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ASYMPTOTIC STUDY FOR STOKES-BRINKMAN MODEL WITH JUMP EMBEDDED TRANSMISSION CONDITIONS

PHILIPPE ANGOT¹, GILLES CARBOU², VICTOR PÉRON^{2,3}

ABSTRACT. In this paper, one considers the coupling of a Brinkman model and Stokes equations with jump embedded transmission conditions. In this model, one assumes that the viscosity in the porous region is very small. Then we derive a Wentzel–Kramers–Brillouin (WKB) expansion in power series of the square root of this small parameter for the velocity and the pressure which are solution of the transmission problem. This WKB expansion is justified rigorously by proving uniform errors estimates.

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1. INTRODUCTION

We address the problem of fluid flow modeling in complex media which combine porous regions and fluid regions with free flow. This issue holds for instance in the study of aquifer media made up of a porous media containing cracks and conduits (see [8]) and also the passive control of the flow around an obstacle covered by a porous thin layer (see [4]).

In this paper the free flow satisfies the linear Stokes equation. In the porous media, we consider two models, the Brinkman and the Darcy models. There are several interface conditions in the literature. In the case of the Stokes-Brinkman coupling, the simpler interface condition is the continuity of the velocity and the normal stress. More accurate models are given by Ochoa-Tapia & Whitaker transmission conditions, by Beavers & Joseph conditions or Beaver, Joseph

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& Saffman conditions (see [8]). In this paper, we deal with the more general jump embedded transmission conditions described in [2]. This condition links the jumps and the averages of both the velocity and the shear stress on the interface. A small parameter appears in this system, in particular the equivalent viscosity in the porous part is small, so that when this parameter tends to zero, we expect that the flow in the porous part will be described by the Darcy law. Our interest lies in the asymptotic study justifying the obtention of the limit model. In particular we describe the boundary layer due to the jump conditions appearing in the porous medium.

Let us describe the Stokes-Brinkman model with jump embedded transmission conditions. The problem is set in the domain $\Omega \subset \mathbb{R}^3$ made of a fluid region Ω_+ and a porous subdomain Ω_- . We assume that the domains Ω_- and Ω_+ are Lipschitz and bounded, and that $\Omega_- \subset \overline{\Omega_-} \subset \Omega$.

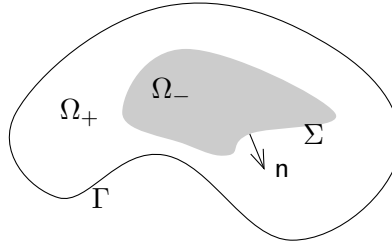


Figure 1 – The domain Ω and the subdomains Ω_- , Ω_+

We denote $\Sigma = \partial\Omega_-$ so that we have $\Omega = \Omega_- \cup \Sigma \cup \Omega_+$ (see Figure 1). We denote by \mathbf{n} the outward unit normal at $\partial\Omega_-$.

On the fluid region Ω_+ , the velocity v_ε^+ and the pressure p_ε^+ satisfy the Stokes equations. In the porous region Ω_- , the velocity v_ε^- and the pressure p_ε^- satisfy a Brinkman model. We couple these models by the Beaver-Joseph conditions at the common boundary Σ . These conditions link the jumps of the velocity and the normal stress vector with the averages of these quantities on Σ .

We denote by $\boldsymbol{\sigma}^+(v_\varepsilon^+, p_\varepsilon^+)$ (resp. $\boldsymbol{\sigma}^-(v_\varepsilon^-, p_\varepsilon^-)$) the stress tensor in the fluid (resp. in the porous) medium :

$$\boldsymbol{\sigma}^+(v_\varepsilon^+, p_\varepsilon^+) = 2\mu\mathbf{d}(v_\varepsilon^+) - p_\varepsilon^+\mathbb{I}, \quad \boldsymbol{\sigma}^-(v_\varepsilon^-, p_\varepsilon^-) = 2\varepsilon\mathbf{d}(v_\varepsilon^-) - p_\varepsilon^-\mathbb{I},$$

with

- $\mathbf{d}(v) = \frac{1}{2}(\nabla v + \nabla^\top v)$, that is $(\mathbf{d}(v))_{ij} = \frac{1}{2}\left(\frac{\partial v^j}{\partial x_i} + \frac{\partial v^i}{\partial x_j}\right)$,
- μ is the viscosity of the fluid and ε is the effective viscosity in the porous medium.

The problem writes

$$(1.1) \quad \begin{cases} -\nabla \cdot \boldsymbol{\sigma}^-(v_\varepsilon^-, p_\varepsilon^-) + \kappa v_\varepsilon^- = \mathbf{g}^- & \text{in } \Omega_- \\ -\nabla \cdot \boldsymbol{\sigma}^+(v_\varepsilon^+, p_\varepsilon^+) = \mathbf{g}^+ & \text{in } \Omega_+ \\ \nabla \cdot v_\varepsilon^- = 0 & \text{in } \Omega_- \\ \nabla \cdot v_\varepsilon^+ = 0 & \text{in } \Omega_+ \\ \boldsymbol{\sigma}^+(v_\varepsilon^+, p_\varepsilon^+) \cdot \mathbf{n} = \boldsymbol{\sigma}^-(v_\varepsilon^-, p_\varepsilon^-) \cdot \mathbf{n} + \mathbf{M}\{v_\varepsilon\} + \mathbf{l} & \text{on } \Sigma \\ \{\boldsymbol{\sigma}(v_\varepsilon, p_\varepsilon) \cdot \mathbf{n}\} = \mathbf{S}(v_\varepsilon^+ - v_\varepsilon^-) + \mathbf{h} & \text{on } \Sigma \\ v_\varepsilon = 0 & \text{on } \Gamma \end{cases}$$

where:

- $\kappa > 0$ is fixed positive constant.
- In the previous equations set on Σ , we denote by $\{w\}$ the mean value across Σ of w .
- The right-hand sides in (1.1) are given data which are defined as follow : $g^\pm : \Omega_\pm \rightarrow \mathbb{R}^3$ are vector fields in $(L^2(\Omega_\pm))^3$, $h \in (H^{-1/2}(\Sigma))^3$, and $l \in (H^{-1/2}(\Sigma))^3$ are vector fields defined on Σ .
- The matrix M is zero on n^\perp and satisfies:

$$M\xi = \beta(\xi \cdot n)n$$

with $\beta > 0$.

- The matrix S satisfies $S|_{n^\perp} = \alpha\mathbb{I}|_{n^\perp}$ and $S(n) = \frac{1}{\varepsilon}n$:

$$S(\xi) = \frac{1}{\varepsilon}(n \cdot \xi)n + \alpha(\xi - (n \cdot \xi)n),$$

with $\alpha > 0$.

We remark that using the divergence free conditions, we can replace the first equation in (1.1) by

$$(1.2) \quad -\varepsilon\Delta v_\varepsilon^- + \nabla p_\varepsilon^- + \kappa v_\varepsilon^- = g^- \quad \text{in } \Omega_-$$

Notation 1.1. For any set $\mathcal{O} \subset \mathbb{R}^3$ we denote $\mathbf{L}^2(\mathcal{O})$ the space $(L^2(\mathcal{O}))^3$ and $\mathbf{H}^s(\mathcal{O})$ the space $(H^s(\mathcal{O}))^3$.

2. STATEMENT OF THE MAIN RESULTS

In order to prove existence of weak solutions for (1.1), let us describe the associated variational formulation.

We denote by $\text{PH}^1(\Omega)$ the space of vector fields $v \in \mathbf{L}^2(\Omega)$ such that $v^+ := v|_{\Omega_+} \in \mathbf{H}^1(\Omega_+)$ and $v^- := v|_{\Omega_-} \in \mathbf{H}^1(\Omega_-)$. We introduce a weak formulation of the problem (1.1) for $v = (v^+, v^-)$ in the space

$$V = \{u \in \text{PH}^1(\Omega) \mid \text{div } u^\pm = 0 \quad \text{in } \Omega_\pm, \quad u^+ = 0 \quad \text{on } \Gamma\},$$

endowed with the piecewise H^1 norm. Such a variational formulation writes : Find $v = (v^-, v^+) \in V$ such that

$$(2.1) \quad \forall u \in V, \quad a_\varepsilon(v, u) = b(u),$$

where

$$\begin{aligned} a_\varepsilon(v, u) &:= 2\varepsilon \int_{\Omega_-} \mathbf{d}(v^-) : \mathbf{d}(u^-) \, d\mathbf{x} + \int_{\Omega_-} \kappa v^- \cdot u^- \, d\mathbf{x} + 2\mu \int_{\Omega_+} \mathbf{d}(v^+) : \mathbf{d}(u^+) \, d\mathbf{x} \\ &+ \langle \beta\{v \cdot n\}, \{u \cdot n\} \rangle_{-1/2, \Sigma} + \langle \alpha[v_\tau], [u_\tau] \rangle_{-1/2, \Sigma} + \varepsilon^{-1} \langle [v \cdot n], [u \cdot n] \rangle_{-1/2, \Sigma}, \end{aligned}$$

with $\langle \cdot, \cdot \rangle_{-1/2, \Sigma}$ being a duality pairing between $H^{-1/2}(\Sigma)^3$ and $H^{1/2}(\Sigma)^3$, $[v] = v^+ - v^-$ denoting the jump of v across Σ , and

$$b(u) = \int_{\Omega} g \cdot u \, d\mathbf{x} + \langle h, [u] \rangle_{-1/2, \Sigma} + \langle l, \{u\} \rangle_{-1/2, \Sigma},$$

compare with [1, Th 1.1], [2, Th 2.1]. The duality pairing $\langle \cdot, \cdot \rangle_{-1/2, \Sigma}$ between $H^{-1/2}(\Sigma)^3$ and $H^{1/2}(\Sigma)^3$ coincides with the duality pairing in $L^2(\Sigma)^3$.

For $\varepsilon > 0$, applying Lax Milgram Theorem, we prove the existence and uniqueness of weak solution for (1.1) and we obtain uniform estimates of the solutions with respect to the parameter ε .

Theorem 2.1. *Let $\mathbf{g}^\pm \in \mathbf{L}^2(\Omega_\pm)^3$ and $\mathbf{h}, \mathbf{l} \in L^2(\Sigma)^3$. Then, the problem (2.1) has a unique solution $v = v_\varepsilon \in V$ for all $\varepsilon > 0$. Moreover, for $\varepsilon > 0$ small enough, the following uniform estimate holds :*

$$(2.2) \quad \begin{aligned} \varepsilon \|\mathbf{d}(v^-)\|_{0, \Omega_-}^2 + \frac{\kappa}{4} \|v^-\|_{0, \Omega_-}^2 + \frac{\mu C^2}{3} \|v^+\|_{1, \Omega_+}^2 + \frac{1}{4\varepsilon} \|[v \cdot \mathbf{n}]\|_{0, \Sigma}^2 + \frac{\alpha}{4} \|v\|_{0, \Sigma}^2 + \frac{\beta}{4} \|\{v \cdot \mathbf{n}\}\|_{0, \Sigma}^2 \\ \leq c (\|\mathbf{g}\|_{0, \Omega}^2 + \|\mathbf{h}\|_{0, \Sigma}^2 + \|\mathbf{l}\|_{0, \Sigma}^2) , \end{aligned}$$

with a constant $c = c(\mu, \kappa, \alpha, \beta)$.

Remark 2.2. This result still holds when the data \mathbf{h} and \mathbf{l} belong to the space $H^{-\frac{1}{2}}(\Sigma)^3$. For the sake of simplicity, we prove this lemma when \mathbf{h} and \mathbf{l} belong to the space $L^2(\Sigma)^3$. We eventually compare this stability result with [1, 2] where the author prove also energy estimates. In [1, Th 1.1] and [2, Th 2.1], estimates are non-necessary uniform, whereas estimates (2.2) are uniform with respect to the parameter ε .

The asymptotic limit of the Stokes-Brinkman model towards the Stokes-Darcy one with Beavers-Joseph interface conditions is studied in [2, 8] in the case of a flat interface when the viscosity $\tilde{\mu}$ in the porous region is very small $\tilde{\mu} = \varepsilon \ll 1$ and when the jump of the normal velocities is penalized. We address in this paper the problem of the convergence of the model (1.1) when the parameter ε tends to zero.

When ε tends to zero, we formally converge to the following Stokes Darcy problem with Beavers & Joseph interface conditions

$$(2.3) \quad \left\{ \begin{array}{ll} \nabla p_0^- + \kappa v_0^- = \mathbf{g}^- & \text{in } \Omega_- \\ -\nabla \cdot \boldsymbol{\sigma}(v_0^+, p_0^+) = \mathbf{g}^+ & \text{in } \Omega_+ \\ \nabla \cdot v_0^- = 0 & \text{in } \Omega_- \\ \nabla \cdot v_0^+ = 0 & \text{in } \Omega_+ \\ \boldsymbol{\sigma}(v_0^+, p_0^+) \cdot \mathbf{n} = -p_0^- \mathbf{n} + \frac{\beta}{2} ((v_0^+ + v_0^-) \cdot \mathbf{n}) \mathbf{n} + \mathbf{l} & \text{on } \Sigma \\ (v_0^+ - v_0^-) \cdot \mathbf{n} = 0 & \text{on } \Sigma \\ v_0^+ = 0 & \text{on } \Gamma . \end{array} \right.$$

This problem is incompatible with the limit (as ε tends to zero) of the tangential part of the interface condition $\{\boldsymbol{\sigma}(v_\varepsilon, p_\varepsilon) \cdot \mathbf{n}\} = \mathcal{S}(v_\varepsilon^+ - v_\varepsilon^-) + \mathbf{h}$ on Σ so that it appears a boundary layer inside the porous medium. We describe this boundary layer by an asymptotic expansion at any order with a WKB method. We derive this expansion in power series of the small parameter $\sqrt{\varepsilon}$ for both the velocity v_ε and the pressure p_ε which are solutions of the transmission problem. This expansion is justified rigorously by proving uniform estimates for remainders of this expansion, Theorem 7.1.

