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Exponential stability of the system of transmission of the wave equation with a delay term in the boundary feedback

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Abstract. We consider a system of transmission of the wave equation with Neumann feedback control that contains a delay term and that acts on the exterior boundary. First, we prove under some assumptions that the closed-loop system generates a C_0 -semigroup of contractions on an appropriate Hilbert space. Then, under further assumptions, we show that the closed-loop system is exponentially stable. To establish this result, we introduce a suitable energy function and use multiplier method together with an estimate taken from [3] (Lemma 7.2) and compactness-uniqueness arguments.

Keywords: Wave equation, transmission problem, time delays, boundary stabilization, exponential stability.

1 Introduction

It is by now well-known that certain infinite-dimensional second-order systems are not robust with respect to arbitrarily small delays in the damping. This lack of stability robustness was first shown to hold for the one-dimensional wave equation ([2]). Later, further examples illustrating this phenomenon were considered in [1]: the two-dimensional wave equation with damping introduced through Neumann-type boundary conditions on one edge of a square boundary and the Euler-Bernoulli beam equation in one dimension with damping introduced through a specific set of boundary conditions on the right end point.

Recently, Xu et al [9] established sufficient conditions that guarantee the exponential stability of the one-dimensional wave equation with a delay term in the boundary feedback. Nicaise and Pignotti [5] extended this result to the multi-dimensional wave equation with a delay term in the boundary or internal feedbacks. The same type of result was obtained by Nicaise and Rebiai [6] for the Schrödinger equation.

Motivated by the references [9], [5] and [6]; we investigate in this paper the problem of exponential stability for the system of transmission of the wave equation with a delay term in the boundary feedback.

Let Ω be an open bounded domain of \mathbb{R}^n with a boundary Γ of class C^2 which consists of two non-empty parts Γ_1 and Γ_2 such that $\overline{\Gamma_1} \cap \overline{\Gamma_2} = \emptyset$. Let Γ_0

with $\overline{\Gamma_0} \cap \overline{\Gamma_1} = \overline{\Gamma_0} \cap \overline{\Gamma_2} = \emptyset$ be a regular hypersurface of class C^2 which separates Ω into two domains Ω_1 and Ω_2 such that $\Gamma_1 \subset \partial\Omega_1$ and $\Gamma_2 \subset \partial\Omega_2$. Furthermore, we assume that there exists a real vector field $h \in (C^2(\overline{\Omega}))^n$ such that:

(H.1) The Jacobian matrix J of h satisfies

$$\int_{\Omega} J(x) \zeta(x) \cdot \zeta(x) d\Omega \geq \alpha \int_{\Omega} |\zeta(x)|^2 d\Omega,$$

for some constant $\alpha > 0$ and for all $\zeta \in L^2(\Omega; \mathbb{R}^n)$;

(H.2) $h(x) \cdot \nu(x) \leq 0$ on Γ_1 ;

(H.3) $h(x) \cdot \nu(x) \geq 0$ on Γ_0 .

where ν is the unit normal on Γ or Γ_0 pointing towards the exterior of Ω or Ω_1 .

Let $a_1, a_2 > 0$ be given. Consider the system of transmission of the wave equation with a delay term in the boundary conditions:

$$y''(x, t) - a(x) \Delta y(x, t) = 0 \quad \text{in } \Omega \times (0, +\infty), \quad (1)$$

$$y(x, 0) = y^0(x), y'(x, 0) = y^1(x, 0) \quad \text{in } \Omega, \quad (2)$$

$$y_1(x, t) = 0 \quad \text{on } \Gamma_1 \times (0, +\infty), \quad (3)$$

$$\frac{\partial y_2(x, t)}{\partial \nu} = -\mu_1 y_2'(x, t) - \mu_2 y_2'(x, t - \tau) \quad \text{on } \Gamma_2 \times (0, +\infty), \quad (4)$$

$$y_1(x, t) = y_2(x, t), \quad \text{on } \Gamma_0 \times (0, +\infty), \quad (5)$$

$$a_1 \frac{\partial y_1(x, t)}{\partial \nu} = a_2 \frac{\partial y_2(x, t)}{\partial \nu} \quad \text{on } \Gamma_0 \times (0, +\infty), \quad (6)$$

$$y_2'(x, t - \tau) = f_0(x, t - \tau) \quad \text{on } \Gamma_2 \times (0, \tau). \quad (7)$$

where:

- $a(x) = \begin{cases} a_1, & x \in \Omega_1 \\ a_2, & x \in \Omega_2 \end{cases}$
- $y(x, t) = \begin{cases} y_1(x, t), & (x, t) \in \Omega_1 \times (0, +\infty) \\ y_2(x, t), & (x, t) \in \Omega_2 \times (0, +\infty) \end{cases}$
- $\frac{\partial}{\partial \nu}$ is the normal derivative.
- μ_1 and μ_2 are positive real numbers.
- τ is the time delay
- y^0, y^1, f_0 are the initial data which belong to suitable spaces.

In the absence of delay, that is $\mu_2 = 0$, Liu and Williams [4] have shown that the solution of (1)-(6) decays exponentially to zero in the energy space $H_{\Gamma_1}^1(\Omega) \times L^2(\Omega)$ provided that

$$a_1 > a_2 \quad (8)$$

and $\{\Omega, \Gamma_0, \Gamma_1, \Gamma_2\}$ satisfies (H.1), (H.2), (H.3), and

(H.4) $h(x) \cdot \nu(x) \geq \gamma > 0$.

The purpose of this paper is to investigate the stability of problem (1) – (7) in the case where both μ_1 and μ_2 are different from zero. To this end, assume as in [5] that

$$\mu_1 > \mu_2. \quad (9)$$

and define the energy of a solution of (1) – (7) by

$$E(t) = \frac{1}{2} \int_{\Omega} \left[|y'(x, t)|^2 + a(x) |\nabla(y(x, t))|^2 \right] dx + \frac{\xi}{2} \int_{\Gamma_2} \int_0^1 |y(x, t - \tau\rho)|^2 d\rho d\sigma(x), \quad (10)$$

where

$$a_2\tau\mu_2 < \xi < a_2\tau(2\mu_1 - \mu_2), \quad (11)$$

We show that if $\{\Omega, \Gamma_0, \Gamma_1, \Gamma_2\}$ satisfies (H.1), (H.2) and (H.3), then there is an exponential decay rate for $E(t)$. The proof of this result combines multipliers technique and compactness-uniqueness arguments.

The main result of this paper can be stated as follows.

Theorem 1. *Assume (H1), (H.2), (H.3), (8) and (9). Then there exist constants $M \geq 1$ and $\omega > 0$ such that*

$$E(t) \leq Me^{-\omega t} E(0).$$

Theorem 1 is proved in Section 3. In Section 2, we investigate the well-posedness of system (1) – (7) using semigroup theory.

2 Well-posedness of problem (1) – (7)

Inspired from [5] and [6], we introduce the auxilliary variable $z(x, \rho, t) = y(x, t - \tau\rho)$. With this new unknown, problem (1) – (7) is equivalent to

$$y''(x, t) - a(x)\Delta y(x, t) = 0 \quad \text{in } \Omega \times (0, +\infty), \quad (12)$$

$$y(x, 0) = y^0(x), y'(x, 0) = y^1(x) \quad \text{in } \Omega, \quad (13)$$

$$y(x, t) = 0 \quad \text{on } \Gamma_1 \times (0, +\infty), \quad (14)$$

$$\frac{\partial z(x, \rho, t)}{\partial t} + \frac{1}{\tau} \frac{\partial z(x, \rho, t)}{\partial \rho} = 0 \quad \text{on } \Gamma_2 \times (0, +\infty) \quad (15)$$

$$\frac{\partial y_2(x, t)}{\partial \nu} = -\mu_1 y_2'(x, t) - \mu_2 z(x, 1, t) \quad \text{on } \Gamma_2 \times (0, +\infty), \quad (16)$$

$$y_1(x, t) = y_2(x, t) \quad \text{on } \Gamma_0 \times (0, +\infty), \quad (17)$$

$$a_1 \frac{\partial y_1(x, t)}{\partial \nu} = a_2 \frac{\partial y_2(x, t)}{\partial \nu} \quad \text{on } \Gamma_0 \times (0, +\infty), \quad (18)$$

$$z(x, 0, t) = y'(x, t) \quad \text{on } \Gamma_2 \times (0, +\infty) \quad (19)$$

$$z(x, \rho, 0) = f_0(x, -\tau\rho) \quad \text{on } \Gamma_2 \times (0, 1) \quad (20)$$

