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Lorentzian Coxeter Groups and Boyd-Maxwell Ball Packings

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Abstract. In the recent study of infinite root systems, fractal patterns of ball packings were observed while visualizing roots in affine space. In fact, the observed fractals are exactly the ball packings described by Boyd and Maxwell. This correspondence is a corollary of a more fundamental result: given a geometric representation of a Coxeter group in Lorentz space, the set of limit directions of weights equals the set of limit roots.

Résumé. Lors de la visualisation des systèmes de racines infinis dans l'espace affine, des formes fractales ressemblant à des empilements de sphères furent observées. En fait, ces fractales sont précisement les empilements de sphères décrits par Boyd et Maxwell. Cette correspondance est une conséquence d'un résultat plus fondamental: étant donné un groupe de Coxeter agissant sur un espace de Lorentz, l'ensemble des directions limites des poids égale l'ensemble des racines limites du système de racines correspondant.

Keywords: Sphere packing, ball packing, infinite Coxeter groups, infinite root systems, limit roots, Coxeter graphs

1 Introduction

Apollonian ball packings form a renowed class of infinite ball packings, see for instance [GLM⁺05, GLM⁺06]. *Boyd–Maxwell ball packings* are generated by inversions, generalizing Apollonian ball packings. We name these packings after Boyd, who first studied them in [Boy74], and Maxwell, who related them to Lorentzian Coxeter groups in [Max82]. Here, we revisit Boyd–Maxwell ball packings. This retrospective is motivated by recent studies on limit roots of infinite Coxeter groups.

Limit roots were introduced and studied in [HLR14]. They are the limit directions of the roots in a geometric representation of an infinite Coxeter system. Properties of limit roots of infinite Coxeter groups were investigated in a recent series of papers, see [HLR14, DHR13, HPR13]. In many examples, the Coxeter group acts as a discrete reflection group on Lorentz space, and patterns of fractal ball packings appear, see Figure 1. A description of this fractal structure is conjectured in [HLR14, Section 3.2] and proved in [DHR13, Theorem 4.10].

While investigating limit roots, we observed that patterns appearing in these examples are similar to Boyd–Maxwell ball packings. The main result here relates the set of limit roots to Boyd–Maxwell ball

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(a) Positive roots of depth \leq 7 for the Coxeter group of rank 4 with a complete Coxeter graph with all edges labeled by 4.

(b) Positive roots of depth ≤ 7 for the Coxeter group of rank 4 with a complete Coxeter graph with all edges labeled by 4 except one dotted edge labeled by -1.1 (see Section 2.3 for Vinberg's convention).

Fig. 1: The pattern of a ball packing and a ball cluster approximated by roots generated by two rank-4 Coxeter groups, seen in the affine space spanned by simple roots.

packings. More specifically, Theorem 3.3 states that the set of limit directions of weights of a Lorentzian Coxeter group W is equal to its set of limit roots. This result allows to conclude that for a Lorentzian Coxeter system (W, S) with Maxwell's condition, namely that (W, S) is of level 2, the set of limit roots is the complement of the interiors of all balls in the Boyd–Maxwell packing, see Corollary 3.4. Moreover, even if Maxwell's condition is not satisfied, the statement is generalized to non-packing clusters of balls, see Corollary 3.5. In [CL13], we also study the tangency graphs of Boyd–Maxwell ball packings and enumerate level-2 Coxeter systems.

In Section 2, we recall the notions of geometric representations of Coxeter system, limit root, and review the work of Boyd and Maxwell. In Section 3, we relate limit weights, limit roots and Boyd–Maxwell ball packings and examine a generalization to ball clusters.

2 Coxeter groups, limit roots and Boyd–Maxwell Packings

2.1 Geometric representation of a Coxeter group

Let (W, S) be a finitely generated *Coxeter system*, where $S = \{s_1, s_2, \ldots, s_n\}$ is a finite set of generators and the *Coxeter group* W is generated by the relations $(s_i s_j)^{m_{ij}} = e$, with $m_{ii} = 1$ and $m_{ij} \ge 2$ or $= \infty$ if $i \ne j$. The cardinality |S| = n is the *rank* of the Coxeter system (W, S). Let V be a real vector space of dimension n, equipped with a basis $\Delta = \{\alpha_1, \alpha_2, \ldots, \alpha_n\}$. Given a Coxeter system (W, S), we define

