_\alpha d\xi = 0, \quad \int_{\mathbb{R}} \xi Q_\alpha d\xi = 0. \quad (33)$$

**Proof of proposition 3.7** *The proof relies on averages w.r.t the variable  $\xi$  of Eq. (32) by using relations (28) and (31). Then using (33), the quantities*

$$\sum_1^N \int_{\mathbb{R}} (\mathcal{B}_\alpha) d\xi, \quad \text{and} \quad \int_{\mathbb{R}} \xi (\mathcal{B}_\alpha) d\xi,$$

respectively gives (17) and (18) that completes the proof. ■

There are a lot of functions that satisfy the integral relations (28). The interest of the kinetic maxwellian defined by (27) lies in its link with the kinetic entropy

$$H_\alpha(f, \xi, z_b) = l_\alpha \left( \frac{\xi^2}{2} f + \frac{g^2 \pi^2}{6} f^3 + g z_b f \right) = l_\alpha H(f, \xi, z_b), \quad (34)$$

where  $f \geq 0$ ,  $\xi \in \mathbb{R}$ ,  $z_b \in \mathbb{R}$ . Indeed one can check the relations

$$\begin{aligned} E_\alpha &= l_\alpha \int_{\mathbb{R}} H(\overline{M}_\alpha, \xi, z_b) d\xi, \\ u_\alpha \left( E_\alpha + \frac{g}{2} h_\alpha h \right) &= l_\alpha \int_{\mathbb{R}} \xi H(\overline{M}_\alpha, \xi, z_b) d\xi. \end{aligned} \quad (35)$$

**Proposition 3.8** *The solutions of the kinetic equation (32) are entropy solutions in the sense they satisfy on each layer the kinetic energy inequality*

$$\begin{aligned}
& \frac{\partial}{\partial t} H(\bar{M}_\alpha, \xi, z_b) + \xi \frac{\partial}{\partial x} H(\bar{M}_\alpha, \xi, z_b) - g \frac{\partial z_b}{\partial x} \frac{\partial}{\partial \xi} H(\bar{M}_\alpha, \xi, z_b) \\
& \leq \frac{G_{\alpha+1/2}}{h} H(\bar{M}_{\alpha+1/2}, \xi, z_b) - \frac{G_{\alpha-1/2}}{h} H(\bar{M}_{\alpha-1/2}, \xi, z_b) \\
& \quad + \frac{g^2 \pi^2}{6} \frac{G_{\alpha+1/2}}{h} (\bar{M}_{\alpha+1/2} + 2\bar{M}_\alpha) (\bar{M}_{\alpha+1/2} - \bar{M}_\alpha)^2 \\
& \quad - \frac{g^2 \pi^2}{6} \frac{G_{\alpha-1/2}}{h} (\bar{M}_{\alpha-1/2} + 2\bar{M}_\alpha) (\bar{M}_{\alpha-1/2} - \bar{M}_\alpha)^2 \\
& \quad + \frac{g^2 \pi^2}{3} \frac{G_{\alpha+1/2} - G_\alpha}{h} \bar{M}_\alpha^3 + H'(\bar{M}_\alpha, \xi, z_b) Q_\alpha
\end{aligned} \tag{36}$$

Integration in  $\xi$  and sum on  $\alpha$  of relations (36) lead to a global macroscopic energy inequality, that is analog to (25)

$$\begin{aligned}
& \frac{\partial}{\partial t} \left( \sum_{\alpha=1}^N E_\alpha \right) + \frac{\partial}{\partial x} \left( \sum_{\alpha=1}^N u_\alpha \left( E_\alpha + \frac{gh}{2} h_\alpha \right) \right) \\
& \leq -\frac{g^2 \pi^2}{6} \sum_{\alpha=1}^{N-1} \frac{|G_{\alpha+1/2}|}{h} \int_{\mathbb{R}} (\bar{M}_{\alpha+1} + 2\bar{M}_\alpha) (\bar{M}_{\alpha+1} - \bar{M}_\alpha)^2 d\xi.
\end{aligned} \tag{37}$$

**Proof of Proposition 3.8** *Kinetic energy inequality (36) is obtained by multiplying the kinetic equation (32) by  $\partial_1 H(\bar{M}_\alpha, \xi, z_b)$ , where  $\partial_1$  denotes the derivative in the first variable. Indeed, it is easy to see that, remember the topography  $z_b$  does not depend on time,*

$$\partial_1 H(\bar{M}_\alpha, \xi, z_b) \frac{\partial M_\alpha}{\partial t} = \frac{\partial}{\partial t} H(\bar{M}_\alpha, \xi, z_b),$$

likewise we have

$$\xi \partial_1 H(\bar{M}_\alpha, \xi, z_b) \frac{\partial M_\alpha}{\partial x} = \xi \frac{\partial}{\partial x} H(\bar{M}_\alpha, \xi, z_b) - \xi \partial_3 H(\bar{M}_\alpha, \xi, z_b) \frac{\partial z_b}{\partial x},$$

and

$$-g \frac{\partial z_b}{\partial x} \partial_1 H(\bar{M}_\alpha, \xi, z_b) \frac{\partial M_\alpha}{\partial \xi} = -g \frac{\partial z_b}{\partial x} \frac{\partial}{\partial \xi} H(\bar{M}_\alpha, \xi, z_b) + g \frac{\partial z_b}{\partial x} \partial_2 H(\bar{M}_\alpha, \xi, z_b),$$

But it follows from definition (34) of the kinetic energy that

$$g \frac{\partial z_b}{\partial x} \partial_2 H(\bar{M}_\alpha, \xi, z_b) - \xi \partial_3 H(\bar{M}_\alpha, \xi, z_b) \frac{\partial z_b}{\partial x} = 0,$$

It remains to obtain a suitable expression for the quantity

$$\partial_1 H(\bar{M}_\alpha, \xi, z_b) N_{\alpha+1/2}.$$

Let's denote

$$N_\alpha = \frac{1}{2}(N_{\alpha+1/2} + N_{\alpha-1/2}), \quad G_\alpha = \frac{1}{2}(G_{\alpha+1/2} + G_{\alpha-1/2}).$$

It follows from (30) that

$$\begin{aligned} \partial_1 H(\bar{M}_\alpha, \xi, z_b)(N_{\alpha+1/2} - N_\alpha) &= \frac{G_{\alpha+1/2}}{h} \partial_1 H(\bar{M}_\alpha, \xi, z_b)(\bar{M}_{\alpha+1/2} - \bar{M}_\alpha) \\ &\quad + \frac{G_{\alpha+1/2} - G_\alpha}{h} \partial_1 H(\bar{M}_\alpha, \xi, z_b) \bar{M}_\alpha, \end{aligned} \quad (38)$$

and the definition (34) of  $H_\alpha(\bar{M}_\alpha, \xi, z_b)$  gives us

$$\begin{aligned} H(\bar{M}_{\alpha+1/2}, \xi, z_b) &= H(\bar{M}_\alpha, \xi, z_b) + \partial_1 H(\bar{M}_\alpha, \xi, z_b)(\bar{M}_{\alpha+1/2} - \bar{M}_\alpha) \\ &\quad + \frac{g^2 \pi^2}{6} (\bar{M}_{\alpha+1/2} + 2\bar{M}_\alpha)(\bar{M}_{\alpha+1/2} - \bar{M}_\alpha)^2. \end{aligned} \quad (39)$$

Relation (38) together with (39) leads to

$$\begin{aligned} \partial_1 H(\bar{M}_\alpha, \xi, z_b)(N_{\alpha+1/2} - N_\alpha) &= \frac{G_{\alpha+1/2}}{h} H(\bar{M}_{\alpha+1/2}, \xi, z_b) - \frac{G_\alpha}{h} H(\bar{M}_\alpha, \xi, z_b) \\ &\quad - \frac{g^2 \pi^2}{6} \frac{G_{\alpha+1/2}}{h} (\bar{M}_{\alpha+1/2} + 2\bar{M}_\alpha)(\bar{M}_{\alpha+1/2} - \bar{M}_\alpha)^2 + \frac{g^2 \pi^2}{3} \frac{G_{\alpha+1/2} - G_\alpha}{h} \bar{M}_\alpha^3, \end{aligned}$$

and the same kind of relation occurs for  $H'_\alpha(\bar{M}_\alpha, \xi, z_b)(N_{\alpha-1/2} - N_\alpha)$ . Kinetic energy inequality per layer (36) obviously follows.

In order to derive the global energy inequality (37), let's now detail the right hand side of relation (36), the left hand side being treated using integral relations (35). The first line involved vertical kinetic exchange terms that vanish when summing on the layers. Second and third lines involved nonpositive terms due to the upwind definition (30) of the interface Maxwellian  $\bar{M}_{\alpha+1/2}$ . They will lead to the nonpositive right hand side in relation (37). Finally, the terms in the last line have no particular sign at the kinetic level. But after integration in  $\xi$ , one observes that

$$\int_{\mathbb{R}} \frac{g^2 \pi^2}{3} \frac{(G_{\alpha+1/2} - G_{\alpha-1/2})}{h} \bar{M}_\alpha^3 d\xi = \frac{g}{2} h (G_{\alpha+1/2} - G_{\alpha-1/2})$$

that can be interpreted as a macroscopic vertical exchange term and then vanishes when summing on the layers. Moreover, due to the particular choice (27) for the kinetic maxwellian  $\bar{M}_\alpha$ , one has

$$\partial_1 H(\bar{M}_\alpha, \xi, z_b) = \frac{\xi^2}{2} + \frac{g^2 \pi^2}{2} \bar{M}_\alpha^2 + g z_b = \frac{\xi^2}{2} + g \eta - \frac{(\xi - u_\alpha)^2}{2} = -\frac{u_\alpha^2}{2} + g \eta + u_\alpha \xi \quad (40)$$

and it hence follows from integral relations (33) on the collision term that the last term on the right hand side of (36) vanishes when integrating in  $\xi$ . ■

We end this section with a last result that extends to the present layer-averaged framework a subdifferential inequality and an energy minimization principle that were exhibited in the classical shallow water framework in [2], see also [6] for the first use of this approach. This result will be used in the next section to extend the entropy inequality to the fully discrete case.

**Lemma 3.9** (i) For any  $h_\alpha \geq 0$ ,  $u_\alpha \in \mathbb{R}$ ,  $f \geq 0$  and  $\xi \in \mathbb{R}$

$$H(f, \xi, z_b) \geq H(\overline{M}_\alpha, \xi, z_b) + E'_\alpha(U_\alpha) \left( \frac{1}{\xi} \right) (f - \overline{M}_\alpha). \quad (41)$$

(ii) For any function  $f(\xi)$  nonnegative, setting  $h_\alpha = \int_{\mathbb{R}} f(\xi) d\xi$  and  $h_\alpha u_\alpha = \int_{\mathbb{R}} \xi f(\xi) d\xi$  (assumed finite), one has

$$E_\alpha = l_\alpha \int_{\mathbb{R}} H(\overline{M}_\alpha, \xi, z_b) d\xi \leq l_\alpha \int_{\mathbb{R}} H(f(\xi), \xi, z_b) d\xi.$$

**Proof** The property (ii) follows from (i) by taking  $f = f(\xi)$  and integrating (41) with respect to  $\xi$  since  $f$  and  $\overline{M}_\alpha$  share the same first two moments. For proving (i), we first notice that (remember that  $h_\alpha = l_\alpha h$ )

$$E'_\alpha(U_\alpha) = (gh + gz_b - u_\alpha^2/2, u_\alpha),$$

and then

$$E'_\alpha(U_\alpha) \left( \frac{1}{\xi} \right) = gh + gz_b - u_\alpha^2/2 + \xi u_\alpha = \frac{\xi^2}{2} + gh + gz_b - \frac{(\xi - u_\alpha)^2}{2}.$$

Now the definition (27) of the maxwellian  $\overline{M}_\alpha$  yields

$$gh - \frac{(\xi - u_\alpha)^2}{2} = \begin{cases} \frac{g^2 \pi^2}{2} \overline{M}_\alpha^2 & \text{if } \overline{M}_\alpha > 0, \\ \text{is nonpositive} & \text{if } \overline{M}_\alpha = 0, \end{cases}$$

Using relation (40), it follows that

$$\partial_1 H(\overline{M}_\alpha, \xi, z_b) = \begin{cases} E'_\alpha \left( \frac{1}{\xi} \right) & \text{if } \overline{M}_\alpha > 0, \\ \geq E'_\alpha \left( \frac{1}{\xi} \right) & \text{if } \overline{M}_\alpha = 0. \end{cases} \quad (42)$$

We conclude using the convexity of  $H_\alpha$  with respect to  $f$ , see definition (34),

$$\begin{aligned} H(f, \xi, z_b) &\geq H(\overline{M}_\alpha, \xi, z_b) + \partial_1 H(\overline{M}_\alpha, \xi, z_b) (f - \overline{M}_\alpha) \\ &\geq H(\overline{M}_\alpha, \xi, z_b) + E'_\alpha \left( \frac{1}{\xi} \right) (f - \overline{M}_\alpha), \end{aligned} \quad (43)$$

which proves the claim. ■

### 3.3.2 Discrete model

The method proposed in [2] in order to solve the Saint-Venant system from its kinetic interpretation can be extended to the system (17)-(18) and its kinetic interpretation given in Proposition 3.7. It is the purpose of this paragraph.