Now, we endow the Hilbert space

$$\mathcal{H} = H_{\Gamma_1}^1(\Omega) \times L^2(\Omega) \times L^2(\Gamma_2; L^2(0, 1))$$

with the inner product

$$\left\langle \begin{pmatrix} u \\ v \\ z \end{pmatrix}; \begin{pmatrix} \bar{u} \\ \bar{v} \\ \bar{z} \end{pmatrix} \right\rangle = \int_{\Omega} (a(x)\nabla u(x)\nabla \bar{u}(x) + v(x)\bar{v}(x)) dx + \xi \int_{\Gamma_2} \int_0^1 z(x, \rho) \bar{z}(x, \rho) d\rho d\sigma(x)$$

and define a linear operator in \mathcal{H} by

$$D(A) = \{(u, v, z)^T \in H^2(\Omega_1, \Omega_2, \Gamma_1) \times H_{\Gamma_1}^1(\Omega) \times L^2(\Gamma_2; H^1(0, 1); \frac{\partial u}{\partial \nu} = -\mu_1 v - \mu_2 z(\cdot, 1), v = z(\cdot, 0) \text{ on } \Gamma_2\} \quad (21)$$

$$A \begin{pmatrix} u \\ v \\ z \end{pmatrix} = \begin{pmatrix} v \\ a(x)\Delta u \\ -\tau^{-1} \frac{\partial z}{\partial \rho} \end{pmatrix} \quad (22)$$

The spaces used for the definition of \mathcal{H} and $D(A)$ are

$$H_{\Gamma_1}^1(\Omega) = \{u \in H^1(\Omega) : u = 0 \text{ on } \Gamma_1\}$$

$$H^2(\Omega_1, \Omega_2, \Gamma_1) = \{u_i \in H^2(\Omega_i) : u = 0 \text{ on } \Gamma_1, u_1 = u_2 \text{ and } a_1 \frac{\partial u_1}{\partial \nu} = a_2 \frac{\partial u_2}{\partial \nu} \text{ on } \Gamma_0\}$$

Then we can rewrite (12) – (20) as an abstract Cauchy problem in \mathcal{H}

$$\begin{cases} \frac{d}{dt} Y(t) = AY(t) \\ Y(0) = Y_0 \end{cases} \quad (23)$$

where

$$Y(t) = (y, y', z)^T \text{ and } Y_0 = (y_0, y_1, f_0(\cdot, -\cdot, \tau))^T$$

Proposition 1. *The operator A defined by (21) and (22) generates a strongly continuous semigroup on \mathcal{H} . Thus, for every $Y_0 \in \mathcal{H}$, problem (23) has a unique solution Y whose regularity depends on the initial datum Y_0 as follows:*

$$Y(\cdot) \in C([0, +\infty); \mathcal{H}) \text{ if } Y_0 \in \mathcal{H},$$

$$Y(\cdot) \in C([0, +\infty); D(A)) \cap C^1([0, +\infty); \mathcal{H}) \text{ if } Y_0 \in D(A).$$

Proof. Let $Y = \begin{pmatrix} u \\ v \\ z \end{pmatrix} \in D(A)$. Then

$$\begin{aligned} \langle AY, Y \rangle &= \int_{\Omega} a(x)\nabla u(x) \cdot \nabla v(x) dx + \int_{\Omega} (a(x)\Delta u(x))v(x) dx - \\ &\quad \frac{\xi}{\tau} \int_{\Gamma_2} \int_0^1 z_{\rho}(x, \rho) z(x, \rho) d\rho d\Gamma \end{aligned} \quad (24)$$