An immediate corollary of our asymptotic expansion is the following convergence theorem:

Theorem 2.3. *We assume that the data in (1.1) satisfy:*

$$\mathbf{g}^- \in \mathbf{H}^5(\Omega_-), \quad \mathbf{g}^+ \in \mathbf{H}^4(\Omega_+), \quad \mathbf{l} \in \mathbf{H}^{\frac{9}{2}}(\Sigma) \quad \text{and} \quad \mathbf{h} \in \mathbf{H}^{\frac{9}{2}}(\Sigma).$$

Then, the solution v_ε for (1.1) given by Theorem 2.1 satisfies:

$$\begin{aligned} v_\varepsilon^+(x) &= v_0^+(x) + r_\varepsilon^+(x) && \text{for } x \in \Omega_+ \\ v_\varepsilon^-(x) &= v_0^-(x) + \tilde{v}_0^-(x, \frac{d(x)}{\sqrt{\varepsilon}}) + r_\varepsilon^-(x) && \text{for } x \in \Omega_- \end{aligned}$$

where

- v_0 is the solution of (2.3),
- \tilde{v}_0^- is a boundary layer term of the form $\tilde{v}_0^-(x, \mathbf{z}) = w_0(x) \exp(-\sqrt{\kappa}\mathbf{z})$,
- $d(x)$ is the euclidean distance to Σ ,
- r_ε is a remainder terms.

This remainder term satisfies the following estimate:

$$\varepsilon \|\mathbf{d}(r_\varepsilon^-)\|_{0,\Omega_-}^2 + \frac{\kappa}{4} \|r_\varepsilon^-\|_{0,\Omega_-}^2 + \frac{\mu C^2}{3} \|r_\varepsilon^+\|_{1,\Omega_+}^2 \leq C\varepsilon^{\frac{1}{2}}.$$

The concept of WKB expansion is rather classical in the modeling of problems arising in fluid mechanics. For instance in [7, 5, 6] the authors derive WKB expansions with boundary layer terms or thin layer asymptotics to describe penalization methods in the context of viscous incompressible flow.

In this work, one difficulty to validate the WKB expansion lies in the proof of both existence and regularity results for one part of the asymptotics which appear in this expansion at any order and which solve Darcy-Stokes problems with non-standard transmission conditions. It is possible to tackle these problems by carefully introducing a Dirichlet-to-Neumann operator which lead us to prove well-posedness results and elliptic regularity results simply for the Stokes operator with mixed boundary conditions.

The outline of the paper proceeds as follows. We prove the well-posedness result for Problem (2.1) together with uniform estimates with respect to the small parameter in Section 3. In Sections 4 and 5, one exhibits a formal WKB expansion for the solution of the transmission problem. The equations satisfied by the asymptotics at any order are explicit in section 5 and existence and regularity results concerning the asymptotics which satisfy Darcy-Stokes problems with non-standard transmission conditions are claimed in Prop. 5.1. The proof of this proposition is postponed to section 6. In section 7, one proves uniform errors estimates to validate this WKB expansion.

3. UNIFORM ESTIMATES

In this section, we prove the well-posedness result for Problem (2.1) together with uniform estimates, Theorem 2.1.

Notation 3.1. *We denote by $\|\cdot\|_{s,\mathcal{O}}$ the norm in the Sobolev space $H^s(\mathcal{O})$.*

Proof of Theorem 2.1. Since Ω_+ is a Lipschitz bounded domain, since $v^+ = 0$ on $\partial\Omega$ for $v \in V$, we have the following Poincaré inequality in Ω_+ :

$$(3.1) \quad \exists C > 0 \quad \text{s.t.} \quad \forall v \in V \quad \|\mathbf{d}(v^+)\|_{0,\Omega_+} \geq C \|v^+\|_{1,\Omega_+},$$

then, there hold $\forall v \in V$,

$$a_\varepsilon(v, v) \geq 2\varepsilon \|\mathbf{d}(v^-)\|_{0,\Omega_-}^2 + \kappa \|v^-\|_{0,\Omega_-}^2 + 2\mu C^2 \|v^+\|_{1,\Omega_+}^2 + \varepsilon^{-1} \|[v \cdot \mathbf{n}]\|_{0,\Sigma}^2 + \beta \|\{v \cdot \mathbf{n}\}\|_{0,\Sigma}^2 + \alpha \|v_\tau\|_{0,\Sigma}^2.$$

Hence, a_ε is V -coercive, and according to the Lax-Milgram Lemma, the problem (2.1) has a unique solution $v_\varepsilon \in V$ for all $\varepsilon > 0$. We also infer : if v satisfies (2.1), then $\forall v \in V$,

$$2\varepsilon \|\mathbf{d}(v^-)\|_{0,\Omega_-}^2 + \kappa \|v^-\|_{0,\Omega_-}^2 + 2\mu C^2 \|v^+\|_{1,\Omega_+}^2 + \varepsilon^{-1} \|[v \cdot \mathbf{n}]\|_{0,\Sigma}^2 + \beta \|\{v \cdot \mathbf{n}\}\|_{0,\Sigma}^2 + \alpha \|v_\tau\|_{0,\Sigma}^2 \leq \|v^+\|_{0,\Omega_+} \|\mathbf{g}^+\|_{0,\Omega_+} + \|v^-\|_{0,\Omega_-} \|\mathbf{g}^-\|_{0,\Omega_-} + \|[v]\|_{0,\Sigma} \|\mathbf{h}\|_{0,\Sigma} + \|\{v\}\|_{0,\Sigma} \|\mathbf{l}\|_{0,\Sigma}.$$

Let $\varepsilon > 0$ such that $\alpha \leq \frac{1}{2\varepsilon}$. Then, for all $\eta_1, \eta_2 > 0$, there holds

$$(3.2) \quad \begin{aligned} & 2\varepsilon \|\mathbf{d}(v^-)\|_{0,\Omega_-}^2 + \kappa \|v^-\|_{0,\Omega_-}^2 + 2\mu C^2 \|v^+\|_{1,\Omega_+}^2 + \frac{1}{2\varepsilon} \|[v \cdot \mathbf{n}]\|_{0,\Sigma}^2 + \beta \|\{v \cdot \mathbf{n}\}\|_{0,\Sigma}^2 + \alpha \|v_\tau\|_{0,\Sigma}^2 \\ & \leq \frac{\eta_1}{2} \|v\|_{0,\Omega}^2 + \frac{1}{2\eta_1} \|\mathbf{g}\|_{0,\Omega}^2 + \frac{\eta_2}{2} \|[v]\|_{0,\Sigma}^2 + \frac{1}{2\eta_2} \|\mathbf{h}\|_{0,\Sigma}^2 + \|\{v \cdot \mathbf{n}\}\|_{0,\Sigma} \|\mathbf{l}\|_{0,\Sigma} + \|\{v_\tau\}\|_{0,\Sigma} \|\mathbf{l}\|_{0,\Sigma}. \end{aligned}$$

We treat hereafter the last two terms : for all $\eta_3 > 0$, there holds

$$\|\{v \cdot \mathbf{n}\}\|_{0,\Sigma} \|\mathbf{l}\|_{0,\Sigma} \leq \frac{\eta_3}{2} \|\{v \cdot \mathbf{n}\}\|_{0,\Sigma}^2 + \frac{1}{2\eta_3} \|\mathbf{l}\|_{0,\Sigma}^2.$$

Since $\{v_\tau\} = v_\tau^+ - \frac{1}{2}[v_\tau]$, there holds

$$\|\{v_\tau\}\|_{0,\Sigma} \|\mathbf{l}\|_{0,\Sigma} \leq (\|v_\tau^+\|_{0,\Sigma} + \frac{1}{2}\|[v_\tau]\|_{0,\Sigma}) \|\mathbf{l}\|_{0,\Sigma},$$

hence, using a trace inequality, there exists $c > 0$ such that

$$\|\{v_\tau\}\|_{0,\Sigma} \|\mathbf{l}\|_{0,\Sigma} \leq (c\|v^+\|_{1,\Omega_+} + \frac{1}{2}\|[v_\tau]\|_{0,\Sigma}) \|\mathbf{l}\|_{0,\Sigma}.$$

We infer : for all $\eta_4, \eta_5 > 0$,

$$\|\{v_\tau\}\|_{0,\Sigma} \|\mathbf{l}\|_{0,\Sigma} \leq \frac{\eta_4 c}{2} \|v^+\|_{1,\Omega_+}^2 + \frac{c}{2\eta_4} \|[v_\tau]\|_{0,\Sigma}^2 + \frac{\eta_5}{4} \|[v_\tau]\|_{0,\Sigma}^2 + \frac{1}{4\eta_5} \|\mathbf{l}\|_{0,\Sigma}^2.$$

According to (3.2), we obtain

$$(3.3) \quad \begin{aligned} & 2\varepsilon \|\mathbf{d}(v^-)\|_{0,\Omega_-}^2 + \kappa \|v^-\|_{0,\Omega_-}^2 + 2\mu C^2 \|v^+\|_{1,\Omega_+}^2 + \frac{1}{2\varepsilon} \|[v \cdot \mathbf{n}]\|_{0,\Sigma}^2 + \alpha \|[v]\|_{0,\Sigma}^2 + \beta \|\{v \cdot \mathbf{n}\}\|_{0,\Sigma}^2 \\ & \leq \frac{\eta_1}{2} \|v\|_{0,\Omega}^2 + \frac{\eta_2}{2} \|[v]\|_{0,\Sigma}^2 + \frac{\eta_3}{2} \|\{v \cdot \mathbf{n}\}\|_{0,\Sigma}^2 + \frac{\eta_4 c}{2} \|v^+\|_{1,\Omega_+}^2 + \frac{\eta_5}{4} \|[v_\tau]\|_{0,\Sigma}^2 \\ & \quad + \frac{1}{2\eta_1} \|\mathbf{g}\|_{0,\Omega}^2 + \frac{1}{2\eta_2} \|\mathbf{h}\|_{0,\Sigma}^2 + \left(\frac{1}{2\eta_3} + \frac{c}{2\eta_4} + \frac{1}{4\eta_5}\right) \|\mathbf{l}\|_{0,\Sigma}^2. \end{aligned}$$

We fix constants η_i such that

$$\frac{\eta_4 c}{2} \leq \frac{\mu C^2}{3}, \quad \frac{\eta_1}{2} \leq \min\left(\frac{\kappa}{2}, \frac{\mu C^2}{3}\right), \quad \frac{\eta_2}{2} + \frac{\eta_5}{4} \leq \frac{\alpha}{2}, \quad \frac{\eta_3}{2} \leq \frac{\beta}{2}.$$

According to (3.3), we infer

$$(3.4) \quad \begin{aligned} \varepsilon \|\mathbf{d}(v^-)\|_{0,\Omega_-}^2 + \frac{\kappa}{4} \|v^-\|_{0,\Omega_-}^2 + \frac{\mu C^2}{3} \|v^+\|_{1,\Omega_+}^2 + \frac{1}{4\varepsilon} \|[v \cdot \mathbf{n}]\|_{0,\Sigma}^2 + \frac{\alpha}{4} \|[v]\|_{0,\Sigma}^2 + \frac{\beta}{4} \|\{v \cdot \mathbf{n}\}\|_{0,\Sigma}^2 \\ \leq C(\mu, \kappa, \alpha, \beta) (\|\mathbf{g}\|_{0,\Omega}^2 + \|\mathbf{h}\|_{0,\Sigma}^2 + \|\mathbf{l}\|_{0,\Sigma}^2). \end{aligned}$$

□

4. FORMAL ASYMPTOTIC EXPANSION

Hypothesis 4.1. *From now on, we assume that the surfaces Σ (interface) and Γ (external boundary) are smooth.*

Notation 4.2. *If X is a vector field defined on Σ , we denote by X_τ the tangent components of X : $X_\tau(x) = X(x) - (X(x) \cdot \mathbf{n}(x))\mathbf{n}(x)$, so that $X = (X \cdot \mathbf{n})\mathbf{n} + X_\tau$. For example, we denote by $(\mathbf{d}(v_\varepsilon) \cdot \mathbf{n})_\tau$ the tangent components of the normal constraint $\mathbf{d}(v_\varepsilon) \cdot \mathbf{n}$ defined on the interface Σ .*

