



# Modelling of seismic waves propagation in harmonic domain by hybridizable discontinuous Galerkin method (HDG)

M. Bonnasse-Gahot<sup>1,2</sup>, H. Calandra<sup>3</sup>, J. Diaz<sup>1</sup> and S. Lanteri<sup>2</sup>

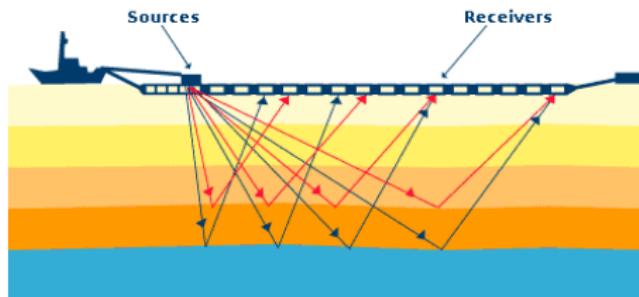
<sup>1</sup> INRIA Bordeaux-Sud-Ouest, team-project Magique 3D

<sup>2</sup> INRIA Sophia-Antipolis-Méditerranée, team-project Nachos

<sup>3</sup> TOTAL Exploration-Production

# Motivation

## Examples of seismic applications



# Motivation

Imaging method : the full wave inversion

- ▶ Iterative procedure
- ▶ Inverse problem requiring to solve a **lot of forward problems**

# Motivation

Imaging method : the full wave inversion

- ▶ Iterative procedure
- ▶ Inverse problem requiring to solve a **lot of forward problems**

Seismic imaging : time-domain or harmonic-domain ?

- ▶ Time-domain : **imaging condition complicated** but **low computational cost**
- ▶ Harmonic-domain : **imaging condition simple** but **huge computational cost**

## Motivation

Imaging method : the full wave inversion

- ▶ Iterative procedure
- ▶ Inverse problem requiring to solve a **lot of forward problems**

Seismic imaging : time-domain or harmonic-domain ?

- ▶ Time-domain : **imaging condition complicated** but **low computational cost**
- ▶ Harmonic-domain : **imaging condition simple** but **huge computational cost**

Forward problem of the inversion process

- ▶ Elastic wave propagation in harmonic domain : **Helmholtz equation**
- ▶ Reduction of the size of the linear system

# Motivation

Seismic imaging in heterogeneous complex media

- ▶ Complex topography
- ▶ High heterogeneities

# Motivation

Seismic imaging in heterogeneous complex media

- ▶ Complex topography
- ▶ High heterogeneities

Use of unstructured meshes with FE methods

# Motivation

Seismic imaging in heterogeneous complex media

- ▶ Complex topography
- ▶ High heterogeneities

Use of unstructured meshes with FE methods

DG method

- ▶ Flexible choice of interpolation orders ( $p$  – adaptativity)
- ▶ Highly parallelizable method
- ▶ Increased computational cost as compared to classical FEM

# Motivation

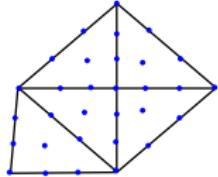
Seismic imaging in heterogeneous complex media

- ▶ Complex topography
- ▶ High heterogeneities

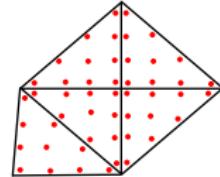
Use of unstructured meshes with FE methods

DG method

- ▶ Flexible choice of interpolation orders ( $p$  – adaptativity)
- ▶ Highly parallelizable method
- ▶ Increased computational cost as compared to classical FEM



DOF of classical FEM



DOF of DGM

# Motivation

## Objective of this work

- ▶ Development of an hybridizable DG (HDG) method
- ▶ Comparison with a reference method : a standard nodal DG method

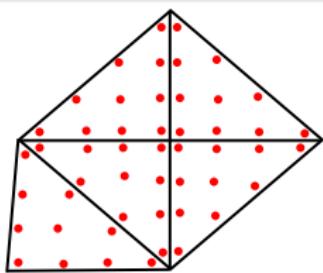


FIGURE : Degrees of freedom of DGM

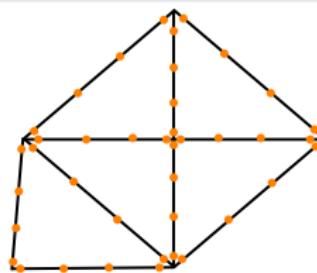


FIGURE : Degrees of freedom of HDGM

# HDG methods

## HDG methods

- ▶ **B. Cockburn, J. Gopalakrishnan, R. Lazarov** *Unified hybridization of discontinuous Galerkin, mixed and continuous Galerkin methods for second order elliptic problems*, SIAM Journal on Numerical Analysis, Vol. 47 (2009)
- ▶ **S. Lanteri, L. Li, R. Perrussel**, *Numerical investigation of a high order hybridizable discontinuous Galerkin method for 2d time-harmonic Maxwell's equations*, COMPEL, Vol. 32 (2013) (time-harmonic domain)
- ▶ **N.C. Nguyen, J. Peraire, B. Cockburn**, *High-order implicit hybridizable discontinuous Galerkin methods for acoustics and elastodynamics*, J. of Comput. Physics, Vol. 230 (2011) (time domain for seismic applications)

# Contents

3D Helmholtz elastic equations

Notations and definitions

Hybridizable Discontinuous Galerkin method

2D Numerical results

3D preliminary numerical results

Conclusions-Perspectives

# 3D Helmholtz elastic equations

First order formulation of Helmholtz wave equations

$$\mathbf{x} = (x, y, z) \in \Omega \subset \mathbb{R}^3,$$

$$\begin{cases} i\omega \rho(\mathbf{x}) \mathbf{v}(\mathbf{x}) = \nabla \cdot \underline{\underline{\sigma}}(\mathbf{x}) + \underline{f}_s(\mathbf{x}) \\ i\omega \underline{\underline{\sigma}}(\mathbf{x}) = \underline{\underline{C}}(\mathbf{x}) \underline{\underline{\varepsilon}}(\mathbf{v}(\mathbf{x})) \end{cases}$$

- ▶ Free surface condition :  $\underline{\underline{\sigma}} \mathbf{n} = 0$  on  $\Gamma_f$
- ▶ Absorbing boundary condition :  $\underline{\underline{\sigma}} \mathbf{n} = P \mathbf{A} \mathbf{P}^T \mathbf{v}$  on  $\Gamma_a$

- ▶  $\mathbf{v}$  : velocity vector
- ▶  $\underline{\underline{\sigma}}$  : stress tensor
- ▶  $\underline{\underline{\varepsilon}}$  : strain tensor

# 3D Helmholtz elastic equations

First order formulation of Helmholtz wave equations

$$\mathbf{x} = (x, y, z) \in \Omega \subset \mathbb{R}^3,$$

$$\begin{cases} i\omega \rho(\mathbf{x}) \mathbf{v}(\mathbf{x}) = \nabla \cdot \underline{\underline{\sigma}}(\mathbf{x}) + \mathbf{f}_s(\mathbf{x}) \\ i\omega \underline{\underline{\sigma}}(\mathbf{x}) = \underline{\underline{C}}(\mathbf{x}) \underline{\underline{\varepsilon}}(\mathbf{v}(\mathbf{x})) \end{cases}$$

- ▶ Free surface condition :  $\underline{\underline{\sigma}} \mathbf{n} = 0$  on  $\Gamma_f$
- ▶ Absorbing boundary condition :  $\underline{\underline{\sigma}} \mathbf{n} = P \mathbf{A} \mathbf{P}^T \mathbf{v}$  on  $\Gamma_a$