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a bilinear form \mathcal{B} as follows:

$$\mathcal{B}(\alpha_i, \alpha_j) = \begin{cases} -\cos(\pi/m_{ij}) & \text{if } m_{ij} < \infty; \\ -c_{ij} & \text{if } m_{ij} = \infty, \end{cases}$$

where c_{ij} are chosen arbitrarily with $c_{ij} = c_{ji} \ge 1$. For a non-isotropic vector $\alpha \in V$, we define the reflection σ_{α}

$$\sigma_{\alpha}(\mathbf{x}) := \mathbf{x} - 2\frac{\mathcal{B}(\mathbf{x},\alpha)}{\mathcal{B}(\alpha,\alpha)}\alpha \quad \text{for all } \mathbf{x} \in V.$$
⁽¹⁾

Then the homomorphism $\rho : W \to \operatorname{GL}(V)$ that sends s_i to σ_{α_i} is a faithful geometric representation of the Coxeter group W as a discrete subgroup of the orthogonal group $O_{\mathcal{B}}(V)$, i.e., the group of linear transformations of V preserving the bilinear form \mathcal{B} . The group W acts on (V, \mathcal{B}) by $w(\mathbf{x}) = \rho(w)(\mathbf{x})$. We refer the readers to [Hum92, Chapter 5] and [HLR14, Section 1] for more details.

If the bilinear form \mathcal{B} is of signature (n - 1, 1), the pair (V, \mathcal{B}) is called a *n*-dimensional *Lorentz space* and the Coxeter system (W, S) with W acting on (V, \mathcal{B}) is said to be *Lorentzian*. By abuse of language, we also say that W is a *Lorentizian Coxeter group*.

Remark 2.1 Being Lorentzian is an extrinsic property of the Coxeter system (W, S) since it depends on the choice of the bilinear form \mathcal{B} associated to the vector space V on which W acts. In the present investigation, we are mainly concerned with Coxeter groups acting on Lorentz space, therefore we call them Lorentzian. In the literature, the term hyperbolic is used, but with different meanings, see for instance [Bou68, Vin71, Max78, Max82, Hum92, DHR13].

In a Lorentz space, a vector \mathbf{x} is *space-like* (resp. *time-like*, *light-like*) if $\mathcal{B}(\mathbf{x}, \mathbf{x})$ is positive (resp. negative, zero). The set of light-like vectors $Q = {\mathbf{x} \in V \mid \mathcal{B}(\mathbf{x}, \mathbf{x}) = 0}$ form a cone called the *light cone*.

Let $\Phi = W(\Delta)$ be the orbit of Δ under the action of W. Then the pair (Φ, Δ) is a *based root system*. The vectors in Δ are called *simple roots*, and the vectors in Φ are called *roots*. Let V^* be the dual vector space of V with dual basis Δ^* . If the bilinear form \mathcal{B} is non-singular, which is the case for Lorentz spaces, V^* can be identified with V, and $\Delta^* = \{\omega_1, \omega_2, \dots, \omega_n\}$ can be identified with a set of vectors in V such that

$$\mathcal{B}(\alpha_i, \omega_j) = \delta_{ij},\tag{2}$$

where δ_{ij} is the Kronecker delta function. Vectors in Δ^* are called *fundamental weights*, and vectors in the orbit

$$\Omega := W(\Delta^*) = \bigcup_{\omega \in \Delta^*} W(\omega)$$

are called *weights*. We refer the readers to [Bou68, Chapter VI, Section 1.10] for more details.

2.2 Limit roots

To study the asymptotic directions of the roots, we pass to the *projective space* $\mathbb{P}V$, i.e., the topological space of 1-dimensional subspaces of V. For a non-zero vector $\mathbf{x} \in V \setminus \{0\}$, let $\hat{\mathbf{x}} \in \mathbb{P}V$ denote the line passing through \mathbf{x} and the origin. The group action of W on V by reflection induces a *projective action* of W on $\mathbb{P}V$:

$$w \cdot \widehat{\mathbf{x}} = \widetilde{w(\mathbf{x})}, \quad w \in W, \quad \mathbf{x} \in V.$$

For a set $X \subset V$, we define the corresponding projective set

$$\widehat{X} := \{ \widehat{\mathbf{x}} \in \mathbb{P}V \mid \mathbf{x} \in X \}.$$

In this sense, we have projective roots $\widehat{\Phi}$, projective weights $\widehat{\Omega}$ and the projective light cone \widehat{Q} .

