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# Deembedding of filters in multiplexers via rational approximation and interpolation

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**ABSTRACT:** In this paper we present a method to recover electrical parameters of filters embedded in a multiplexer for which scattering measurements are given. Unlike other approaches proposed for this problem, this method does not require a priori knowledge of the scattering parameters of the junction. This feature renders the procedure well suited for tuning purposes or for fault diagnosis. Technically the algorithm starts with a rational approximation step, in order to derive a rational representation of certain scattering parameters of the multiplexer. This representation is then used in a second step to identify the electrical model of each filter. This second step relies on a rational interpolation technique used to extract the filter's responses.

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**KEYWORDS:** Diplexer, Filter tuning, Fault diagnosis, Deembedding, Rational approximation, Padé interpolation

## I. Introduction

MICROWAVE multiplexers are present in nearly every transmission or reception unit of communication systems. The complex reciprocal loading effects of the filters connected via the junction (see Fig. 1) make their synthesis a difficult task: the latter commonly relies on computer-driven simulation, in the circuitual and full wave domain, that are coupled to optimization methods [1]. The practical realization of such devices remains however a delicate matter, because of the inevitable dimension mismatch between the synthesized multiplexer and the realized one. The filters are therefore equipped with tuning elements (e.g. screws, irises) that need to be adjusted in the final manufacturing phase. When filters can be accessed and measured at their two ports, methods based on rational approximation have been developed ([2], [3], [4]) to extract the electrical parameters of the measured filter, and have led to efficient tuning procedures. By contrast, when tuning a multiplexer, detaching filters from the common junction is however hardly a possible or suitable option. Note also that in the early stages of a fault diagnosis procedure on a damaged device, the ability to identify precisely the malfunctioning filter(s), without disassembling the entire

multiplexer, might also be of great value. Concerning instead the design phase, moreover, a method to obtain the electrical parameters of the filters from a whole simulated multiplexer or manufactured prototype can be of great value to recognize mismatches with respect to the ideal model and thus improve the design.

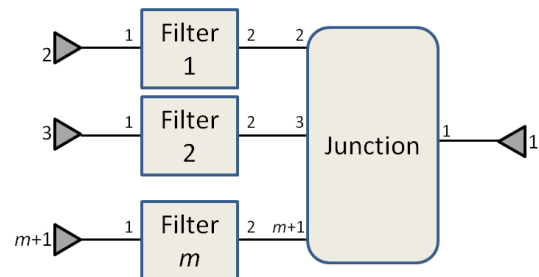


Fig. 1: Model of the multiplexer and ports/filters numbering

For all these reasons, important efforts have been spent mostly on tuning techniques for multiplexers which rely only on external scattering measurements. Most of them are based on neural networks or on the minimization of a tuning criterion ([5], [6]), which however might suffer from the presence of local minima and usually give no real insight about the internal state of the hardware. Other approaches ([7]) have therefore been proposed in order to derive the filter's electrical parameters: these however necessitate the additional knowledge of the junction's scattering parameters, that might not be at hand in practice.

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We describe here a method requiring the sole knowledge of the multiplexer's external scattering measurements, the filter's order and the coupling geometries (in-line, box section, triplets etc...) they implement. The procedure outputs are most of the electrical parameters of the filters up to the resonating frequency offsets of the resonators closest to the junction and their associated output couplings. It is also analytically proven that this is the best that can be done when starting from external measurements. From the latter our algorithm derives polynomial models of certain of the device's scattering parameters, from which the rational scattering matrices of the filters can be extracted. Eventually the filter's scattering matrices extracted in this way are synthesized in terms of electrical circuits of coupled resonators with the specified topologies, yielding the coupling matrix of each filter.

The core concepts of our method are described in the conference paper [8] whereas the present paper is enriched with a technical appendix which is essential for who wants to implement in practice the proposed procedure or understand in some details its functioning. Practical examples are expanded and one has also been added dealing with filters with finite transmission zeros embedded in a triplexer in order to demonstrate the versatility of the proposed approach.

## II. Filter's Response Extraction

We consider the situation depicted in Figure 2:

- a reciprocal filter is loaded on its second port by an unknown load
- we have access to the reflection parameter  $S_{1,1}$  of this system (filter+load) at port one of the filter
- we know the filter's order  $N$ , its coupling topology and the location of its transmission zeros.