- |  |  |
|--|--|
| <ul style="list-style-type: none"> <li>▶ <math>\rho</math> : mass density</li> <li>▶ <math>\underline{\underline{C}}</math> : tensor of elasticity coefficients</li> </ul> | <ul style="list-style-type: none"> <li>▶ <math>v_p</math> : P-wave velocity</li> <li>▶ <math>v_s</math> : S-wave velocity</li> <li>▶ <math>f_s</math> : source term, <math>f_s \in L^2(\Omega)</math></li> </ul> |
|--|--|

# Contents

3D Helmholtz elastic equations

Notations and definitions

Hybridizable Discontinuous Galerkin method

2D Numerical results

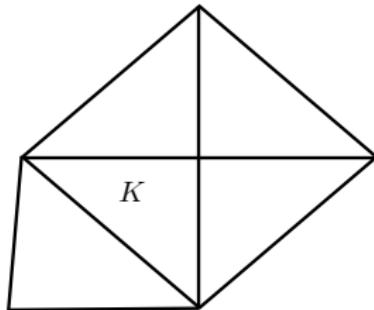
3D preliminary numerical results

Conclusions-Perspectives

# Notations and definitions

## Notations

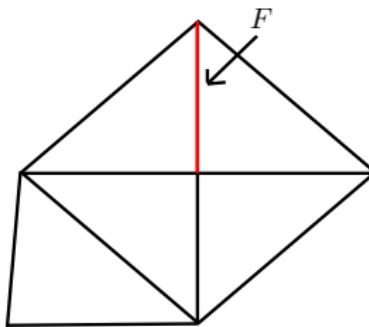
- $\mathcal{T}_h$  mesh of  $\Omega$  composed of tetrahedrons  $K$



# Notations and definitions

## Notations

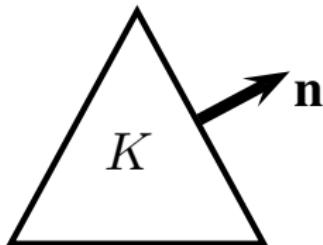
- ▶  $\mathcal{T}_h$  mesh of  $\Omega$  composed of tetrahedrons  $K$
- ▶  $\mathcal{F}_h$  set of all faces  $F$  of  $\mathcal{T}_h$



# Notations and definitions

## Notations

- ▶  $\mathcal{T}_h$  mesh of  $\Omega$  composed of tetrahedrons  $K$
- ▶  $\mathcal{F}_h$  set of all faces  $F$  of  $\mathcal{T}_h$
- ▶  $\mathbf{n}$  the normal outward vector of an element  $K$



# Notations and definitions

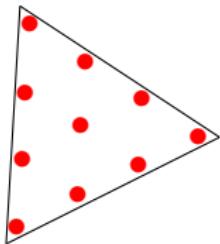
## Approximations spaces

- ▶  $P_p(K)$  set of polynomials of degree at most  $p$  on  $K$

# Notations and definitions

## Approximations spaces

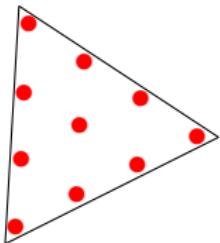
- ▶  $P_p(K)$  set of polynomials of degree at most  $p$  on  $K$
- ▶  $\mathbf{V}_h^p = \{\mathbf{v} \in (L^2(\Omega))^3 : \mathbf{v}|_K \in \mathbf{V}^p(K) = (P_p(K))^3, \forall K \in \mathcal{T}_h\}$



# Notations and definitions

## Approximations spaces

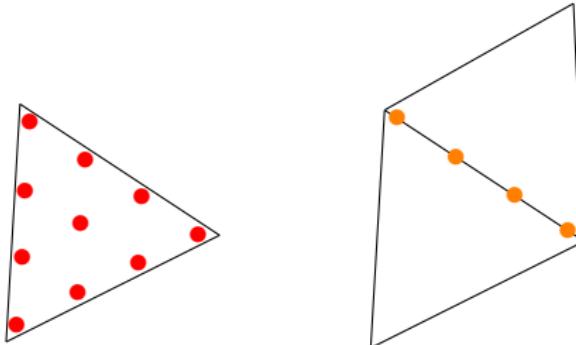
- ▶  $P_p(K)$  set of polynomials of degree at most  $p$  on  $K$
- ▶  $\mathbf{V}_h^p = \{\mathbf{v} \in (L^2(\Omega))^3 : \mathbf{v}|_K \in \mathbf{V}^p(K) = (P_p(K))^3, \forall K \in \mathcal{T}_h\}$
- ▶  $\Sigma_h^p = \{\underline{\underline{\sigma}} \in (L^2(\Omega))^6 : \underline{\underline{\sigma}}|_K \in \Sigma^p(K) = (P_p(K))^6, \forall K \in \mathcal{T}_h\}$



# Notations and definitions

## Approximations spaces

- ▶  $P_p(K)$  set of polynomials of degree at most  $p$  on  $K$
- ▶  $\mathbf{V}_h^p = \{\mathbf{v} \in (L^2(\Omega))^3 : \mathbf{v}|_K \in \mathbf{V}^p(K) = (P_p(K))^3, \forall K \in \mathcal{T}_h\}$
- ▶  $\Sigma_h^p = \{\underline{\underline{\sigma}} \in (L^2(\Omega))^6 : \underline{\underline{\sigma}}|_K \in \Sigma^p(K) = (P_p(K))^6, \forall K \in \mathcal{T}_h\}$
- ▶  $\mathbf{M}_h = \{\eta \in (L^2(\mathcal{F}_h))^3 : \eta|_F \in (P_p(F))^3, \forall F \in \mathcal{F}_h\}$



# Notations and definitions

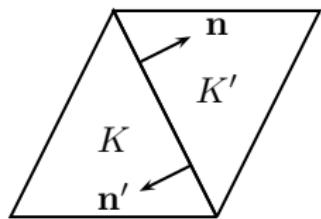
## Definitions

- ▶ Jump  $[\![\cdot]\!]$  of a vector  $\mathbf{v}$  through  $F$  :

$$[\![\mathbf{v}]\!] = \mathbf{v}^+ \cdot \mathbf{n}^+ + \mathbf{v}^- \cdot \mathbf{n}^- = \mathbf{v}^+ \cdot \mathbf{n}^+ - \mathbf{v}^- \cdot \mathbf{n}^+$$

- ▶ Jump of a tensor  $\underline{\underline{\sigma}}$  through  $F$  :

$$[\![\underline{\underline{\sigma}}]\!] = \underline{\underline{\sigma}}^+ \mathbf{n}^+ + \underline{\underline{\sigma}}^- \mathbf{n}^- = \underline{\underline{\sigma}}^+ \mathbf{n}^+ - \underline{\underline{\sigma}}^- \mathbf{n}^+$$



# Contents

3D Helmholtz elastic equations

Notations and definitions

**Hybridizable Discontinuous Galerkin method**

Formulation

Algorithm

2D Numerical results

3D preliminary numerical results

Conclusions-Perspectives

# HDG formulation of the equations

## Local HDG formulation

$$\begin{cases} i\omega \rho \mathbf{v} - \nabla \cdot \underline{\underline{\sigma}} = 0 \\ i\omega \underline{\sigma} - \underline{\underline{C}_\varepsilon}(\mathbf{v}) = 0 \end{cases}$$