Rewriting the transmission conditions set on Σ in (1.1), we use Formulation (1.2) to obtain the following equivalent problem,

$$(4.1) \quad -\varepsilon \Delta v_\varepsilon^- + \nabla p_\varepsilon^- + \kappa v_\varepsilon^- = \mathbf{g}^- \quad \text{in } \Omega_-$$

$$(4.2) \quad -\nabla \cdot \boldsymbol{\sigma}^+(v_\varepsilon^+, p_\varepsilon^+) = \mathbf{g}^+ \quad \text{in } \Omega_+$$

$$(4.3) \quad \nabla \cdot v_\varepsilon^- = 0 \quad \text{in } \Omega_-$$

$$(4.4) \quad \nabla \cdot v_\varepsilon^+ = 0 \quad \text{in } \Omega_+$$

$$(4.5) \quad 2\mu \mathbf{d}(v_\varepsilon^+) \cdot \mathbf{n} - p_\varepsilon^+ \mathbf{n} = 2\varepsilon \mathbf{d}(v_\varepsilon^-) \cdot \mathbf{n} - p_\varepsilon^- \mathbf{n} + \frac{\beta}{2} ((v_\varepsilon^+ + v_\varepsilon^-) \cdot \mathbf{n}) \mathbf{n} + \mathbf{l} \quad \text{on } \Sigma$$

$$(4.6) \quad \alpha (v_\varepsilon^+ - v_\varepsilon^-)_\tau = 2\mu (\mathbf{d}(v_\varepsilon^+) \cdot \mathbf{n})_\tau - (\mathbf{h} + \frac{1}{2}\mathbf{l})_\tau \quad \text{on } \Sigma$$

$$(4.7) \quad \varepsilon^{-1} (v_\varepsilon^+ - v_\varepsilon^-) \cdot \mathbf{n} = 2\mu (\mathbf{d}(v_\varepsilon^+) \cdot \mathbf{n}) \cdot \mathbf{n} - p_\varepsilon^+ - \frac{\beta}{4} (v_\varepsilon^+ + v_\varepsilon^-) \cdot \mathbf{n} - (\mathbf{h} + \frac{1}{2}\mathbf{l}) \cdot \mathbf{n} \quad \text{on } \Sigma$$

$$(4.8) \quad v_\varepsilon^+ = 0 \quad \text{on } \Gamma$$

(we recall that $\nabla \cdot \boldsymbol{\sigma}^+(v_\varepsilon^+, p_\varepsilon^+) = \mu \Delta v_\varepsilon^+ - \nabla p_\varepsilon^+$)

As already said, when ε tends to zero, we formally converge to Problem (2.3). The well posedness of this limit system is establish in Section 6. This problem is incompatible with the

limit of the jump condition (4.6) on the tangential velocity:

$$\alpha(v_0^+ - v_0^-)_\tau = 2\mu (\mathbf{d}(v_0^+) \cdot \mathbf{n})_\tau - (\mathbf{h} + \frac{1}{2}\mathbf{l})_\tau \quad \text{on } \Sigma.$$

Hence, it appears a boundary layer that we describe with a multiscale method, namely a WKB expansion: we exhibit series expansions in powers of $\sqrt{\varepsilon}$ for the flow v_ε , and the pressure p_ε . In the fluid part, we do not expect the formation of a boundary layer, therefore we look for an *Ansatz* on the form:

$$(4.9) \quad v_\varepsilon^+(x) \approx \sum_{j \geq 0} \varepsilon^{\frac{j}{2}} v_j^+(x),$$

$$(4.10) \quad p_\varepsilon^+(x) \approx \sum_{j \geq 0} \varepsilon^{\frac{j}{2}} p_j^+(x).$$

In the porous part, we denote by $d : \Omega_- \rightarrow \mathbb{R}^+$ the euclidean distance to Σ , $d(x) = \text{dist}(x, \Sigma)$. We describe the velocity and the pressure on the following way:

$$(4.11) \quad v_\varepsilon^-(x) \approx \sum_{j \geq 0} \varepsilon^{\frac{j}{2}} v_j^-(x, \frac{d(x)}{\sqrt{\varepsilon}}), \quad \text{with} \quad v_j^-(x, z) = v_j^-(x) + \tilde{v}_j^-(x, z),$$

$$(4.12) \quad p_\varepsilon^-(x) \approx \sum_{j \geq 0} \varepsilon^{\frac{j}{2}} p_j^-(x, \frac{d(x)}{\sqrt{\varepsilon}}), \quad \text{with} \quad p_j^-(x, z) = p_j^-(x) + \tilde{p}_j^-(x, z).$$

The terms \tilde{v}_j^- and \tilde{p}_j^- are boundary layer terms and defined on $\Sigma \times \mathbb{R}^+$. They are required to tend to 0 (such as their derivatives) when $z \rightarrow \infty$.

We use the following notations (see also [7]) : $x = (x_1, x_2, x_3)$ denotes the cartesian coordinates in \mathbb{R}^3 ,

$$\partial_i = \frac{\partial}{\partial x_i}, \quad \partial_z = \frac{\partial}{\partial z}, \quad \partial_{zz} = \frac{\partial^2}{\partial z^2}, \quad \nabla \tilde{p}^- = (\partial_1 \tilde{p}^-, \partial_2 \tilde{p}^-, \partial_3 \tilde{p}^-)^T.$$

For $\tilde{v}^- = (\tilde{v}^{-,1}, \tilde{v}^{-,2}, \tilde{v}^{-,3})^T$,

$$\nabla \cdot \tilde{v}^- = \partial_1 \tilde{v}^{-,1} + \partial_2 \tilde{v}^{-,2} + \partial_3 \tilde{v}^{-,3}, \quad (\mathbf{d}(\tilde{v}^-))_{ij} = \frac{1}{2} \left(\frac{\partial \tilde{v}^{-,j}}{\partial x_i} + \frac{\partial \tilde{v}^{-,i}}{\partial x_j} \right),$$

$$(\nabla d \cdot \nabla) \partial_z \tilde{v}^- = \sum_{j=1}^3 \frac{\partial d}{\partial x_j} \frac{\partial^2 \tilde{v}^-}{\partial x_j \partial z}, \quad \Delta \tilde{v}^- = \partial_1^2 \tilde{v}^- + \partial_2^2 \tilde{v}^- + \partial_3^2 \tilde{v}^-.$$

We recall that \mathbf{n} is the outward unit normal on $\partial\Omega_- = \Sigma$. Since we assume that Σ is a regular manifold, then d is smooth in a neighborhood of Σ . In this neighborhood, $|\nabla d| = 1$. In addition, for $x \in \Sigma$, $\nabla d(x) = -\mathbf{n}(x)$ so that $\partial_n d = -1$ on Σ . We extend \mathbf{n} by setting $\mathbf{n}(x) = -\nabla d(x)$. If X is a vector field defined in Ω_- , in a neighborhood of Σ , we define the tangential part of X by $X_\tau(x) = X(x) - (X(x) \cdot \mathbf{n}(x))\mathbf{n}(x)$.

By simple calculations there holds

$$\begin{aligned} \Delta \left(x \mapsto \tilde{v}^-(x, \frac{d(x)}{\sqrt{\varepsilon}}) \right) &= \varepsilon^{-1} \partial_{zz} \tilde{v}^-(x, \frac{d(x)}{\sqrt{\varepsilon}}) + \varepsilon^{-\frac{1}{2}} \left(2(\nabla d \cdot \nabla) \partial_z \tilde{v}^- + \Delta d \partial_z \tilde{v}^- \right) \left(x, \frac{d(x)}{\sqrt{\varepsilon}} \right) \\ &\quad + \Delta \tilde{v}^-(x, \frac{d(x)}{\sqrt{\varepsilon}}) \end{aligned}$$

$$\begin{aligned} \nabla \left(x \mapsto \tilde{p}^-(x; \frac{d(x)}{\sqrt{\varepsilon}}) \right) &= \varepsilon^{-\frac{1}{2}} \nabla d \partial_z \tilde{p}^-(x, \frac{d(x)}{\sqrt{\varepsilon}}) + \nabla \tilde{p}^-(x, \frac{d(x)}{\sqrt{\varepsilon}}) \\ \nabla \cdot \left(x \mapsto \tilde{v}^-(x; \frac{d(x)}{\sqrt{\varepsilon}}) \right) &= \varepsilon^{-\frac{1}{2}} \nabla d \cdot \partial_z \tilde{v}^-(x, \frac{d(x)}{\sqrt{\varepsilon}}) + \nabla \cdot \tilde{v}^-(x, \frac{d(x)}{\sqrt{\varepsilon}}) \\ \mathbf{d} \left(x \mapsto \tilde{v}^-(x, \frac{d(x)}{\sqrt{\varepsilon}}) \right) \cdot \mathbf{n} &= -\varepsilon^{-\frac{1}{2}} \frac{1}{2} (\partial_z \tilde{v}^- + (\partial_z \tilde{v}^- \cdot \mathbf{n}) \mathbf{n}) (x; \frac{d(x)}{\sqrt{\varepsilon}}) + \mathbf{d}(\tilde{v}^-) \cdot \mathbf{n} (x; \frac{d(x)}{\sqrt{\varepsilon}}) \quad \text{on } \Sigma. \end{aligned}$$

We insert the *Ansatz* (4.9)-(4.10)-(4.11)-(4.12) in equations (4.1)-(4.8) and we perform the identification of terms with the same power in $\sqrt{\varepsilon}$.

4.1. Asymptotic expansion of (4.1).

Order $\varepsilon^{-\frac{1}{2}}$. There holds $\nabla d \partial_z \tilde{p}_0^- = 0$, and since $\tilde{p}_0^- \rightarrow 0$ when $z \rightarrow \infty$, we infer

$$(4.13) \quad \tilde{p}_0^- = 0 \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

Order ε^0 . According to (4.13), (v_0^-, p_0^-) solves

$$(4.14) \quad -\partial_{zz} \tilde{v}_0^- + \kappa \tilde{v}_0^- + \kappa v_0^- + \nabla p_0^- + \nabla d \partial_z \tilde{p}_1^- = \mathbf{g}^- \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

Hence, taking the limit of this equation when $z \rightarrow \infty$, since \tilde{v}_0^- and all its derivatives tends to zero when z tends to $+\infty$, we infer

$$(4.15) \quad \kappa v_0^- + \nabla p_0^- = \mathbf{g}^- \quad \text{in } \Omega_- ,$$

and by difference with the previous equation, we obtain

$$(4.16) \quad -\partial_{zz} \tilde{v}_0^- + \kappa \tilde{v}_0^- + \nabla d \partial_z \tilde{p}_1^- = 0 \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

Order $\varepsilon^{\frac{1}{2}}$. (v_1^-, p_1^-) solves

$$(4.17) \quad -\partial_{zz} \tilde{v}_1^- + \kappa \tilde{v}_1^- + \kappa v_1^- + \nabla p_1^-(x) + \nabla \tilde{p}_1^- + \nabla d \partial_z \tilde{p}_2^- = (2(\nabla d \cdot \nabla) \partial_z \tilde{v}_0^- + \Delta d \partial_z \tilde{v}_0^-),$$

and taking the limit of this equation when $z \rightarrow \infty$, we infer

$$(4.18) \quad \kappa v_1^-(x) + \nabla p_1^-(x) = 0 \quad \text{in } \Omega_- .$$

By difference with the previous equation, we obtain

$$(4.19) \quad -\partial_{zz} \tilde{v}_1^- + \kappa \tilde{v}_1^- + \nabla \tilde{p}_1^- + \nabla d \partial_z \tilde{p}_2^- = 2(\nabla d \cdot \nabla) \tilde{v}_{0,Z}^- + \Delta d \partial_z \tilde{v}_0^- , \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

Order $\varepsilon^{\frac{j}{2}}$, $j \geq 2$. (v_j^-, p_j^-) solves

$$(4.20) \quad \begin{aligned} -\partial_{zz} \tilde{v}_j^- + \kappa \tilde{v}_j^- + \kappa v_j^- + \nabla p_j^- + \nabla \tilde{p}_j^- + \nabla d \partial_z \tilde{p}_{j+1}^- \\ = (2(\nabla d \cdot \nabla) \partial_z \tilde{v}_{j-1}^- + \Delta d \partial_z \tilde{v}_{j-1}^-) + \Delta (v_{j-2}^- + \tilde{v}_{j-2}^-), \end{aligned}$$

hence, we infer

$$(4.21) \quad \kappa v_j^- + \nabla p_j^- = \Delta v_{j-2}^- \quad \text{in } \Omega_- ,$$

$$(4.22) \quad \begin{aligned} -\partial_{zz} \tilde{v}_j^- + \kappa \tilde{v}_j^- + \nabla \tilde{p}_j^- + \nabla d \partial_z \tilde{p}_{j+1}^- \\ = (2(\nabla d \cdot \nabla) \partial_z \tilde{v}_{j-1}^- + \Delta d \partial_z \tilde{v}_{j-1}^-) + \Delta \tilde{v}_{j-2}^- \quad \text{in } \Omega_- \times \mathbb{R}^+ . \end{aligned}$$

4.2. Asymptotic expansion of (4.2).

Order ε^0 . v_0^+ satisfies

$$(4.23) \quad -\nabla \cdot \boldsymbol{\sigma}^+(v_0^+, p_0^+) = \mathbf{g}^+ \quad \text{in } \Omega_+ .$$

Order $\varepsilon^{\frac{j}{2}}$, $j \geq 1$. For all $j \in \mathbb{N}$, v_j^+ satisfies

$$(4.24) \quad -\nabla \cdot \sigma^+(v_j^+, p_j^+) = 0 \quad \text{in } \Omega_+.$$

4.3. Asymptotic expansion of (4.3).

Order $\varepsilon^{-\frac{1}{2}}$. There holds

$$(4.25) \quad \nabla d \cdot \partial_z \tilde{v}_0^- = 0 \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

Since $\nabla d = -n$, since \tilde{v}_0^- tends to zero when z tends to $+\infty$, we obtain

$$(4.26) \quad \tilde{v}_0^- \cdot n = 0 \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

Order $\varepsilon^{\frac{j}{2}}$, $j \geq 0$. There holds

$$(4.27) \quad \nabla d \cdot \partial_z \tilde{v}_{j+1}^- + \nabla \cdot v_j^- + \nabla \cdot \tilde{v}_j^- = 0 \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