We would like to approximate the solution  $U(t, x)$ , see (26),  $x \in \mathbb{R}$ ,  $t \geq 0$  of the system (17)-(18) by discrete values  $U_i^n$ ,  $i \in \mathbb{Z}$ ,  $n \in \mathbb{N}$ . In order to do so, we consider a grid of points  $x_{i+1/2}$ ,  $i \in \mathbb{Z}$ ,

$$\dots < x_{i-1/2} < x_{i+1/2} < x_{i+3/2} < \dots,$$

and we define the cells (or finite volumes) and their lengths

$$C_i = ]x_{i-1/2}, x_{i+1/2}[ , \quad \Delta x_i = x_{i+1/2} - x_{i-1/2}.$$

We consider discrete times  $t^n$  with  $t^{n+1} = t^n + \Delta t^n$ , and we define the piecewise constant functions  $U^n(x)$  corresponding to time  $t^n$  and  $z(x)$  as

$$U^n(x) = U_i^n, \quad z_b(x) = z_{b,i}, \quad \text{for } x_{i-1/2} < x < x_{i+1/2}.$$

A finite volume scheme for solving (17)-(18) is a formula of the form

$$U_i^{n+1} = U_i^n - \sigma_i(F_{i+1/2-} - F_{i-1/2+}) + \Delta t^n S_i, \quad (44)$$

where  $\sigma_i = \Delta t^n / \Delta x_i$ , telling how to compute the values  $U_i^{n+1}$  knowing  $U_i^n$  and discretized values  $z_{b,i}$  of the topography. The quantity  $S_i$  is a source terms accounting for the discrete momentum exchange terms between each layer in (18). Here we consider first-order explicit three points schemes where

$$F_{i+1/2-} = \mathcal{F}_l(U_i^n, U_{i+1}^n, z_{i+1} - z_i), \quad F_{i+1/2+} = \mathcal{F}_r(U_i^n, U_{i+1}^n, z_{i+1} - z_i).$$

The functions  $\mathcal{F}_{l/r}(U_l, U_r, \Delta z) \in \mathbb{R}^2$  are the numerical fluxes, see [7].

The proposed discrete scheme is based on the equivalence between the kinetic and the macroscopic levels stated in Proposition 3.7 and can be divided into three steps

- To construct the discrete kinetic maxwellian  $M_{\alpha,i}^n$  starting from the macroscopic quantities  $U_i^n$  and definition (27)

$$M_{\alpha,i}^n(\xi) = \frac{l_\alpha}{g\pi} (2gh_{\alpha,i}^n - (\xi-))_+^{1/2} \quad (45)$$

- To update the kinetic quantities through a finite volume scheme that will be precised hereafter to compute the quantities  $M_{\alpha,i}^{n+1,-}$ . This step is in general performed without considering the collision term and it follows that  $M_{\alpha,i}^{n+1,-}$  is no more of a maxwellian.



- To compute the new macroscopic quantities  $U_i^{n+1}$  as the integral of the kinetic quantities  $M_{\alpha,i}^{n+1,-}$

$$U_{\alpha,i}^{n+1} = \int_{\mathbb{R}} \begin{pmatrix} 1 \\ \xi \end{pmatrix} M_{\alpha,i}^{n+1,-}(\xi) d\xi.$$

The difference between  $M_{\alpha,i}^{n+1,-}$  and  $M_{\alpha,i}^{n+1}$ , computed using (45) at time  $t^{n+1}$ , can be seen as an instantaneous relaxation on the maxwellian, see [7].

Such a kind of kinetic scheme was presented for the classical shallow water system in [3], see also [7], and for the layer-averaged system (17)-(18) in [4]. In these works, the kinetic step was fully explicit. Here, and in order to demonstrate stability properties, we propose a implicit-explicit variant that is presented in details in the next section. Note that for practical computations, the integration processes is not performed on the cell unknowns but directly to compute macroscopic fluxes, as it will be explained hereafter. It follows the presented scheme can be entirely written at the macroscopic level, avoiding expensive computations at the kinetic level. Nevertheless, the kinetic interpretation is an efficient way to demonstrate the properties of the scheme.

### 3.3.3 Discrete kinetic equation

Let us now detail the kinetic scheme we propose. It can be written in a one-step version

$$M_{\alpha,i}^{n+1-} = M_{\alpha,i} - \sigma_i \left( \xi M_{\alpha,i+1/2} + \delta M_{\alpha,i+1/2-} - \xi M_{\alpha,i-1/2} - \delta M_{\alpha,i-1/2+} \right) + \Delta t^n (N_{\alpha+1/2,i}^{n+1-} - N_{\alpha-1/2,i}^{n+1-}) \quad (46)$$

or divided into an explicit and an implicit steps

$$M_{\alpha,i}^{n*} = M_{\alpha,i} - \sigma_i \left( \xi M_{\alpha,i+1/2} + \delta M_{\alpha,i+1/2-} - \xi M_{\alpha,i-1/2} - \delta M_{\alpha,i-1/2+} \right) \quad (47)$$

$$M_{\alpha,i}^{n+1-} = M_{\alpha,i}^{n*} + \Delta t^n (N_{\alpha+1/2,i}^{n+1-} - N_{\alpha-1/2,i}^{n+1-}) \quad (48)$$

with  $\sigma_i = \Delta t^n / \Delta x_i$ . To simplify the notations, we omit the variable  $\xi$  and the superscript  $n$ . The quantities  $M_{\alpha,i\pm 1/2}$  and  $\delta M_{\alpha,i+1/2\pm}$  respectively account for the conservative part and the topography source term, their definitions will be precised later. The definition of the quantity  $N_{\alpha+1/2,i}^{n+1-}$  requires a discrete extension of relation (30)

$$N_{\alpha+1/2,i}^{n+1-} = \frac{G_{\alpha+1/2,i}^{n*}}{h_i^{n*}} \overline{M}_{\alpha+1/2,i}^{n+1-}, \quad \overline{M}_{\alpha+1/2,i}^{n+1-} = \begin{cases} \overline{M}_{\alpha,i}^{n+1-} & \text{if } G_{\alpha+1/2,i}^{n*} \leq 0 \\ \overline{M}_{\alpha+1}^{n+1-} & \text{if } G_{\alpha+1/2,i}^{n*} \geq 0 \end{cases} \quad (49)$$

where the discrete mass exchange term  $G_{\alpha+1/2,i}$  is computed using a discrete version of relation (23)

$$\Delta x_i G_{\alpha+1/2,i} = \sum_{j=1}^{\alpha} \left( \int_{\mathbb{R}} \xi (M_{j,i+1/2} - M_{j,i-1/2}) d\xi \right)$$

$$-l_j \sum_{p=1}^N \int_{\mathbb{R}} \xi (M_{p,i+1/2} - M_{p,i-1/2}) d\xi \Big). \quad (50)$$

The explicit step (47) is very similar to the kinetic scheme proposed in [3] and analysed in [2] for the classical shallow water problem. We first prove hereafter that the implicit step (48) leads to a well posed problem. Then, in the next section, we prove the stability properties of the whole scheme (46). Using (49) and (50), the implicit step (48) can be written

$$-\Delta t^n \frac{|G_{\alpha+1/2,i}|_+}{h_{\alpha+1,i}^{n+1}} M_{\alpha+1,i}^{n+1-} + \left( 1 - \Delta t^n \frac{|G_{\alpha+1/2,i}|_- - |G_{\alpha-1/2,i}|_+}{h_{\alpha,i}^{n+1}} \right) M_{\alpha,i}^{n+1-} + \Delta t^n \frac{|G_{\alpha-1/2,i}|_-}{h_{\alpha-1,i}^{n+1}} M_{\alpha-1,i}^{n+1-} = M_{\alpha,i}^{n*} \quad (51)$$

that is equivalent to solve the linear system

$$(\mathbf{I}_N + \Delta t \mathbf{G}_{N,i}) M_i^{n+1-} = M_i^{n*}, \quad M_i^k = (M_{1,i}^k, \dots, M_{N,i}^k)^T \quad (52)$$

with

$$G_{N,i} = \begin{pmatrix} -\frac{|G_{3/2,i}|_-}{h_{1,i}^{n+1}} & -\frac{|G_{3/2,i}|_+}{h_{1,i}^{n+1}} & 0 & 0 & \dots & 0 \\ \frac{|G_{3/2,i}|_-}{h_{2,i}^{n+1}} & \ddots & \ddots & 0 & \dots & 0 \\ 0 & \ddots & \ddots & \ddots & 0 & 0 \\ \vdots & 0 & \frac{|G_{\alpha-1/2,i}|_-}{h_{\alpha,i}^{n+1}} & -\frac{|G_{\alpha+1/2,i}|_- - |G_{\alpha-1/2,i}|_+}{h_{\alpha,i}^{n+1}} & -\frac{|G_{\alpha+1/2,i}|_+}{h_{\alpha,i}^{n+1}} & 0 \\ \vdots & \ddots & 0 & \ddots & \ddots & -\frac{|G_{N-1/2,i}|_+}{h_{N-1,i}^{n+1}} \\ 0 & \dots & 0 & 0 & \frac{|G_{N-1/2,i}|_-}{h_{N,i}^{n+1}} & \frac{|G_{N-1/2,i}|_+}{h_{N,i}^{n+1}} \end{pmatrix}.$$

**Lemma 3.10** *The matrix  $\mathbf{A}_{N,i} = \mathbf{I}_N + \Delta t \mathbf{G}_{N,i}$  satisfies the following properties*

- (i) *The matrix  $\mathbf{A}_{N,i}$  is invertible for any  $h_i^{n+1} > 0$  and then the linear system (52) has a unique solution.*
- (ii) *Its inverse  $\mathbf{A}_{N,i}^{-1}$  has only positive coefficients and then the kinetic density  $M_i^{n+1-}$  is positive if  $M_i^{n*}$  is.*
- (iii) *For any vector  $T$  with non negative entries i.e.  $T_\alpha \geq 0$ , for  $1 \leq \alpha \leq N$ , one has*

$$\|\mathbf{A}_{N,i}^{-1} T\|_\infty \leq \|T\|_\infty.$$

*and then the solution of the linear system (52) does not raise difficulties for any  $h_i > 0$  even if  $h_i$  is arbitrarily small.*

**Proof of lemma 3.10** (i) For any  $h_i^{n+1} > 0$ , the matrix  $\mathbf{A}_{N,i}^T$  is a strictly dominant diagonal matrix. It follows that  $\mathbf{A}_{N,i}$  is invertible.

