

# Modeling of elastic Helmholtz equations by hybridizable discontinuous Galerkin method (HDG) for geophysical applications 

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## Motivation

## Examples of seismic applications


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## Motivation

Imaging method : the full wave inversion

- Iterative procedure
- Inverse problem requiring to solve many direct problems
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## Motivation

Imaging method: the full wave inversion

- Iterative procedure
- Inverse problem requiring to solve many direct problems

Seismic imaging : time-domain or harmonic-domain?

- Time-domain : imaging condition complicated but low computational cost
- Harmonic-domain : imaging condition simple but huge computational cost


## Motivation

Imaging method : the full wave inversion

- Iterative procedure
- Inverse problem requiring to solve many direct problems

Seismic imaging : time-domain or harmonic-domain?

- Time-domain : imaging condition complicated but low computational cost
- Harmonic-domain : imaging condition simple but huge computational cost

Forward problem of the inversion process

- Elastic wave propagation in harmonic domain: Helmholtz equation
- Reduction of the size of the linear system


## Motivation

Seismic imaging in heterogeneous complex media

- Complex topography
- High heterogeneities


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Use of unstructured meshes with FE methods

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Use of unstructured meshes with FE methods
DG method

- Flexible choice of interpolation orders ( $p$ - adaptativity)
- Highly parallelizable method
- Increased computational cost as compared to classical FEM


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DOF of classical FEM


DOF of DGM
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## Motivation

Objective of this work

- Development of an hybridizable DG (HDG) method
- Comparison with a reference method : a standard nodal DG method


Figure: Degrees of freedom of DGM


Figure: Degrees of freedom of HDGM

## HDG methods

HDG methods

- B. Cockburn, J. Gopalakrishnan, R. Lazarov Unified hybridization of discontinuous Galerkin, mixed and continuous Galerkin methods for second order elliptic problems, SIAM Journal on Numerical Analysis, Vol. 47 (2009)
- S. Lanteri, L. Li, R. Perrussel, Numerical investigation of a high order hybridizable discontinuous Galerkin method for 2d time-harmonic Maxwell's equations, COMPEL, Vol. 32 (2013) (time-harmonic domain)
- N.C. Nguyen, J. Peraire, B. Cockburn, High-order implicit hybridizable discontinuous Galerkin methods for acoustics and elastodynamics, J. of Comput. Physics, Vol. 230 (2011) (time domain for seismic applications)


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## 3D Helmholtz elastic equations

## Notations and definitions

## Hybridizable Discontinuous Galerkin method

2D Numerical results

3D preliminary numerical results

Conclusions-Perspectives
M. Bonnasse-Gahot - Modeling of Helmholtz equations by HDGm

## 3D Helmholtz elastic equations

First order formulation of Helmholtz wave equations
$\mathbf{x}=(x, y, z) \in \Omega \subset \mathbb{R}^{3}$,

$$
\left\{\begin{array}{l}
i \omega \rho(\mathbf{x}) \mathbf{v}(\mathbf{x})=\nabla \cdot \underline{\underline{\sigma}}(\mathbf{x})+f_{s}(\mathbf{x}) \\
i \omega \underline{\underline{\sigma}}(\mathbf{x})=\underline{\underline{C}}(\mathbf{x}) \underline{\underline{\varepsilon}}(\mathbf{v}(\mathbf{x}))
\end{array}\right.
$$

- Free surface condition : $\underline{\underline{\sigma}}=0$ on $\Gamma_{\text {, }}$
- Absorbing boundary condition : $\underline{\underline{\sigma}}=P A P^{T} v$ on $\Gamma_{a}$
- v : velocity vector
- $\underline{\underline{\sigma}}$ : stress tensor
- $\underline{\underline{\varepsilon}}$ : strain tensor