We infer

$$(4.28) \quad \nabla \cdot v_j^- = 0 \quad \text{in } \Omega_- ,$$

$$(4.29) \quad \nabla d \cdot \partial_z \tilde{v}_{j+1}^- + \nabla \cdot \tilde{v}_j^- = 0 \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

4.4. Asymptotic expansion of (4.4).

Order $\varepsilon^{\frac{j}{2}}$. For all $j \in \mathbb{N}$, v_j^+ satisfies

$$(4.30) \quad \nabla \cdot v_j^+ = 0 \quad \text{in } \Omega_+ .$$

4.5. Asymptotic expansion of (4.5).

Order ε^0 .

$$(4.31) \quad 2\mu \mathbf{d}(v_0^+) \cdot n - p_0^+ n = -p_0^- n - \tilde{p}_0^-(\cdot, 0)n + \frac{\beta}{2} ((v_0^+ + v_0^- + \tilde{v}_0^-(\cdot, 0)) \cdot n) n + l \quad \text{on } \Sigma$$

Order $\varepsilon^{\frac{1}{2}}$.

$$(4.32) \quad 2\mu \mathbf{d}(v_1^+) \cdot n - p_1^+ n = -\frac{1}{2} (\partial_z \tilde{v}_0^-(\cdot, 0) + (\partial_z \tilde{v}_0^-(\cdot, 0) \cdot n)n) - p_1^- n - \tilde{p}_1^-(\cdot, 0)n \\ + \frac{\beta}{2} ((v_1^+ + v_1^- + \tilde{v}_1^-(\cdot, 0)) \cdot n) n \quad \text{on } \Sigma$$

Order $\varepsilon^{\frac{j}{2}}$, $j \geq 2$.

$$(4.33) \quad 2\mu \mathbf{d}(v_j^+) \cdot n - p_j^+ n = -\frac{1}{2} (\partial_z \tilde{v}_{j-1}^-(\cdot, 0) + (\partial_z \tilde{v}_{j-1}^-(\cdot, 0) \cdot n)n) + 2\mathbf{d}(v_{j-2}^-) \cdot n \\ + 2\mathbf{d}(\tilde{v}_{j-2}^-)(\cdot, 0) \cdot n - p_j^- n - \tilde{p}_j^-(\cdot, 0)n \\ + \frac{\beta}{2} ((v_j^+ + v_j^- + \tilde{v}_j^-(\cdot, 0)) \cdot n) n \quad \text{on } \Sigma.$$

4.6. Asymptotic expansion of (4.6).

Order ε^0 .

$$(4.34) \quad \alpha(v_0^+ - v_0^- - \tilde{v}_0^-(\cdot, 0))_\tau = 2\mu (\mathbf{d}(v_0^+) \cdot n)_\tau - \left(h + \frac{1}{2} l \right)_\tau \quad \text{on } \Sigma$$

Order $\varepsilon^{\frac{j}{2}}, j \geq 1$.

$$(4.35) \quad \alpha(v_j^+ - v_j^- - \tilde{v}_j^-(\cdot, 0))_\tau = 2\mu \left(\mathbf{d}(v_j^+) \cdot \mathbf{n} \right)_\tau \quad \text{on } \Sigma$$

4.7. Asymptotic expansion of (4.7).

Order ε^0 .

$$(4.36) \quad (v_0^+ - v_0^- - \tilde{v}_0^-(\cdot, 0)) \cdot \mathbf{n} = 0 \quad \text{on } \Sigma$$

Order $\varepsilon^{\frac{1}{2}}$.

$$(4.37) \quad (v_1^+ - v_1^- - \tilde{v}_1^-(\cdot, 0)) \cdot \mathbf{n} = 0 \quad \text{on } \Sigma$$

Order ε .

(4.38)

$$(v_2^+ - v_2^- - \tilde{v}_2^-(\cdot, 0)) \cdot \mathbf{n} = 2\mu \left(\mathbf{d}(v_0^+) \cdot \mathbf{n} \right) \cdot \mathbf{n} - p_0^+ - \frac{\beta}{4} (v_0^+ + v_0^- + \tilde{v}_0^-(\cdot, 0)) \cdot \mathbf{n} - \left(\mathbf{h} + \frac{1}{2} \mathbf{l} \right) \cdot \mathbf{n} \quad \text{on } \Sigma$$

Order $\varepsilon^{\frac{j}{2}}, j > 2$.

(4.39)

$$(v_j^+ - v_j^- - \tilde{v}_j^-(\cdot, 0)) \cdot \mathbf{n} = 2\mu \left(\mathbf{d}(v_{j-2}^+) \cdot \mathbf{n} \right) \cdot \mathbf{n} - p_j^+ - \frac{\beta}{4} \left(v_{j-2}^+ + v_{j-2}^- + \tilde{v}_{j-2}^-(\cdot, 0) \right) \cdot \mathbf{n} \quad \text{on } \Sigma.$$

4.8. Asymptotic expansion of (4.8).

Order δ^j . For all $j \in \mathbb{N}$, v_j^+ satisfies the external boundary condition

$$(4.40) \quad v_j^+ = 0 \quad \text{in } \Gamma.$$

5. DETERMINATION, EXISTENCE AND REGULARITY OF THE ASYMPTOTICS

Let $k \geq 1$. We assume that the data in (4.1)-(4.8) satisfy:

$$(5.1) \quad \mathbf{g}^- \in \mathbf{H}^k(\Omega_-), \quad \mathbf{g}^+ \in \mathbf{H}^{k-1}(\Omega_+), \quad \mathbf{l} \in \mathbf{H}^{k-\frac{1}{2}}(\Sigma) \quad \text{and} \quad \mathbf{h} \in \mathbf{H}^{k-\frac{1}{2}}(\Sigma).$$

5.1. Equations satisfied by asymptotics of order 0.

5.1.1. *Determination of \tilde{p}_0^- and $\tilde{v}_0^- \cdot \mathbf{n}$.* We have already obtained in (4.13) and (4.26) that

$$(5.2) \quad \tilde{p}_0^- = 0 \quad \text{and} \quad \tilde{v}_0^- \cdot \mathbf{n} = 0 \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

5.1.2. *Determination of v_0^- , p_0^- , v_0^+ and p_0^+ .* According to (4.36) and (4.26), we infer

$$(v_0^+ - v_0^-) \cdot \mathbf{n} = 0 \quad \text{on } \Sigma.$$

According to and (4.13) and (4.31)

$$2\mu \mathbf{d}(v_0^+) \cdot \mathbf{n} - p_0^+ \mathbf{n} = -p_0^- \mathbf{n} + \frac{\beta}{2} \left((v_0^+ + v_0^-) \cdot \mathbf{n} \right) \mathbf{n} + \mathbf{l} \quad \text{on } \Sigma.$$

Finally, according to (4.15), (4.28), (4.31), and according to (4.23) and (4.30) when $j = 0$, (v_0^-, v_0^+) , and (p_0^-, p_0^+) satisfy the transmission problem :

$$(5.3) \quad \begin{cases} \kappa v_0^- + \nabla p_0^- = g^- & \text{in } \Omega_- \\ \nabla \cdot v_0^- = 0 & \text{in } \Omega_- \\ -\nabla \cdot \boldsymbol{\sigma}^+(v_0^+, p_0^+) = g^+ & \text{in } \Omega_+ \\ \nabla \cdot v_0^+ = 0 & \text{in } \Omega_+ \\ (v_0^+ - v_0^-) \cdot \mathbf{n} = 0 & \text{on } \Sigma \\ 2\mu \mathbf{d}(v_0^+) \cdot \mathbf{n} - p_0^+ \mathbf{n} = -p_0^- \mathbf{n} + \frac{\beta}{2} ((v_0^+ + v_0^-) \cdot \mathbf{n}) \mathbf{n} + l & \text{on } \Sigma \\ v_0^+ = 0 & \text{on } \Gamma. \end{cases}$$

Let us claim the following existence and regularity result concerning such a problem

Proposition 5.1. *Let $k \geq 1$. Let $g^- \in \mathbf{H}^k(\Omega_-)$, $g^+ \in \mathbf{H}^{k-1}(\Omega_+)$, $l \in \mathbf{H}^{k-\frac{1}{2}}(\Sigma)$ and $h \in \mathbf{H}^{k-\frac{1}{2}}(\Sigma)$. We assume that h satisfies the compatibility condition*

$$\int_{\Sigma} h \, d\sigma = 0.$$

We consider the following problem:

$$(5.4) \quad \begin{cases} \kappa v^- + \nabla p^- = g^- & \text{in } \Omega_- \\ \nabla \cdot v^- = 0 & \text{in } \Omega_- \\ -\nabla \cdot \boldsymbol{\sigma}^+(v^+, p^+) = g^+ & \text{in } \Omega_+ \\ \nabla \cdot v^+ = 0 & \text{in } \Omega_+ \\ (v^+ - v^-) \cdot \mathbf{n} = h & \text{on } \Sigma \\ 2\mu \mathbf{d}(v^+) \cdot \mathbf{n} - p^+ \mathbf{n} = -p^- \mathbf{n} + \frac{\beta}{2} ((v^+ + v^-) \cdot \mathbf{n}) \mathbf{n} + l & \text{on } \Sigma \\ v^+ = 0 & \text{on } \Gamma. \end{cases}$$

Then (5.4) admits a weak solution unique up to additive constants for p^+ and p^- . It satisfies the regularity properties:

$$v^+ \in \mathbf{H}^{k+1}(\Omega_+), \quad v^- \in \mathbf{H}^k(\Omega_-), \quad p^+ \in \mathbf{H}^k(\Omega_+) \quad \text{and} \quad p^- \in \mathbf{H}^{k+1}(\Omega_-).$$

This proposition will be proved in Section 6.

Using Proposition 5.1 with $g^\pm = g^\pm$, $l = l$ and $h = 0$, using the regularity of the data (5.1), we obtain the existence and uniqueness (up to an additive constant for the pressures) of the asymptotics $v_0^+ \in \mathbf{H}^{k+1}(\Omega_+)$, $v_0^- \in \mathbf{H}^k(\Omega_-)$, $p_0^+ \in \mathbf{H}^k(\Omega_+)$ and $p_0^- \in \mathbf{H}^{k+1}(\Omega_-)$.

5.1.3. *Determination of $(\tilde{v}_0^-)_\tau$.* The tangential component of (4.16) reduce to

$$-\partial_{zz}(\tilde{v}_0^-)_\tau + \kappa(\tilde{v}_0^-)_\tau = 0 \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

Hence, since $\tilde{v}_0^- \rightarrow 0$ when $z \rightarrow \infty$, we obtain

$$\tilde{v}_0^-(x, z)_\tau = \tilde{v}_0^-(x, 0)_\tau \exp(-\sqrt{\kappa}z) \quad \text{in } \Omega_- \times \mathbb{R}^+,$$

According to (4.34), we infer

$$\tilde{v}_0^-(x, 0)_\tau = (v_0^+ - v_0^-)_\tau - \frac{2\mu}{\alpha} (\mathbf{d}(v_0^+) \cdot \mathbf{n})_\tau + \frac{1}{\alpha} \left(h + \frac{1}{2} l \right)_\tau \quad \text{on } \Sigma.$$

Hereafter, we introduce w_0^0 which is a tangential extension of $(v_0^+ - v_0^-)_\tau - \frac{2\mu}{\alpha}(\mathbf{d}(v_0^+) \cdot \mathbf{n})_\tau + \frac{1}{\alpha}(h + \frac{1}{2}l)_\tau$ in the domain Ω_- . We choose this extension w_0^0 such that it has a support in a tubular neighborhood of the interface Σ . Since $v_0^+ \in \mathbf{H}^{k+1}(\Omega_+)$, since $v_0^- \in \mathbf{H}^k(\Omega_-)$, since l and h belong to the space $\mathbf{H}^{k-\frac{1}{2}}(\Sigma)$, we can take this extension satisfying $w_0^0 \in \mathbf{H}^k(\Omega_-)$, and using (5.2) we obtain

$$(5.5) \quad \tilde{v}_0^-(x, z) = w_0^0(x) \exp(-\sqrt{\kappa} z) \quad \text{in } \Omega_- \times \mathbb{R}^+,$$

and then \tilde{v}_0^- is completely defined by (5.5).