# HDG formulation of the equations

## Local HDG formulation

$$\left\{ \begin{array}{l} \int_K i\omega \rho^K \mathbf{v}^K \cdot \mathbf{w} + \int_K \underline{\underline{\sigma}}^K : \nabla \mathbf{w} - \int_{\partial K} \hat{\underline{\underline{\sigma}}}^{\partial K} \cdot \mathbf{n} \cdot \mathbf{w} = 0 \\ \int_K i\omega \underline{\underline{\sigma}}^K : \underline{\underline{\xi}} + \int_K \mathbf{v}^K \cdot \nabla \cdot (\underline{\underline{C}}^K \underline{\underline{\xi}}) - \int_{\partial K} \hat{\mathbf{v}}^{\partial K} \cdot \underline{\underline{C}}^K \underline{\underline{\xi}} \cdot \mathbf{n} = 0 \end{array} \right.$$

$\hat{\underline{\underline{\sigma}}}^K$  and  $\hat{\mathbf{v}}^K$  are numerical traces of  $\underline{\underline{\sigma}}^K$  and  $\mathbf{v}^K$  respectively on  $\partial K$

# HDG formulation of the equations

## Local HDG formulation

$$\begin{cases} \int_K i\omega \rho^K \mathbf{v}^K \cdot \mathbf{w} + \int_K \underline{\underline{\sigma}}^K : \nabla \mathbf{w} - \int_{\partial K} \hat{\underline{\sigma}}^{\partial K} \cdot \mathbf{n} \cdot \mathbf{w} = 0 \\ \int_K i\omega \underline{\underline{\sigma}}^K : \underline{\underline{\xi}} + \int_K \mathbf{v}^K \cdot \nabla \cdot (\underline{\underline{C}}^K \underline{\underline{\xi}}) - \int_{\partial K} \hat{\mathbf{v}}^{\partial K} \cdot \underline{\underline{C}}^K \underline{\underline{\xi}} \cdot \mathbf{n} = 0 \end{cases}$$

We define :

$$\begin{aligned} \hat{\mathbf{v}}^F &= \lambda^F, & \forall F \in \mathcal{F}_h, \\ \hat{\underline{\sigma}}^{\partial K} \cdot \mathbf{n} &= \underline{\underline{\sigma}}^K \cdot \mathbf{n} - \tau \mathbf{I} (\mathbf{v}^K - \lambda^{\partial K}), & \text{on } \partial K \end{aligned}$$

where  $\tau$  is the stabilization parameter ( $\tau > 0$ )

# HDG formulation of the equations

## Local HDG formulation

We replace  $\hat{\mathbf{v}}^K$  and  $(\underline{\underline{\sigma}}^K \cdot \mathbf{n})$  by their definitions into the local equations

$$\left\{ \begin{array}{l} \int_K i\omega \rho^K \mathbf{v}^K \cdot \mathbf{w} + \int_K \underline{\underline{\sigma}}^K : \nabla \mathbf{w} - \int_{\partial K} \underline{\underline{\sigma}}^K \cdot \mathbf{n} \cdot \mathbf{w} \\ \qquad \qquad \qquad + \int_{\partial K} \tau \mathbf{I} (\mathbf{v}^K - \lambda^{\partial K}) \cdot \mathbf{w} = 0 \\ \int_K i\omega \underline{\underline{\sigma}}^K : \underline{\underline{\xi}} + \int_K \mathbf{v}^K \cdot \nabla \cdot (\underline{\underline{C}}^K \underline{\underline{\xi}}) - \int_{\partial K} \lambda^{\partial K} \cdot \underline{\underline{C}}^K \underline{\underline{\xi}} \cdot \mathbf{n} = 0 \end{array} \right.$$

# HDG formulation of the equations

## Local HDG formulation

$$\left\{ \begin{array}{l} \int_K i\omega \rho^K \mathbf{v}^K \cdot \mathbf{w} - \int_K (\nabla \cdot \underline{\underline{\sigma}}^K) \cdot \mathbf{w} + \int_{\partial K} \tau \mathbf{I} (\mathbf{v}^K - \lambda^{\partial K}) \cdot \mathbf{w} = 0 \\ \int_K i\omega \underline{\underline{\sigma}}^K : \underline{\underline{\xi}} + \int_K \mathbf{v}^K \cdot \nabla \cdot (\underline{\underline{C}}^K \underline{\underline{\xi}}) - \int_{\partial K} \lambda^{\partial K} \cdot \underline{\underline{C}}^K \underline{\underline{\xi}} \cdot \mathbf{n} = 0 \end{array} \right.$$

# HDG formulation of the equations

## Local HDG formulation

$$\left\{ \begin{array}{l} \int_K i\omega \underline{\rho}^K \underline{\mathbf{v}}^K \cdot \mathbf{w} - \int_K (\nabla \cdot \underline{\underline{\sigma}}^K) \cdot \mathbf{w} + \int_{\partial K} \tau \mathbf{I} (\underline{\mathbf{v}}^K - \underline{\lambda}^{\partial K}) \cdot \mathbf{w} = 0 \\ \int_K i\omega \underline{\underline{\sigma}}^K : \underline{\underline{\xi}} + \int_K \underline{\mathbf{v}}^K \cdot \nabla \cdot (\underline{\underline{\mathbf{C}}}^K \underline{\underline{\xi}}) - \int_{\partial K} \underline{\lambda}^{\partial K} \cdot \underline{\underline{\mathbf{C}}}^K \underline{\underline{\xi}} \cdot \mathbf{n} = 0 \end{array} \right.$$

We define :

$$\underline{\mathbf{W}}^K = \left( \underline{\underline{\mathbf{v}}}_x^K, \underline{\underline{\mathbf{v}}}_y^K, \underline{\underline{\mathbf{v}}}_z^K, \underline{\underline{\sigma_{xx}}}^K, \underline{\underline{\sigma_{yy}}}^K, \underline{\underline{\sigma_{zz}}}^K, \underline{\underline{\sigma_{xy}}}^K, \underline{\underline{\sigma_{xz}}}^K, \underline{\underline{\sigma_{yz}}}^K \right)^T$$

$$\underline{\Lambda} = \left( \underline{\Lambda}^{F_1}, \underline{\Lambda}^{F_2}, \dots, \underline{\Lambda}^{F_{n_f}} \right)^T, \text{ where } n_f = \text{card}(\mathcal{F}_h)$$

## Discretization of the local HDG formulation

$$\underline{\mathbb{A}}^K \underline{\mathbf{W}}^K + \underline{\mathbb{C}}^K \underline{\Lambda} = 0$$

# HDG formulation of the equations

## Transmission condition

In order to determine  $\lambda^K$ , the continuity of the normal component of  $\hat{\underline{\sigma}}^K$  is weakly enforced, rendering this numerical trace conservative :

$$\int_F \llbracket \hat{\underline{\sigma}}^K \cdot \mathbf{n} \rrbracket \cdot \eta = 0$$

# HDG formulation of the equations

## Transmission condition

In order to determine  $\lambda^K$ , the continuity of the normal component of  $\hat{\underline{\sigma}}^K$  is weakly enforced, rendering this numerical trace conservative :

$$\int_F \llbracket \hat{\underline{\sigma}}^K \cdot \mathbf{n} \rrbracket \cdot \eta = 0$$

Replacing  $(\hat{\underline{\sigma}}^K \cdot \mathbf{n})$  and summing over all faces, the transmission condition becomes :

$$\sum_{K \in \mathcal{T}_h} \int_{\partial K} (\underline{\underline{\sigma}}^K \cdot \mathbf{n}) \cdot \eta - \sum_{K \in \mathcal{T}_h} \int_{\partial K} \tau \mathbf{I} (\mathbf{v}^K - \lambda^{\partial K}) \cdot \eta = 0$$