## 3D Helmholtz elastic equations

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i \omega \underline{\underline{\underline{\sigma}}(\mathbf{x})=\underline{\underline{C}}(\mathbf{x}) \underline{\underline{\varepsilon}}(\mathbf{v}(\mathbf{x}))} \text {. }
\end{array}\right.
$$

- Free surface condition : $\underline{\underline{\sigma}} \mathbf{n}=0$ on $\Gamma_{\text {, }}$
- Absorbing boundary condition : $\underline{\underline{n}}=P A P^{T} v$ on $\Gamma_{a}$
- $\rho$ : mass density
- C : tensor of elasticity coefficients
- $f_{s}$ : source term, $f_{s} \in L^{2}(\Omega)$


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## Notations and definitions

Notations

- $\mathcal{T}_{h}$ mesh of $\Omega$ composed of tetrahedrons $K$

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- $\mathcal{T}_{h}$ mesh of $\Omega$ composed of tetrahedrons $K$
- $\mathcal{F}_{h}$ set of all faces F of $\mathcal{T}_{h}$

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## Notations and definitions

Notations

- $\mathcal{T}_{h}$ mesh of $\Omega$ composed of tetrahedrons $K$
- $\mathcal{F}_{h}$ set of all faces F of $\mathcal{T}_{h}$
- $\mathbf{n}$ the normal outward vector of an element $K$



## Notations and definitions

Approximations spaces

- $P_{p}(K)$ set of polynomials of degree at most $p$ on $K$
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## Notations and definitions

Approximations spaces

- $P_{p}(K)$ set of polynomials of degree at most $p$ on $K$
- $\mathbf{V}_{h}^{p}=\left\{\mathbf{v} \in\left(L^{2}(\Omega)\right)^{3}:\left.\quad \mathbf{v}\right|_{K} \in \mathbf{V}^{p}(K)=\left(P_{p}(K)\right)^{3}, \forall K \in \mathcal{T}_{h}\right\}$



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Approximations spaces

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- $\Sigma_{h}^{p}=\left\{\underline{\underline{\sigma}} \in\left(L^{2}(\Omega)\right)^{6}: \quad \underline{\underline{\sigma}} \mid K \in \Sigma^{p}(K)=\left(P_{p}(K)\right)^{6}, \forall K \in \mathcal{T}_{h}\right\}$



## Notations and definitions

Approximations spaces

- $P_{p}(K)$ set of polynomials of degree at most $p$ on $K$
- $\mathbf{V}_{h}^{p}=\left\{\mathbf{v} \in\left(L^{2}(\Omega)\right)^{3}:\left.\quad \mathbf{v}\right|_{K} \in \mathbf{V}^{p}(K)=\left(P_{p}(K)\right)^{3}, \forall K \in \mathcal{T}_{h}\right\}$
- $\Sigma_{h}^{p}=\left\{\underline{\underline{\sigma}} \in\left(L^{2}(\Omega)\right)^{6}: \quad \underline{\underline{\sigma}} \mid K \in \Sigma^{p}(K)=\left(P_{p}(K)\right)^{6}, \forall K \in \mathcal{T}_{h}\right\}$
- $\mathbf{M}_{h}=\left\{\eta \in\left(L^{2}\left(\mathcal{F}_{h}\right)\right)^{3}:\left.\quad \eta\right|_{F} \in\left(P_{p}(F)\right)^{3}, \forall F \in \mathcal{F}_{h}\right\}$



## Notations and definitions

## Definitions

- Jump $\llbracket \rrbracket$ of a vector $v$ through $F$ :

$$
\llbracket v \rrbracket=v^{+} \cdot \mathbf{n}^{+}+v^{-} \cdot \mathbf{n}^{-}=v^{+} \cdot \mathbf{n}^{+}-v^{-} \cdot \mathbf{n}^{+}
$$