5.2. Equations satisfied by asymptotics of order 1.

5.2.1. *Determination of \tilde{p}_1^- .* According to (4.26), by taking the normal components in (4.16), we obtain $\tilde{p}_{1,z}^- = 0$. Hence, since $\tilde{p}_1^- \rightarrow 0$ when $z \rightarrow \infty$, we obtain

$$(5.6) \quad \tilde{p}_1^- = 0 \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

5.2.2. *Determination of $\tilde{v}_1^- \cdot \mathbf{n}$.* According to (4.29)

$$\partial_z \tilde{v}_1^- \cdot \nabla d = -\nabla \cdot \tilde{v}_0^- = -(\nabla \cdot w_0^0) \exp(-\sqrt{\kappa} z) \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

Hence, since \tilde{v}_1^- tends to zero when z tends to $+\infty$,

$$(5.7) \quad \tilde{v}_1^- \cdot \mathbf{n} = \nu_1^0 \exp(-\sqrt{\kappa} z) \quad \text{with } \nu_1^0 = \frac{-1}{\sqrt{\kappa}}(\nabla \cdot w_0^0) \in \mathbf{H}^{k-1}(\Omega_-).$$

5.2.3. *Determination of (v_1^-, v_1^+) , and (p_1^-, p_1^+) .* According to (4.26) and (5.6), the transmission condition (4.32) writes :

$$2\mu \mathbf{d}(v_1^+) \cdot \mathbf{n} - p_1^+ \mathbf{n} = -\frac{1}{2} \partial_z \tilde{v}_0^- - p_1^- \mathbf{n} + \frac{\beta}{2} ((v_1^+ + v_1^- + \tilde{v}_1^-) \cdot \mathbf{n}) \mathbf{n} \quad \text{on } \Sigma$$

Hence, according to (4.18), (4.28), and to (4.24) and (4.30) when $j = 1$, and according to (4.32), (4.37), (5.5) and (5.7), (v_1^-, v_1^+) , and (p_1^-, p_1^+) satisfy the transmission problem:

$$(5.8) \quad \begin{cases} \kappa v_1^- + \nabla p_1^- = 0 & \text{in } \Omega_- \\ \nabla \cdot v_1^- = 0 & \text{in } \Omega_- \\ -\nabla \cdot \boldsymbol{\sigma}^+(v_1^+, p_1^+) = 0 & \text{in } \Omega_+ \\ \nabla \cdot v_1^+ = 0 & \text{in } \Omega_+ \\ (v_1^+ - v_1^-) \cdot \mathbf{n} = \nu_1^0 & \text{on } \Sigma \\ 2\mu \mathbf{d}(v_1^+) \cdot \mathbf{n} - p_1^+ \mathbf{n} = -p_1^- \mathbf{n} + \frac{\beta}{2} ((v_1^+ + v_1^-) \cdot \mathbf{n}) \mathbf{n} + \frac{1}{2} \sqrt{\kappa} w_0^0 + \frac{\beta}{2} \nu_1^0 \mathbf{n} & \text{on } \Sigma \\ v_1^+ = 0 & \text{on } \Gamma. \end{cases}$$

We apply Proposition 5.1 with

$$g^\pm = 0, \quad l = \frac{1}{2} \sqrt{\kappa} w_0^0 + \frac{\beta}{2} \nu_1^0 \mathbf{n} \in \mathbf{H}^{k-\frac{3}{2}}(\Sigma), \quad \text{and } h = \nu_1^0 \in \mathbf{H}^{k-\frac{3}{2}}(\Sigma),$$

and $k \geq 2$. We remark that ν_1^0 satisfies the compatibility condition $\int_\Sigma \nu_1^0 d\sigma = 0$ since ν_1^0 is the divergence of a tangential vector field. Indeed, the divergence of w_0^0 on the surface Σ is the divergence operator on Σ applied to a tangent vector field on Σ ; hence by the Stokes formula, the integral of $\nabla \cdot w_0^0$ vanishes since Σ does not have boundary.

Hence we obtain the existence and uniqueness (up to an additive constant for the pressures) for the asymptotics v_1^-, v_1^+, p_1^- and p_1^+ which satisfy:

$$(5.9) \quad v_1^+ \in \mathbf{H}^k(\Omega_+), \quad v_1^- \in \mathbf{H}^{k-1}(\Omega_-), \quad p_1^+ \in H^{k-1}(\Omega_+) \quad \text{and} \quad p_1^- \in H^k(\Omega_-).$$

5.2.4. *Determination of $(\tilde{v}_1^-)_\tau$.* Using $\tilde{p}_1^- = 0$, then (4.19) writes now

$$(5.10) \quad -\partial_{zz}\tilde{v}_1^- + \kappa\tilde{v}_1^- + \nabla d \partial_z \tilde{p}_2^- = -\sqrt{\kappa} (2(\nabla d \cdot \nabla)w_0^0 + \Delta d w_0^0) \exp(-\sqrt{\kappa} z), \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

Taking the tangential components of the previous equation, we obtain that $\tilde{v}_1^-(x, z)_\tau$ solves

$$(5.11) \quad -\partial_{zz}(\tilde{v}_1^-)_\tau + \kappa(\tilde{v}_1^-)_\tau = -\sqrt{\kappa} (2(\nabla d \cdot \nabla)w_0^0 + \Delta d w_0^0)_\tau \exp(-\sqrt{\kappa} z) \quad \text{in } \Omega_- \times \mathbb{R}^+,$$

with (see (4.35)):

$$(5.12) \quad (\tilde{v}_1^-)_\tau(x, 0) = (v_1^+ - v_1^-)_\tau - \frac{2\mu}{\alpha} (\mathbf{d}(v_1^+) \cdot \mathbf{n})_\tau \quad \text{on } \Sigma.$$

On the one hand, we introduce w_1^0 which is a tangential extension of $(v_1^+ - v_1^-)_\tau - \frac{2\mu}{\alpha} (\mathbf{d}(v_1^+) \cdot \mathbf{n})_\tau$ in the domain Ω_- . We choose this extension w_1^0 such that it has a support in a tubular neighborhood of the interface Σ . Since $v_1^+ \in \mathbf{H}^k(\Omega_+)$ and $v_1^- \in \mathbf{H}^{k-1}(\Omega_-)$, we can take an extension satisfying

$$w_1^0 \in \mathbf{H}^{k-1}(\Omega_-).$$

On the other hand, we denote

$$w_1^1 = -\frac{1}{2} (2\nabla d \cdot \nabla w_0^0 + \Delta d w_0^0)_\tau \exp(-\sqrt{\kappa} z).$$

We remark that $w_1^1 \in \mathbf{H}^{k-1}(\Omega_-)$ since $w_0^0 \in \mathbf{H}^k(\Omega_-)$.

By solving (5.11) with the boundary condition (5.12) and we obtain that

$$(5.13) \quad (\tilde{v}_1^-)_\tau(x, z) = (w_1^0(x) + z w_1^1(x)) \exp(-\sqrt{\kappa} z) \quad \text{with } w_1^0, w_1^1 \text{ in } \mathbf{H}^{k-1}(\Omega_-).$$

5.3. **Characterization of the asymptotics of order $j \geq 2$.** First let us write the equations satisfied by the asymptotics at order j .

At order 2, from the normal components of (4.19), using (5.5), we have

$$(-\partial_{zz}\tilde{v}_1^- + \kappa\tilde{v}_1^-) \cdot \mathbf{n} - \partial_z \tilde{p}_2^- = -\sqrt{\kappa} (2(\nabla d \cdot \nabla)w_0^0 + \Delta d w_0^0) \cdot \mathbf{n} \exp(-\sqrt{\kappa} z).$$

Hence, according to (5.7), \tilde{p}_2^- satisfies

$$\partial_z \tilde{p}_2^- = \sqrt{\kappa} (2(\nabla d \cdot \nabla)w_0^0) \cdot \mathbf{n} \exp(-\sqrt{\kappa} z), \quad \text{in } \Omega_- \times \mathbb{R}^+.$$

We infer

$$(5.14) \quad \tilde{p}_2^- = q_2^0 \exp(-\sqrt{\kappa} z) \quad \text{with } q_2^0 = -(2(\nabla d \cdot \nabla)w_0^0) \cdot \mathbf{n} \in \mathbf{H}^{k-1}(\Omega_-).$$

At order $j \geq 3$, from the normal part of (4.22), replacing j by $j - 1$, we obtain that

$$\partial_z \tilde{p}_j^- = \left(-\partial_{zz}\tilde{v}_{j-1}^- + \kappa\tilde{v}_{j-1}^- + \nabla \tilde{p}_{j-1}^- - 2\nabla d \cdot \nabla \partial_z \tilde{v}_{j-2}^- - \Delta d \partial_z \tilde{v}_{j-2}^- - \Delta \tilde{v}_{j-3}^- \right) \cdot \mathbf{n},$$

so that

$$(5.15) \quad \tilde{p}_j^-(x, z) = - \int_z^{+\infty} \left(-\partial_{zz}\tilde{v}_{j-1}^- + \kappa\tilde{v}_{j-1}^- + \nabla\tilde{p}_{j-1}^- - 2\nabla d \cdot \nabla\partial_z\tilde{v}_{j-2}^- - \Delta d\partial_z\tilde{v}_{j-2}^- - \Delta\tilde{v}_{j-3}^- \right) \cdot n \, dz.$$

From (4.29), there holds

$$\partial_z\tilde{v}_j^- \cdot n = \nabla \cdot \tilde{v}_{j-1}^- \quad \text{in } \Omega_- \times \mathbb{R}^+,$$

hence

$$(5.16) \quad \tilde{v}_j^- \cdot n(x, z) = - \int_z^{+\infty} \nabla \cdot \tilde{v}_{j-1}^-(\cdot, \xi) \, d\xi.$$

From (4.21), (4.24), (4.28), (4.30), (4.33), (4.40) and (4.39), we obtain the following system characterizing v_j^+ , p_j^+ , v_j^- and p_j^- :

$$(5.17) \quad \begin{cases} \kappa v_j^- + \nabla p_j^- = \Delta v_{j-2}^- & \text{in } \Omega_- \\ \nabla \cdot v_j^- = 0 & \text{in } \Omega_- \\ -\nabla \cdot \sigma^+(v_j^+, p_j^+) = 0 & \text{in } \Omega_+ \\ \nabla \cdot v_j^+ = 0 & \text{in } \Omega_+ \\ (v_j^+ - v_j^-) \cdot n = h_j & \text{on } \Sigma \\ 2\mu \mathbf{d}(v_j^+) \cdot n - p_j^+ n = -p_j^- n + \frac{\beta}{2} \left((v_j^+ + v_j^-) \cdot n \right) n + l_j & \text{on } \Sigma \\ v_j^+ = 0 & \text{on } \Gamma, \end{cases}$$

where the data are given by:

$$\begin{aligned} l_j &= - \left(\partial_z \tilde{v}_{j-1}^-(\cdot, 0) + (\partial_z \tilde{v}_{j-1}^- \cdot n)n \right) + 2(\mathbf{d}(v_{j-2}^-) \cdot n + \mathbf{d}(\tilde{v}_{j-2}^-)(\cdot, 0) \cdot n) + \left(\frac{\beta}{2} \tilde{v}_j^-(\cdot, 0) \cdot n - \tilde{p}_j^-(\cdot, 0) \right) n, \\ h_2 &= \tilde{v}_2^- \cdot n(\cdot, 0) + 2\mu (\mathbf{d}(v_0^+) \cdot n) \cdot n - p_0^+ - \frac{\beta}{4} (v_0^+ + v_0^- + \tilde{v}_0^-) \cdot n - (h + \frac{1}{2}l) \cdot n, \\ h_j &= \tilde{v}_j^- \cdot n(\cdot, 0) + 2\mu (\mathbf{d}(v_{j-2}^+) \cdot n) \cdot n - p_{j-2}^+ - \tilde{p}_{j-2}^- - \frac{\beta}{4} (v_{j-2}^+ + v_{j-2}^- + \tilde{v}_{j-2}^-) \cdot n \quad \text{for } j \geq 3. \end{aligned}$$

Taking the tangential part of (4.22) and writing (4.35), we obtain that $(\tilde{v}_j^-)_\tau$ satisfies

$$(5.18) \quad \begin{cases} -\partial_{zz}(\tilde{v}_j^-)_\tau + \kappa(\tilde{v}_j^-)_\tau = \left(-\nabla\tilde{p}_j^- + 2(\nabla d \cdot \nabla)\partial_z\tilde{v}_{j-1}^- + \Delta d\partial_z\tilde{v}_{j-1}^- + \Delta\tilde{v}_{j-2}^- \right)_\tau & \text{for } z \geq 0, \\ (\tilde{v}_j^-)_\tau(\cdot, 0) = (v_j^+ - v_j^-)_\tau - \frac{2\mu}{\alpha} (\mathbf{d}(v_j^+) \cdot n)_\tau. \end{cases}$$

5.4. Existence and regularity of the asymptotics.

Proposition 5.2. *Let $k \geq 1$. We assume that the data in (4.1)-(4.8) satisfy:*

$$g^- \in \mathbf{H}^k(\Omega_-), \quad g^+ \in \mathbf{H}^{k-1}(\Omega_+), \quad l \in \mathbf{H}^{k-\frac{1}{2}}(\Sigma) \quad \text{and} \quad h \in \mathbf{H}^{k-\frac{1}{2}}(\Sigma).$$

Then, for all $j \in \{0, \dots, k-1\}$, there exists asymptotics v_j^+ , v_j^- , \tilde{v}_j^- , p_j^+ , p_j^- and \tilde{p}_j^- , satisfying

- (5.2), (5.3) and (5.5), when $j = 0$,
- (5.6), (5.7), (5.8), and (5.13), when $j = 1$,
- (5.14), (5.15), (5.16), (5.17) and (5.18), when $j \geq 2$.

There holds

$$v_j^+ \in \mathbf{H}^{k-j+1}(\Omega_+), \quad v_j^- \in \mathbf{H}^{k-j}(\Omega_-), \quad p_j^+ \in \mathbf{H}^{k-j}(\Omega_+) \quad \text{and} \quad p_j^- \in \mathbf{H}^{k-j+1}(\Omega_-).$$

In addition, the boundary layer terms \tilde{v}_j^- are on the following form: $\tilde{v}_0^- \cdot \mathbf{n} = 0$ and

$$\begin{aligned} \tilde{v}_j^- \cdot \mathbf{n}(x, \mathbf{z}) &= \sum_{l=0}^{j-1} \nu_j^l(x) \mathbf{z}^l \exp(-\sqrt{\kappa} \mathbf{z}) \quad \text{with} \quad \nu_j^l \in \mathbf{H}^{k-j}(\Omega_-), \\ (\tilde{v}_j^-)_\tau(x, \mathbf{z}) &= \sum_{l=0}^j w_j^l(x) \mathbf{z}^l \exp(-\sqrt{\kappa} \mathbf{z}) \quad \text{with} \quad w_j^l \in \mathbf{H}^{k-j}(\Omega_-). \end{aligned}$$

Concerning the boundary layer terms for the pressure, $\tilde{p}_0^- = \tilde{p}_1^- = 0$ and for $j \geq 2$,

$$\tilde{p}_j^-(x, \mathbf{z}) = \sum_{l=0}^{j-2} q_j^l(x) \mathbf{z}^l \exp(-\sqrt{\kappa} \mathbf{z}) \quad \text{with} \quad q_j^l \in \mathbf{H}^{k-j+1}(\Omega_-)$$

Proof. From Sections 5.1 and 5.2, the property is true for $j = 0$ and $j = 1$.