(ii) Denoting  $\mathbf{G}_{N,i}^d$  (resp.  $\mathbf{G}_{N,i}^{nd}$ ) the diagonal (resp. non diagonal) part of  $\mathbf{G}_{N,i}$  we can write

$$\mathbf{A}_{N,i} = \mathbf{I}_N + \Delta t \mathbf{G}_{N,i} = (\mathbf{I}_N + \Delta t \mathbf{G}_{N,i}^d) (\mathbf{I}_N - (\mathbf{I}_N + \Delta t \mathbf{G}_{N,i}^d)^{-1} (-\Delta t \mathbf{G}_{N,i}^{nd})),$$

where all the entries of the matrix

$$\mathbf{J}_{N,i} = (\mathbf{I}_N + \Delta t \mathbf{G}_{N,i}^d)^{-1} (-\Delta t \mathbf{G}_{N,i}^{nd}),$$

are non negative and less than 1. And hence, we can write

$$(\mathbf{I}_N + \Delta t \mathbf{G}_{N,i})^{-1} = \sum_{k=0}^{\infty} \mathbf{J}_{N,i}^k,$$

proving all the entries of  $(\mathbf{I}_N + \Delta t \mathbf{G}_{N,i})^{-1}$  are non negative.

(iii) Let us consider the vector  $\mathbf{1}$  whose entries are all equal to 1. Since we have

$$(\mathbf{I}_N + \Delta t \mathbf{G}_{N,i})^t \mathbf{1} = \mathbf{1},$$

we also have

$$\mathbf{1} = (\mathbf{I}_N + \Delta t \mathbf{G}_{N,i})^{-t} \mathbf{1}.$$

Now let  $\mathbf{T}$  be a vector whose entries  $\{T_\alpha\}_{1 \leq \alpha \leq N}$  are non negative, then

$$(\mathbf{I}_N + \Delta t \mathbf{G}_{N,i})^{-t} \mathbf{T} \leq (\mathbf{I}_N + \Delta t \mathbf{G}_{N,i})^{-t} \mathbf{1} \|\mathbf{T}\|_\infty = \mathbf{1} \|\mathbf{T}\|_\infty,$$

that completes the proof. ■

## 4 Properties of the scheme

In this section, we examine the properties of the scheme (46).

### 4.1 Without topography

We first consider the problem without topography. The scheme (46) reduces to

$$\begin{cases} M_{\alpha,i}^{n*} = M_{\alpha,i} - \sigma_i \xi (M_{\alpha,i+1/2} - M_{\alpha,i-1/2}) \\ M_{\alpha,i}^{n+1-} = M_{\alpha,i}^{n*} + \Delta t^n (N_{\alpha+1/2,i}^{n+1-} - N_{\alpha-1/2,i}^{n+1-}) \end{cases} \quad (53)$$

with

$$M_{\alpha,i+1/2} = \mathbf{1}_{\xi>0} M_{\alpha,i} + \mathbf{1}_{\xi<0} M_{\alpha,i+1}, \quad (54)$$

$$M_{\alpha,i-1/2} = \mathbf{1}_{\xi>0}M_{\alpha,i-1} + \mathbf{1}_{\xi<0}M_{\alpha,i}. \quad (55)$$

The discrete kinetic equations (53) allow to precise the numerical fluxes in (44) having the form

$$F_{i+1/2-} = \left( \sum_{\alpha=1}^N F_{h_{\alpha},i+1/2-}, F_{q_1,i+1/2-}, \dots, F_{q_N,i+1/2-} \right)^T, \quad (56)$$

with

$$\begin{aligned} F_{h_{\alpha},i+1/2-} &= \int_{\mathbb{R}} \xi M_{\alpha,i+1/2} d\xi = \int_{\xi>0} \xi M_{\alpha,i} d\xi + \int_{\xi<0} \xi M_{\alpha,i+1} d\xi, \\ F_{q_{\alpha},i+1/2-} &= \int_{\mathbb{R}} \xi^2 M_{\alpha,i+1/2} d\xi + \Delta x_i u_{\alpha+1/2,i} G_{\alpha+1/2,i} \\ &= \int_{\xi>0} \xi^2 M_{\alpha,i} d\xi + \int_{\xi<0} \xi^2 M_{\alpha,i+1} d\xi + \Delta x_i u_{\alpha+1/2,i} G_{\alpha+1/2,i}. \end{aligned}$$

The source  $S_i$  in (44) is defined by

$$S_i = u_{\alpha+1/2,i} G_{\alpha+1/2,i} - u_{\alpha-1/2,i} G_{\alpha-1/2,i},$$

where  $u_{\alpha+1/2,i}$  is defined by (21) and  $G_{\alpha+1/2,i}$  is given by expression (50) that can be rewritten under the form

$$\Delta x_i G_{\alpha+1/2,i} = \sum_{j=1}^{\alpha} \left( F_{h_j,i+1/2-} - F_{h_j,i-1/2+} - l_j \sum_{p=1}^N (F_{h_p,i+1/2-} - F_{h_p,i-1/2+}) \right).$$

In the following proposition, we prove fundamental stability properties for the numerical scheme (46).

**Theorem 4.1** *Under the CFL condition*

$$\Delta t^n \leq \frac{1}{2} \min_{1 \leq \alpha \leq N} \min_{i \in I} \frac{\Delta x_i}{|u_{\alpha,i}| + \sqrt{2gh_i}}. \quad (57)$$

the scheme (46) satisfies the following properties

- (i) The kinetic functions remain nonnegative  $M_{\alpha,i}^{n+1-} \geq 0, \forall \alpha, i$
- (ii) One has the kinetic energy equality

$$\begin{aligned} H(\overline{M}_{\alpha,i}^{n+1-}, \xi, z_{b,i}) &= H(\overline{M}_{\alpha,i}, \xi, z_{b,i}) - \sigma_i \left( \tilde{H}_{\alpha,i+1/2} - \tilde{H}_{\alpha,i-1/2} \right) \\ &\quad - \Delta t^n \left( \hat{H}_{\alpha+1/2,i}^{n+1-} - \hat{H}_{\alpha-1/2,i}^{n+1-} \right) + d_{\alpha,i} + e_{\alpha,i}, \end{aligned} \quad (58)$$

where  $\tilde{H}_{\alpha,i\pm 1/2}, \hat{H}_{\alpha\pm 1/2,i}$  are defined by

$$\tilde{H}_{\alpha,i\pm 1/2} = \xi H(\overline{M}_{\alpha,i\pm 1/2}, \xi, z_{b,i\pm 1/2}) - \xi H(\overline{M}_{\alpha,i}, \xi, z_{b,i}),$$

$$\widehat{H}_{\alpha\pm 1/2,i}^{n+1-} = \frac{G_{\alpha\pm 1/2,i}}{h_i^{n+1}} H(\overline{M}_{\alpha\pm 1/2,i}^{n+1-}, \xi, z_{b,i}),$$

and  $d_{\alpha,i}, e_{\alpha,i}$  are given by

$$\begin{aligned} d_{\alpha,i} &= \frac{g^2 \pi^2}{6} \sigma_i \xi \left( \overline{M}_{\alpha,i+1} + 2\overline{M}_{\alpha,i} + \sigma_i \xi (\overline{M}_{\alpha,i}^{n*} + 2\overline{M}_{\alpha,i}) \right) (\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i})^2 \\ &\quad - \frac{g^2 \pi^2}{6} \sigma_i \xi \left( \overline{M}_{\alpha,i-1} + 2\overline{M}_{\alpha,i} + \sigma_i \xi (\overline{M}_{\alpha,i}^{n*} + 2\overline{M}_{\alpha,i}) \right) (\overline{M}_{\alpha,i-1/2} - \overline{M}_{\alpha,i})^2 \\ &\quad - \frac{g^2 \pi^2}{6} \Delta t^n \frac{G_{\alpha+1/2,i}}{h_i^{n+1}} (\overline{M}_{\alpha+1/2,i}^{n+1-} + 2\overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha+1/2,i}^{n+1-} - \overline{M}_{\alpha,i}^{n+1-})^2 \\ &\quad + \frac{g^2 \pi^2}{6} \Delta t^n \frac{G_{\alpha-1/2,i}}{h_i^{n+1}} (\overline{M}_{\alpha-1/2,i}^{n+1-} + 2\overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha-1/2,i}^{n+1-} - \overline{M}_{\alpha,i}^{n+1-})^2, \\ e_{\alpha,i} &= \frac{g^2 \pi^2}{3} \Delta t^n \frac{(G_{\alpha+1/2,i} - G_{\alpha-1/2,i})}{h_i^{n+1}} (\overline{M}_{\alpha,i}^{n+1-})^3. \end{aligned}$$

**Corollary 4.2** *Under the CFL condition (57), one has the macroscopic energy inequality*

$$\sum_{\alpha=1}^N \overline{E}_{\alpha,i}^{n+1} \leq \sum_{\alpha=1}^N \overline{E}_{\alpha,i} - \sigma_i \left( \sum_{\alpha=1}^N \int_{\mathbb{R}} \tilde{H}_{i+1/2} d\xi - \sum_{\alpha=1}^N \int_{\mathbb{R}} \tilde{H}_{i-1/2} d\xi \right).$$

**Remark 4.3** *Even if we considered the system without topography, we keep the notations  $z_{b,j}$ ,  $j = i, i + 1/2, i - 1/2, \dots$  so that the obtained formula can be easily extended to the case of a non flat topography.*

**Remark 4.4** *We have seen in remark 3.6 that, at the continuous level, the quantities*

$$\frac{G_{\alpha\pm 1/2}}{h},$$

*remain bounded when  $h \rightarrow 0$ . At the discrete level, the quantities  $N_{\alpha\pm 1/2,i}^{n+1-}$  and  $\widehat{H}_{\alpha\pm 1/2,i}^{n+1-}$  involve the discrete mass exchange terms*

$$\frac{G_{\alpha\pm 1/2,i}}{h_i^{n+1}}.$$

*For any  $h_i^{n+1}$  and even if  $h_i^{n+1} \rightarrow 0$ , from Eq. (53), we have*

$$\Delta t^n N_{3/2,i}^{n+1-} = M_{1,i}^{n+1-} - M_{1,i}^{n*},$$

*and hence after integration in  $\xi$  we get*

$$\Delta t^n l_1 G_{3/2,i} = l_1 h_i^{n+1} - h_{1,i}^{n+1/2},$$

and since  $\sum_{j=1}^N h_{j,i}^{n+1/2} = h_i^{n+1}$ , we have  $h_{1,i}^{n+1/2} \leq h_i^{n+1}$  and this gives us the estimate

$$1 - \frac{1}{l_1} \leq \Delta t^n \frac{G_{3/2,i}}{h_i^{n+1}} \leq 1.$$

Using the process, one can prove

$$\frac{G_{\alpha+1/2,i}}{h_i^{n+1}},$$

is bounded for any  $\alpha$ .

**Proof of theorem 4.1** (i) The scheme (51) also writes

$$\begin{aligned} - \Delta t^n \frac{|G_{\alpha+1/2,i}|_+}{h_{\alpha+1,i}^{n+1}} M_{\alpha+1,i}^{n+1-} + \left( 1 - \Delta t^n \frac{|G_{\alpha+1/2,i}|_- - |G_{\alpha-1/2,i}|_+}{h_{\alpha,i}^{n+1}} \right) M_{\alpha,i}^{n+1-} \\ + \Delta t^n \frac{|G_{\alpha-1/2,i}|_-}{h_{\alpha-1,i}^{n+1}} M_{\alpha-1,i}^{n+1-} = M_{\alpha,i} - \sigma_i(\xi M_{\alpha,i+1/2} - M_{\alpha,i-1/2}), \end{aligned}$$

leading to

$$\begin{aligned} - \Delta t^n \frac{|G_{\alpha+1/2,i}|_+}{h_{\alpha+1,i}^{n+1}} M_{\alpha+1,i}^{n+1-} + \left( 1 - \Delta t^n \frac{|G_{\alpha+1/2,i}|_- - |G_{\alpha-1/2,i}|_+}{h_{\alpha,i}^{n+1}} \right) M_{\alpha,i}^{n+1-} \\ + \Delta t^n \frac{|G_{\alpha-1/2,i}|_-}{h_{\alpha-1,i}^{n+1}} M_{\alpha-1,i}^{n+1-} \geq (1 - \sigma_i(|\xi|)) M_{\alpha,i}. \quad (59) \end{aligned}$$

Since  $M_{\alpha,i} \geq 0$  and (57) holds, using the property (ii) of lemma 3.10, the inequality (59) proves (i).

(ii) The proof of Theorem 3.6 in [2] and the proof of the inequality (58) shares common points, namely the linear dissipation of the scheme is, in both cases, based on the convexity of the kinetic entropy (34) and the form of the Maxwellian (27). But the proof of inequality (58) is more complex because of the momentum exchange terms along the vertical axis and their implicit treatment. Notice that, compared to Theorem 3.6 in [2], the derivation of the horizontal linear dissipation is obtained in a different way.