Applying Green's first theorem, we obtain

$$\begin{aligned}
\int_{\Omega} (a(x)\Delta u(x))v(x)dx &= a_1 \int_{\Gamma_1} v(x) \frac{\partial u(x)}{\partial \nu} d\Gamma - a_1 \int_{\Omega_1} \nabla u(x) \cdot \nabla v(x) dx + \\
& a_2 \int_{\Gamma_2} v(x) \frac{\partial u(x)}{\partial \nu} d\Gamma - a_2 \int_{\Omega_2} \nabla u(x) \cdot \nabla v(x) dx \\
&= a_2 \int_{\Gamma_2} v(x) \{-\mu_1 v(x) - \mu_2 z(x, 1)\} d\Gamma - \int_{\Omega} a(x) \nabla u(x) \cdot \nabla v(x) dx
\end{aligned} \tag{25}$$

Integrating by parts in ρ , we get

$$\int_{\Gamma_2} \int_0^1 z_{\rho}(x, \rho) z(x, \rho) d\rho d\Gamma = \frac{1}{2} \int_{\Gamma_2} \{z^2(x, 1) - z^2(x, 0)\} d\Gamma \tag{26}$$

Inserting (25) and (26) into (24) results in

$$\begin{aligned}
\langle AY, Y \rangle &= -a_2 \mu_1 \int_{\Gamma_2} v^2(x) d\Gamma - a_2 \mu_2 \int_{\Gamma_2} v(x) z(x, 1) d\Gamma - \\
& \frac{\xi}{2\tau} \int_{\Gamma_2} z^2(x, 1) d\Gamma + \frac{\xi}{2\tau} \int_{\Gamma_2} v^2(x) d\Gamma
\end{aligned}$$

from which follows using the Cauchy-Schwarz inequality

$$\langle AY, Y \rangle \leq -(a_2 \mu_1 - \frac{a_2 \mu_2}{2} + \frac{\xi}{2\tau}) \int_{\Gamma_2} v^2(x) d\Gamma - (\frac{\xi}{2\tau} - \frac{a_2 \mu_2}{2}) \int_{\Gamma_2} z^2(x, 1) d\Gamma \tag{27}$$

(27) implies that

$$\langle AY, Y \rangle \leq 0$$

Thus A is dissipative.

Now we show that for a fixed $\lambda > 0$ and $(g, h, k)^T \in \mathcal{H}$, there exists $Y = (u, v, z)^T \in D(A)$ such that

$$(\lambda I - A)Y = (g, h, k)^T$$

or equivalently

$$\lambda u - v = g \tag{28}$$

$$\lambda v - a(x)\Delta u = h \tag{29}$$

$$\lambda z + \frac{1}{\tau} z_{\rho} = k \tag{30}$$

Suppose that we have found u with the appropriate regularity, then we can determine z . Indeed, from (21) and (30) we have

$$\begin{cases} z_{\rho}(x, \rho) = -\lambda \tau z(x, \rho) + \tau k(x, \rho) \\ z(x, 0) = v(x) \end{cases}$$

The unique solution of the above initial value problem is

$$z(x, \rho) = e^{-\lambda\tau\rho}v(x) + \tau e^{-\lambda\tau\rho} \int_0^\rho e^{\lambda\tau s} k(x, s) ds$$

and in particular

$$z(x, 1) = \lambda e^{-\lambda\tau} u(x) + z_0(x), \quad x \in \Gamma_2$$

where

$$z_0(x) = -e^{-\lambda\tau} g(x) + \tau e^{-\lambda\tau} \int_0^1 e^{\lambda\tau s} k(x, s) ds$$

By (28) and (29), the function u satisfies

$$\lambda^2 u - a(x) \Delta u = h + \lambda g \quad (31)$$

Problem (31) can be reformulated as

$$\int_{\Omega} (\lambda^2 u - a(x) \Delta u) w dx = \int_{\Omega} (h + \lambda g) w dx, \quad w \in H_{\Gamma_1}^1(\Omega) \quad (32)$$

Using Green's first theorem and recalling (21), we express the right-hand side of (32) as follows

$$\begin{aligned} \int_{\Omega} (\lambda^2 u - a(x) \Delta u) w dx &= \int_{\Omega} (\lambda^2 u w + a(x) \nabla u \cdot \nabla w) dx + a_2 \int_{\Gamma_2} \{\mu_1 (\lambda u - g) w \\ &+ \mu_2 (\lambda e^{-\lambda\tau} u(x) + z_0(x)) w\} d\Gamma \end{aligned}$$