Let $h(\mathbf{x})$ denote the sum of the coordinates of \mathbf{x} in the basis Δ , and call it the *height* of the vector \mathbf{x} . We say that \mathbf{x} is *future-directed* (resp. *past-directed*) if $h(\mathbf{x})$ is positive (resp. negative). The hyperplane $\{\mathbf{x} \in V \mid h(\mathbf{x}) = 1\}$ is the affine subspace $\operatorname{aff}(\Delta)$ spanned by the simple roots. It is useful to identify the projective space $\mathbb{P}V$ with the affine subspace $\operatorname{aff}(\Delta)$ with a *projective hyperplane added at infinity*. For a vector $\mathbf{x} \in V$, if $h(\mathbf{x}) \neq 0$, $\hat{\mathbf{x}}$ is identified with the *projective vector*

$$\mathbf{x}/h(\mathbf{x}) \in \operatorname{aff}(\Delta).$$

Otherwise, if $h(\mathbf{x}) = 0$, the direction $\hat{\mathbf{x}}$ is identified to a point on the projective hyperplane at infinity. In fact, if $h(\mathbf{x}) \neq 0$, $\hat{\mathbf{x}}$ is identified with the intersection of $\operatorname{aff}(\Delta)$ with the straight line passing through \mathbf{x} and the origin. In this sense, the projective roots $\hat{\Phi}$, projective weights $\hat{\Omega}$ and projective light cone \hat{Q} are respectively identified with the intersection of $\operatorname{aff}(\Delta)$ with the 1-subspaces spanned by the roots Φ , weights Ω and light-like vectors. The projective light cone \hat{Q} is projectively equivalent to a sphere. In Figure 2, simple roots, fundamental weights and some positive roots are represented in $\operatorname{aff}(\Delta)$.



Fig. 2: Simple roots, fundamental weights, and positive roots of depth ≤ 6 of a geometric Coxeter system of rank 3 seen in the affine space spanned by the simple roots. The Coxeter graph is shown in the upper-left corner.

Definition 2.2 (Hohlweg–Labbé–Ripoll [HLR14, Definition 2.12]) The set $E(\Phi)$ of limit roots is the set of accumulation points of $\widehat{\Phi}$, in other words,

$$E(\Phi) = \{ \widehat{\mathbf{x}} \in \mathbb{P}V \mid \text{ there is an injective sequence } (\gamma_i)_{i \in \mathbb{N}} \in \Phi \text{ such that } \lim_{i \to \infty} \widehat{\gamma}_i = \widehat{\mathbf{x}} \}$$

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Theorem 2.7 of [HLR14] asserts that

$$E(\Phi) \subseteq \widehat{Q} \cap \operatorname{cone}(\Delta).$$

Consequently, there is no limit root in the set $\widehat{Q} \setminus \operatorname{cone}(\Delta)$. If the set $\widehat{Q} \setminus \operatorname{cone}(\Delta)$ consists of open spherical caps, it was conjectured that $E(\Phi)$ is equal to the complement of the *W*-orbit of these caps, see [HLR14, Conjecture 3.9]. This conjecture is proved in [DHR13, Theorem 4.10], and presented here in Theorem 2.3. The aim here is to relate this result to the one of Maxwell, see Theorem 2.4.

Theorem 2.3 ([DHR13, Theorem 4.10], [HMN13, Theorem 1.2]) Let $(W, S)_B$ be an irreducible Lorentzian Coxeter system. Then

 $E(\Phi) = \widehat{Q} \setminus (W \cdot (\widehat{Q} \setminus \operatorname{conv}(\Delta))).$

In particular, if $\widehat{Q} \subset \operatorname{conv}(\Delta)$, then $E(\Phi) = \widehat{Q}$.

2.3 Boyd–Maxwell packing

In a metric space, a *ball packing* is a set of balls with disjoint interiors. A famous example is the Apollonian ball packing, see for instance [Boy73, BdPP94, GLM⁺05, GLM⁺06, Che13]. The *residual set* of a ball packing is the complement of the interiors of all balls in the packing. An Apollonian packing can be generated by inversions. In [Boy74], Boyd presents a new class of infinite ball packings generalizing this construction process. Moreover, he notices a connection to reflection groups. In [Max82], Maxwell revisits these packings, and interprets them using Lorentzian Coxeter groups.