Let us assume that the property is true up to order $j - 1$, with $j \leq k - 1$. We claim without proof the following lemma:

Lemma 5.3. For all $l \in \mathbb{N}$,

$$\int_{\mathbf{z}}^{+\infty} \xi^l \exp(-\sqrt{\kappa} \xi) d\xi = \left(\sum_{j=0}^l \frac{l!}{j!} \frac{1}{\kappa^{\frac{l-j+1}{2}}} \mathbf{z}^j \right) \exp(-\sqrt{\kappa} \mathbf{z}).$$

First step: construction of \tilde{p}_j^- .

Concerning \tilde{p}_j^- , if $j = 2$, we have Property (5.14). If $j > 2$, then from (5.15) and with the induction property at order $j - 1$ and $j - 2$, $\partial_{\mathbf{z}} \tilde{p}_j^-$ writes:

$$(5.19) \quad \partial_{\mathbf{z}} \tilde{p}_j^-(x, \mathbf{z}) = \left(\sum_{l=0}^{j-2} \psi_l(x) \mathbf{z}^l \right) \exp(-\sqrt{\kappa} \mathbf{z})$$

where the terms ψ_l are linear combinations of the terms ν_{j-1}^s , ∇q_{j-1}^s , ν_{j-2}^s , w_{j-2}^s , $\nabla \nu_{j-2}^s$, ∇w_{j-2}^s , $\Delta \nu_{j-3}^s$, and Δw_{j-3}^s . In particular, all these terms belong to the space $\mathbf{H}^{k-j+1}(\Omega_-)$.

Hence, by using the previous lemma, we obtain by (5.19) that \tilde{p}_j^- writes:

$$\tilde{p}_j^-(x, \mathbf{z}) = \sum_{l=0}^{j-2} q_j^l(x) \mathbf{z}^l \exp(-\sqrt{\kappa} \mathbf{z}) \quad \text{with} \quad q_j^l \in \mathbf{H}^{k-j+1}(\Omega_-).$$

Second step: construction of $\tilde{v}_j^- \cdot \mathbf{n}$.

From (5.16), with the induction property at order $j - 1$, we remark that $\partial_{\mathbf{z}} \tilde{v}_j^- \cdot \mathbf{n}$ writes

$$\partial_{\mathbf{z}} \tilde{v}_j^- \cdot \mathbf{n} = \left(\sum_{l=0}^{j-1} \phi_l(x) \mathbf{z}^l \right) \exp(-\sqrt{\kappa} \mathbf{z})$$

where the terms ϕ_l are linear combinations of the $\nabla \cdot \nu_{j-1}^s$ and the $\nabla \cdot w_{j-1}^s$, so that the ϕ_i are in $\mathbf{H}^{k-j}(\Omega_-)$. Using the Lemma 5.3, we obtain that $\tilde{v}_j^- \cdot \mathbf{n}$ writes:

$$\tilde{v}_j^- \cdot \mathbf{n}(x, z) = \sum_{l=0}^{j-1} \nu_j^l(x) z^l \exp(-\sqrt{\kappa} z) \quad \text{with} \quad \nu_j^l \in \mathbf{H}^{k-j}(\Omega_-).$$

Third step: construction of the terms v_j^+ , p_j^+ , v_j^- and p_j^- .

The terms v_j^+ , p_j^+ , v_j^- and p_j^- satisfy (5.17), and the data satisfy

$$\Delta v_{j-2}^- \in \mathbf{H}^{k-j}(\Omega_-), \quad l_j \in \mathbf{H}^{k-j-\frac{1}{2}}(\Sigma), \quad h_j \in \mathbf{H}^{k-j-\frac{1}{2}}(\Sigma).$$

We remark that we can add an arbitrary constant to p_{j-2}^+ in order to obtain the compatibility condition:

$$\int_{\Sigma} h_j d\sigma = 0.$$

Hence, we can apply Proposition 5.1 with $k = k - j \geq 1$, $g^- = \Delta v_{j-2}^-$, $g^+ = 0$, $l = l_j$, $h = h_j$. Therefore we obtain the existence and the uniqueness (up to an additive constant for the pressures) of the solution of (5.17) and the solution satisfies:

$$v_j^+ \in \mathbf{H}^{k-j+1}(\Omega_+), \quad v_j^- \in \mathbf{H}^{k-j}(\Omega_-), \quad p_j^+ \in \mathbf{H}^{k-j}(\Omega_+) \quad \text{and} \quad p_j^- \in \mathbf{H}^{k-j+1}(\Omega_-).$$

Fourth step: construction of $(\tilde{v}_j^-)_\tau$.

On the interface Σ , from the previous results, $\alpha(v_j^+ - v_j^-)_\tau - 2\mu(\partial_n v_j^+)_\tau \in \mathbf{H}^{k-j-\frac{1}{2}}(\Sigma)$. We introduce $w_j^0 \in \mathbf{H}^{k-j}(\Omega_+)$ a tangential extension of this boundary data. We choose this extension w_j^0 such that it has a support in a tubular neighborhood of the interface Σ .

We have now the following lemma:

Lemma 5.4. *Let $K \geq 0$. The solution of the ODE*

$$\begin{cases} -\partial_{zz} f + \kappa f = \left(\sum_{i=0}^K \gamma_i z^i \right) \exp(-\sqrt{\kappa} z) & \text{for all } z \geq 0 \\ f(0) = f_0 \\ f(z) \rightarrow 0 & \text{when } z \rightarrow +\infty, \end{cases}$$

is $f(z) = \left(f_0 + \sum_{i=0}^K \beta_i z^{i+1} \right) \exp(-\sqrt{\kappa} z)$. Here,

$$\begin{pmatrix} \gamma_0 \\ \vdots \\ \gamma_K \end{pmatrix} = M_K \begin{pmatrix} \beta_0 \\ \vdots \\ \beta_K \end{pmatrix}$$

where M_K is a $(K+1) \times (K+1)$ nonsingular matrix with entries m_{ij} such that $m_{ii} = 2i\sqrt{\kappa}$, $m_{i,i+1} = -i(i+1)$, and with vanishing other entries.

From (5.18), using the induction hypothesis on \tilde{p}_j^- , \tilde{v}_{j-1}^- and on \tilde{v}_{j-2}^- , $(\tilde{v}_j^-)_\tau$ satisfies the system:

$$\begin{cases} -\partial_{zz}(\tilde{v}_j^-)_\tau(x, z) + \kappa(\tilde{v}_j^-)_\tau(x, z) = \left(\sum_{i=0}^{j-1} \gamma_i(x) z^i \right) \exp(-\sqrt{\kappa}z), & \text{for } x \in \Omega_-, z \geq 0, \\ (\tilde{v}_j^-)_\tau(x, 0) = w_j^0(x) & \text{for } x \in \Omega_-, \end{cases}$$

where $\gamma_i \in \mathbf{H}^{k-j}(\Omega_-)$ and $w_j^0 \in \mathbf{H}^{k-j}(\Omega_-)$. Hence, by applying Lemma 5.4, we obtain that

$$(\tilde{v}_j^-)_\tau(x, z) = \sum_{l=0}^j w_j^l(x) z^l \exp(-\sqrt{\kappa}z) \quad \text{with} \quad w_j^l \in \mathbf{H}^{k-j}(\Omega_-).$$

This concludes the proof of the property at order j . □

6. REGULARITY PROOF FOR THE ASYMPTOTICS

Following the previous section, we want to solve a collection of elementary problems satisfying : Find $v = (v^-, v^+)$, and $p = (p^-, p^+)$ such that

$$(6.1) \quad \begin{cases} \kappa v^- + \nabla p^- = g^- & \text{in } \Omega_- \\ \nabla \cdot v^- = 0 & \text{in } \Omega_- \\ -\nabla \cdot \boldsymbol{\sigma}^+(v^+, p^+) = g^+ & \text{in } \Omega_+ \\ \nabla \cdot v^+ = 0 & \text{in } \Omega_+ \\ \boldsymbol{\sigma}^+(v^+, p^+) \cdot \mathbf{n} = l - p^- \mathbf{n} + \frac{\beta}{2}((v^+ + v^-) \cdot \mathbf{n}) \mathbf{n} & \text{on } \Sigma \\ (v^+ - v^-) \cdot \mathbf{n} = h & \text{on } \Sigma \\ v^+ = 0 & \text{on } \Gamma, \end{cases}$$

where $\boldsymbol{\sigma}^+(v^+, p^+) = 2\mu \mathbf{d}(v^+) - p^+ \mathbf{n}$, associated with the data g, h, l . We remark that because of the divergence free condition, we need the compatibility condition $\int_{\Sigma} h d\sigma = 0$.

Remark 6.1. The first asymptotic $v_0 = (v_0^+, v_0^-)$ satisfies the problem (6.1) with $g^\pm = \mathbf{g}^\pm$, $l = l$ and $h = 0$. The term $v_1 = (v_1^+, v_1^-)$ satisfies the problem (6.1) with $g^\pm = 0$, $h = \nu_1^0$, $l_\tau = \sqrt{\kappa} w_0^0$ and $l \cdot \mathbf{n} = \frac{\beta}{2} \nu_1^0$.

We first address this problem when $h = 0$.

6.1. Elementary problem without jump for the normal components. We consider the following problem:

$$(6.2) \quad \begin{cases} \kappa v^- + \nabla p^- = g^- & \text{in } \Omega_- \\ \nabla \cdot v^- = 0 & \text{in } \Omega_- \\ -\nabla \cdot \boldsymbol{\sigma}^+(v^+, p^+) = g^+ & \text{in } \Omega_+ \\ \nabla \cdot v^+ = 0 & \text{in } \Omega_+ \\ \boldsymbol{\sigma}^+(v^+, p^+) \cdot \mathbf{n} = l - p^- \mathbf{n} + \frac{\beta}{2}((v^+ + v^-) \cdot \mathbf{n})\mathbf{n} & \text{on } \Sigma \\ (v^+ - v^-) \cdot \mathbf{n} = 0 & \text{on } \Sigma \\ v^+ = 0 & \text{on } \Gamma . \end{cases}$$

Then, we introduce a variational problem for $v = (v^-, v^+)$ associated with the problem (6.2) in the space

$$W = \{u \in \mathbf{L}^2(\Omega) \mid \nabla u^+ \in \mathbf{L}^2(\Omega_+), \quad \operatorname{div} u = 0 \quad \text{in } \Omega, \quad u^+ = 0 \quad \text{on } \Gamma\} .$$

Such a variational formulation writes : Find $v = (v^-, v^+) \in W$ such that

$$(6.3) \quad \forall u \in W, \quad a(v, u) = b(u),$$

where

$$a(v, u) := \mu \int_{\Omega_+} \mathbf{d}(v^+) : \mathbf{d}(u^+) \, d\mathbf{x} + \kappa \int_{\Omega_-} v^- \cdot u^- \, d\mathbf{x} + \beta \int_{\Sigma} (v \cdot \mathbf{n})(u \cdot \mathbf{n}) \, ds ,$$

and

$$b(u) = \int_{\Omega} g \cdot u \, d\mathbf{x} - \int_{\Sigma} l \cdot u^+ \, ds .$$

Endowed with the norm

$$\|v\|_W = \|v^+\|_{1, \Omega_+} + \|v^-\|_{0, \Omega_-} + \|v \cdot \mathbf{n}\|_{0, \Sigma} ,$$

the functional space W is a Hilbert space since W is a closed subspace of the Hilbert space $\{u \in \mathbf{L}^2(\Omega) \mid \nabla u^+ \in \mathbf{L}^2(\Omega_+), \quad u^+ = 0 \quad \text{on } \Gamma\}$.

As a consequence of both the Poincaré inequality in $\mathbf{H}_{0, \Gamma}^1(\Omega_+)$ and the Lax-Milgram lemma, we infer the well-posedness of problem (6.3) :

Proposition 6.2. *For given data $g \in \mathbf{L}^2(\Omega)$ and $l \in \mathbf{L}^2(\Sigma)$, the problem (6.3) is well posed in W .*

If we assume in addition that $\operatorname{curl} g^- \in \mathbf{L}^2(\Omega_-)$, then the solution v of the problem (6.3) belongs to the space V (we remind that V is introduced in the beginning of Section 3).

Proposition 6.3. *For given data $g \in \mathbf{L}^2(\Omega)$ with $\operatorname{curl} g^- \in \mathbf{L}^2(\Omega_-)$, and $l \in \mathbf{L}^2(\Sigma)$, the solution v of the problem (6.3) belongs to the space V .*

Proof. Taking the curl in the first equation in (6.2), $\kappa \operatorname{curl} v^- = \operatorname{curl} g^- \in \mathbf{L}^2(\Omega_-)$. Moreover $\nabla \cdot v^- = 0$ and $v^- \cdot \mathbf{n} = v^+ \cdot \mathbf{n} \in \mathbf{H}^{\frac{1}{2}}(\Sigma)$ since v^+ belongs to $\mathbf{H}^1(\Omega_+)$. Since Ω_- is a smooth domain, we infer $v^- \in \mathbf{H}^1(\Omega_-)$. \square

The next proposition ensures a regularity result in Sobolev spaces for the solutions of problem (6.2). It is the main result of this section.