In order to prove (58) we multiply the first equation of (51) by  $H'(\bar{M}_{\alpha,i}, \xi, z_{b,i})$  and the second equation of (51) by  $H'(\bar{M}_{\alpha,i}^{n+1-}, \xi, z_{b,i})$ .

First, using the identity

$$\bar{M}_{\alpha,i+1}^3 = \bar{M}_{\alpha,i}^3 + 3\bar{M}_{\alpha,i}^2(\bar{M}_{\alpha,i+1} - \bar{M}_{\alpha,i}) + (\bar{M}_{\alpha,i+1} + 2\bar{M}_{\alpha,i})(\bar{M}_{\alpha,i+1} - \bar{M}_{\alpha,i})^2,$$

we obtain an expression of the linear dissipation

$$\begin{aligned} H_\alpha(\bar{M}_{\alpha,i+1}, \xi, z_{b,i+1}) - H_\alpha(\bar{M}_{\alpha,i}, \xi, z_{b,i}) - H'_\alpha(\bar{M}_{\alpha,i}, \xi, z_{b,i})(\bar{M}_{\alpha,i+1} - \bar{M}_{\alpha,i}) \\ = \frac{g^2 \pi^2}{6} (\bar{M}_{\alpha,i+1} + 2\bar{M}_{\alpha,i})(\bar{M}_{\alpha,i+1} - \bar{M}_{\alpha,i})^2 \end{aligned}$$

that can also be written using the definitions (54)-(55) under the form

$$H_\alpha(\overline{M}_{\alpha,i+1/2}, \xi, z_{b,i+1/2}) = H_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i}) + H'_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i})(\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i}) + \frac{g^2\pi^2}{6}(\overline{M}_{\alpha,i+1} + 2\overline{M}_{\alpha,i})(\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i})^2. \quad (60)$$

Notice that at the interface  $i - 1/2$ , we obtain an expression similar to (60) having the form

$$H_\alpha(\overline{M}_{\alpha,i-1/2}, \xi, z_{b,i-1/2}) = H_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i}) + H'_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i})(\overline{M}_{\alpha,i-1/2} - \overline{M}_{\alpha,i}) + \frac{g^2\pi^2}{6}(\overline{M}_{\alpha,i-1} + 2\overline{M}_{\alpha,i})(\overline{M}_{\alpha,i-1/2} - \overline{M}_{\alpha,i})^2. \quad (61)$$

For  $\xi \leq 0$ , the first equation of (51) writes

$$M_{\alpha,i}^{n*} = M_{\alpha,i} - \sigma_i \xi (M_{\alpha,i+1} - M_{\alpha,i}), \quad (62)$$

whereas for  $\xi \geq 0$ , we have

$$M_{\alpha,i}^{n*} = M_{\alpha,i} - \sigma_i \xi (M_{\alpha,i} - M_{\alpha,i-1}). \quad (63)$$

Now let us multiply Eqs. (62),(63) by  $H'_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i})$  and using the expressions (60),(61) for the quantities

$$\begin{aligned} H'_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i})(M_{\alpha,i+1} - M_{\alpha,i}) &= \frac{1}{l_\alpha} H'_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i})(\overline{M}_{\alpha,i+1} - \overline{M}_{\alpha,i}), \\ H'_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i})(M_{\alpha,i} - M_{\alpha,i-1}) &= \frac{1}{l_\alpha} H'_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i})(\overline{M}_{\alpha,i} - \overline{M}_{\alpha,i-1}), \end{aligned}$$

we obtain the relation

$$\begin{aligned} l_\alpha H'_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i})(M_{\alpha,i}^{n*} - M_{\alpha,i}) &= -\sigma_i (\tilde{H}_{\alpha,i+1/2} - \tilde{H}_{\alpha,i-1/2}) \\ &+ \frac{g^2\pi^2}{6} \sigma_i \xi (\overline{M}_{\alpha,i+1} + 2\overline{M}_{\alpha,i})(\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i})^2 \\ &- \frac{g^2\pi^2}{6} \sigma_i \xi (\overline{M}_{\alpha,i-1} + 2\overline{M}_{\alpha,i})(\overline{M}_{\alpha,i-1/2} - \overline{M}_{\alpha,i})^2 \end{aligned} \quad (64)$$

The identities

$$\begin{aligned} H'(f, \xi, z_b) f &= H(f, \xi, z_b) + \frac{\pi^2 g^2}{3} f^3, \\ H'(f, \xi, z_b) \tilde{f} &= H(\tilde{f}, \xi, z_b) + \frac{\pi^2 g^2}{2} (f^2 - \tilde{f}^2) \tilde{f} + \frac{\pi^2 g^2}{3} \tilde{f}^3, \end{aligned}$$

help us to write

$$H'_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i})(M_{\alpha,i}^{n*} - M_{\alpha,i}) = \frac{1}{l_\alpha} H'_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i})(\overline{M}_{\alpha,i}^{n*} - \overline{M}_{\alpha,i})$$

$$= \frac{1}{l_\alpha} H_\alpha(\overline{M}_{\alpha,i}^{n*}, \xi, z_{b,i}) - \frac{1}{l_\alpha} H_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i}) - L_\alpha,$$

where as in (27),  $\overline{M}_{\alpha,i}^{n*} = l_\alpha M_{\alpha,i}^{n*}$  with

$$L_\alpha = \frac{\pi^2 g^2}{6} (\overline{M}_{\alpha,i}^{n*} + 2\overline{M}_{\alpha,i}) (\overline{M}_{\alpha,i}^{n*} - \overline{M}_{\alpha,i})^2.$$

Notice that from relations (62)-(63) we have

$$L_\alpha = \frac{\pi^2 g^2}{6} (\sigma_i |\xi|)^2 (\overline{M}_{\alpha,i}^{n*} + 2\overline{M}_{\alpha,i}) ((\overline{M}_{\alpha,i+1} - \overline{M}_{\alpha,i})^2 \mathbf{1}_{\xi \leq 0} + (\overline{M}_{\alpha,i} - \overline{M}_{\alpha,i-1})^2 \mathbf{1}_{\xi \geq 0}).$$

Therefore we are able to rewrite (64) under the form

$$\begin{aligned} H_\alpha(\overline{M}_{\alpha,i}^{n*}, \xi, z_{b,i}) &= H_\alpha(\overline{M}_{\alpha,i}, \xi, z_{b,i}) - \sigma_i \left( \tilde{H}_{\alpha,i+1/2} - \tilde{H}_{\alpha,i-1/2} \right) \\ &\quad + l_\alpha \frac{g^2 \pi^2}{6} \sigma_i \xi (\overline{M}_{\alpha,i+1} + 2\overline{M}_{\alpha,i} + \sigma_i \xi (\overline{M}_{\alpha,i}^{n*} + 2\overline{M}_{\alpha,i})) (\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i})^2 \\ &\quad - l_\alpha \frac{g^2 \pi^2}{6} \sigma_i \xi (\overline{M}_{\alpha,i-1} + 2\overline{M}_{\alpha,i} + \sigma_i \xi (\overline{M}_{\alpha,i}^{n*} + 2\overline{M}_{\alpha,i})) (\overline{M}_{\alpha,i-1/2} - \overline{M}_{\alpha,i})^2 \end{aligned}$$

In a second step, we multiply the second equation of (53) by  $H'_\alpha(\overline{M}_{\alpha,i}^{n+1-}, \xi, z_{b,i})$ . The proof of proposition 3.8 gives us an expression for the quantity

$$H'_\alpha(\overline{M}_{\alpha,i}^{n+1-}, \xi, z_{b,i}) (N_{\alpha+1/2,i}^{n+1-} - N_{\alpha-1/2,i}^{n+1-}),$$

under the form

$$\begin{aligned} H'_\alpha(\overline{M}_{\alpha,i}^{n+1-}, \xi, z_{b,i}) (N_{\alpha+1/2,i}^{n+1-} - N_{\alpha-1/2,i}^{n+1-}) &= \frac{G_{\alpha+1/2,i}}{h_i^{n+1}} H_\alpha(\overline{M}_{\alpha+1/2,i}^{n+1-}, \xi, z_{b,i}) - \frac{G_{\alpha-1/2,i}}{h_i^{n+1}} H_\alpha(\overline{M}_{\alpha-1/2,i}^{n+1-}, \xi, z_{b,i}) \\ &\quad - l_\alpha \frac{g^2 \pi^2}{6} \frac{G_{\alpha+1/2,i}}{h_i^{n+1}} (\overline{M}_{\alpha+1/2,i}^{n+1-} + 2\overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha+1/2,i}^{n+1-} - \overline{M}_{\alpha,i}^{n+1-})^2 \\ &\quad + l_\alpha \frac{g^2 \pi^2}{6} \frac{G_{\alpha-1/2,i}}{h_i^{n+1}} (\overline{M}_{\alpha-1/2,i}^{n+1-} + 2\overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha-1/2,i}^{n+1-} - \overline{M}_{\alpha,i}^{n+1-})^2 \\ &\quad + l_\alpha \frac{g^2 \pi^2}{3} \frac{(G_{\alpha+1/2,i} - G_{\alpha-1/2,i})}{h_i^{n+1}} (\overline{M}_{\alpha,i}^{n+1-})^3. \end{aligned}$$

and we also have

$$\begin{aligned} H'_\alpha(\overline{M}_{\alpha,i}^{n+1-}, \xi, z_{b,i}) (\overline{M}_{\alpha,i}^{n+1-} - \overline{M}_{\alpha,i}^{n*}) &= H_\alpha(\overline{M}_{\alpha,i}^{n+1-}, \xi, z_{b,i}) - H_\alpha(\overline{M}_{\alpha,i}^{n*}, \xi, z_{b,i}) \\ &\quad - l_\alpha \frac{g^2 \pi^2}{6} (\overline{M}_{\alpha,i}^{n*} + 2\overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha,i}^{n+1-} - \overline{M}_{\alpha,i}^{n*})^2. \end{aligned}$$

Using the two previous expressions, we are able to rewrite the second equation of (53) multiplied by  $H'_\alpha(\overline{M}_{\alpha,i}^{n+1-}, \xi, z_{b,i})$  under the form

$$H_\alpha(\overline{M}_{\alpha,i}^{n+1-}, \xi, z_{b,i}) = H_\alpha(\overline{M}_{\alpha,i}^{n*}, \xi, z_{b,i}) - \sigma_i \left( \widehat{H}_{\alpha+1/2,i}^{n+1-} - \widehat{H}_{\alpha-1/2,i}^{n+1-} \right)$$



$$\begin{aligned}
& -l_\alpha \frac{g^2 \pi^2}{6} \Delta t^n \frac{G_{\alpha+1/2,i}}{h_i^{n+1}} (\overline{M}_{\alpha+1/2,i}^{n+1-} + 2\overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha+1/2,i}^{n+1-} - \overline{M}_{\alpha,i}^{n+1-})^2 \\
& + l_\alpha \frac{g^2 \pi^2}{6} \Delta t^n \frac{G_{\alpha-1/2,i}}{h_i^{n+1}} (\overline{M}_{\alpha-1/2,i}^{n+1-} + 2\overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha-1/2,i}^{n+1-} - \overline{M}_{\alpha,i}^{n+1-})^2 \\
& + l_\alpha \frac{g^2 \pi^2}{3} \Delta t^n \frac{(G_{\alpha+1/2,i} - G_{\alpha-1/2,i})}{h_i^{n+1}} (\overline{M}_{\alpha,i}^{n+1-})^3 \\
& - l_\alpha \frac{g^2 \pi^2}{6} (\overline{M}_{\alpha,i}^{n*} + 2\overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha,i}^{n+1-} - \overline{M}_{\alpha,i}^{n*})^2.
\end{aligned} \tag{66}$$

The sum of relations (65), (66) divided by  $l_\alpha$  gives the result. ■

**Proof of corollary 4.2** An integration in  $\xi$  of relation (58) and a sum of the obtained relation for  $\alpha = 1, \dots, N$  gives

$$\begin{aligned}
\sum_{\alpha=1}^N \overline{E}_{\alpha,i}^{n+1} &= \sum_{\alpha=1}^N \overline{E}_{\alpha,i} - \sigma_i \left( \sum_{\alpha=1}^N \int_{\mathbb{R}} \tilde{H}_{i+1/2} d\xi - \sum_{\alpha=1}^N \int_{\mathbb{R}} \tilde{H}_{i-1/2} d\xi \right) \\
&+ \sum_{\alpha=1}^N \int_{\mathbb{R}} d_{\alpha,i} d\xi + \sum_{\alpha=1}^N \int_{\mathbb{R}} e_{\alpha,i} d\xi.
\end{aligned}$$

Using (62), the first line of  $d_{\alpha,i}$  writes

$$\frac{g^2 \pi^2}{6} \sigma_i \xi \left( (1 - \sigma_i \xi) \overline{M}_{\alpha,i+1} + (2 + 3\sigma_i \xi - (\sigma_i \xi)^2) \overline{M}_{\alpha,i} \right) (\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i})^2,$$

that is non positive under the CFL condition (57) and likewise, we obtain that the second line of  $d_{\alpha,i}$  is non positive. Moreover, it is obvious from the definition of  $G_{\alpha+1/2,i}$  given by (21) that the other lines of  $d_i$  are also non positive.