- Jump of a tensor $\underline{\underline{\sigma}}$ through $F$ :


$$
\llbracket \underline{\underline{\sigma}} \rrbracket=\underline{\underline{\sigma}}^{+} \mathbf{n}^{+}+\underline{\underline{\sigma}}^{-} \mathbf{n}^{-}=\underline{\underline{\sigma}}^{+} \mathbf{n}^{+}-\underline{\underline{\sigma}}^{-} \mathbf{n}^{+}
$$

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## HDG formulation of the equations

Local HDG formulation

$$
\left\{\begin{array}{l}
i \omega \rho \mathbf{v}-\nabla \cdot \underline{\underline{\sigma}}=0 \\
i \omega \underline{\underline{\sigma}}-\underline{\underline{C_{\varepsilon}}}(\mathbf{v})=0
\end{array}\right.
$$

## HDG formulation of the equations

Local HDG formulation

$$
\left\{\begin{array}{l}
\int_{K} i \omega \rho^{K} \mathbf{v}^{K} \cdot \mathbf{w}+\int_{K} \underline{\underline{\sigma}}^{K}: \nabla \mathbf{w}-\int_{\partial K} \underline{\underline{\widehat{\sigma}}}^{\partial K} \cdot \mathbf{n} \cdot \mathbf{w}=0 \\
\int_{K} i \omega \underline{\underline{\sigma}}^{K}: \underline{\underline{\xi}}+\int_{K} \mathbf{v}^{K} \cdot \nabla \cdot\left(\underline{\underline{C}}^{K} \underline{\underline{\xi}}\right)-\int_{\partial K} \widehat{\mathbf{v}}^{\partial K} \cdot \underline{\underline{C}}^{K} \underline{\underline{\xi}} \cdot \mathbf{n}=0
\end{array}\right.
$$

$\underline{\underline{\sigma}}^{K}$ and $\widehat{\mathbf{v}}^{K}$ are numerical traces of $\underline{\underline{\sigma}}^{K}$ and $\mathbf{v}^{K}$ respectively on $\partial K$

## HDG formulation of the equations

Local HDG formulation

$$
\left\{\begin{array}{l}
\int_{K} i \omega \rho^{K} \mathbf{v}^{K} \cdot \mathbf{w}+\int_{K} \underline{\underline{\sigma}}^{K}: \nabla \mathbf{w}-\int_{\partial K}{\underline{\underline{\widehat{\sigma}^{2}}}}^{\partial K} \cdot \mathbf{n} \cdot \mathbf{w}=0 \\
\int_{K} i \omega \underline{\underline{\sigma}}^{K}: \underline{\underline{\xi}}+\int_{K} \mathbf{v}^{K} \cdot \nabla \cdot\left(\underline{\underline{C}}^{K} \underline{\underline{\xi}}\right)-\int_{\partial K} \widehat{\mathbf{v}}^{\partial K} \cdot \underline{\underline{C}}^{K} \underline{\underline{\xi}} \cdot \mathbf{n}=0
\end{array}\right.
$$

We define :

$$
\begin{array}{lll}
\widehat{\mathbf{v}}^{\partial K} & =\widehat{\mathbf{v}}^{F}, & \\
\widehat{\mathbf{v}}^{F} & =\lambda^{F}, & \forall F \in \mathcal{F}_{h}, \\
\underline{\underline{\sigma}}^{\partial K} \cdot \mathbf{n} & =\underline{\underline{\sigma}}^{K} \cdot \mathbf{n}-\tau \mathbf{I}\left(\mathbf{v}^{K}-\lambda^{F}\right), & \text { on } \partial K
\end{array}
$$

where $\tau$ is the stabilization parameter $(\tau>0)$
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## HDG formulation of the equations