Proposition 6.4. *Let $k \geq 1$. We assume that $g^+ \in \mathbf{H}^{k-1}(\Omega_+)$, $g^- \in \mathbf{H}^k(\Omega_-)$, and $l \in H^{k-\frac{1}{2}}(\Sigma)$. Then the solution of the problem (6.2) satisfies $v^+ \in \mathbf{H}^{k+1}(\Omega_+)$, $v^- \in \mathbf{H}^k(\Omega_-)$, $p^+ \in H^k(\Omega_+)$ and $p^- \in H^{k+1}(\Omega_-)$.*

Proof. We remark that if we know $v^- \cdot n$ on Σ , then p^- and v^- are completely determined. Indeed, taking the divergence of the first equation in (6.2) we obtain that

$$\Delta p^- = \nabla \cdot g^-.$$

In addition, taking the scalar product of the same equation with n , we obtain that

$$\kappa v^- \cdot n + \frac{\partial p^-}{\partial n} = g^- \cdot n.$$

Thus, p^- satisfies:

$$\begin{cases} \Delta p^- = \nabla \cdot g^- & \text{in } \Omega_- \\ \frac{\partial p^-}{\partial n} = g^- \cdot n - \kappa v^- \cdot n & \text{on } \Sigma \end{cases}$$

With an additional condition on the mean of p^- , this characterizes completely p^- .

We fix $j \in \{1, \dots, k\}$. We introduce the following Dirichlet to Neumann operator $T : H^{j-\frac{1}{2}}(\Sigma) \rightarrow H^{j+\frac{1}{2}}(\Sigma)$ in the following way: for $\varphi \in H^{j-\frac{1}{2}}(\Sigma)$ we solve

$$(6.4) \quad \begin{cases} \Delta p = \nabla \cdot g^- & \text{in } \Omega_- \\ \frac{\partial p}{\partial n} = g^- \cdot n - \kappa \varphi & \text{on } \Sigma \\ \int_{\Omega_-} p \, d\mathbf{x} = 0 \end{cases}$$

and we denote by $T(\varphi)$ the trace of the obtained p on Σ .

We remark that since $g^- \in \mathbf{H}^k(\Omega_-)$, if $\varphi \in H^{j-\frac{1}{2}}(\Sigma)$ then by classical elliptic regularity results, $p \in H^{j+1}(\Omega_-)$ so that there exists a constant C independent on φ and g^- such that

$$(6.5) \quad \|p\|_{H^{j+1}(\Omega_-)} + \|T(\varphi)\|_{H^{j+\frac{1}{2}}(\Sigma)} \leq C(\|g^-\|_{\mathbf{H}^k(\Omega_-)} + \|\varphi\|_{H^{j-\frac{1}{2}}(\Sigma)}).$$

In addition, we have $v^- = \frac{1}{\kappa}(g^- - \nabla p^-)$, and we infer

$$\|v^-\|_{\mathbf{H}^j(\Omega_-)} \leq C(\|g^-\|_{\mathbf{H}^k(\Omega_-)} + \|\varphi\|_{H^{j-\frac{1}{2}}(\Sigma)}).$$

Now on the domain Ω_+ , we rewrite the boundary conditions: we remark first that $v^+ \cdot n = v^- \cdot n$ on Σ . In addition, from the equations on Ω_- , $p^- = T(v^- \cdot n)$. Therefore, (v^+, p^+) satisfies the problem:

$$(6.6) \quad \begin{cases} -\nabla \cdot \sigma^+(v^+, p^+) = g^+ & \text{in } \Omega_+ \\ \nabla \cdot v^+ = 0 & \text{in } \Omega_+ \\ \sigma^+(v^+, p^+) \cdot n = l + \beta(v^+ \cdot n)n - T(v^+ \cdot n)n & \text{on } \Sigma \\ v^+ = 0 & \text{on } \Gamma. \end{cases}$$

Hereafter we use the next proposition which ensures a well-posedness result together with elliptic regularity result for the Stokes operator with mixed boundary conditions (namely with a stress boundary condition on Σ and a Dirichlet boundary condition on Γ):

Proposition 6.5. *Let $g \in \mathbf{L}^2(\Omega_+)$ and $\gamma \in \mathbf{H}^{-\frac{1}{2}}(\Sigma)$. Then the problem*

$$(6.7) \quad \begin{cases} -\nabla \cdot \boldsymbol{\sigma}^+(w, p) = g & \text{in } \Omega_+ \\ \nabla \cdot w = 0 & \text{in } \Omega_+ \\ w = 0 & \text{in } \Gamma \\ \boldsymbol{\sigma}^+(w, p) \cdot \mathbf{n} = \gamma & \text{in } \Sigma \end{cases}$$

has a unique solution (w, p) in the space $\mathbf{H}_{0,\Gamma}^1(\Omega_+) \times \mathbf{L}^2(\Omega_+)$.

Let $s \geq 1$. We assume that there exists a solution (w, p) of the problem (6.7) in the space $\mathbf{H}^s(\Omega_+) \times \mathbf{H}^{s-1}(\Omega_+)$. If $g \in \mathbf{H}^{s-1}(\Omega_+)$ and $\gamma \in \mathbf{H}^{s-\frac{1}{2}}(\Sigma)$, then $(w, p) \in \mathbf{H}^{s+1}(\Omega_+) \times \mathbf{H}^s(\Omega_+)$ and there exists a constant $C > 0$ such that

$$\|w\|_{\mathbf{H}^{s+1}(\Omega_+)} + \|p\|_{\mathbf{H}^s(\Omega_+)} \leq C \left(\|g\|_{\mathbf{H}^{s-1}(\Omega_+)} + \|\gamma\|_{\mathbf{H}^{s-\frac{1}{2}}(\Sigma)} \right).$$

This proposition is a consequence of elliptic regularity result for the Stokes operator with normal stress boundary conditions (see [3, Th. III.5.7] on page 192).

Using Prop. 6.5, let us prove now by induction on j that for all $j \in \{1, \dots, k\}$, we have the property $\mathcal{P}(j)$:

$$\mathcal{P}(j): \quad \mathbf{v}^+ \in \mathbf{H}^{j+1}(\Omega_+), \quad \mathbf{p}^+ \in \mathbf{H}^j(\Omega_+).$$

Proof of $\mathcal{P}(1)$: we already know that $\mathbf{v}^+ \in \mathbf{H}^1(\Omega_+)$ using Proposition 6.4. Hence, $\mathbf{v}^+ \cdot \mathbf{n} \in \mathbf{H}^{\frac{1}{2}}(\Sigma)$. Thus, by (6.5), we obtain that the right-hand-side of the third equation in (6.6) is in $\mathbf{H}^{\frac{1}{2}}(\Sigma)$ and thus, by Proposition 6.5, we obtain that

$$\mathbf{v}^+ \in \mathbf{H}^2(\Omega_+) \quad \text{and} \quad \mathbf{p}^+ \in \mathbf{H}^1(\Omega_+).$$

Induction: we assume that $j < k$ and that $\mathcal{P}(j)$ is satisfied, i.e. $\mathbf{v}^+ \in \mathbf{H}^{j+1}(\Omega_+)$. Hence $\mathbf{v}^+ \cdot \mathbf{n} \in \mathbf{H}^{j+\frac{1}{2}}(\Sigma)$. Thus, by (6.5), we obtain that the right-hand-side of the third equation in (6.6) is in $\mathbf{H}^{j+\frac{1}{2}}(\Sigma)$ and thus, by Proposition 6.5, we obtain that

$$\mathbf{v}^+ \in \mathbf{H}^{j+2}(\Omega_+) \quad \text{and} \quad \mathbf{p}^+ \in \mathbf{H}^{j+1}(\Omega_+).$$

This complete the proof of Proposition 6.4. □

6.2. General existence and regularity result for the elementary problems. We address now the elementary problem with non vanishing jump for the normal velocity (6.1). We prove the following result:

Proposition 6.6. *Let $k \geq 1$. We consider $g^- \in \mathbf{H}^k(\Omega_-)$, $g^+ \in \mathbf{H}^{k-1}(\Omega_+)$, $l \in \mathbf{H}^{k-\frac{1}{2}}(\Sigma)$ and $h \in \mathbf{H}^{k-\frac{1}{2}}(\Sigma)$. Assume that h satisfies the compatibility condition*

$$\int_{\Sigma} h \, d\sigma = 0.$$

Then the solution of problem (6.1) satisfies $v^+ \in \mathbf{H}^{k+1}(\Omega_+)$, $v^- \in \mathbf{H}^k(\Omega_-)$, $p^+ \in \mathbf{H}^k(\Omega_+)$ and $p^- \in \mathbf{H}^{k+1}(\Omega_-)$.

Proof. Let $w \in \mathbf{H}^k(\Omega_-)$ such that $\nabla \cdot w = 0$ in Ω_- with $w \cdot n = -h$ on Σ . We denote $w = v^- - w$. Then (v^+, p^+, v^-, p^-) satisfies (6.1) if and only if (v^+, p^+, w, p^-) satisfies:

$$(6.8) \quad \begin{cases} \kappa w + \nabla p^- = g^- - \kappa w & \text{in } \Omega_- \\ \nabla \cdot w = 0 & \text{in } \Omega_- \\ -\nabla \cdot \sigma^+(v^+, p^+) = g^+ & \text{in } \Omega_+ \\ \nabla \cdot v^+ = 0 & \text{in } \Omega_+ \\ \sigma^+(v^+, p^+) \cdot n = -p^- n + \frac{\beta}{2} ((v^+ + w) \cdot n) n + l - \frac{\beta}{2} h n & \text{on } \Sigma \\ (v^+ - w) \cdot n = 0 & \text{on } \Sigma \\ v^+ = 0 & \text{on } \Gamma. \end{cases}$$

We remark that $g^- - \kappa w \in \mathbf{H}^k(\Omega_-)$ and that $l - \frac{\beta}{2} h n \in \mathbf{H}^{k-\frac{1}{2}}(\Sigma)$ therefore, by Proposition 6.4, we obtain that $v^+ \in \mathbf{H}^{k+1}(\Omega_+)$, $p^+ \in \mathbf{H}^k(\Omega_+)$ and $p^- \in \mathbf{H}^{k+1}(\Omega_-)$. We obtain in addition that $w \in \mathbf{H}^k(\Omega_-)$ and since $w \in \mathbf{H}^k(\Omega_-)$, we obtain the same regularity for v^- . \square

7. ESTIMATES OF REMAINDERS

We address the validation of the WKB expansion found before. We claim the following theorem:

Theorem 7.1. *Let $k \geq 5$. We assume that the data in (1.1) satisfy hypothesis (5.1) that we recall here:*

$$g^- \in \mathbf{H}^k(\Omega_-), \quad g^+ \in \mathbf{H}^{k-1}(\Omega_+), \quad l \in \mathbf{H}^{k-\frac{1}{2}}(\Sigma) \quad \text{and} \quad h \in \mathbf{H}^{k-\frac{1}{2}}(\Sigma).$$

We fix $k = k - 2$. Let v_j^+ , v_j^- , \tilde{v}_j^- , p_j^+ , p_j^- and \tilde{p}_j^- , $j \in \{0, \dots, k-1\}$ given by Proposition 5.2. We define $(r_{k,\varepsilon}^\pm, q_{k,\varepsilon}^\pm)$ by removing to the solution $(v_\varepsilon, p_\varepsilon)$ of (1.1) the truncated expansion up to the order $\varepsilon^{\frac{k}{2}}$:

$$(7.1) \quad r_{k,\varepsilon}^+(x) = v_\varepsilon^+(x) - \sum_{j=0}^k \varepsilon^{\frac{j}{2}} v_j^+(x),$$

$$(7.2) \quad r_{k,\varepsilon}^-(x) = v_\varepsilon^-(x) - \sum_{j=0}^k \varepsilon^{\frac{j}{2}} \left(v_j^-(x) + \tilde{v}_j^-(x, \frac{d(x)}{\sqrt{\varepsilon}}) \right),$$

$$(7.3) \quad q_{k,\varepsilon}^+(x) = p_\varepsilon^+(x) - \sum_{j=0}^k \varepsilon^{\frac{j}{2}} p_j^+(x),$$

$$(7.4) \quad q_{k,\varepsilon}^-(x) = p_\varepsilon^-(x) - \sum_{j=0}^k \varepsilon^{\frac{j}{2}} \left(p_j^-(x) + \tilde{p}_j^-(x, \frac{d(x)}{\sqrt{\varepsilon}}) \right).$$

We have the following estimate :

$$\varepsilon \|\nabla r_{k,\varepsilon}^-\|_{0,\Omega_-}^2 + \frac{\kappa}{4} \|r_{k,\varepsilon}^-\|_{0,\Omega_-}^2 + \frac{\mu C^2}{3} \|r_{k,\varepsilon}^+\|_{1,\Omega_+}^2 \leq C \varepsilon^{\frac{k-2}{2}}.$$