Now we intend to prove that for any  $i$ , one has

$$\sum_{\alpha=1}^N \int_{\mathbb{R}} e_{\alpha,i} d\xi = \sum_{\alpha=1}^N \int_{\mathbb{R}} \frac{g^2 \pi^2}{3} \Delta t^n \frac{(G_{\alpha+1/2,i} - G_{\alpha-1/2,i})}{h_i^{n+1}} (\overline{M}_{\alpha,i}^{n+1-})^3 d\xi = 0. \tag{67}$$

Since  $\overline{M}_{\alpha,i}^{n+1-}$  is not a Maxwellian, it is not possible to conclude, as in the proof of Proposition 3.8, that

$$\int_{\mathbb{R}} (\overline{M}_{\alpha,i}^{n+1-})^3 d\xi = \frac{g}{2} h_{\alpha,i}^2,$$

and that the sum (67) is zero. So we proceed as follows.

Let us rewrite Eq. (46) under the equivalent form

$$\begin{cases} M_{\alpha,i}^{n+1/2} = M_{\alpha,i} - \sigma_i \left( \xi M_{\alpha,i+1/2} + \delta M_{\alpha,i+1/2-} - \xi M_{\alpha,i-1/2} - \delta M_{\alpha,i-1/2+} \right) + \Delta t^n Q_{\alpha,i} \\ M_{\alpha,i}^{n+1} = M_{\alpha,i}^{n+1/2} + \Delta t^n (N_{\alpha+1/2,i}^{n+1} - N_{\alpha-1/2,i}^{n+1}) + \Delta t^n Q_{\alpha,i}^{n+1/2} \end{cases} \tag{68}$$

where  $Q_{\alpha,i}$ ,  $Q_{\alpha,i}^{n+1/2}$  are two collision terms satisfying

$$\int_{\mathbb{R}} \begin{pmatrix} 1 \\ \xi \end{pmatrix} Q_{\alpha,i} d\xi = \int_{\mathbb{R}} \begin{pmatrix} 1 \\ \xi \end{pmatrix} Q_{\alpha,i}^{n+1/2} d\xi = \begin{pmatrix} 0 \\ 0 \end{pmatrix},$$

We are now able to reuse the arguments of the proof of Theorem 4.1-(ii) with the discretization (68) but since

$$H'_{\alpha}(\overline{M}_{\alpha,i}^{n+1}, \xi, z_{b,i}) = l_{\alpha} \left( g(h_i^{n+1} + z_{b,i}) + \frac{(u_{\alpha,i}^{n+1})^2}{2} + \xi u_{\alpha,i}^{n+1} \right),$$

we have

$$\int_{\mathbb{R}} H'_{\alpha}(\overline{M}_{\alpha,i}^{n+1}, \xi, z_{b,i}) Q_{\alpha,i}^{n+1/2} d\xi = 0,$$

and the integration in  $\xi$  of the second line of Eq. (68) multiplied by  $H'_{\alpha}(\overline{M}_{\alpha,i}^{n+1}, \xi, z_{b,i})$  gives exactly the relation (66) but where the superscript  $n^*$  (resp.  $n+1^-$ ) becomes  $n+1/2$  (resp.  $n+1$ ). And hence, we have

$$\begin{aligned} \sum_{\alpha=1}^N \int_{\mathbb{R}} e_{\alpha,i} d\xi &= \sum_{\alpha=1}^N \int_{\mathbb{R}} \frac{g^2 \pi^2}{3} \Delta t^n \frac{(G_{\alpha+1/2,i} - G_{\alpha-1/2,i})}{h_i^{n+1}} (\overline{M}_{\alpha,i}^{n+1-})^3 d\xi \\ &= \sum_{\alpha=1}^N \int_{\mathbb{R}} \frac{g^2 \pi^2}{3} \Delta t^n \frac{(G_{\alpha+1/2,i} - G_{\alpha-1/2,i})}{h_i^{n+1}} (\overline{M}_{\alpha,i}^{n+1})^3 d\xi = 0, \end{aligned}$$

that concludes the proof. ■

## 4.2 With topography

In this paragraph we examine the properties of the discrete scheme (46) when the topography source term is no more neglected.

The hydrostatic reconstruction scheme (HR scheme for short) is a general method giving, with any solver, a robust and efficient discretization of the source terms in conservation laws. It has been initially proposed for the Saint-Venant system leading to a consistent, well-balanced, positive scheme satisfying a semi-discrete entropy inequality [1]. Here we use the HR technique to discretize the topography source term appearing in (18) and we prove the kinetic scheme (46) coupled with the HR technique leads to a consistent, well-balanced, positive scheme satisfying a fully discrete entropy inequality.

The hydrostatic reconstruction scheme uses reconstructed states

$$U_{\alpha,i+1/2-} = (h_{\alpha,i+1/2-}, h_{\alpha,i+1/2-} u_i), \quad U_{\alpha,i+1/2+} = (h_{\alpha,i+1/2+}, h_{\alpha,i+1/2+} u_{i+1}), \quad (69)$$

defined by

$$\begin{aligned} h_{i+1/2-} &= (h_i + z_i - z_{b,i+1/2})_+, \\ h_{i+1/2+} &= (h_{i+1} + z_{b,i+1} - z_{b,i+1/2})_+, \\ h_{\alpha,i+1/2\pm} &= l_{\alpha} h_{i+1/2\pm}, \\ M_{\alpha,i+1/2\pm} &= M(U_{\alpha,i+1/2\pm}, \xi) \end{aligned} \quad (70)$$

and

$$z_{b,i+1/2} = \max(z_{b,i}, z_{b,i+1}). \quad (71)$$

We note that the definitions of  $h_{i+1/2\pm}$  in (70)-(71) ensure that  $h_{i+1/2-} \leq h_i$ , and  $h_{i+1/2+} \leq h_{i+1}$ . Therefore, because of (27) one has

$$0 \leq M_{\alpha,i+1/2-} \leq M_{\alpha,i}, \quad 0 \leq M_{\alpha,i+1/2+} \leq M_{\alpha,i+1}, \quad (72)$$

thus

$$M(U_{\alpha,i}, \xi) = 0 \Rightarrow M(U_{\alpha,i+1/2-}, \xi) = 0 \text{ and } M(U_{\alpha,i-1/2+}, \xi) = 0.$$

The quantities  $\delta M_{\alpha,i+1/2\pm}$  depend on  $\xi$ ,  $U_{\alpha,i}$ ,  $U_{\alpha,i+1}$ ,  $\Delta z_{i+1/2} = z_{i+1} - z_i$ , and satisfy the moment relations

$$\int_{\mathbb{R}} \delta M_{\alpha,i+1/2-} d\xi = 0, \quad \int_{\mathbb{R}} \xi \delta M_{\alpha,i+1/2-} d\xi = g \frac{h_{\alpha,i}^2}{2} - g \frac{h_{\alpha,i+1/2-}^2}{2}, \quad (73)$$

$$\int_{\mathbb{R}} \delta M_{\alpha,i-1/2+} d\xi = 0, \quad \int_{\mathbb{R}} \xi \delta M_{\alpha,i-1/2+} d\xi = g \frac{h_{\alpha,i}^2}{2} - g \frac{h_{\alpha,i-1/2+}^2}{2}. \quad (74)$$

We also assume that,

$$M(U_{\alpha,i}, \xi) = 0 \Rightarrow \delta M_{\alpha,i+1/2-}(\xi) = 0 \text{ and } \delta M_{\alpha,i-1/2+}(\xi) = 0. \quad (75)$$

For reasons that will appear later during the derivation of the entropy inequality, we make the choice

$$\begin{aligned} \delta M_{\alpha,i+1/2-} &= (\xi - u_{\alpha,i})(M_{\alpha,i} - M_{\alpha,i+1/2-}), \\ \delta M_{\alpha,i-1/2+} &= (\xi - u_{\alpha,i})(M_{\alpha,i} - M_{\alpha,i-1/2+}), \end{aligned} \quad (76)$$

that satisfies the assumptions (73), (74) and (75). This allows to precise the numerical fluxes in (44) having the form

$$F_{i+1/2-} = \left( \sum_{\alpha=1}^N F_{h_{\alpha,i+1/2-}}, F_{q_{1,i+1/2-}}, \dots, F_{q_{N,i+1/2-}} \right)^T, \quad (77)$$

with

$$\begin{aligned} F_{h_{\alpha,i+1/2-}} &= \int_{\mathbb{R}} \xi M_{\alpha,i+1/2} d\xi = \int_{\xi>0} \xi M_{\alpha,i+1/2-} d\xi + \int_{\xi<0} \xi M_{\alpha,i+1/2+} d\xi, \\ F_{q_{\alpha,i+1/2-}} &= \int_{\mathbb{R}} (\xi^2 M_{\alpha,i+1/2} + \xi \delta M_{\alpha,i+1/2-}) d\xi + \Delta x_i u_{\alpha+1/2,i} G_{\alpha+1/2,i} \\ &= \int_{\xi>0} \xi^2 M_{\alpha,i+1/2-} d\xi + \int_{\xi<0} \xi^2 M_{\alpha,i+1/2+} d\xi \\ &\quad + \frac{gh_{\alpha,i}^2}{2} - \frac{gh_{\alpha,i+1/2-}^2}{2} + \Delta x_i u_{\alpha+1/2,i} G_{\alpha+1/2,i}. \end{aligned} \quad (78)$$

Now we prove some properties of the scheme (46) with the choice (76) Notice that the implicit part of the scheme (46) requires to invert the matrix  $\mathbf{I}_N + \Delta t \mathbf{G}_{N,i}$  whose properties have been studied in lemma 3.10.

In order to take into account the topography source term, we consider a modified version of the CFL condition (57) having the form

$$\sigma_i(v_m + |u_{\alpha,i}|) \leq \beta < 1 \quad \text{for all } \alpha, i, \quad (79)$$

where  $\sigma_i = \Delta t^n / \Delta x_i$ , and  $\beta$  is a given constant. Notice that the CFL condition (79) is equivalent to

$$\Delta t^n \leq \beta \min_{1 \leq \alpha \leq N} \min_{i \in I} \frac{\Delta x_i}{|u_{\alpha,i}| + 2\sqrt{2gh_i}}$$

and we have the following theorem.