Therefore (32), can be rewritten as

$$\begin{aligned} \int_{\Omega} (\lambda^2 u w + a(x) \nabla u \cdot \nabla w) dx + a_2 \int_{\Gamma_2} (\mu_1 + \mu_2 e^{-\lambda\tau}) \lambda u w d\Gamma &= \int_{\Omega} (h + \lambda g) w d\Gamma \\ + a_2 \mu_1 \int_{\Gamma_2} g w d\Gamma - a_2 \mu_2 \int_{\Gamma_2} z_0 w d\Gamma, \quad \forall w \in H_{\Gamma_1}^1(\Omega). \end{aligned} \quad (33)$$

Since the left-hand side of (33) is coercive on $H_{\Gamma_1}^1(\Omega)$, the Lax-Milgram Theorem guarantees the existence and uniqueness of a solution $y \in H_{\Gamma_1}^1(\Omega)$ of (31). If we consider $w \in \mathcal{D}(\Omega)$ in (28), then y is a solution in $\mathcal{D}'(\Omega)$ of

$$\lambda^2 u - a(x) \Delta u = h + \lambda g \quad (34)$$

and thus $\Delta u \in L^2(\Omega)$.

Combining (33) together with (34), we obtain after using Green's first theorem

$$a_2 \int_{\Gamma_2} (\mu_1 + \mu_2 e^{-\lambda\tau}) \lambda u w d\Gamma + a_2 \int_{\Omega} \frac{\partial u}{\partial \nu} w d\Gamma = a_2 \mu_1 \int_{\Gamma_2} g w d\Gamma - a_2 \mu_2 \int_{\Gamma_2} z_0 w d\Gamma$$

which implies that

$$\frac{\partial u(x)}{\partial \nu} = -\mu_1 v(x) - \mu_2 z(x, 1)$$

So, we have found $(u, v, z)^T \in D(A)$ which satisfies (28) – (30). Thus, by the Lumer-Phillips Theorem (see for instance [8], Theorem 1.4.3), generates a strongly continuous semigroup of contractions on \mathcal{H} .

3 Proof of Theorem 1

We prove Theorem 1 for smooth initial data. The general case follows by a standard density argument.

We proceed in several steps.

Step 1.

Since

$$E(t) = \frac{1}{2} \|(y, y', z)\|_{\mathcal{H}}^2$$

Then, we deduce from the proof of Proposition 1 that $E(t)$ is non-increasing and

$$\frac{d}{dt}E(t) \leq -C \int_{\Gamma_2} \{y^2(x, t) + y'^2(x, t)\} d\Gamma \quad (35)$$

where

$$C = \min\{a_2\mu_1 - \frac{a_2\mu_2}{2} + \frac{\xi}{2\tau}, \frac{\xi}{2\tau} - \frac{a_2\mu_2}{2}\}$$

Step 2.

Set

$$E(t) = \mathcal{E}(t) + E_d(t)$$

where

$$\mathcal{E}(t) = \frac{1}{2} \int_{\Omega} \{a(x) |\nabla y(x, t)|^2 + |y'(x, t)|^2\} dx$$

and

$$E_d(t) = \frac{\xi}{2\tau} \int_{\Gamma_2} \int_0^1 |y'(x, t - \tau\rho)|^2 d\rho d\Gamma$$

$E_d(t)$ can be rewritten via a change of variable as

$$E_d(t) = \frac{\xi}{2\tau^2} \int_t^{t+\tau} \int_{\Gamma_2} y'^2(x, s - \tau) d\Gamma ds \quad (36)$$

From (36), we obtain

$$E_d(t) \leq C_1 \int_0^T \int_{\Gamma_2} y'^2(x, s - \tau) d\Gamma ds \quad (37)$$

for $0 \leq t \leq T$ and T large enough.

Step 3.