Local HDG formulation
We replace $\widehat{v}^{K}$ and $\left(\hat{\underline{\sigma}}^{K} \cdot \mathbf{n}\right)$ by their definitions into the local equations

$$
\left\{\begin{aligned}
& \int_{K} i \omega \rho^{K} \mathbf{v}^{K} \cdot \mathbf{w}+\int_{K} \underline{\underline{\sigma}}^{K}: \nabla \mathbf{w}-\int_{\partial K} \underline{\underline{\sigma}}^{K} \cdot \mathbf{n} \cdot \mathbf{w} \\
&+\int_{\partial K} \tau \mathbf{l}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \mathbf{w}=0 \\
& \int_{K} i \omega_{\underline{\sigma^{K}}}: \underline{\underline{\xi}}+\int_{K} \mathbf{v}^{K} \cdot \nabla \cdot\left(\underline{\underline{C}}^{K} \underline{\underline{\xi}}\right)-\int_{\partial K} \lambda^{F} \cdot \underline{\underline{C}}^{K} \underline{\underline{\xi}} \cdot \mathbf{n}=0
\end{aligned}\right.
$$

## HDG formulation of the equations

Local HDG formulation

$$
\left\{\begin{array}{l}
\int_{K} i \omega \rho^{K} \mathbf{v}^{K} \cdot \mathbf{w}-\int_{K}\left(\nabla \cdot \underline{\underline{\sigma}}^{K}\right) \cdot \mathbf{w}+\int_{\partial K} \tau \mathbf{l}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \mathbf{w}=0 \\
\int_{K} i \omega \underline{\underline{\sigma}}^{K}: \underline{\underline{\xi}}+\int_{K} \mathbf{v}^{K} \cdot \nabla \cdot\left(\underline{\underline{C}}^{K} \underline{\underline{\xi}}\right)-\int_{\partial K} \lambda^{F} \cdot \underline{\underline{C}}^{K} \underline{\underline{\xi}} \cdot \mathbf{n}=0
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\int_{K} i \omega \underline{\underline{\sigma}}^{K}: \underline{\underline{\xi}}+\int_{K} \mathbf{v}^{K} \cdot \nabla \cdot\left(\underline{\underline{C}}^{K} \underline{\underline{\xi}}\right)-\int_{\partial K} \lambda^{F} \cdot \underline{\underline{C}}^{K} \underline{\underline{\xi}} \cdot \mathbf{n}=0
\end{array}\right.
$$

We define :
$\underline{W}^{k}=\left(\underline{V_{x}}{ }^{k},{\underline{V_{y}}}^{k},{\underline{V_{z}}}^{k},{\underline{\sigma_{x x}}}^{k},{\underline{\sigma_{y y}}}^{k},{\underline{\sigma_{z z}}}^{k}, \underline{\sigma_{x y}}{ }^{k},{\underline{\sigma_{x z}}}^{k},{\underline{\sigma_{y z}}}^{k}\right)^{T}$
$\underline{\Lambda}=\left(\underline{\Lambda}^{F_{1}}, \underline{\Lambda}^{F_{2}}, \ldots, \underline{\Lambda}^{F_{n_{f}}}\right)^{T}$, where $n_{f}=\operatorname{card}\left(\mathcal{F}_{h}\right)$
Discretization of the local HDG formulation

$$
\mathbb{A}^{K} \underline{W}^{K}+\sum_{F \in \partial K} \mathbb{C}^{K, F} \underline{\Lambda}=0
$$

## HDG formulation of the equations

Local HDG formulation

$$
\left\{\begin{array}{l}
\int_{K} i \omega \rho^{K} \mathbf{v}^{K} \cdot \mathbf{w}-\int_{K}\left(\nabla \cdot \underline{\underline{\sigma}}^{K}\right) \cdot \mathbf{w}+\int_{\partial K} \tau \mathbf{l}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \mathbf{w}=0 \\
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\end{array}\right.
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$\underline{\Lambda}=\left(\underline{\Lambda}^{F_{1}}, \underline{\Lambda}^{F_{2}}, \ldots, \underline{\Lambda}^{F_{n f}}\right)^{T}$, where $n_{f}=\operatorname{card}\left(\mathcal{F}_{h}\right)$
Discretization of the local HDG formulation

$$
\mathbb{A}^{K} \underline{W}^{K}+\mathbb{C}^{K} \underline{\Lambda}=0
$$

## HDG formulation of the equations

Transmission condition
In order to determine $\lambda^{K}$, the continuity of the normal component of $\underline{\underline{\sigma}}^{K}$ is weakly enforced, rendering this numerical trace conservative :

$$
\int_{F} \llbracket \underline{\underline{\sigma}}^{K} \cdot \mathbf{n} \rrbracket \cdot \eta=0
$$