# HDG formulation of the equations

## Transmission condition

In order to determine  $\lambda^K$ , the continuity of the normal component of  $\hat{\underline{\sigma}}^K$  is weakly enforced, rendering this numerical trace conservative :

$$\int_F [\hat{\underline{\sigma}}^K \cdot \mathbf{n}] \cdot \eta = 0$$

Replacing  $(\hat{\underline{\sigma}}^K \cdot \mathbf{n})$  and summing over all faces, the transmission condition becomes :

$$\sum_{K \in \mathcal{T}_h} \int_{\partial K} (\underline{\sigma}^K \cdot \mathbf{n}) \cdot \eta - \sum_{K \in \mathcal{T}_h} \int_{\partial K} \tau \mathbf{I} (\mathbf{v}^K - \lambda^{\partial K}) \cdot \eta = 0$$

## Discretization of the transmission condition

$$\sum_{K \in \mathcal{T}_h} [\mathbb{B}^K \underline{W}^K + \mathbb{L}^K \underline{\Lambda}] = 0$$

# HDG formulation of the equations

## Global HDG formulation

$$\left\{ \begin{array}{l} \int_K i\omega \underline{\rho}^K \underline{\mathbf{v}}^K \cdot \mathbf{w} - \int_K (\nabla \cdot \underline{\underline{\sigma}}^K) \cdot \mathbf{w} + \int_{\partial K} \tau \mathbf{I} (\underline{\mathbf{v}}^K - \underline{\lambda}^{\partial K}) \cdot \mathbf{w} = 0 \\ \int_K i\omega \underline{\underline{\sigma}}^K : \underline{\underline{\xi}} + \int_K \underline{\mathbf{v}}^K \cdot \nabla \cdot (\underline{\underline{C}}^K \underline{\underline{\xi}}) - \int_{\partial K} \underline{\lambda}^{\partial K} \cdot \underline{\underline{C}}_K \underline{\underline{\xi}} \cdot \mathbf{n} = 0 \\ \sum_{K \in \mathcal{T}_h} \int_{\partial K} (\underline{\underline{\sigma}}^K \cdot \mathbf{n}) \cdot \eta - \sum_{K \in \mathcal{T}_h} \int_{\partial K} \tau \mathbf{I} (\underline{\mathbf{v}}^K - \underline{\lambda}^{\partial K}) \cdot \eta = 0 \end{array} \right.$$

# HDG formulation of the equations

## Global HDG formulation

$$\left\{ \begin{array}{l} \int_K i\omega \underline{\rho}^K \underline{\mathbf{v}}^K \cdot \mathbf{w} - \int_K (\nabla \cdot \underline{\underline{\sigma}}^K) \cdot \mathbf{w} + \int_{\partial K} \tau \mathbf{I} (\underline{\mathbf{v}}^K - \underline{\lambda}^{\partial K}) \cdot \mathbf{w} = 0 \\ \int_K i\omega \underline{\underline{\sigma}}^K : \underline{\underline{\xi}} + \int_K \underline{\mathbf{v}}^K \cdot \nabla \cdot (\underline{\underline{\mathcal{C}}}^K \underline{\underline{\xi}}) - \int_{\partial K} \underline{\lambda}^{\partial K} \cdot \underline{\underline{\mathcal{C}}}^K \underline{\underline{\xi}} \cdot \mathbf{n} = 0 \\ \sum_{K \in \mathcal{T}_h} \int_{\partial K} (\underline{\underline{\sigma}}^K \cdot \mathbf{n}) \cdot \eta - \sum_{K \in \mathcal{T}_h} \int_{\partial K} \tau \mathbf{I} (\underline{\mathbf{v}}^K - \underline{\lambda}^{\partial K}) \cdot \eta = 0 \end{array} \right.$$

## Global HDG discretization

$$\left\{ \begin{array}{l} \underline{\mathbb{A}}^K \underline{\mathbb{W}}^K + \underline{\mathbb{C}}^K \underline{\Lambda} = 0 \\ \sum_{K \in \mathcal{T}_h} [\underline{\mathbb{B}}^K \underline{\mathbb{W}}^K + \underline{\mathbb{L}}^K \underline{\Lambda}] = 0 \end{array} \right.$$

# HDG formulation of the equations

## Global HDG formulation

$$\left\{ \begin{array}{l} \int_K i\omega \rho^K \mathbf{v}^K \cdot \mathbf{w} - \int_K (\nabla \cdot \underline{\underline{\sigma}}^K) \cdot \mathbf{w} + \int_{\partial K} \tau \mathbf{I} (\mathbf{v}^K - \lambda^{\partial K}) \cdot \mathbf{w} = 0 \\ \int_K i\omega \underline{\underline{\sigma}}^K : \underline{\underline{\xi}} + \int_K \mathbf{v}^K \cdot \nabla \cdot (\underline{\underline{C}}^K \underline{\underline{\xi}}) - \int_{\partial K} \lambda^{\partial K} \cdot \underline{\underline{C}}_K \underline{\underline{\xi}} \cdot \mathbf{n} = 0 \\ \sum_{K \in \mathcal{T}_h} \int_{\partial K} (\underline{\underline{\sigma}}^K \cdot \mathbf{n}) \cdot \eta - \sum_{K \in \mathcal{T}_h} \int_{\partial K} \tau \mathbf{I} (\mathbf{v}^K - \lambda^{\partial K}) \cdot \eta = 0 \end{array} \right.$$

## Global HDG discretization

$$\left\{ \begin{array}{l} \underline{\underline{W}}^K = -(\mathbb{A}^K)^{-1} \mathbb{C}^K \underline{\underline{\Lambda}} \\ \sum_{K \in \mathcal{T}_h} [\mathbb{B}^K \underline{\underline{W}}^K + \mathbb{L}^K \underline{\underline{\Lambda}}] = 0 \end{array} \right.$$

# HDG formulation of the equations

## Global HDG formulation

$$\left\{ \begin{array}{l} \int_K i\omega \underline{\rho}^K \underline{\mathbf{v}}^K \cdot \mathbf{w} - \int_K (\nabla \cdot \underline{\underline{\sigma}}^K) \cdot \mathbf{w} + \int_{\partial K} \tau \mathbf{I} (\underline{\mathbf{v}}^K - \underline{\lambda}^{\partial K}) \cdot \mathbf{w} = 0 \\ \int_K i\omega \underline{\underline{\sigma}}^K : \underline{\underline{\xi}} + \int_K \underline{\mathbf{v}}^K \cdot \nabla \cdot (\underline{\underline{\mathcal{C}}}^K \underline{\underline{\xi}}) - \int_{\partial K} \underline{\lambda}^{\partial K} \cdot \underline{\underline{\mathcal{C}}}_K \underline{\underline{\xi}} \cdot \mathbf{n} = 0 \\ \sum_{K \in \mathcal{T}_h} \int_{\partial K} (\underline{\underline{\sigma}}^K \cdot \mathbf{n}) \cdot \eta - \sum_{K \in \mathcal{T}_h} \int_{\partial K} \tau \mathbf{I} (\underline{\mathbf{v}}^K - \underline{\lambda}^{\partial K}) \cdot \eta = 0 \end{array} \right.$$

## Global HDG discretization

$$\sum_{K \in \mathcal{T}_h} [-\mathbb{B}^K (\mathbb{A}^K)^{-1} \mathbb{C}^K + \mathbb{L}^K] \underline{\Lambda} = 0$$