By applying energy methods (multiplier $2h \cdot \nabla y + (\operatorname{div} h - \alpha)y$) (see the appendix) to problem (1) – (7), we obtain for all $T > 0$.

$$\begin{aligned} \int_0^T \mathcal{E}(t) dt &\leq C_2(\mathcal{E}(0) + \mathcal{E}(T)) + C_3 \int_0^T \int_{\Gamma_2} \{(\frac{\partial y(x, t)}{\partial \nu})^2 + y'^2(x, t)\} d\Gamma dt + \\ &C_4 \int_0^T \int_{\Gamma_2} |\nabla_{\sigma} y(x, t)|^2 d\Gamma dt + C_5 \int_0^T \int_{\Omega} |y(x, t)|^2 d\Omega dt \end{aligned} \quad (38)$$

where $\nabla_\sigma y$ is the tangential gradient of y .

Step 4.

We eliminate the tangential gradient from (38) by using the following estimate due to Lasiecka and Triggiani (Lemma 7.2 in [3])

$$\int_\epsilon^{T-\epsilon} \int_{\Gamma_2} |\nabla_\sigma y(x, t)|^2 d\Gamma dt \leq C_6 \left\{ \int_0^T \int_{\Gamma_2} \left\{ \left(\frac{\partial y(x, t)}{\partial \nu} \right)^2 + y'^2(x, t) \right\} d\Gamma dt + \|y\|_{L^2(0, T; H^{1/2+\delta}(\Omega))}^2 \right\}$$

where ϵ and δ are arbitrary positive constants. We obtain

$$\begin{aligned} \int_0^T \mathcal{E}(t) dt &\leq C_2(\mathcal{E}(0) + \mathcal{E}(T)) + C_7 \int_0^T \int_{\Gamma_2} \left\{ \left(\frac{\partial y(x, t)}{\partial \nu} \right)^2 + y'^2(x, t) \right\} d\Gamma dt + \\ &\quad C_8 \|y\|_{L^2(0, T; H^{1/2+\delta}(\Omega))}^2 \end{aligned} \quad (39)$$

Step 5.

We differentiate $\mathcal{E}(t)$ with respect to t and apply Green's first theorem. We obtain

$$\frac{d}{dt} \mathcal{E}(t) = a_2 \int_{\Gamma_2} y'(x, t) \frac{\partial y(x, t)}{\partial \nu} d\Gamma dt \quad (40)$$

From (40), we get via the Cauchy-Schwarz inequality

$$\mathcal{E}(0) \leq \mathcal{E}(T) + \frac{a_2}{2} \int_0^T \int_{\Gamma_2} \left\{ y'^2(x, t) + \left(\frac{\partial y(x, t)}{\partial \nu} \right)^2 \right\} d\Gamma dt \quad (41)$$

Insertion of (41) into (39) yields

$$\begin{aligned} \int_0^T \mathcal{E}(t) dt &\leq 2C_2 \mathcal{E}(T) + C_9 \int_0^T \int_{\Gamma_2} \left\{ \left(\frac{\partial y(x, t)}{\partial \nu} \right)^2 + y'^2(x, t) \right\} d\Gamma dt + \\ &\quad C_8 \|y\|_{L^2(0, T; H^{1/2+\delta}(\Omega))}^2 \end{aligned} \quad (42)$$

Step 6.

Since $E(t)$ is non-increasing and $E(t) = \mathcal{E}(t) + E_d(t)$, then (42) together with (37) implies that

$$\begin{aligned} TE(T) &\leq 2C_2 \mathcal{E}(T) + C_9 \int_0^T \int_{\Gamma_2} \left\{ \left(\frac{\partial y(x, t)}{\partial \nu} \right)^2 + y'^2(x, t) \right\} d\Gamma dt + \\ &\quad C_8 \|y\|_{L^2(0, T; H^{1/2+\delta}(\Omega))}^2 + TC_1 \int_0^T \int_{\Gamma_2} y'^2(x, t - \tau) d\Gamma dt \end{aligned} \quad (43)$$

Thus invoking again the identity $E(t) = \mathcal{E}(t) + E_d(t)$ and recalling the boundary condition (4), we obtain from (43)

$$E(T) \leq C_{10} \int_0^T \int_{\Gamma_2} \{ y'^2(x, t) + y'^2(x, t - \tau) \} d\Gamma dt + C_{11} \|y\|_{L^2(0, T; H^{1/2+\delta}(\Omega))}^2 \quad (44)$$

for T large enough.

Step 7.