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$$

Replacing ( $\hat{\underline{\sigma}}^{K} \cdot \mathbf{n}$ ) and summing over all faces, the transmission condition becomes :

$$
\sum_{K \in \mathcal{T}_{h}} \int_{\partial K}\left(\underline{\underline{\sigma}}^{K} \cdot \mathbf{n}\right) \cdot \eta-\sum_{K \in \mathcal{T}_{h}} \int_{\partial K} \tau \mathbf{I}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \eta=0
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\sum_{K \in \mathcal{T}_{h}} \int_{\partial K}\left(\underline{\underline{\sigma}}^{K} \cdot \mathbf{n}\right) \cdot \eta-\sum_{K \in \mathcal{T}_{h}} \int_{\partial K} \tau \mathbf{I}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \eta=0
$$

Discretization of the transmission condition

$$
\sum_{K \in \mathcal{T}_{h}}\left[\mathbb{B}^{K} \underline{W}^{K}+\mathbb{L}^{K} \underline{\Lambda}\right]=0
$$

## HDG formulation of the equations

Global HDG formulation

$$
\left\{\begin{array}{l}
\int_{K} i \omega \rho^{K} \mathbf{v}^{K} \cdot \mathbf{w}-\int_{K}\left(\nabla \cdot \underline{\underline{\sigma}}^{K}\right) \cdot \mathbf{w}+\int_{\partial K} \tau \mathbf{l}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \mathbf{w}=0 \\
\int_{K} i \omega \underline{\underline{\sigma}}^{K}: \underline{\underline{\xi}}+\int_{K} \mathbf{v}^{K} \cdot \nabla \cdot\left(\underline{\underline{C}}^{K} \underline{\underline{\xi}}\right)-\int_{\partial K} \lambda^{F} \cdot \underline{\underline{C}} K \underline{\underline{\xi}} \cdot \mathbf{n}=0 \\
\sum_{K \in \mathcal{T}_{h}} \int_{\partial K}\left(\underline{\underline{\sigma}}^{K} \cdot \mathbf{n}\right) \cdot \eta-\sum_{K \in \mathcal{T}_{h}} \int_{\partial K} \tau \mathbf{l}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \eta=0
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\int_{K} i \omega_{\underline{\sigma^{K}}}: \underline{\underline{\xi}}+\int_{K} v^{K} \cdot \nabla \cdot\left(\underline{\underline{C}}^{K} \underline{\underline{\xi}}\right)-\int_{\partial K} \lambda^{F} \cdot \underline{\underline{C_{K}}} \underline{\underline{\xi}} \cdot \mathbf{n}=0 \\
\sum_{K \in \mathcal{T}_{h}} \int_{\partial K}\left(\underline{\underline{\sigma}}^{K} \cdot \mathbf{n}\right) \cdot \eta-\sum_{K \in \mathcal{T}_{h}} \int_{\partial K} \tau \mathbf{l}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \eta=0
\end{array}\right.
$$

Global HDG discretization

$$
\left\{\begin{array}{l}
\mathbb{A}^{K} \underline{W}^{K}+\mathbb{C}^{K} \underline{\Lambda}=0 \\
\sum_{K \in \mathcal{T}_{h}}\left[\mathbb{B}^{K} \underline{W}^{K}+\mathbb{L}^{K} \underline{\Lambda}\right]=0
\end{array}\right.
$$