Given a space-like vector x in the Lorentz space (V, \mathcal{B}) , the *normalized vector* \overline{x} of x is given by

$$\overline{\mathbf{x}} = \mathbf{x} / \sqrt{\mathcal{B}(\mathbf{x}, \mathbf{x})}$$

The normalized vector $\overline{\mathbf{x}}$ lies on the one-sheet hyperboloid $\mathcal{H} = {\mathbf{x} \in V | \mathcal{B}(\mathbf{x}, \mathbf{x}) = 1}$. For n > 2, there is a classical correspondence between (n-2)-dimensional balls and space-like directions in *n*-dimensional Lorentz space, see for example [Max82, Section 2], [Cec08, Section 2.2] or [HJ03, Section 1.1]. Given a space-like vector \mathbf{x} , let $H_{\mathbf{x}}$ be the *orthogonal hyperplane* $H_{\mathbf{x}} = {\mathbf{x}' \in V | \mathcal{B}(\mathbf{x}, \mathbf{x}') = 0}$. The intersection of \hat{Q} with the projective half-space $\hat{H}_{\mathbf{x}}^- = {\hat{\mathbf{x}}' \in \mathbb{P}V | \mathcal{B}(\mathbf{x}, \mathbf{x}') \leq 0}$ is a closed ball (spherical cap) on \hat{Q} . We denote this ball by Ball(\mathbf{x}). After a stereographic projection, Ball(\mathbf{x}) becomes a ball in an (n-2)-dimensional Euclidean space. For two space-like vectors \mathbf{x} and \mathbf{x}' , if they are *not both* future-directed, we have

- Ball(\mathbf{x}) and Ball(\mathbf{x}') are disjoint if $\mathcal{B}(\overline{\mathbf{x}}, \overline{\mathbf{x}}') < -1$;
- Ball(\mathbf{x}) is tangent to Ball(\mathbf{x}') if $\mathcal{B}(\overline{\mathbf{x}}, \overline{\mathbf{x}}') = -1$;
- The boundary of $Ball(\mathbf{x})$ and $Ball(\mathbf{x}')$ intersect transversally if $\mathcal{B}(\overline{\mathbf{x}}, \overline{\mathbf{x}}') > -1$;

Therefore, if a set of space-like vectors represents a ball packing, we must have $\mathcal{B}(\overline{\mathbf{x}}, \overline{\mathbf{x}}') \leq -1$ for any two vectors. The packing corresponding to a pair of opposite vectors $\{\mathbf{x}, -\mathbf{x}\}$ is said to be trivial.

To encode both the Coxeter system (W, S) and the bilinear form \mathcal{B} , we adopt Vinberg's convention for Coxeter graphs. Namely, if $c_{ij} > 1$ the edge ij is dotted and labeled by $-c_{ij}$. A graph G is said to be of *level 0* if it represents a finite or affine Coxeter system (W, S). The list of level-0 Coxeter graphs can be

found in [Hum92, Chapter 2]. A graph G is of level $\leq r$ if every induced subgraph of G on n-r vertices is of level 0. A graph G is of level r if it is of level $\leq r$ but not of level $\leq r-1$. Correspondingly, a Coxeter system (W, S) with a Coxeter graph of level r is said to be of level r.

For a Lorentzian Coxeter system (W, S), let Ω_r be the set of *space-like weights*. Maxwell proved that Coxeter groups of level 2 are Lorentzian [Max82, Proposition 1.6] and the following theorem.

Theorem 2.4 (Maxwell [Max82, Theorem 3.2]) Let (W, S) be a Lorentzian Coxeter system and Ω_r its set of space-like weights. The set $\{Ball(\omega) \mid \omega \in \Omega_r\}$ is a ball packing if and only if W is of level 2.

We call the ball packing in Theorem 2.4 the *Boyd–Maxwell ball packing* generated by the Coxeter system (W, S), and denote it by $\mathcal{P}(W, S)$. Maxwell manually enumerated the Coxeter graphs representing irreducible Coxeter groups of level 2, and suggested a computer verification. This is done in [CL13, Section 5].

3 Relation between limit roots, limit weights and Boyd–Maxwell Packings

From now on, we assume that (W, S) is a Coxeter system with W acting on a non-singular space (V, \mathcal{B}) . We define the set of limit weights $E(\Omega)$ analogously to limit roots.