We drop the lower order term on the right-hand side of (44) by a compactness-uniqueness argument to obtain

$$E(T) \leq C_{12} \int_0^T \int_{\Gamma_2} \{y'^2(x, t) + y'^2(x, t - \tau)\} d\Gamma dt \quad (45)$$

Step 8.

The estimate (45) together with (35) yields

$$E(T) \leq \frac{C_{13}}{1 + C_{13}} E(0) \quad (46)$$

The desired conclusion follows now from (46) since the system (1) – (7) is invariant by translation.

References

1. Datko, D.: Not all feedback stabilized hyperbolic systems are robust with respect to small time delays in their feedbacks. *SIAM J. Control Optim.* 26, 697-713 (1988)
2. Datko, R., Lagnese, J., Polis, M.P.: An example on the effect of time delays in boundary feedback stabilization of wave equations. *SIAM J. Control Optim.* 24, 152-156 (1986)
3. Lasiecka, I., Triggiani, R.: Uniform stabilization of the wave equation with Dirichlet or Neumann feedback control without geometrical conditions. *Appl. Math. Optim.* 25, 189-244 (1992)
4. Liu, W., Williams, G.H.: The exponential stability of the problem of transmission of the wave equation. *Bull. Austr. Math. Soc.* 97, 305-327 (1998)
5. Nicaise, S., Pignotti, C.: Stability and instability results of the wave equation with a delay term in the boundary or internal feedbacks. *SIAM J. Control Optim.* 45, 1561-1585 (2006)
6. Nicaise, N., Rebiai, S.E.: Stabilization of the Schrödinger equation with a delay term in boundary feedback or internal feedback. *Portugal. Math.* 68, 19-39 (2011).
7. Nicaise, S., Valein, J.: Stabilization of second order evolution equations with unbounded feedback with delay. *ESAIM Control Optim. Calc. Var.* 16, 420-456 (2010)
8. Pazy, A.: *Semigroups of linear operators and applications to partial differential equations.* Springer-Verlag, New York (1983)
9. Xu, G.Q., Yung, S.P., Li, L.K.: Stabilization of wave systems with input delay in the boundary control, *ESAIM Control Optim. Calc. Var.* 12, 770-785 (2006)

Appendix: Sketch of Proof of (38)

We multiply both sides of (1) by $2h.\nabla y + (divh - \alpha)y$ and integrate over $(0, T) \times \Omega$. We obtain

$$\begin{aligned}
& 2 \int_0^T \int_{\Omega} a(x) J \nabla y . \nabla y d\Omega dt + \alpha \int_0^T \int_{\Omega} \{y'^2 - a(x) |\nabla y|^2\} d\Omega dt = \\
& - \int_{\Omega} \{2y'h.\nabla y + (divh - \alpha)y'y\}_0^T d\Omega - \int_0^T \int_{\Omega} a(x) y \nabla y . \nabla (divh - \alpha) d\Omega dt + \\
& a_1 \int_0^T \int_{\Gamma_1} \left| \frac{\partial y_1}{\partial \nu} \right|^2 h.\nu d\Gamma dt - (a_1 - a_2) \int_0^T \int_{\Gamma_0} |\nabla y_1|^2 h.\nu d\Gamma dt - \\
& \frac{(a_1 - a_2)^2}{a_2} \int_0^T \int_{\Gamma_0} \left| \frac{\partial y_1}{\partial \nu} \right|^2 h.\nu d\Gamma dt + \int_0^T \int_{\Gamma_2} |y_2'|^2 h.\nu d\Gamma dt + \\
& 2a_2 \int_0^T \int_{\Gamma_2} \left| \frac{\partial y_2}{\partial \nu} \right|^2 h.\nabla y_2 d\Gamma dt - a_2 \int_0^T \int_{\Gamma_2} |\nabla y_2|^2 h.\nu d\Gamma dt + \\
& a_2 \int_0^T \int_{\Gamma_2} \left| \frac{\partial y_2}{\partial \nu} \right|^2 (divh - \alpha) d\Gamma dt
\end{aligned} \tag{47}$$

after using the boundary conditions (3) and (5). Identity (47) is used together with (H.1), (H.2), (H.3) and (8) to obtain estimate (38).