Proof. By construction of the profiles, we derive :

$$\left\{ \begin{array}{ll} -\varepsilon \Delta r_{k,\varepsilon}^- + \nabla q_{k,\varepsilon}^- + \kappa r_{k,\varepsilon}^- = \mathbf{g}_{k,\varepsilon} & \text{in } \Omega_- \\ -\nabla \cdot \boldsymbol{\sigma}^+(r_{k,\varepsilon}^+, \mathbf{q}_{k,\varepsilon}^+) = 0 & \text{in } \Omega_+ \\ \nabla \cdot r_{k,\varepsilon}^- = \mathbf{f}_{k,\varepsilon} & \text{in } \Omega_- \\ \nabla \cdot r_{k,\varepsilon}^+ = 0 & \text{in } \Omega_+ \\ 2\mu \mathbf{d}(r_{k,\varepsilon}^+) \cdot \mathbf{n} - \mathbf{q}_{k,\varepsilon}^+ \cdot \mathbf{n} = 2\varepsilon \mathbf{d}(r_{k,\varepsilon}^-) \cdot \mathbf{n} - \mathbf{q}_{k,\varepsilon}^+ \cdot \mathbf{n} + \frac{\beta}{2} \left((r_{k,\varepsilon}^+ + r_{k,\varepsilon}^-) \cdot \mathbf{n} \right) \cdot \mathbf{n} + l_{k,\varepsilon} & \text{on } \Sigma \\ \alpha (r_{k,\varepsilon}^+ - r_{k,\varepsilon}^-)_\tau = 2\mu \left(\mathbf{d}(r_{k,\varepsilon}^+) \cdot \mathbf{n} \right)_\tau & \text{on } \Sigma \\ \frac{1}{\varepsilon} (r_{k,\varepsilon}^+ - r_{k,\varepsilon}^-) \cdot \mathbf{n} = 2\mu \left(\mathbf{d}(r_{k,\varepsilon}^+) \cdot \mathbf{n} \right) \cdot \mathbf{n} - \mathbf{q}_{k,\varepsilon}^+ - \frac{\beta}{4} \left(r_{k,\varepsilon}^+ + r_{k,\varepsilon}^- \right) \cdot \mathbf{n} + h_{k,\varepsilon} & \text{on } \Sigma \\ r_{k,\varepsilon}^+ = 0 & \text{on } \Gamma . \end{array} \right.$$

with

$$\mathbf{g}_{k,\varepsilon} = \varepsilon^{\frac{k+1}{2}} \left(\Delta \tilde{v}_{k-1}^- + 2(\nabla d \cdot \nabla) \partial_z \tilde{v}_k^- + \Delta d \partial_z \tilde{v}_k^- + \Delta \tilde{v}_{k-1}^- \right) + \varepsilon^{\frac{k+2}{2}} \left(\Delta v_k^- + \Delta \tilde{v}_k^- \right),$$

$$\mathbf{f}_{k,\varepsilon} = -\varepsilon^{\frac{k}{2}} \nabla \cdot \tilde{v}_k^-,$$

$$l_{k,\varepsilon} = \varepsilon^{\frac{k+1}{2}} \left(2\mathbf{d}(v_{k-1}^-) \cdot \mathbf{n} + 2\mathbf{d}(\tilde{v}_{k-1}^-) \cdot \mathbf{n} - \partial_z \tilde{v}_k^- - (\partial_z \tilde{v}_k^- \cdot \mathbf{n}) \mathbf{n} \right) + \varepsilon^{\frac{k+2}{2}} \left(\mathbf{d}(v_k^-) \cdot \mathbf{n} + \mathbf{d}(\tilde{v}_k^-) \cdot \mathbf{n} \right),$$

$$h_{\varepsilon,k} = \varepsilon^{\frac{k-1}{2}} \left(\boldsymbol{\sigma}^+(v_{k-1}^+, p_{k-1}^+) \cdot \mathbf{n} - \frac{\beta}{4} (v_{k-1}^+ + v_{k-1}^- + \tilde{v}_{k-1}^-) \right) \cdot \mathbf{n} \\ + \varepsilon^{\frac{k}{2}} \left(\boldsymbol{\sigma}^+(v_k^+, p_k^+) \cdot \mathbf{n} - \frac{\beta}{4} (v_k^+ + v_k^- + \tilde{v}_k^-) \right) \cdot \mathbf{n}.$$

Let us introduce $\psi_{k,\varepsilon} : \Omega_- \rightarrow \mathbb{R}^3$ such that

$$\nabla \cdot \psi_{k,\varepsilon} = \mathbf{f}_{k,\varepsilon} \quad \text{in } \Omega_- .$$

From the expression of \tilde{v}_k^- given by Proposition 5.2, since $z \mapsto z^l \exp(-\sqrt{\kappa} z)$ is bounded on \mathbb{R}^+ , we obtain that

$$\|\mathbf{f}_{k,\varepsilon}\|_{L^2(\Omega_-)} \leq C \varepsilon^{\frac{k}{2}} \max_l \{ \|w_k^l\|_{H^1(\Omega_-)}, \|\nu_k^l\|_{H^1(\Omega_-)} \}.$$

In addition,

$$\partial_i \mathbf{f}_{k,\varepsilon} = -\varepsilon^{\frac{k}{2}} \left(\nabla \cdot \partial_i \tilde{v}_k^- \right) \left(x, \frac{d(x)}{\sqrt{\varepsilon}} \right) - \varepsilon^{\frac{k-1}{2}} \partial_i d(x) \left(\nabla \cdot \partial_z \tilde{v}_k^- \right) \left(x, \frac{d(x)}{\sqrt{\varepsilon}} \right).$$

Hence, there exists a constant $C > 0$ such that

$$\|\nabla \mathbf{f}_{k,\varepsilon}\|_{L^2(\Omega_-)} \leq C \varepsilon^{\frac{k-1}{2}} \max_l \{ \|w_k^l\|_{H^2(\Omega_-)}, \|\nu_k^l\|_{H^2(\Omega_-)} \}.$$

Therefore we can assume that there exists $C > 0$ such that for all k and ε ,

$$(7.5) \quad \|\psi_{k,\varepsilon}\|_{\mathbf{H}^1(\Omega_-)} \leq C \varepsilon^{\frac{k}{2}} \max_l \{ \|w_k^l\|_{H^1(\Omega_-)}, \|\nu_k^l\|_{H^1(\Omega_-)} \}$$

and

$$\|\psi_{k,\varepsilon}\|_{\mathbf{H}^2(\Omega_-)} \leq C \varepsilon^{\frac{k-1}{2}} \max_l \{ \|w_k^l\|_{H^2(\Omega_-)}, \|\nu_k^l\|_{H^2(\Omega_-)} \}.$$

Now we denote $\bar{r}_{k,\varepsilon}^- = r_{k,\varepsilon}^- - \psi_{k,\varepsilon}$ and we obtain that $\bar{r}_{k,\varepsilon}^-$, $r_{k,\varepsilon}^+$ and $q_{k,\varepsilon}^\pm$ satisfy:

$$(7.6) \quad \begin{cases} -\varepsilon \Delta \bar{r}_{k,\varepsilon}^- + \nabla q_{k,\varepsilon}^- + \kappa \bar{r}_{k,\varepsilon}^- = \bar{g}_{k,\varepsilon} & \text{in } \Omega_- \\ -\nabla \cdot \boldsymbol{\sigma}^+(r_{k,\varepsilon}^+, q_{k,\varepsilon}^+) = 0 & \text{in } \Omega_+ \\ \nabla \cdot \bar{r}_{k,\varepsilon}^- = 0 & \text{in } \Omega_- \\ \nabla \cdot r_{k,\varepsilon}^+ = 0 & \text{in } \Omega_+ \\ 2\mu \mathbf{d}(r_{k,\varepsilon}^+) \cdot \mathbf{n} - q_{k,\varepsilon}^+ \mathbf{n} = 2\varepsilon \mathbf{d}(\bar{r}_{k,\varepsilon}^-) \cdot \mathbf{n} - q_{k,\varepsilon}^+ \mathbf{n} + \frac{\beta}{2} \left((r_{k,\varepsilon}^+ + \bar{r}_{k,\varepsilon}^-) \cdot \mathbf{n} \right) \mathbf{n} + \bar{l}_{k,\varepsilon} & \text{on } \Sigma \\ \alpha (r_{k,\varepsilon}^+ - \bar{r}_{k,\varepsilon}^-)_\tau = 2\mu \left(\mathbf{d}(r_{k,\varepsilon}^+) \cdot \mathbf{n} \right)_\tau + \alpha (\psi_{k,\varepsilon})_\tau & \text{on } \Sigma \\ \frac{1}{\varepsilon} (r_{k,\varepsilon}^+ - \bar{r}_{k,\varepsilon}^-) \cdot \mathbf{n} = 2\mu \left(\mathbf{d}(r_{k,\varepsilon}^+) \cdot \mathbf{n} \right) \cdot \mathbf{n} - q_{k,\varepsilon}^+ - \frac{\beta}{4} \left(r_{k,\varepsilon}^+ + \bar{r}_{k,\varepsilon}^- \right) \cdot \mathbf{n} + \bar{h}_{k,\varepsilon} & \text{on } \Sigma \\ r_{k,\varepsilon}^+ = 0 & \text{on } \Gamma, \end{cases}$$

where

$$\begin{aligned} \bar{g}_{k,\varepsilon} &= \mathbf{g}_{k,\varepsilon} + \varepsilon \Delta \psi_{k,\varepsilon} - \kappa \psi_{k,\varepsilon}, \\ \bar{l}_{k,\varepsilon} &= l_{k,\varepsilon} + \varepsilon (2\mathbf{d}(\psi_{k,\varepsilon}) \cdot \mathbf{n}) + \frac{\beta}{2} (\psi_{k,\varepsilon} \cdot \mathbf{n}) \mathbf{n}, \\ \bar{h}_{k,\varepsilon} &= h_{k,\varepsilon} + \frac{1}{\varepsilon} \psi_{k,\varepsilon} \cdot \mathbf{n} - \frac{\beta}{4} \psi_{k,\varepsilon} \cdot \mathbf{n}. \end{aligned}$$

By assumption (5.1), since $k = k - 2$, the terms $w_k^l, w_{k-1}^l, \nu_k^l, \nu_{k-1}^l, v_k^-, v_{k-1}^-$ are bounded in $H^2(\Omega_-)$, and v_k^+, v_{k-1}^+, p_k^+ and p_{k-1}^+ are bounded in $H^2(\Omega_+)$.

Hence, we obtain the following estimates: there exists C such that for all ε ,

$$(7.7) \quad \begin{aligned} \|\bar{g}_{k,\varepsilon}\|_{L^2(\Omega_-)} &\leq C\varepsilon^{\frac{k}{2}} \\ \|\bar{l}_{k,\varepsilon}\|_{L^2(\Sigma)} &\leq C\varepsilon^{\frac{k}{2}} \\ \|\bar{h}_{k,\varepsilon}\|_{L^2(\Sigma)} &\leq C\varepsilon^{\frac{k-2}{2}}. \end{aligned}$$

Then, the remainder terms $(\bar{r}_{k,\varepsilon}^-, r_{k,\varepsilon}^+, q_{k,\varepsilon}^-, q_{k,\varepsilon}^+)$ satisfy Problem (1.1) with the following data:

$$\mathbf{g}^- = \bar{g}_{k,\varepsilon}, \quad \mathbf{g}^+ = 0, \quad l = \bar{l}_{k,\varepsilon}, \quad \mathbf{h} = -\alpha (\psi_{k,\varepsilon})_\tau - \frac{1}{2} \bar{l}_{k,\varepsilon} - \bar{h}_{k,\varepsilon} \mathbf{n}$$

and by (7.7), using Theorem 2.1, we obtain that

$$\varepsilon \|\mathbf{d}(\bar{r}_{k,\varepsilon}^-)\|_{0,\Omega_-}^2 + \frac{\kappa}{4} \|\bar{r}_{k,\varepsilon}^-\|_{0,\Omega_-}^2 + \frac{\mu C^2}{3} \|r_{k,\varepsilon}^+\|_{1,\Omega_+}^2 \leq C\varepsilon^{\frac{k-2}{2}}.$$

Since $r_{k,\varepsilon}^- = \bar{r}_{k,\varepsilon}^- + \psi_{k,\varepsilon}$, this concludes the proof of Theorem 7.1 by using estimate (7.5). \square

Applying Theorem 7.1 for $k = 5$, by straightforward estimates of the order 1, order 2 and order 3 terms of the asymptotic expansion, we conclude the proof of Theorem 2.3.

REFERENCES

- [1] P. Angot. A fictitious domain model for the Stokes/Brinkman problem with jump embedded boundary conditions. *C. R. Math. Acad. Sci. Paris*, 348(11-12):697–702, 2010.
- [2] P. Angot. On the well-posed coupling between free fluid and porous viscous flows. *Appl. Math. Lett.*, 24(6):803–810, 2011.
- [3] F. Boyer and P. Fabrie. *Éléments d'analyse pour l'étude de quelques modèles d'écoulements de fluides visqueux incompressibles*, volume 52. Springer, 2005.
- [4] C.-H. Bruneau and I. Mortazavi. Contrôle passif d'écoulements incompressibles autour d'obstacles à l'aide de milieux poreux. *Comptes Rendus de l'Académie des Sciences-Series IIB-Mechanics*, 329(7):517–521, 2001.
- [5] G. Carbou. Penalization method for viscous incompressible flow around a porous thin layer. *Nonlinear Anal. Real World Appl.*, 5(5):815–855, 2004.
- [6] G. Carbou. Brinkmann model and double penalization method for the flow around a porous thin layer. *J. Math. Fluid Mech.*, 10(1):126–158, 2008.
- [7] G. Carbou and P. Fabrie. Boundary layer for a penalization method for viscous incompressible flow. *Adv. Differential Equations*, 8(12):1453–1480, 2003.
- [8] N. Chen, M. Gunzburger, and X. Wang. Asymptotic analysis of the differences between the Stokes-Darcy system with different interface conditions and the Stokes-Brinkman system. *J. Math. Anal. Appl.*, 368(2):658–676, 2010.