## HDG formulation of the equations

Global HDG formulation

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\left\{\begin{array}{l}
\int_{K} i \omega \rho^{K} \mathbf{v}^{K} \cdot \mathbf{w}-\int_{K}\left(\nabla \cdot \underline{\underline{\sigma}}^{K}\right) \cdot \mathbf{w}+\int_{\partial K} \tau \mathbf{l}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \mathbf{w}=0 \\
\int_{K} i \omega_{\underline{\sigma^{K}}}: \underline{\underline{\xi}}+\int_{K} v^{K} \cdot \nabla \cdot\left(\underline{\underline{C}}^{K} \underline{\underline{\xi}}\right)-\int_{\partial K} \lambda^{F} \cdot \underline{\underline{C_{K}}} \underline{\underline{\xi}} \cdot \mathbf{n}=0 \\
\sum_{K \in \mathcal{T}_{h}} \int_{\partial K}\left(\underline{\underline{\sigma}}^{K} \cdot \mathbf{n}\right) \cdot \eta-\sum_{K \in \mathcal{T}_{h}} \int_{\partial K} \tau \mathbf{l}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \eta=0
\end{array}\right.
$$

Global HDG discretization

$$
\left\{\begin{array}{l}
\underline{W}^{K}=-\left(\mathbb{A}^{K}\right)^{-1} \mathbb{C}^{K} \Lambda \\
\sum_{K \in \mathcal{T}_{b}}\left[\mathbb{B}^{\kappa} \underline{W}^{K}+\mathbb{L}^{K} \Lambda\right]=0
\end{array}\right.
$$

## HDG formulation of the equations

Global HDG formulation

$$
\left\{\begin{array}{l}
\int_{K} i \omega \rho^{K} \mathbf{v}^{K} \cdot \mathbf{w}-\int_{K}\left(\nabla \cdot \underline{\underline{\sigma}}^{K}\right) \cdot \mathbf{w}+\int_{\partial K} \tau \mathbf{l}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \mathbf{w}=0 \\
\int_{K} i \omega \underline{\underline{\sigma}}^{K}: \underline{\underline{\xi}}+\int_{K} \mathbf{v}^{K} \cdot \nabla \cdot\left(\underline{\underline{C}}^{K} \underline{\underline{\xi}}\right)-\int_{\partial K} \lambda^{F} \cdot \underline{\underline{C_{K}}} K \underline{\underline{\xi}} \cdot \mathbf{n}=0 \\
\sum_{K \in \mathcal{T}_{h}} \int_{\partial K}\left(\underline{\underline{\sigma}}^{K} \cdot \mathbf{n}\right) \cdot \eta-\sum_{K \in \mathcal{T}_{h}} \int_{\partial K} \tau \mathbf{I}\left(\mathbf{v}^{K}-\lambda^{F}\right) \cdot \eta=0
\end{array}\right.
$$

Global HDG discretization

$$
\sum_{K \in \mathcal{T}_{h}}\left[-\mathbb{B}^{K}\left(\mathbb{A}^{K}\right)^{-1} \mathbb{C}^{K}+\mathbb{L}^{K}\right] \underline{\Lambda}=0
$$

## Main idea of the algorithm using the HDG formulation

1. Construction of the linear system $\mathbb{M} \Lambda$
with $\mathbb{M}=\sum_{K \in \mathcal{T}_{h}}\left[-\mathbb{B}^{K}\left(\mathbb{A}^{K}\right)^{-1} \mathbb{C}^{K}+\mathbb{L}^{K}\right]$
for $K=1$ to $N b_{\text {tri }}$ do
Compute matrices $\mathbb{B}^{K},\left(\mathbb{A}^{K}\right)^{-1}, \mathbb{C}^{K}$ and $\mathbb{L}^{K}$
Construction of the corresponding section of $\mathbb{M}$ end for
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## Main idea of the algorithm using the HDG formulation