Definition 3.1 The set of limit weights $E(\Omega)$ is the set of accumulation points of the projective weights $\overline{\Omega}$. That is

 $E(\Omega) = \{ \widehat{\mathbf{x}} \in \mathbb{P}V \mid \text{there is an injective sequence } (\omega_i)_{i \in \mathbb{N}} \in \Omega \text{ such that } \lim_{i \to \infty} \widehat{\omega}_i = \widehat{\mathbf{x}} \}.$

As we know, limit roots lie in the isotropic cone Q. Here is an analogous result for limit weights.

Theorem 3.2 Consider an injective sequence of weights $(\omega_k)_{k \in \mathbb{N}}$ and suppose that $(\widehat{\omega}_k)_{k \in \mathbb{N}}$ converges to a limit $\widehat{\psi}$. Then

- (i) $h(\omega_k)$ tends to $-\infty$,
- (ii) $\widehat{\psi}$ lies in \widehat{Q} .

When (V, \mathcal{B}) is a Lorentz space, even more is true.

Theorem 3.3 The set of limit weights of a Lorentzian Coxeter system (W, S) is equal to its set of limit roots. That is, $E(\Omega) = E(\Phi)$.

Figure 3 illustrates the result of the previous theorem.

We denote the set of space-like weights by Ω_r . It is the union of the orbits of space-like fundamental weights. Space-like weights in Ω_r correspond to balls in the Boyd–Maxwell ball packing $\mathcal{P}(W, S)$. From this correspondence, we get the following corollary of Theorem 3.2.

Corollary 3.4 The set $E(\Phi)$ of limit roots of an irreducible Lorentzian Coxeter system (W, S) of level 2 is equal to the residual set of the Boyd-Maxwell ball packing $\mathcal{P}(W, S)$.

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Fig. 3: Positive roots and space-like weights for a Lorentzian Coxeter system of rank 4 seen in the affine space spanned by the simple roots. The Coxeter graph is shown on the upper-left corner.

Let us now explain the relation between ball packings studied by Boyd and Maxwell and ball packings observed in the study of limit roots. In the point of view of [HLR14] and [DHR13], (W, S) is of level 2 if and only if $\hat{Q} \setminus \text{conv}(\Delta)$ is not empty and consists of a union of disjoint spherical caps. We notice from Equation (2) that

$$\operatorname{aff}(\Delta \setminus \{\alpha\}) = \operatorname{aff}(\Delta) \cap H_{\omega}, \quad \forall \alpha \in \Delta,$$

where ω is the fundamental weights corresponding to the simple root α . In other words, the supporting hyperplane $\operatorname{aff}(\Delta \setminus \{\alpha\})$ of the simplex $\operatorname{conv}(\Delta)$ is exactly the intersection of $\operatorname{aff}(\Delta)$ and the orthogonal hyperplane for the fundamental weight ω . Therefore, the spherical caps obtained by the space-like fundamental weights are exactly the spherical caps in $\widehat{Q} \setminus \operatorname{conv}(\Delta)$. Consequently, if the Coxeter system is Lorentzian of level 2, the fractal structure described in Theorem 2.3 is the residual set of the Boyd–Maxwell ball packing described in Theorem 2.4.

For a Lorentzian Coxeter system of level ≤ 1 , every facet of $\operatorname{conv}(\Delta)$ is disjoint from, or tangent to Q. In this case, as observed in [HPR13, DHR13], $E(\Phi) = \hat{Q}$, so no pattern of ball packing or ball cluster appears. In the framework of Maxwell, it can be explained by the absence of space-like weights. For a Lorentzian Coxeter system of level higher than 2, the space-like weights still represent (n - 2)dimensional balls, but some balls will intersect each other. We name this configuration *Boyd–Maxwell ball cluster* generated by the Coxeter group, see Figure 1(b) for an example. The notion of residual set naturally extends to ball clusters. To generalize Corollary 3.4 to Boyd–Maxwell ball clusters, most of the arguments and discussions in the previous part applies. However, slight modifications are necessary. Finally, we obtain the following corollary, which completes the connection between Theorem 2.3 and 2.4. **Corollary 3.5** The set $E(\Phi)$ of limit roots of an irreducible Lorentzian Coxeter system (W, S) is equal to the residual set of the Boyd-Maxwell ball cluster generated by (W, S).

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