H. Yahia ¹

B. Xu 1,3

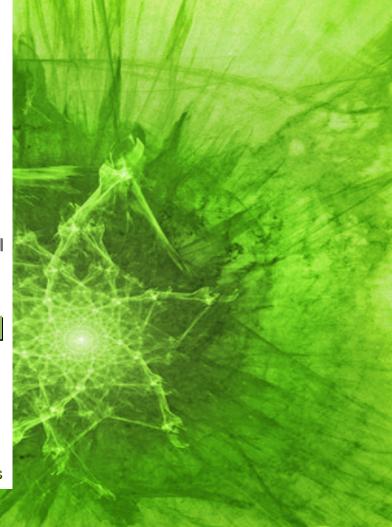
Oriol PONT 1,2

¹ INRIA Bordeaux - ² IRIT-CNRS - ³ Ch. U. Hospital Basel

Complexity in Electrophysiological Dynamics

Emergence and measures of organization

2nd BCAM Workshop on Nonlinear Dynamics in Biological Systems



Cosmos and chaos

Classic view:

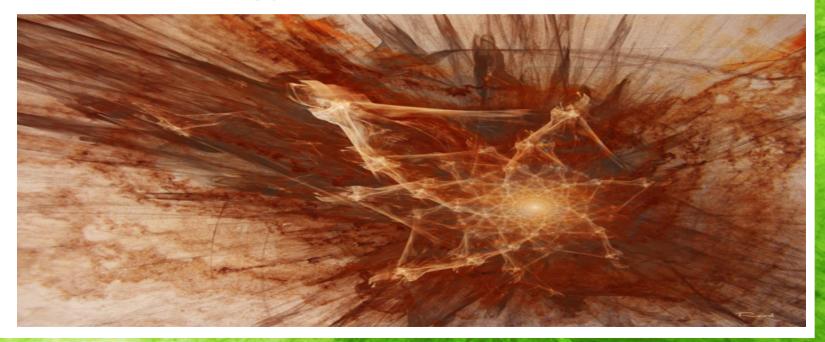
- Hesiod's cosmogony: gods created from chaos
- Anaximander: Apeiron is the principle of everything
- Ovid: in the origin, mix of the four elements
- Genesis 1:2: The earth was formless and void; darkness covered the abyss and the Spirit of God moved upon the waters



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Cosmos and chaos

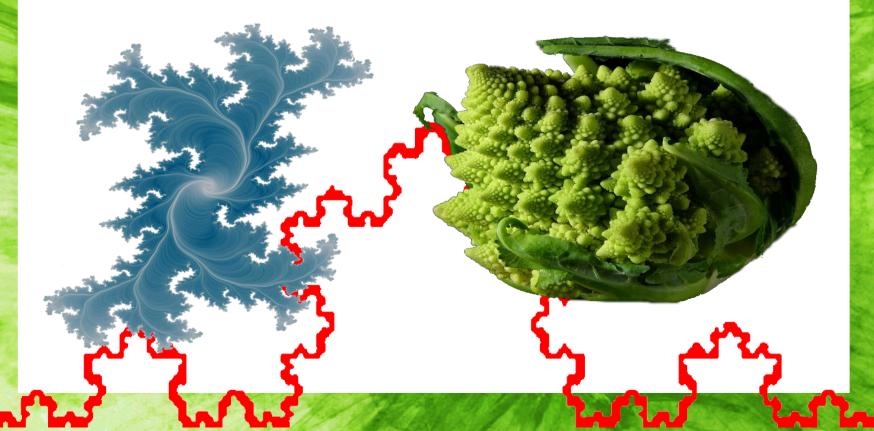
Modern view: entropy (from order to chaos)



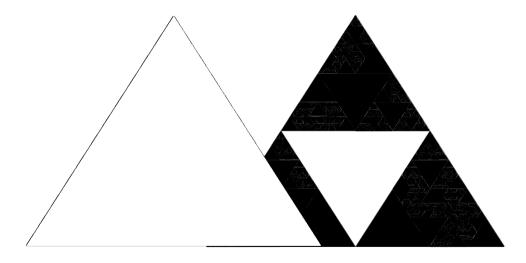
Cosmos and chaos Order from chaos?



Scale invariance

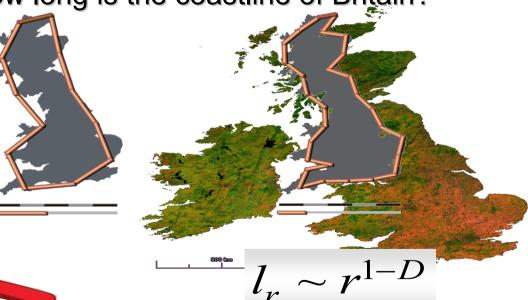


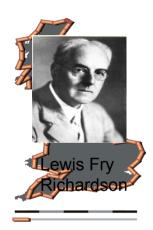
Scale invariance



Measure for scale invariance

How long is the coastline of Britain?