1. Construction of the linear system $\mathbb{M} \wedge$
2. Construction of the right hand side $\mathbb{S}$

## Main idea of the algorithm using the HDG formulation

1. Construction of the linear system $\mathbb{M} \wedge$
2. Construction of the right hand side $\mathbb{S}$
3. Resolution $\mathbb{M} \underline{\Lambda}=\mathbb{S}$

## Main idea of the algorithm using the HDG formulation

1. Construction of the linear system $\mathbb{M} \underline{\wedge}$
2. Construction of the right hand side $\mathbb{S}$
3. Resolution $\mathbb{M} \underline{\wedge}=\mathbb{S}$
4. Computation of the solutions of the initial problem
for $K=1$ to $N b_{t r i}$ do
Compute $\underline{W}^{K}=-\left(\mathbb{A}^{K}\right)^{-1} \mathbb{C}^{K} \underline{\Lambda}$
end for

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## Plane wave



Computational domain $\Omega$ setting

- Physical parameters:
- $\rho=2000 \mathrm{~kg} \cdot \mathrm{~m}^{-3}$
- $\lambda=16 \mathrm{GPa}$
- $\mu=8 \mathrm{GPa}$
- Plane wave :

$$
u=\nabla e^{i(k \cos \theta x+k \sin \theta y)}
$$

where $k=\frac{\omega}{v_{p}}$

- $\theta=0$
- Three meshes :
- 3000 elements
- 10000 elements
- 45000 elements


## Plane wave



Convergence order of the HDG scheme
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## Plane wave

Memory used

| \# Elements | Degree | Memory (MB) |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  |  | HDG | UDG | IPDG |
| 3100 | 1 | 44 | 288 | 58 |
| 10300 | 1 | 161 | 1076 | 221 |
| 45000 | 1 | 797 | 5492 | 1156 |
| 3100 | 2 | 97 | 804 | 215 |
| 10300 | 2 | 355 | 3097 | 852 |
| 45000 | 2 | 1746 | 15965 | 4454 |
| 3100 | 3 | 170 | 1656 | 598 |
| 10300 | 3 | 624 | 6600 | 2394 |
| 45000 | 3 | 3080 | 34597 | 12362 |
| 3100 | 4 | 254 | 2749 | 1324 |
| 10300 | 4 | 947 | 10098 | 5251 |
| 45000 | 4 | 4653 | 50297 | 27314 |

## Plane wave

Memory used

| \# Elements | Degree | Memory ratio |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  |  | HDG | UDG | IPDG |
| 3100 | 1 | 1 | 6.55 | 1.32 |
| 10300 | 1 | 1 | 6.68 | 1.37 |
| 45000 | 1 | 1 | 6.89 | 1.45 |
| 3100 | 2 | 1 | 8.29 | 2.22 |
| 10300 | 2 | 1 | 8.72 | 2.4 |
| 45000 | 2 | 1 | 9.14 | 2.55 |
| 3100 | 3 | 1 | 9.74 | 3.52 |
| 10300 | 3 | 1 | 10.58 | 3.84 |
| 45000 | 3 | 1 | 11.23 | 4.01 |
| 3100 | 4 | 1 | 10.82 | 5.21 |
| 10300 | 4 | 1 | 10.66 | 5.54 |
| 45000 | 4 | 1 | 10.81 | 5.87 |

## Marmousi test-case



Computational domain $\Omega$ composed of 235000 triangles

## Parallel results for the Marmousi test-case with the

 HDG-P3 scheme, $f=2 \mathrm{~Hz}$

## Caracteristics of the computing processors used

- Plafrim platform
- Hardware specification : 16 nodes, 12 cores by nodes
- Caracteristics of computing nodes:
- 2 Hexa-core Westmere Intel ${ }^{\circledR}$ Xeon ${ }^{\circledR}$ X5670
- Frequency : $2,93 \mathrm{GHz}$
- Cache L3 : 12 Mo
- RAM : 96 Go
- Infiniband DDR: 20Gb/s
- Ethernet : 1Gb/s