**Theorem 4.5** *Under the CFL condition (79), the scheme (46) with the choice (76) verifies the following properties.*

- (i) *The kinetic functions remain nonnegative  $M_{\alpha,i}^{n+1-} \geq 0$ .*
- (ii) *The scheme (46) is kinetic well-balanced.*
- (iii) *One has the kinetic relation*

$$\begin{aligned} H(\overline{M}_{\alpha,i}^{n+1-}, \xi, z_{b,i}) &= H(\overline{M}_{\alpha,i}, \xi, z_{b,i}) - \sigma_i \left( \tilde{H}_{\alpha,i+1/2-} - \tilde{H}_{\alpha,i-1/2+} \right) \\ &\quad - \Delta t^n \left( \hat{H}_{i,\alpha+1/2}^{n+1-} - \hat{H}_{i,\alpha-1/2}^{n+1-} \right) + d_{\alpha,i} + e_{\alpha,i}, \end{aligned} \quad (80)$$

where  $\tilde{H}_{\alpha,i+1/2-}$ ,  $\tilde{H}_{\alpha,i-1/2+}$ ,  $\hat{H}_{i,\alpha+1/2}^{n+1-}$ ,  $\hat{H}_{i,\alpha-1/2}^{n+1-}$  are defined in Theorem 4.1 and the terms  $d_{\alpha,i}$ ,  $e_{\alpha,i}$  satisfy the estimates

$$\begin{aligned} d_{\alpha,i} &\leq \sigma_i \xi \frac{g^2 \pi^2}{6} \left( \overline{M}_{\alpha,i+1/2+} + 2\overline{M}_{\alpha,i+1/2-} + \sigma_i \xi (2\overline{M}_{\alpha,i} + \overline{M}_{\alpha,i}^{n*}) \right) (\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i+1/2-})^2 \\ &\quad - \sigma_i \xi \frac{g^2 \pi^2}{6} \left( \overline{M}_{\alpha,i-1/2-} + 2\overline{M}_{\alpha,i-1/2+} - \sigma_i \xi (2\overline{M}_{\alpha,i} + \overline{M}_{\alpha,i}^{n*}) \right) (\overline{M}_{\alpha,i-1/2} - \overline{M}_{\alpha,i-1/2+})^2 \\ &\quad - \Delta t^n \frac{g^2 \pi^2}{6} \frac{G_{\alpha+1/2,i}}{h_i^{n+1}} (\overline{M}_{\alpha+1/2,i}^{n+1-} + 2\overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha+1/2,i}^{n+1-} - \overline{M}_{\alpha,i}^{n+1-})^2 \\ &\quad + \Delta t^n \frac{g^2 \pi^2}{6} \frac{G_{\alpha-1/2,i}}{h_i^{n+1}} (\overline{M}_{\alpha-1/2,i}^{n+1-} + 2\overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha-1/2,i}^{n+1-} - \overline{M}_{\alpha,i}^{n+1-})^2, \\ e_{\alpha,i} &\leq \sigma_i^2 \frac{g^2 \pi^2}{3} (u_{\alpha,i}^{n+1})^2 (2\overline{M}_{\alpha,i}^{n+1-} + \overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha,i-1/2+}^{n+1-} - \overline{M}_{\alpha,i+1/2-}^{n+1-})^2 \\ &\quad + \Delta t^n \frac{g^2 \pi^2}{3} \frac{(G_{\alpha+1/2,i} - G_{\alpha-1/2,i})}{h_i^{n+1}} (\overline{M}_{\alpha,i}^{n+1-})^3. \end{aligned}$$

**Remark 4.6** *The result of Remark 4.4 concerning the boundedness of the quantities*

$$\frac{G_{\alpha \pm 1/2,i}}{h_i^{n+1}},$$

*remains valid.*

In relation (80), the first two lines of  $d_{\alpha,i}$  are not necessarily non positive and the first line of  $e_{\alpha,i}$  is non negative. Similar estimates have been obtained in [2] in the context of the classical Saint-Venant system and using the same arguments as in [2, Theorem 3.6], the following corollary holds.

**Corollary 4.7** *Under the CFL condition (79), the scheme (46) with the choice (76) has the kinetic entropy inequality*

$$H(\overline{M}_{\alpha,i}^{n+1-}, \xi, z_{b,i}) = H(\overline{M}_{\alpha,i}, \xi, z_{b,i}) - \sigma_i \left( \tilde{H}_{\alpha,i+1/2-} - \tilde{H}_{\alpha,i-1/2+} \right) - \Delta t^n \left( \hat{H}_{i,\alpha+1/2}^{n+1-} - \hat{H}_{i,\alpha-1/2}^{n+1-} \right) + d_{\alpha,i}^1 + e_{\alpha,i}^1, \quad (81)$$

where  $\tilde{H}_{i+1/2-}$ ,  $\tilde{H}_{i-1/2+}$ ,  $\hat{H}_{i+1/2-}^{n+1-}$ ,  $\hat{H}_{i-1/2+}^{n+1-}$  are defined in Theorem 4.5,  $\nu_\beta > 0$  is a dissipation constant depending only on  $\beta$ , and  $C_\beta \geq 0$  is a constant depending only on  $\beta$ . The term proportional to  $C_\beta$  is an error, while the term proportional to  $\nu_\beta$  is a dissipation that reinforces the inequality. The quantities  $d_{\alpha,i}^1$ ,  $e_{\alpha,i}^1$  are defined by

$$\begin{aligned} d_{\alpha,i}^1 &\leq \nu_\beta \sigma_i \xi \frac{g^2 \pi^2}{6} (\overline{M}_{\alpha,i+1/2+} + \overline{M}_{\alpha,i+1/2-}) (\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i+1/2-})^2 \\ &\quad - \nu_\beta \sigma_i \xi \frac{g^2 \pi^2}{6} (\overline{M}_{\alpha,i-1/2-} + \overline{M}_{\alpha,i-1/2+}) (\overline{M}_{\alpha,i-1/2} - \overline{M}_{\alpha,i-1/2+})^2 \\ &\quad - \Delta t^n \frac{g^2 \pi^2}{6} \frac{G_{\alpha+1/2,i}}{h_i^{n+1}} (\overline{M}_{\alpha+1/2,i}^{n+1-} + 2\overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha+1/2,i}^{n+1-} - \overline{M}_{\alpha,i}^{n+1-})^2 \\ &\quad + \Delta t^n \frac{g^2 \pi^2}{6} \frac{G_{\alpha-1/2,i}}{h_i^{n+1}} (\overline{M}_{\alpha-1/2,i}^{n+1-} + 2\overline{M}_{\alpha,i}^{n+1-}) (\overline{M}_{\alpha-1/2,i}^{n+1-} - \overline{M}_{\alpha,i}^{n+1-})^2, \\ e_{\alpha,i}^1 &\leq C_\beta (\sigma_i v_m)^2 \frac{g^2 \pi^2}{6} \overline{M}_i \left( (\overline{M}_{\alpha,i} - \overline{M}_{\alpha,i+1/2-})^2 + (\overline{M}_{\alpha,i} - \overline{M}_{\alpha,i-1/2+})^2 \right) \\ &\quad + \Delta t^n \frac{g^2 \pi^2}{3} \frac{(G_{\alpha+1/2,i} - G_{\alpha-1/2,i})}{h_i^{n+1}} (\overline{M}_{\alpha,i}^{n+1-})^3. \end{aligned}$$

**Corollary 4.8** *Under the CFL condition (79), integrating the relation (81) with respect to  $\xi$  and summing for  $\alpha = 1, \dots, N$ , yields that*

$$\sum_{\alpha=1}^N \overline{E}_{\alpha,i}^{n+1} \leq \sum_{\alpha=1}^N \overline{E}_{\alpha,i} - \sigma_i \left( \sum_{\alpha=1}^N \int_{\mathbb{R}} \tilde{H}_{i+1/2} d\xi - \int_{\mathbb{R}} \tilde{H}_{i-1/2} d\xi \right) \quad (82)$$

$$+ C_\beta (\sigma_i v_m)^2 \left( g(h_i - h_{i+1/2-})^2 + g(h_i - h_{i-1/2+})^2 \right), \quad (83)$$

where  $C_\beta \geq 0$  depends only on  $\beta$ . As in [2, Corollary 3.7], we conclude that relation (83) is the discrete entropy inequality associated to the HR scheme (44),(70),(71) with kinetic numerical flux (77)-(78). With (69)-(71) one has

$$0 \leq h_i - h_{i+1/2-} \leq |z_{b,i+1} - z_{b,i}|, \quad 0 \leq h_i - h_{i-1/2+} \leq |z_{b,i} - z_{b,i-1}|.$$

We conclude that the quadratic error terms proportional to  $C_\beta$  in the right-hand side of (83) (divide (83) by  $\Delta t^n$  to be consistent with (14)) has the following key properties: it vanishes identically when  $z_b = \text{cst}$  (no topography) or when  $\sigma_i \rightarrow 0$  (semi-discrete limit), and as soon as the topography is Lipschitz continuous, it tends to zero strongly when the grid size tends to 0 (consistency with the continuous entropy inequality (14)), even if the solution contains shocks.

**Proof of theorem 4.5** (i) The scheme (46) also writes

$$\begin{aligned} & -\Delta t^n \frac{|G_{\alpha+1/2,i}|_+}{h_{\alpha+1,i}^{n+1}} M_{\alpha+1,i}^{n+1-} + \left( 1 - \Delta t^n \frac{|G_{\alpha+1/2,i}|_- - |G_{\alpha-1/2,i}|_+}{h_{\alpha,i}^{n+1}} \right) M_{\alpha,i}^{n+1-} \\ & + \Delta t^n \frac{|G_{\alpha-1/2,i}|_-}{h_{\alpha-1,i}^{n+1}} M_{\alpha-1,i}^{n+1-} = M_{\alpha,i} - \sigma_i (\xi M_{\alpha,i+1/2} + \delta M_{\alpha,i+1/2} - M_{\alpha,i-1/2} - \delta M_{\alpha,i-1/2}), \end{aligned}$$

and using (72),(76) we have

$$\begin{aligned} & -\Delta t^n \frac{|G_{\alpha+1/2,i}|_+}{h_{\alpha+1,i}^{n+1}} M_{\alpha+1,i}^{n+1-} + \left( 1 - \Delta t^n \frac{|G_{\alpha+1/2,i}|_- - |G_{\alpha-1/2,i}|_+}{h_{\alpha,i}^{n+1}} \right) M_{\alpha,i}^{n+1-} \\ & + \Delta t^n \frac{|G_{\alpha-1/2,i}|_-}{h_{\alpha-1,i}^{n+1}} M_{\alpha-1,i}^{n+1-} \geq (1 - \sigma_i (|\xi| + |u_{\alpha,i}|)) M_{\alpha,i}. \end{aligned}$$

The CFL condition (79) ensures the r.h.s. of the previous equation is non negative and using the property (ii) of lemma 3.10, this proves (i).

(ii) When  $u_{\alpha,i} = 0$ ,  $h_i + z_{b,i} = \text{cst}$  for any  $\alpha, i$  then for any  $\xi$  we have  $M_{\alpha,i+1/2+} = M_{\alpha,i+1/2-} = M_{\alpha,i-1/2+} = M_{\alpha,i-1/2-}$ ,  $G_{\alpha+1/2,i} = G_{\alpha-1/2,i} = 0$  and therefore  $M_{\alpha,i}^{n+1-} = M_{\alpha,i}$  proving (ii).