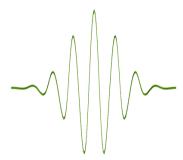


Scale invariance

Singularity exponent $h(\mathbf{x})$ degree of singularity / regularity of signal s around point \mathbf{x}

$$\mu_{\mathbf{S}}[\Psi_{r}(\mathbf{x})] \sim r^{\dim(\mathbf{x}) + h(\mathbf{x})} \qquad d\mu_{\mathbf{S}} \sim ||\nabla s|| d\mathbf{x}$$

 μ is a measure defined from gradient modulus and $\Psi_r = \Psi (\mathbf{x} / r)$ is a wavelet kernel



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Ocean Circulation

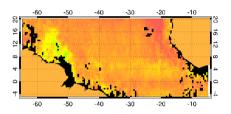
Singularity and local predictability: $\mu_s [\Psi(\mathbf{x})] = A(\mathbf{x}) ||\Psi||^{h(\mathbf{x})} + o(||\Psi||^{h(\mathbf{x})}) d\mu_s = |\nabla s| d\mathbf{x}$

Derived fields, fusion: $d\mu_s = S(\mathbf{x}) d\mu_r$

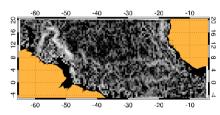
Ocean Circulation

Singularity and local predictability:

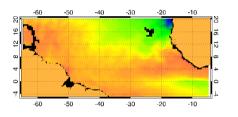
$$\mu_s[\Psi(\mathbf{x})] = A(\mathbb{X}) ||\Psi||^{h(\mathbf{x})} + o(||\Psi||^{h(\mathbb{X})}) \quad d\mu_s = |\nabla s| d\mathbf{x}$$



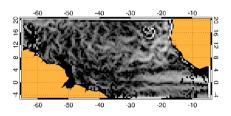
salinity



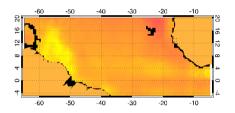
Derived fields, fusion:



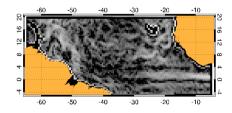
temperature



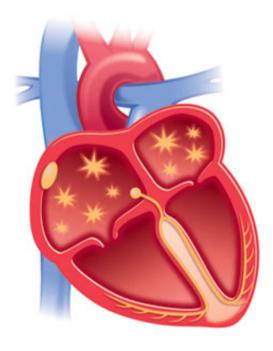
$$d\mu_s = S(\mathbf{x}) d\mu_r$$



salinity fused



Atrial fibrillation



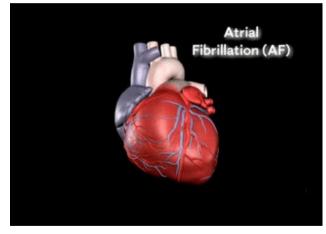
The most extended arrhythmia

Ventricles still pump but:

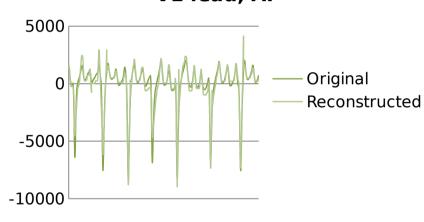
- Can lead to heart failure
- Forms blood clots (trombosis) that can provoke embolism

Signal well reconstructed from most singular manifold

$$\hat{s}(\mathbf{f}) = (i \mathbf{f} / f^2) \cdot F(\nabla s \delta_{IIP})$$



V1 lead, AF

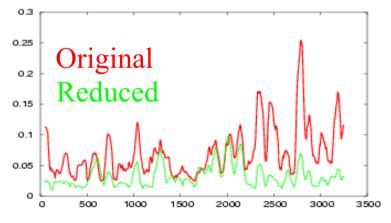


signal s, frequency \mathbf{f} , most singular manifold $\delta_{_{\mathrm{UP}}}$

Signal well reconstructed from most singular manifold

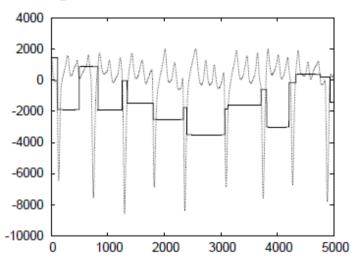
Fast dynamics: MSM orientation defines a reduced signal

Slow dynamics: modulation over the reduced reconstruction



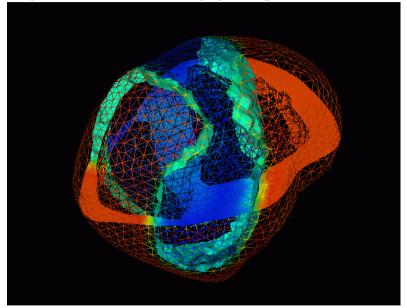
Source field, slow dynamics

$$SF(\mathbf{x}) = d\mu_S / d\mu_r$$



Radon-Nikodym dervative between original signal and reduced signal diff measures

Epicardial mapping of electric potential

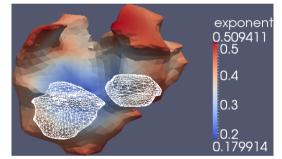


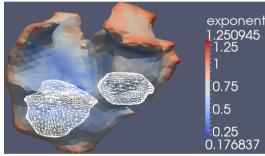


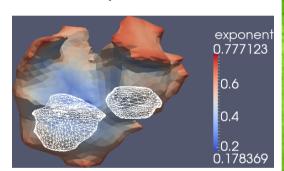
CardioInsightTM projects on the heart surface (inverse-problem), real time