# Main idea of the algorithm using the HDG formulation

---

1. Construction of the linear system  $\underline{\mathbf{M}}\underline{\Lambda}$

with  $\underline{\mathbf{M}} = \sum_{K \in \mathcal{T}_h} \left[ -\underline{\mathbb{B}}^K (\underline{\mathbb{A}}^K)^{-1} \underline{\mathbb{C}}^K + \underline{\mathbb{L}}^K \right]$

**for**  $K = 1$  to  $Nb_{tri}$  **do**

    Compute matrices  $\underline{\mathbb{B}}^K, (\underline{\mathbb{A}}^K)^{-1}, \underline{\mathbb{C}}^K$  and  $\underline{\mathbb{L}}^K$

    Construction of the corresponding section of  $\underline{\mathbf{M}}$

**end for**

---

# Main idea of the algorithm using the HDG formulation

- 
1. Construction of the linear system  $\underline{M}\underline{\Lambda}$
  2. Construction of the right hand side  $\underline{S}$
-

# Main idea of the algorithm using the HDG formulation

- 
1. Construction of the linear system  $\mathbf{M}\boldsymbol{\Lambda}$
  2. Construction of the right hand side  $\mathbf{S}$
  3. Resolution  $\mathbf{M}\boldsymbol{\Lambda} = \mathbf{S}$
-

# Main idea of the algorithm using the HDG formulation

- 
1. Construction of the linear system  $\underline{\mathbf{M}}\underline{\Lambda}$
  2. Construction of the right hand side  $\mathbb{S}$
  3. Resolution  $\underline{\mathbf{M}}\underline{\Lambda} = \mathbb{S}$
  4. Computation of the solutions of the initial problem

```
for  $K = 1$  to  $Nb_{tri}$  do
    Compute  $\underline{\mathbf{W}}^K = -(\mathbb{A}^K)^{-1}\mathbb{C}^K\underline{\Lambda}$ 
end for
```

---

# Contents

3D Helmholtz elastic equations

Notations and definitions

Hybridizable Discontinuous Galerkin method

## 2D Numerical results

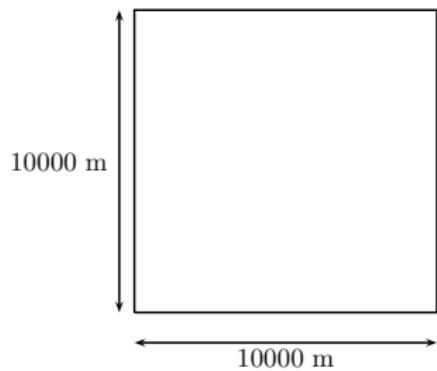
Plane wave in an homogeneous medium

Marmousi test-case

3D preliminary numerical results

Conclusions-Perspectives

# Plane wave



Computational domain  $\Omega$  setting

- ▶ Physical parameters :

- ▶  $\rho = 2000 \text{ kg.m}^{-3}$
- ▶  $\lambda = 16 \text{ GPa}$
- ▶  $\mu = 8 \text{ GPa}$

- ▶ Plane wave :

$$u = \nabla e^{i(k \cos \theta x + k \sin \theta y)}$$

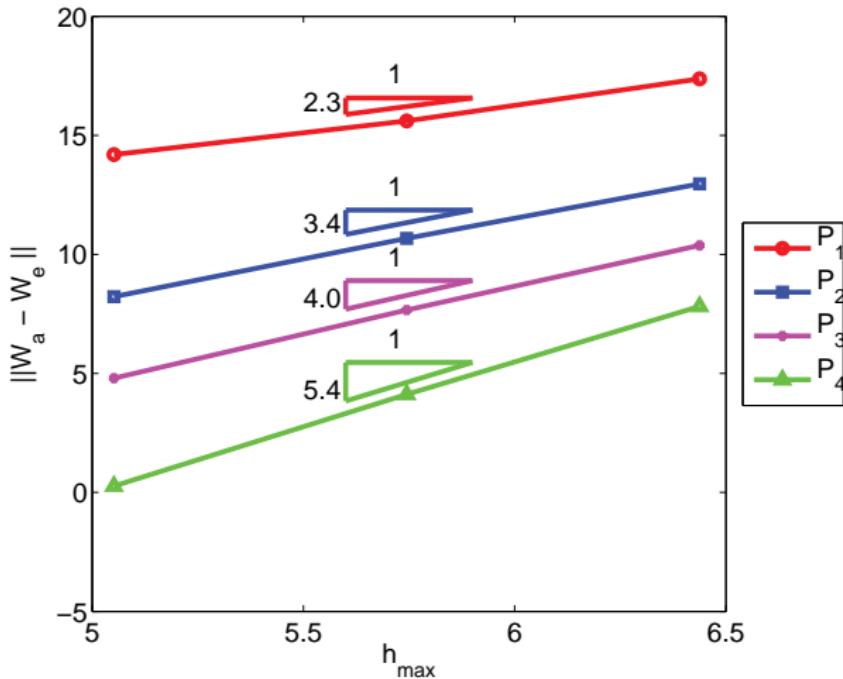
$$\text{where } k = \frac{\omega}{v_p}$$

- ▶  $\theta = 0$

- ▶ Three meshes :

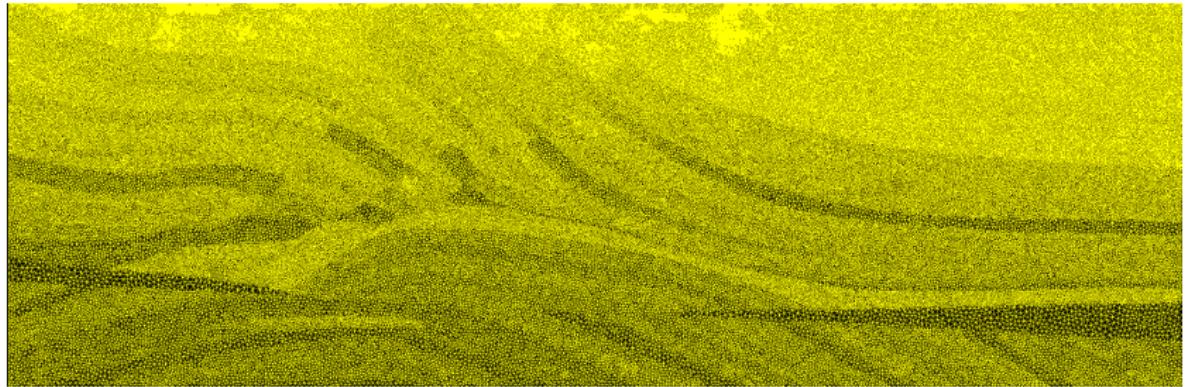
- ▶ 3000 elements
- ▶ 10000 elements
- ▶ 45000 elements