(iii) In order to prove (iii) we proceed as in the proof of Theorem 4.1 but the computations are more complex because of the topography source terms. First, we multiply the first line of (46) with (76) by  $H'(\bar{M}_{\alpha,i}, \xi, z_{b,i})$  and we get

$$\begin{aligned} H(\bar{M}_{\alpha,i}^{n*}, \xi, z_{b,i}) &= H(\bar{M}_{\alpha,i}, \xi, z_{b,i}) \\ & - \sigma_i \xi (H(\bar{M}_{\alpha,i+1/2}, \xi, z_{b,i+1/2}) - H(\bar{M}_{\alpha,i-1/2}, \xi, z_{b,i-1/2})) + R_{\alpha,i}^x + R_{\alpha,i}^t, \end{aligned} \quad (84)$$

where  $R_{\alpha,i}^x$  (resp.  $R_{\alpha,i}^t$ ) is an error term coming from the space (resp. time) discretization and with

$$\begin{aligned} R_{\alpha,i}^x &= \sigma_i \xi (H(\bar{M}_{\alpha,i+1/2}, \xi, z_{b,i+1/2}) - H'(\bar{M}_{\alpha,i}, \xi, z_{b,i}) \bar{M}_{\alpha,i+1/2}) \\ & - \sigma_i \xi (H(\bar{M}_{\alpha,i-1/2}, \xi, z_{b,i-1/2}) - H'(\bar{M}_{\alpha,i}, \xi, z_{b,i}) \bar{M}_{\alpha,i-1/2}) \\ & + \sigma_i (\xi - u_{\alpha,i}) H'(\bar{M}_{\alpha,i}, \xi, z_{b,i}) (\bar{M}_{\alpha,i+1/2-} - \bar{M}_{\alpha,i-1/2+}), \\ R_{\alpha,i}^t &= \frac{\pi^2 g^2}{6} (\bar{M}_{\alpha,i}^{n*} + 2\bar{M}_{\alpha,i}) (\bar{M}_{\alpha,i}^{n*} - \bar{M}_{\alpha,i})^2. \end{aligned}$$

For the quantity  $R_{\alpha,i}^t$ , using (46) we have

$$\begin{aligned} R_{\alpha,i}^t &= \frac{g^2\pi^2}{6}\sigma_i^2(2\overline{M}_{\alpha,i} + \overline{M}_{\alpha,i}^{n*}) (\xi\overline{M}_{\alpha,i+1/2} - \xi\overline{M}_{\alpha,i-1/2} + \delta\overline{M}_{\alpha,i+1/2-} - \delta\overline{M}_{\alpha,i-1/2+})^2 \\ &\leq \frac{2g^2\pi^2}{3}\sigma_i^2(2\overline{M}_{\alpha,i} + \overline{M}_{\alpha,i}^{n*}) \left( \xi^2(\overline{M}_{i+1/2+} - \overline{M}_{i+1/2-})^2 \mathbf{1}_{\xi < 0} \right. \\ &\quad \left. + \xi^2(\overline{M}_{i-1/2-} - \overline{M}_{i-1/2+})^2 \mathbf{1}_{\xi > 0} + u_i^2(\overline{M}_{i+1/2-} - \overline{M}_{i-1/2+})^2 \right). \end{aligned}$$

For the quantity  $R_{\alpha,i}^x$ , we can write  $R_{\alpha,i}^x = R_{\alpha,i+}^x + R_{\alpha,i-}^x$  with

$$\begin{aligned} R_{\alpha,i+}^x &= \sigma_i\xi \left( H(\overline{M}_{\alpha,i+1/2}, \xi, z_{b,i+1/2}) - H'(\overline{M}_{\alpha,i}, \xi, z_{b,i})(\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i+1/2-}) \right) \\ &\quad - \sigma_i u_{\alpha,i} H'(\overline{M}_{\alpha,i}, \xi, z_{b,i}) \overline{M}_{\alpha,i+1/2-} - \sigma_i \xi H(\overline{M}_{\alpha,i}, \xi, z_{b,i}), \end{aligned}$$

and

$$\begin{aligned} R_{\alpha,i-}^x &= -\sigma_i\xi \left( H(\overline{M}_{\alpha,i-1/2}, \xi, z_{b,i-1/2}) - H'(\overline{M}_{\alpha,i}, \xi, z_{b,i})(\overline{M}_{\alpha,i-1/2} - \overline{M}_{\alpha,i-1/2+}) \right) \\ &\quad + \sigma_i u_{\alpha,i} H'(\overline{M}_{\alpha,i}, \xi, z_{b,i}) \overline{M}_{\alpha,i-1/2+} + \sigma_i \xi H(\overline{M}_{\alpha,i}, \xi, z_{b,i}), \end{aligned}$$

and we can rewrite  $R_{\alpha,i+}^x$  under the form

$$\begin{aligned} R_{\alpha,i+}^x &= \sigma_i\xi \left( H(\overline{M}_{\alpha,i+1/2}, \xi, z_{b,i+1/2}) - H'(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2})(\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i+1/2-}) \right) \\ &\quad - \sigma_i\xi \left( H'(\overline{M}_{\alpha,i}, \xi, z_{b,i}) - H'(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2}) \right) (\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i+1/2-}) \\ &\quad - \sigma_i u_{\alpha,i} H'(\overline{M}_{\alpha,i}, \xi, z_{b,i}) \overline{M}_{\alpha,i+1/2-} - \sigma_i \xi H(\overline{M}_{\alpha,i}, \xi, z_{b,i}). \end{aligned}$$

Using the identity

$$\begin{aligned} \overline{M}_{\alpha,i+1/2+}^3 &= \overline{M}_{\alpha,i+1/2-}^3 + 3\overline{M}_{\alpha,i+1/2-}^2(\overline{M}_{\alpha,i+1/2+} - \overline{M}_{\alpha,i+1/2-}) \\ &\quad + (\overline{M}_{\alpha,i+1/2+} + 2\overline{M}_{\alpha,i+1/2-})(\overline{M}_{\alpha,i+1/2+} - \overline{M}_{\alpha,i+1/2-})^2, \end{aligned}$$

we obtain an expression of the linear dissipation under the form

$$\begin{aligned} &H(\overline{M}_{\alpha,i+1/2+}, \xi, z_{b,i+1/2}) - H(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2}) \\ &\quad - H'(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2})(\overline{M}_{\alpha,i+1/2+} - \overline{M}_{\alpha,i+1/2-}) \\ &= \frac{g^2\pi^2}{6}(\overline{M}_{\alpha,i+1/2+} + 2\overline{M}_{\alpha,i+1/2-})(\overline{M}_{\alpha,i+1/2+} - \overline{M}_{\alpha,i+1/2-})^2 \end{aligned}$$

that can also be written under the form

$$\begin{aligned} H(\overline{M}_{\alpha,i+1/2}, \xi, z_{b,i+1/2}) &= H(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2}) \\ &\quad + H'(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2})(\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i+1/2-}) \\ &\quad + \frac{g^2\pi^2}{6}(\overline{M}_{\alpha,i+1/2+} + 2\overline{M}_{\alpha,i+1/2-})(\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i+1/2-})^2. \end{aligned} \tag{85}$$

Therefore, relation (85) allows to write  $L_{2+}$  under the form

$$R_{\alpha,i+}^x = \sigma_i\xi \left( H(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2}) + \frac{g^2\pi^2}{6}(\overline{M}_{\alpha,i+1/2+} + 2\overline{M}_{\alpha,i+1/2-})(\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i+1/2-})^2 \right)$$

$$\begin{aligned}
& -\sigma_i \xi (H'(\overline{M}_{\alpha,i}, \xi, z_{b,i}) - H'(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2})) (\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i+1/2-}) \\
& -\sigma_i u_{\alpha,i} H'(\overline{M}_{\alpha,i}, \xi, z_{b,i}) \overline{M}_{\alpha,i+1/2-} - \sigma_i \xi H(\overline{M}_{\alpha,i}, \xi, z_{b,i}).
\end{aligned}$$

Next, if  $\overline{M}_{\alpha,i}(\xi) > 0$ , one has

$$\begin{aligned}
& H'(\overline{M}_{\alpha,i}, \xi, z_{b,i}) (\overline{M}_{\alpha,i+1/2+} - \overline{M}_{\alpha,i+1/2-}) \\
& = (\eta'(U_{\alpha,i}) \left( \frac{1}{\xi} \right) + g z_{b,i}) (\overline{M}_{\alpha,i+1/2+} - \overline{M}_{\alpha,i+1/2-}),
\end{aligned} \tag{86}$$

and as in (43)

$$\begin{aligned}
& H'(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2}) (\overline{M}_{\alpha,i+1/2+} - \overline{M}_{\alpha,i+1/2-}) \\
& \geq (\eta'(U_{\alpha,i+1/2-}) \left( \frac{1}{\xi} \right) + g z_{b,i+1/2}) (\overline{M}_{\alpha,i+1/2+} - \overline{M}_{\alpha,i+1/2-}).
\end{aligned} \tag{87}$$

Taking the difference between (87) and (86), we obtain

$$\begin{aligned}
& H'(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2}) (\overline{M}_{\alpha,i+1/2+} - \overline{M}_{\alpha,i+1/2-}) - H'(\overline{M}_{\alpha,i}, \xi, z_{b,i}) (\overline{M}_{\alpha,i+1/2+} - \overline{M}_{\alpha,i+1/2-}) \\
& \geq l_{\alpha} (g h_{i+1/2-} - g h_i + g z_{b,i+1/2} - g z_{b,i}) (\overline{M}_{\alpha,i+1/2+} - \overline{M}_{\alpha,i+1/2-}) \geq 0,
\end{aligned} \tag{88}$$

Using (88), it comes

$$\begin{aligned}
L_{\alpha,2+} & \leq \sigma_i \xi (H(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2}) + \frac{g^2 \pi^2}{6} (\overline{M}_{\alpha,i+1/2+} + 2\overline{M}_{\alpha,i+1/2-}) (\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i+1/2-})^2) \\
& \quad -\sigma_i u_{\alpha,i} H'(\overline{M}_{\alpha,i}, \xi, z_{b,i}) \overline{M}_{\alpha,i+1/2-} - \sigma_i \xi H(\overline{M}_{\alpha,i}, \xi, z_{b,i}) \\
& = \sigma_i \xi (H(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2}) + \frac{g^2 \pi^2}{6} (\overline{M}_{\alpha,i+1/2+} + 2\overline{M}_{\alpha,i+1/2-}) (\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i+1/2-})^2) \\
& \quad -\sigma_i H'(\overline{M}_{\alpha,i}, \xi, z_{b,i}) (\xi \overline{M}_{\alpha,i+1/2-} - \xi \overline{M}_{\alpha,i} + \delta \overline{M}_{\alpha,i+1/2-}) - \sigma_i \xi H(\overline{M}_{\alpha,i}, \xi, z_{b,i}).
\end{aligned}$$

Then with (42) we get

$$\begin{aligned}
& \left( \eta'(U_{\alpha,i}) \left( \frac{1}{\xi} \right) + g z_{b,i} \right) (\xi \overline{M}_{\alpha,i+1/2-} - \xi \overline{M}_{\alpha,i} + \delta \overline{M}_{\alpha,i+1/2-}) \\
& = H'(\overline{M}_{\alpha,i}, \xi, z_{b,i}) (\xi \overline{M}_{\alpha,i+1/2-} - \xi \overline{M}_{\alpha,i} + \delta \overline{M}_{\alpha,i+1/2-}),
\end{aligned} \tag{89}$$

We use the identity (89) giving

$$\begin{aligned}
R_{\alpha,i+}^x & \leq \sigma_i \xi (H(\overline{M}_{\alpha,i+1/2-}, \xi, z_{b,i+1/2}) + \frac{g^2 \pi^2}{6} (\overline{M}_{\alpha,i+1/2+} + 2\overline{M}_{\alpha,i+1/2-}) (\overline{M}_{\alpha,i+1/2} - \overline{M}_{\alpha,i+1/2-})^2) \\
& \quad -\sigma_i \left( \eta'(U_{\alpha,i}) \left( \frac{1}{\xi} \right) + g z_{b,i} \right) (\xi \overline{M}_{\alpha,i+1/2-} - \xi \overline{M}_{\alpha,i} + \delta \overline{M}_{\alpha,i+1/2-}) \\
& \quad -\sigma_i \xi H(\overline{M}_{\alpha,i}, \xi, z_{b,i}),
\end{aligned}$$

and an analogous inequality can be obtained for  $L_{2-}$  under the form

$$R_{\alpha,i-}^x \leq -\sigma_i \xi (H(\overline{M}_{\alpha,i-1/2+}, \xi, z_{b,i-1/2}) + \frac{g^2 \pi^2}{6} (\overline{M}_{\alpha,i-1/2-} + 2\overline{M}_{\alpha,i-1/2+}) (\overline{M}_{\alpha,i-1/2} - \overline{M}_{\alpha,i-1/2+})^2)$$



$$\begin{aligned}
& +\sigma_i \left( \eta'(U_{\alpha,i}) \left( \frac{1}{\xi} \right) + gz_{b,i} \right) (\xi \bar{M}_{\alpha,i-1/2+} - \xi \bar{M}_{\alpha,i} + \delta \bar{M}_{\alpha,i-1/2+}) \\
& +\sigma_i \xi H(\bar{M}_{\alpha,i}, \xi, z_{b,i}),
\end{aligned}$$

In a second step, we multiply the second equation of (46) by  $H'_\alpha(\bar{M}_{\alpha,i}^{n+1-}, \xi, z_{b,i})$ . After the same computations as in the proof of Theorem 4.1 – since the implicit (in time) step of the scheme (46) do not depend on the topography source term – we obtain the relation (66).