## Efficiency of the parallelism of the global matrix construction

Efficiency $=\frac{\text { sequential computational time }}{\mathrm{nb} \_ \text {cores } \times \text { computational time with one core }}$

Efficiency of the parallelism of the global matrix construction


## Efficiency of the parallelism of the resolution part



## Speed up for the global matrix construction


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## Speed up (Total simulation time)



Memory required (GB) for the simulation


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Plane wave in an homogeneous medium

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## Plane wave

- Physical parameters:
- $\rho=2000 \mathrm{~kg} . \mathrm{m}^{-3}$
- $\lambda=16 \mathrm{GPa}$
- $\mu=8 \mathrm{GPa}$
- Plane wave :

$$
u=\nabla e^{i(k \cos \theta x+k \sin \theta y)}
$$

where $k=\frac{\omega}{v_{p}}$
Computational domain $\Omega$ setting

- $\theta=0$
- Two meshes :
- 2600 elements
- 6400 elements


## 3D preliminary numerical results



Figure: Exact solution

## 3D preliminary numerical results



Figure: Numerical solution computed with $P_{3}$-HDG scheme on mesh composed of 6400 tetrahedrons

| Elements | Order | CPU Time (s) |  | Memory (MB) |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | HDG | IPDG | HDG | IPDG |
| 2600 | 1 | 11 |  | 572 |  |
| 6400 | 1 | 59 |  | 1981 |  |
| 2600 | 2 | 91 |  | 2254 |  |
| 6400 | 2 | 498 |  | 7894 |  |
| 2600 | 3 | 508 |  | 6242 |  |
| 6400 | 3 | 2386 |  | 21899 |  |


| Elements | Order | CPU Time (s) |  | Memory (MB) |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | HDG | IPDG | HDG | IPDG |
| 2600 | 1 | 11 | 12 | 572 | 461 |
| 6400 | 1 | 59 | 54 | 1981 | 1630 |
| 2600 | 2 | 91 | 123 | 2254 | 2687 |
| 6400 | 2 | 498 | 910 | 7894 | 10666 |
| 2600 | 3 | 508 | 1409 | 6242 | 13359 |
| 6400 | 3 | 2386 | 10370 | 21899 | 47922 |


| Elements | Order | CPU Time (s) |  | Memory (MB) |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | HDG | IPDG | HDG | IPDG |
| 2600 | 1 | 1 |  | 1 |  |
| 6400 | 1 | 1 |  | 1 |  |
| 2600 | 2 | 1 |  | 1 |  |
| 6400 | 2 | 1 |  | 1 |  |
| 2600 | 3 | 1 |  | 1 |  |
| 6400 | 3 | 1 |  | 1 |  |


| Elements | Order | CPU Time (s) |  | Memory (MB) |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | HDG | IPDG | HDG | IPDG |
| 2600 | 1 | 1 | 1.09 | 1 | 0.8 |
| 6400 | 1 | 1 | 0.9 | 1 | 0.8 |
| 2600 | 2 | 1 | 1.35 | 1 | 1.19 |
| 6400 | 2 | 1 | 1.83 | 1 | 1.35 |
| 2600 | 3 | 1 | 2.77 | 1 | 2.14 |
| 6400 | 3 | 1 | 4.35 | 1 | 2.19 |

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## Conclusions-Perspectives

Conclusions
On a same mesh, with the 2D HDG method:

- Memory gain
- Computational time gain


## Conclusions-Perspectives

Conclusions
On a same mesh, with the 2D HDG method:

- Memory gain
- Computational time gain


## Perspectives

- Study of the 3D HDG algorithm for Helmholtz equations
- Solution strategy for the HDG linear system


## Thank you!