## Plane wave



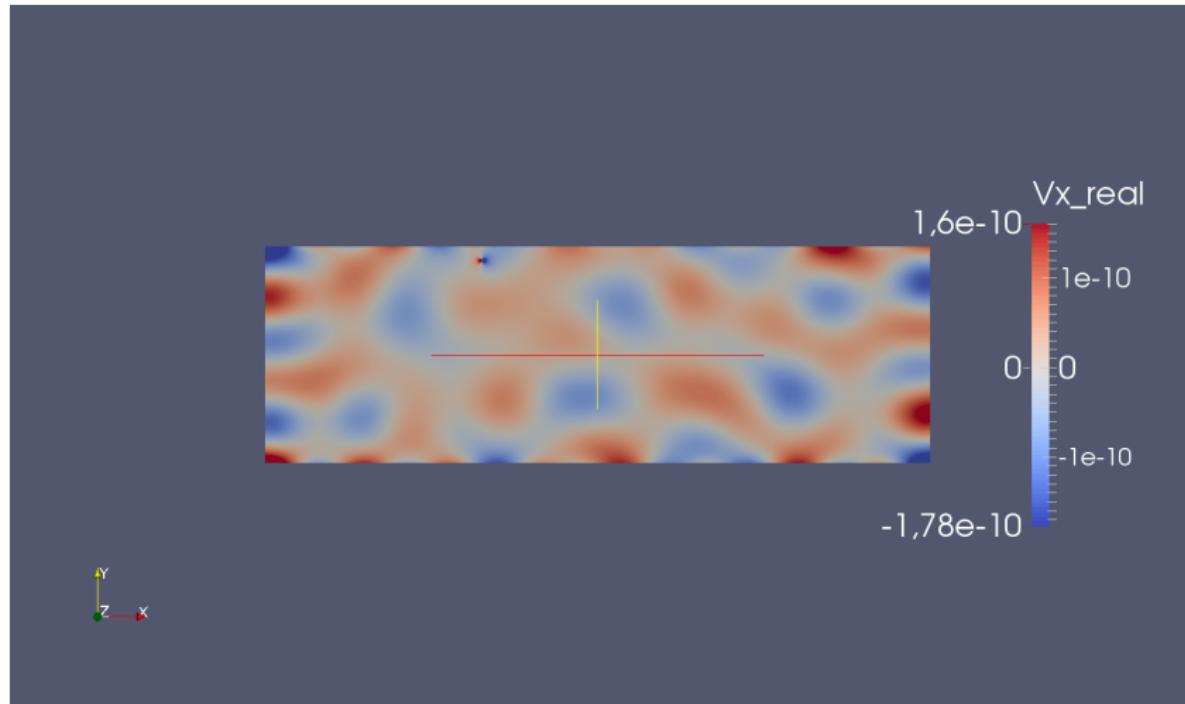
Convergence order of the HDG scheme

# Marmousi test-case



Computational domain  $\Omega$  composed of 235000 triangles

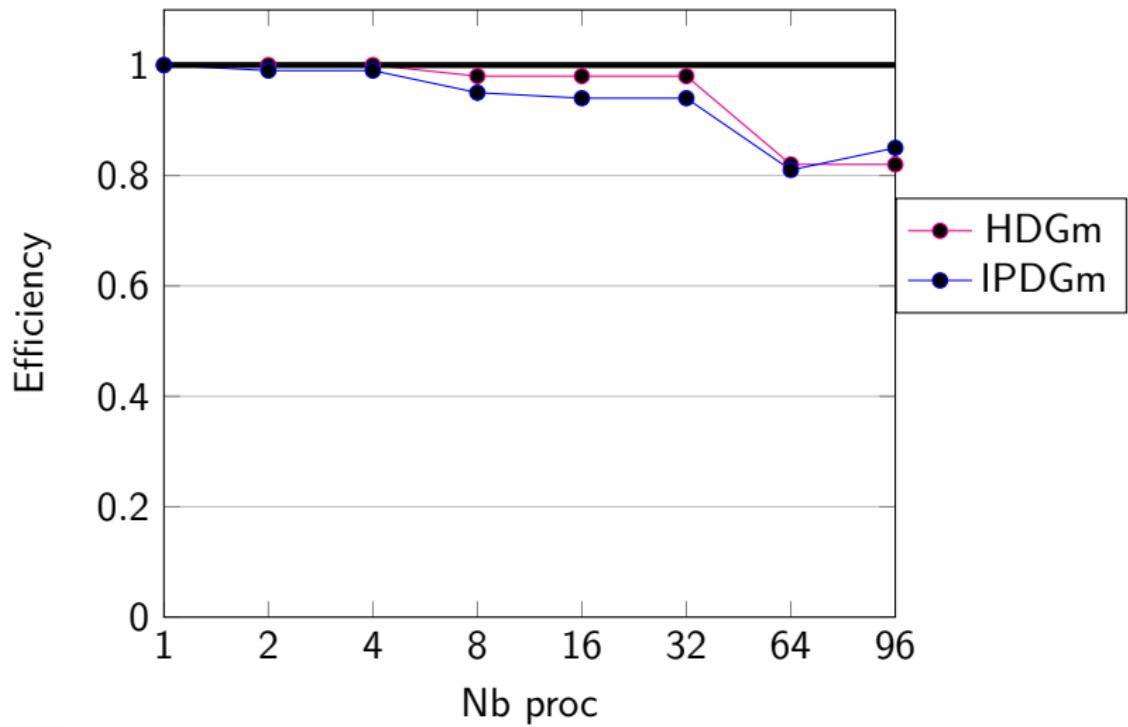
# Parallel results for the Marmousi test-case with the HDG-P3 scheme, $f = 2\text{Hz}$



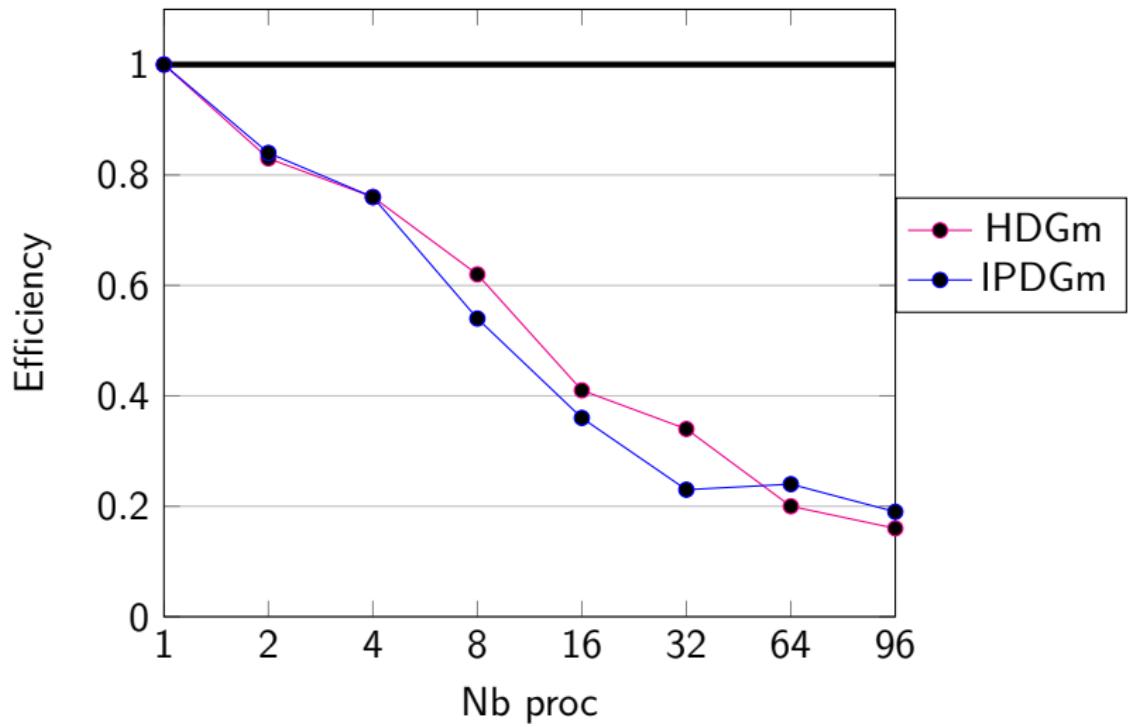
# Characteristics of the computing processors used