Therefore adding the two relations (84),(66) gives (80) proving the result. ■

**Proof of Corollary 4.7** The proof of the result is similar to the one given by some of the authors in [2, Theorem 3.6]. ■

**Proof of corollary 4.8** As in the proof of Corollary (4.2), it is possible to prove that

$$\int_{\mathbb{R}} (\bar{M}_{\alpha,i}^{n+1-})^3 d\xi = \int_{\mathbb{R}} (\bar{M}_{\alpha,i}^{n+1})^3 d\xi,$$

and under the CFL condition (57) the quantity  $d_{\alpha,i}^1$  is non positive then the sum for  $i = 1, \dots, N$  of relations (81) integrated in  $\xi$  gives the result. ■

## 5 Fully discrete entropy inequality for the layer-averaged Navier-Stokes system

The layer-averaging applied to the Euler system in Section 3 can also be applied to the Navier-Stokes system, see [11]. Considering a simplified Newtonian rheology, the Navier-Stokes system (1)-(3) can be written under the form

$$\frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0, \tag{90}$$

$$\frac{\partial u}{\partial t} + \frac{\partial u^2}{\partial x} + \frac{\partial uw}{\partial z} + \frac{\partial p}{\partial x} = \mu \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial z^2}, \tag{91}$$

$$\frac{\partial p}{\partial z} = -g, \tag{92}$$

where  $\mu$  is a viscosity coefficient. The system (90)-(92) is completed with the kinematic boundary conditions (4),(6) and suitable dynamic boundary conditions and its layer-averaged version is given by

$$\frac{\partial}{\partial t} h + \frac{\partial}{\partial x} \sum_{j=1}^N h_\alpha u_\alpha = 0, \tag{93}$$

$$\frac{\partial}{\partial t} (h_\alpha u_\alpha) + \frac{\partial}{\partial x} \left( h_\alpha u_\alpha^2 + \frac{g}{2} h_\alpha h \right) = -gh_\alpha \frac{\partial z_b}{\partial x} + u_{\alpha+1/2} G_{\alpha+1/2} - u_{\alpha-1/2} G_{\alpha-1/2}$$

$$+\frac{\partial}{\partial x} \left( 4\mu h_\alpha \frac{\partial u_\alpha}{\partial x} \right) + 2\mu \frac{u_{\alpha+1} - u_\alpha}{h_{\alpha+1} + h_\alpha} - 2\mu \frac{u_\alpha - u_{\alpha-1}}{h_\alpha + h_{\alpha-1}}, \quad \alpha = 2, \dots, N-1 \quad (94)$$

$$\begin{aligned} \frac{\partial}{\partial t} (h_1 u_1) + \frac{\partial}{\partial x} \left( h_1 u_1^2 + \frac{g}{2} h_1 h \right) &= -g h_1 \frac{\partial z_b}{\partial x} + u_{3/2} G_{3/2} \\ &+ \frac{\partial}{\partial x} \left( 4\mu h_1 \frac{\partial u_1}{\partial x} \right) + 2\mu \frac{u_2 - u_1}{h_2 + h_1} - \kappa u_1, \end{aligned} \quad (95)$$

$$\begin{aligned} \frac{\partial}{\partial t} (h_N u_N) + \frac{\partial}{\partial x} \left( h_N u_N^2 + \frac{g}{2} h_N h \right) &= -g h_N \frac{\partial z_b}{\partial x} - u_{N-1/2} G_{N-1/2} \\ &+ \frac{\partial}{\partial x} \left( 4\mu h_N \frac{\partial u_N}{\partial x} \right) - 2\mu \frac{u_N - u_{N-1}}{h_N + h_{N-1}}, \end{aligned} \quad (96)$$

$$\begin{aligned} \frac{\partial}{\partial t} \left( \frac{z_{\alpha+1/2}^2 - z_{\alpha-1/2}^2}{2} \right) + \frac{\partial}{\partial x} \left( \frac{z_{\alpha+1/2}^2 - z_{\alpha-1/2}^2}{2} u_\alpha \right) dz &= h_\alpha \omega_\alpha \\ &+ z_{\alpha+1/2} G_{\alpha+1/2} - z_{\alpha-1/2} G_{\alpha-1/2}, \end{aligned} \quad (97)$$

where  $\kappa$  is a Navier friction coefficient at the bottom and the exchange terms  $G_{\alpha+1/2}$  are given by (19)-(20).

The system (93)-(97) is rewritten under the compact form

$$\frac{\partial X}{\partial t} + \frac{\partial F(X)}{\partial x} = S_e(X, \partial_t X, \partial_x X) + S_b(X) + S_{v,f}(X), \quad (98)$$

with  $X = \left( \sum_{\alpha=1}^N h_\alpha, q_1, \dots, q_N \right)^T$  and  $q_\alpha = l_\alpha h$ ,  $q_\alpha = l_\alpha h u_\alpha$ . We denote with  $F(X)$  the flux of the conservative part, and with  $S_e(X, \partial_t X, \partial_x X)$ ,  $S_b(X)$  and  $S_{v,f}(X)$  the source terms, representing respectively the mass transfer, the topography, and the viscous and friction effects.

For the time discretization, we apply a time splitting technique to the equations (98) and we write

$$\frac{\tilde{X}^{n+1} - X^n}{\Delta t^n} + \frac{\partial F(X^n)}{\partial x} = S_e(X^n, \tilde{X}^{n+1}) + S_b(X^n), \quad (99)$$

$$\frac{X^{n+1} - \tilde{X}^{n+1}}{\Delta t^n} - S_{v,f}(X^n, X^{n+1}) = 0. \quad (100)$$

Equation (99) corresponds to the semi-discrete in time version of the layer-averaged Euler system (17)-(18) whose discretization has studied in Section 4. It remains to propose a discrete in space scheme for Eq. (100).

We first solve the conservative part of (99) by the proposed explicit kinetic scheme then the second step taking into account the viscous and friction effects is performed using an explicit time.

More precisely, since the viscous and friction terms  $S_{v,f}$  in (100) are dissipative, they are treated via a semi-implicit scheme for stability reasons. This gives using a finite differences discretization

$$(hu)_{\alpha,i}^{n+1} = (\widetilde{hu})_{\alpha,i}^{n+1} + \frac{8\mu}{\Delta x_i} \left( h_{\alpha,i+1/2} \frac{u_{\alpha,i+1}^{1/2} - u_{\alpha,i}^{1/2}}{\Delta x_{i+1} + \Delta x_i} - h_{\alpha,i-1/2} \frac{u_{\alpha,i}^{1/2} - u_{\alpha,i-1}^{1/2}}{\Delta x_i + \Delta x_{i-1}} \right)$$

$$+2\mu \frac{u_{\alpha+1,i}^{\frac{1}{2}} - u_{\alpha,i}^{\frac{1}{2}}}{h_{\alpha+1,i} + h_{\alpha,i}} - 2\mu \frac{u_{\alpha,i}^{\frac{1}{2}} - u_{\alpha-1,i}^{\frac{1}{2}}}{h_{\alpha,i} + h_{\alpha-1,i}} - \kappa \delta_{1,\alpha} u_{\alpha,i}^{\frac{1}{2}}, \quad (101)$$

for  $\alpha = 1, \dots, N$ . The superscript  $y^{\frac{1}{2}}$  means  $y^{\frac{1}{2}} = \frac{y^n + \tilde{y}^{n+1}}{2}$  and  $\delta_{1,\alpha}$  is the Kronecker symbol.

And the following proposition holds.

**Proposition 5.1** *Defining  $E_{\alpha,i}$  by*

$$E_{\alpha,i} = \frac{h_{\alpha,i}}{2} (h_{\alpha,i})^2 + \frac{g}{2} h_{\alpha,i} h_i + g z_{b,i} h_{\alpha,i},$$

with the notations of theorem 4.5 we have

$$\begin{aligned} E_{\alpha,i}^{n+1} &\leq E_{\alpha,i} - \sigma_i \left( \int_{\mathbb{R}} \tilde{H}_{\alpha,i+1/2-} d\xi - \int_{\mathbb{R}} \tilde{H}_{\alpha,i-1/2+} d\xi \right) \\ &\quad - \Delta t^n \left( \int_{\mathbb{R}} \hat{H}_{i,\alpha+1/2}^{n+1-} d\xi - \int_{\mathbb{R}} \hat{H}_{i,\alpha-1/2}^{n+1-} d\xi \right) \\ &\quad + \int_{\mathbb{R}} d_{\alpha,i} d\xi + \int_{\mathbb{R}} e_{\alpha,i} d\xi + \Delta t^n \mathcal{D}_{\alpha,i}, \end{aligned}$$

with

$$\begin{aligned} \mathcal{D}_{\alpha,i} &= \frac{8\mu}{\Delta x_i} \left( h_{\alpha,i+1/2} \frac{u_{\alpha,i+1}^{\frac{1}{2}} + u_{\alpha,i}^{\frac{1}{2}}}{2} \frac{u_{\alpha,i+1}^{\frac{1}{2}} - u_{\alpha,i}^{\frac{1}{2}}}{\Delta x_{i+1} + \Delta x_i} \right. \\ &\quad \left. - h_{\alpha,i-1/2} \frac{u_{\alpha,i}^{\frac{1}{2}} + u_{\alpha,i-1}^{\frac{1}{2}}}{2} \frac{u_{\alpha,i}^{\frac{1}{2}} - u_{\alpha,i-1}^{\frac{1}{2}}}{\Delta x_i + \Delta x_{i-1}} \right) \\ &\quad + 2\mu \frac{u_{\alpha+1,i}^{\frac{1}{2}} + u_{\alpha,i}^{\frac{1}{2}}}{2} \frac{u_{\alpha+1,i}^{\frac{1}{2}} - u_{\alpha,i}^{\frac{1}{2}}}{h_{\alpha+1,i} + h_{\alpha,i}} - 2\mu \frac{u_{\alpha,i}^{\frac{1}{2}} + u_{\alpha-1,i}^{\frac{1}{2}}}{2} \frac{u_{\alpha,i}^{\frac{1}{2}} - u_{\alpha-1,i}^{\frac{1}{2}}}{h_{\alpha,i} + h_{\alpha-1,i}} \\ &\quad - \frac{4\mu}{\Delta x_i} \left( h_{\alpha,i+1/2} \frac{\left( u_{\alpha,i+1}^{\frac{1}{2}} - u_{\alpha,i}^{\frac{1}{2}} \right)^2}{\Delta x_{i+1} + \Delta x_i} - h_{\alpha,i-1/2} \frac{\left( u_{\alpha,i}^{\frac{1}{2}} - u_{\alpha,i-1}^{\frac{1}{2}} \right)^2}{\Delta x_i + \Delta x_{i-1}} \right) \\ &\quad - \mu \frac{\left( u_{\alpha+1,i}^{\frac{1}{2}} - u_{\alpha,i}^{\frac{1}{2}} \right)^2}{h_{\alpha+1,i} + h_{\alpha,i}} - \mu \frac{\left( u_{\alpha,i}^{\frac{1}{2}} - u_{\alpha-1,i}^{\frac{1}{2}} \right)^2}{h_{\alpha,i} + h_{\alpha-1,i}} - \kappa \delta_{1,\alpha} \left( u_{\alpha,i}^{\frac{1}{2}} \right)^2, \end{aligned}$$

**Proof of prop. 5.1** *First we notice that the semi-implicit step does not modify the water depth i.e.  $h_i^{n+1} = \tilde{h}_i^{n+1}$ . Multiplying (101) by  $u_{\alpha,i}^{\frac{1}{2}}$ , it comes*

$$\frac{h_{\alpha,i}^{n+1}}{2} (u_{\alpha,i}^{n+1})^2 = \frac{\tilde{h}_{\alpha,i}^{n+1}}{2} (\tilde{u}_{\alpha,i}^{n+1})^2 + \Delta t^n \mathcal{D}_{\alpha,i}, \quad (102)$$

and  $\mathcal{D}_{\alpha,i}$  is a consistent discretization of the viscous and friction terms appearing in (8). The sum of relation (80) integrated in  $\xi$  over  $\mathbb{R}$  and relation (102) gives the result. ■

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