Singularity exponent epicardial maps

space time space-time



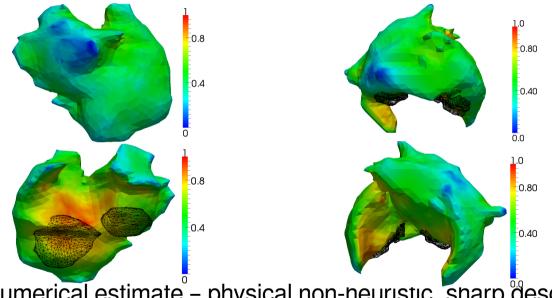




versatility in the domain considered

$$\mu_{s}[\Psi_{r}(x)] \sim r \dim(x) + h(x)$$

Source field in atrial flutter and fibrillation

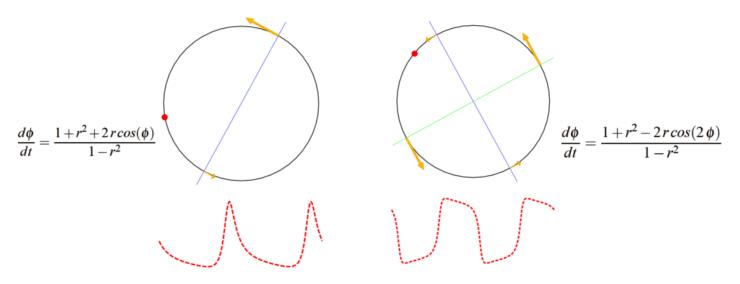


robust numerical estimate - physical non-neuristic, snarp descriptor

A basis for smooth anharmonic signals

P. Hanusse has developed a nonlinear signal-processing technique centred on the characterization of anharmonicity (as smooth perturbations of harmonicity) in periodic signals. Such characterization is rather sparse in multiple physiological systems.

Illustration of the concept



First (left) and second (right) breakings of rotational symmetry framing the notion of *anharmonic* trigonometry

Details

Generally, let x(t) be a periodic side

First symmetry breaking can be re



hcos:
$$(t,r) \to \frac{(1+r^2)\cos(t) - 2r}{1+r^2 - 2r\cos(t)}$$

hsin: $(t,r) \to \frac{(1-r^2)\sin(t)}{1+r^2 - 2r\cos(t)}$

hsin:
$$(t,r) \to \frac{(1-r^2)\sin(t)}{1+r^2-2r\cos(t)}$$

where r is the anharmonicity parameter.

Extension

The higher orders require these generalizations:

$$\operatorname{hpcos}_n : (t, r) \to \sum_{k=1}^{\infty} \frac{\cos(kt)r^{k-1}}{k^n}$$
$$\operatorname{hpsin}_n : (t, r) \to \sum_{k=1}^{\infty} \frac{\sin(kt)r^{k-1}}{k^n}$$

Order 1 is still treatable and corresponds to phase equations that are fractions of trigonometric polynomials.

$$t(\phi) = \phi + \sum_{k=1}^{n} a_k \text{hpsin}_1(\phi - p_k, r_k) - b_k \text{hpcos}_1(\phi - p_k, r_k)$$

$$p\cos_1(t, r) = -\frac{1}{2} \ln(1 + r^2 - 2r\cos(t))$$

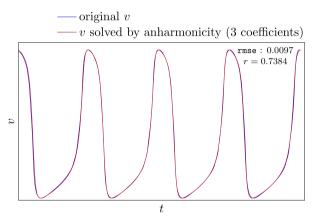
$$p\sin_1(t, r) = \tan^{-1}\left(\frac{r\sin(t)}{1 - r\cos(t)}\right)$$

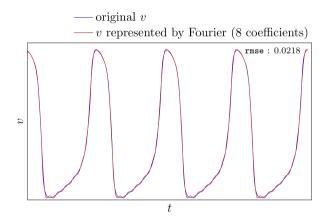
where we re-use *n* as the polynomial order since we already fixed 1 for the family of functions. Typical real-world systems are well described at low orders.

Analysis of FitzHugh-Nagumo action potentials

$$\dot{v} = v(v - \alpha)(1 - v) - w,$$

$$\dot{w} = \varepsilon(v - \gamma w),$$





Only 3 anharmonic terms reconstruct a nearly perfect representation of the signal In contrast, with 8 Fourier terms the signal is ill-represented and plagued with Gibbs phenomena The presented decomposition thus appears far more comact than Fourier for FitzHugh-Nagumo impulse trains

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Thanks for your attention!