- **Plafrim** platform
- Hardware specification : 16 nodes, 12 cores by nodes
- Characteristics of computing nodes :
  - ▶ 2 Hexa-core Westmere Intel® Xeon® X5670
  - ▶ Frequency : 2,93 GHz
  - ▶ Cache L3 : 12 Mo
  - ▶ RAM : 96 Go
  - ▶ Infiniband DDR : 20Gb/s
  - ▶ Ethernet : 1Gb/s

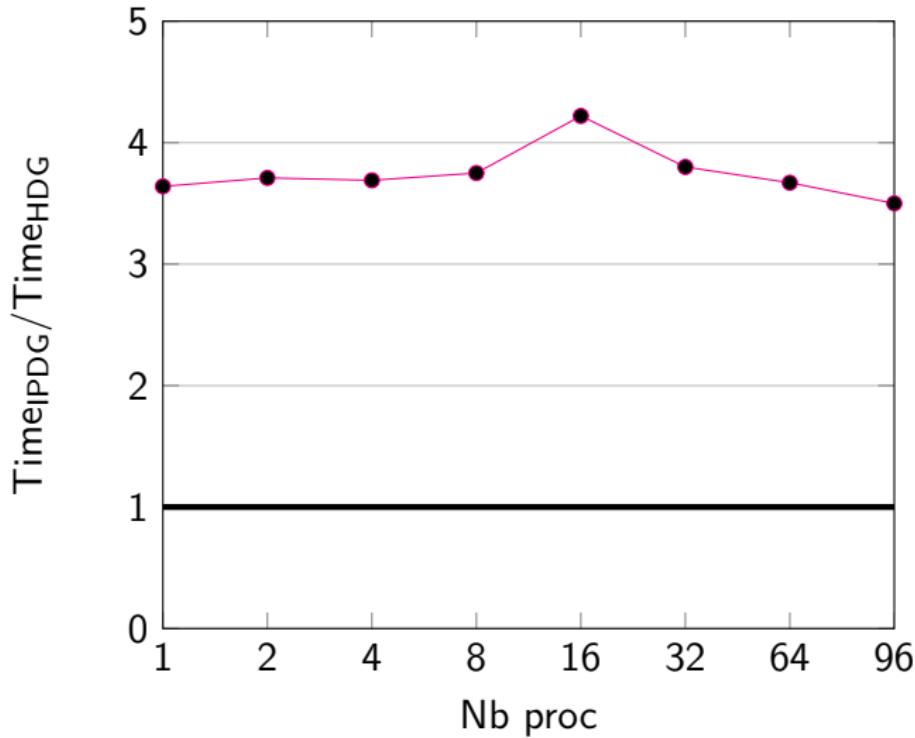
# Efficiency of the parallelism of the global matrix construction



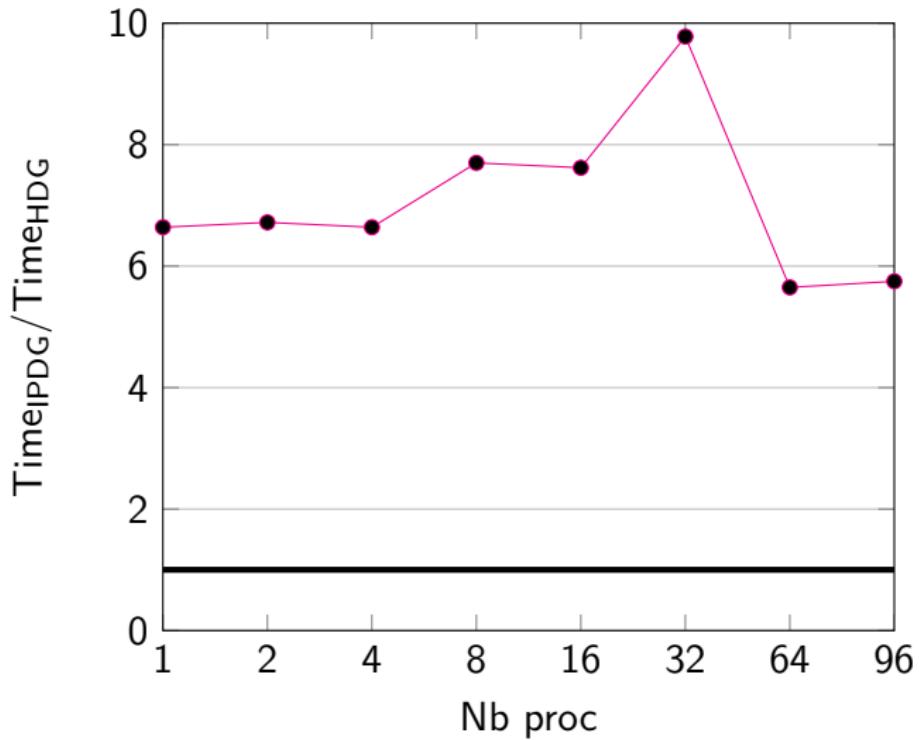
# Efficiency of the parallelism of the whole algorithm



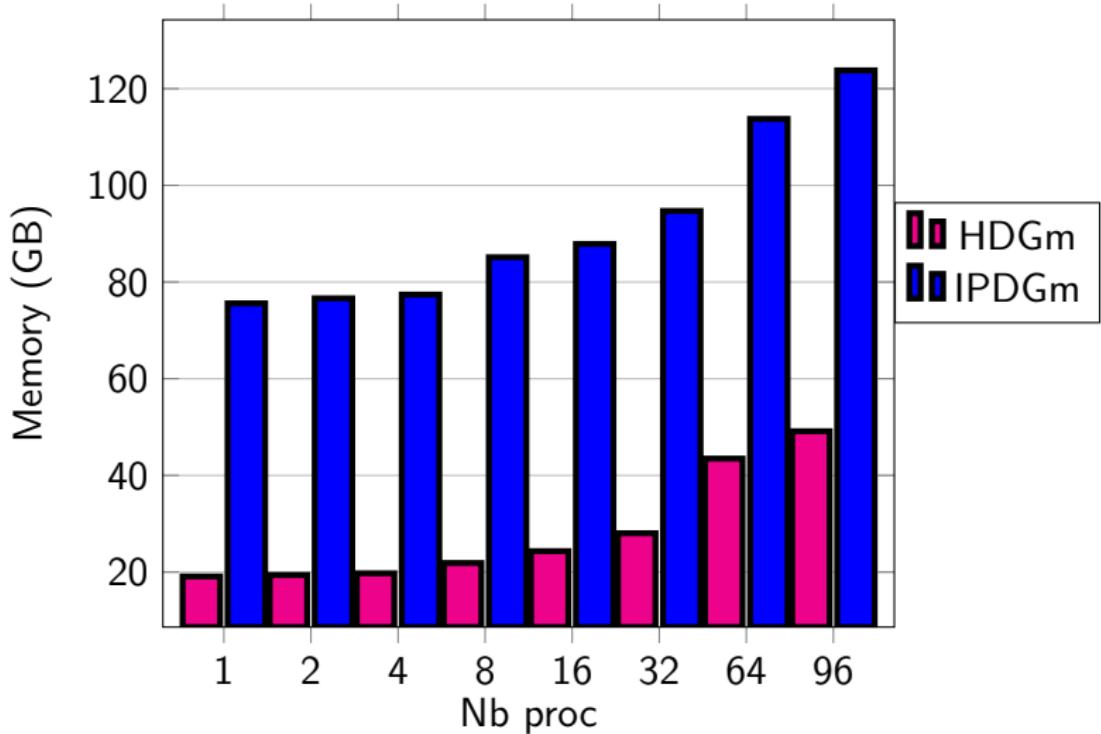
# Speed up for the global matrix construction



## Speed up (Total simulation time)



# Memory required (GB) for the simulation



# Contents

3D Helmholtz elastic equations

Notations and definitions

Hybridizable Discontinuous Galerkin method

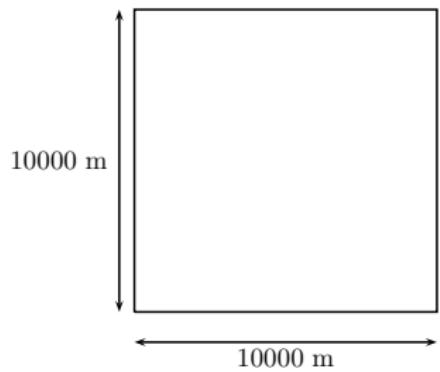
2D Numerical results

3D preliminary numerical results

Plane wave in an homogeneous medium

Conclusions-Perspectives

# Plane wave



Computational domain  $\Omega$   
setting

- ▶ Physical parameters :

- ▶  $\rho = 2000 \text{ kg.m}^{-3}$
- ▶  $\lambda = 16 \text{ GPa}$
- ▶  $\mu = 8 \text{ GPa}$

- ▶ Plane wave :

$$u = \nabla e^{i(k \cos \theta x + k \sin \theta y)}$$

$$\text{where } k = \frac{\omega}{v_p}$$

- ▶  $\theta = 0$
- ▶ Two meshes :
  - ▶ 2600 elements
  - ▶ 6400 elements

# 3D preliminary numerical results

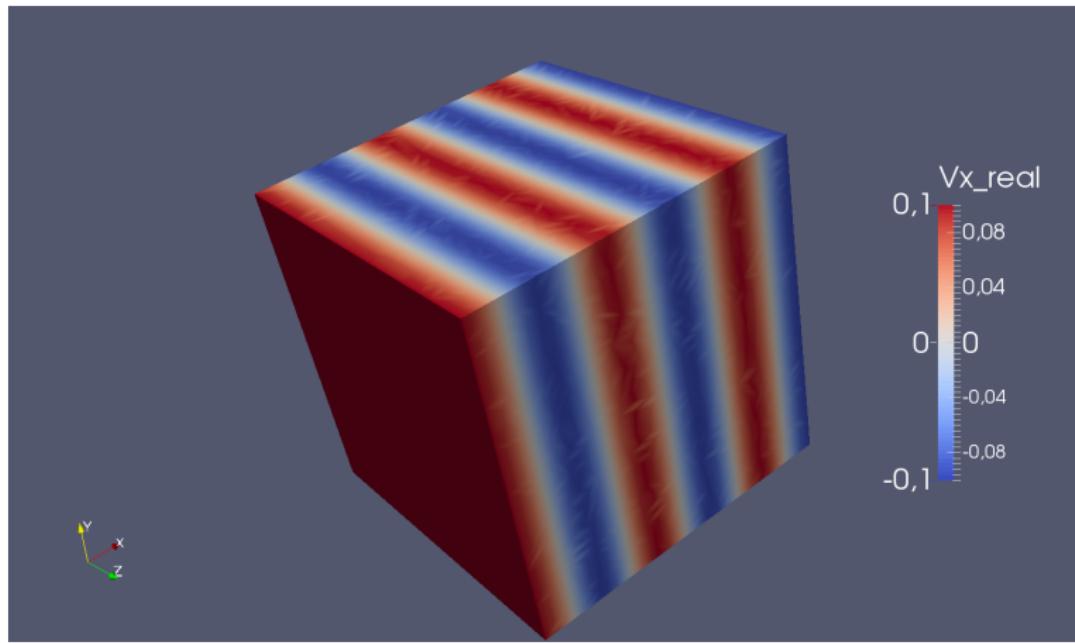


FIGURE : Exact solution

# 3D preliminary numerical results

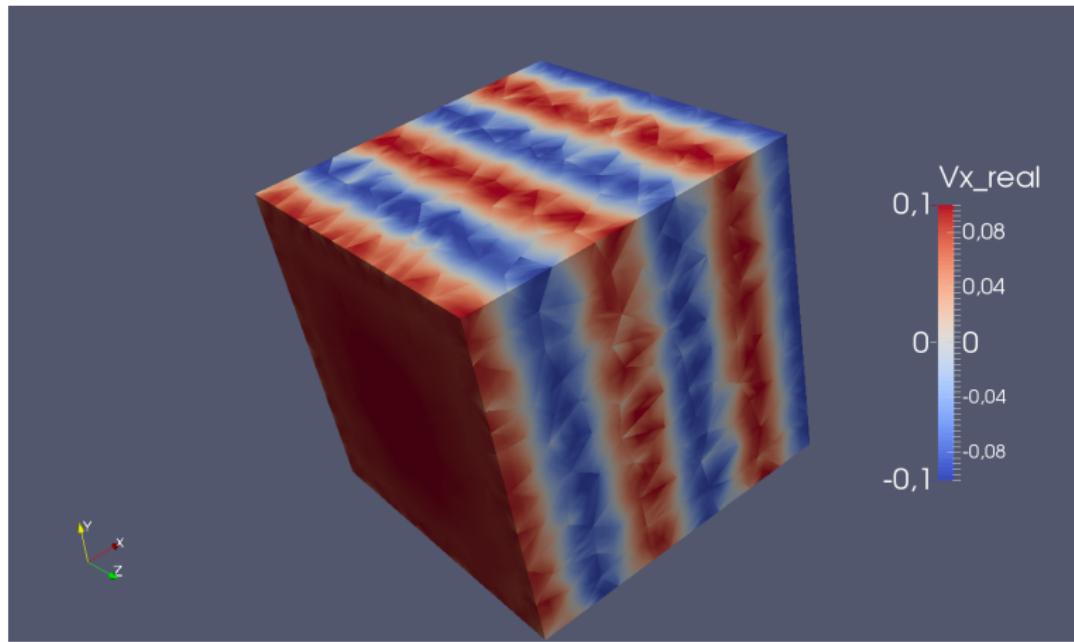


FIGURE : Numerical solution computed with  $P_3$ -HDG scheme on mesh composed of 6400 tetrahedrons

Elements	Order	CPU Time (s)		Memory (MB)	
		HDG	IPDG	HDG	IPDG
2600	1	11		572	
6400	1	59		1981	
2600	2	91		2254	
6400	2	498		7894	
2600	3	508		6242	
6400	3	2386		21899	

Elements	Order	CPU Time (s)		Memory (MB)	
		HDG	IPDG	HDG	IPDG
2600	1	1	1.09	1	0.8
6400	1	1	0.9	1	0.8
2600	2	1	1.35	1	1.19
6400	2	1	1.83	1	1.35
2600	3	1	2.77	1	2.14
6400	3	1	4.35	1	2.19

# Contents

3D Helmholtz elastic equations

Notations and definitions

Hybridizable Discontinuous Galerkin method

2D Numerical results

3D preliminary numerical results

Conclusions-Perspectives

# Conclusions-Perspectives

## Conclusions

On a same mesh, with the 2D HDG method :

- ▶ Memory gain
- ▶ Computational time gain

# Conclusions-Perspectives

## Conclusions

On a same mesh, with the 2D HDG method :

- ▶ Memory gain
- ▶ Computational time gain

## Perspectives

- ▶ Study of the 3D HDG algorithm for Helmholtz equations
- ▶ Solution strategy for the HDG linear system

Thank you !