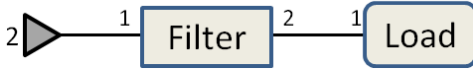


Fig. 2: Filter with an arbitrary load on port 2

This section will try to answer the following natural question: under these circumstances, what can we say about the filter's response itself? We call  $F_{1,1}$ ;  $F_{1,2}$ ;  $F_{2,2}$  the filter's scattering parameters, and  $L_{1,1}$  the reflection parameter of the load. With these notations, the reflection parameter of the overall system expresses as:

$$S_{1,1} = F_{1,1} + \frac{(F_{1,2})^2 L_{1,1}}{1 - L_{1,1} F_{2,2}} \quad (1)$$

Evaluating the preceding expression at a transmission zero  $\omega_c$  of the filter is of interest in our situation, as it readily yields:

$$\begin{cases} F_{1,1}(\omega_c) = S_{1,1}(\omega_c) \\ \dot{F}_{1,1}(\omega_c) = \dot{S}_{1,1}(\omega_c) \end{cases} \quad (2)$$

where the dot notation is used to indicate first derivatives with respect to  $\omega$ , and the second equality can be easily proven by observing that  $F_{1,2}$  appears squared in (1) and vanishes at  $\omega_c$  thus nulling any further contribution to  $\dot{S}_{1,1}(\omega_c)$  other than from  $\dot{F}_{1,1}(\omega_c)$ . Eq. (2) further indicates that the values and the derivatives of the reflection

coefficient of the filter at one of its finite or infinite transmission zeros can be both inferred from those of the system's reflection at the same point. Note that if  $\omega_c$  is a zero of multiplicity  $m$  then the list of equalities (2) can be continued up to a derivative order  $2m-1$ . This remark will be useful for transmission zeros at infinity: for example a filter of order  $N$  with an in-line coupling topology has a zero of order  $N$  at infinity. The rationality of  $F_{1,1} = p/q$  leads us to consider the following interpolation problem: supposing  $(\omega_1, \dots, \omega_l)$  are all the filter's transmission zeros of respective order  $(m_1, \dots, m_l)$  ( $\sum m_k = N$ ), find all the polynomial pairs  $(p, q)$  satisfying:

$$\forall k \in \{1 \dots l\}, \forall i \in \{0 \dots 2m_k - 1\} \left( \frac{p}{q} \right)^{(i)}(\omega_k) = S_{1,1}^{(i)}(\omega_k) \quad (3)$$

The system (3) defines a rational interpolation problem, of Padé multipoint type [9], and can be solved classically using elementary methods from linear algebra. The system of equations (3) contains  $2N$  equations compared to the  $2N+2$  unknowns corresponding to the coefficients of polynomials  $p$  and  $q$ . It is therefore no surprise that the following holds:

*Proposition 2.1:* All the polynomial pairs  $(p, q)$  of degree at most  $N$  that solve (3) form a two-dimensional space<sup>2</sup>. In particular there exist  $(p_1, q_1)$  and  $(p_2, q_2)$ , such that, for all complex numbers  $(\alpha, \beta)$ , the pair  $(p, q)$  defined as:

$$\begin{aligned} p &= \alpha p_1 + \beta p_2 \\ q &= \alpha q_1 + \beta q_2 \end{aligned} \quad (4)$$

is a solution to (3). Moreover, any two pairs of distinct solutions satisfy

$$p_1 q_2 - q_1 p_2 = \gamma^2 r^2 \quad (5)$$

with  $\gamma$  a complex number (depending on the special choice of the solution pair)  $r$  the transmission polynomial

$$r = \prod_{\omega_k \neq \infty} (s - \omega_k)^{m_k}.$$

The lack of uniqueness of the identified filter's reflection coefficient is coherent with the following remark: one can always intercalate between the filter and the load (see Fig. 2) an arbitrary constant (independent on the frequency) chain matrix followed by its inverse, and this without changing the reflection parameter  $S_{1,1}$  of the system. The filter's response is therefore at most recoverable up to the chaining of a constant chain matrix at its port two. The following proposition shows that this is the sole uncertainty on the filter's response:

*Proposition 2.2:* In the notations of proposition (2.1), and for any two pairs  $a=(p_1, q_1)$  and  $b=(p_2, q_2)$  solutions of (3), the scattering matrix:

$$F_{a,b} = \frac{1}{q_1} \begin{pmatrix} p_1 & \mathcal{R} \\ q_1 & q_2 \end{pmatrix} \quad (6)$$

is a possible solution to our de-embedding problem, in the sense that:

- $F_{a,b}$  has  $\omega_k$  as transmission zeros
- $F_{a,b}$  is of McMillan degree at most  $N$  (its determinant is a rational function of degree at most  $N$ )
- It satisfies equation (2) at every transmission zero  $\omega_k$ .

<sup>2</sup>Strictly speaking, the assertion is true only generically (for almost all), with respect to interpolation data  $S_{1,1}(\omega_k)$

Moreover, every solution in this sense to the de-embedding problem can be obtained like that. Eventually, for any two choices of solutions  $(a, b)$  and  $(c, d)$ ,  $F_{a,b}$  and  $F_{c,d}$  differ only by a constant chain matrix connected on their second port.

In order to access the electrical parameters of the filter we need now to consider circuitual realizations of the derived rational scattering matrices. The considered circuitual realization consists in the classical low-pass prototype [10], made of coupled resonators with a specified coupling topology. The following proposition indicates that most of the electrical parameters of the filter are recovered.

**Proposition 2.3:** Suppose that the filter has at most  $N-2$  transmission zeros at finite frequencies, in order to admit a circuitual realization with no source-load coupling [10]. Suppose that  $F$  is a possible solution of the de-embedding problem (in the sense of proposition 2.2) and admits a circuitual representation characterized by a  $(N+2) \times (N+2)$  coupling matrix  $M$ . Then any other solution  $F'$ , also solution of the de-embedding problem, admits a circuit representation with coupling matrix  $M'$  differing from  $M$  only for:

- the frequency offset of the last cavity of the filter (nearest to the load), that is the value  $M_{N,N}$  of the coupling matrix;
- the value of the output coupling represented by  $M_{N,L}$  in the coupling matrix (see [10]).

The preceding proposition, for which a proof is sketched in Sect. VI-D, indicates that the uncertainty on the scattering matrix of the filter has a very localized impact on its circuitual representation. This is coherent with the fact that this uncertainty weighs on the value of a constant chain matrix plugged at its output port.

### III. De-embedding Filters from a Multiplexer

The problem of de-embedding filters from a multiplexer, when starting from external measurements, is in many aspects similar to the simplified situation considered in Sect. II. If we call  $G$  the scattering matrix of the multiplexer:

- Each filter appears to be loaded at its second port by a component regrouping the remaining filters connected via the junctions
- For all these sub-systems, composed of a filter loaded by an unknown component, the reflection at port 1 is known. For the first filter, it is for example given by the scattering parameter  $G_{2,2}$ , for the second by  $G_{3,3}$  etc...
- The transmission zeros, for example of filter 1, are also present in the transmission parameters  $G_{2,1}$ ,  $G_{2,3}$ ,  $G_{2,4}$ ... As we will see, this property can be used to locate the transmission zeros.

The core idea of the de-embedding algorithm is to perform a line-wise rational approximation of the scattering measurement of  $G$ : we will detail it from the perspective of the de-embedding problem of filter 1.

We start with a rational approximation of the row  $[G_{2,1}, G_{2,2}, G_{2,3}, \dots]$ . The order of the rational approximation is chosen so as to get a proper fitting of the measurements and is in general higher than the filter's order: an additional increase in degree is introduced to approximate the effects of the loading. If the filter is expected to have finite transmission zeros, they can be identified using the following remark: the finite transmission zeros of filter 1 are the common zeros of the transmission parameters  $G_{2,1}, G_{2,3}, \dots$  (apart from unlikely situations where all the

filters have the same transmission zeros, or the junction itself has a transmission zero). Their locations can therefore be obtained by analyzing the common zeros of the numerators of the rational approximations of  $G_{2,1}, G_{2,3}, \dots$ . Eventually, the values and derivatives of the right-hand term of (3) will be provided by evaluating the rational approximation of  $G_{2,2}$ . Extraction of coupling parameters is then obtained following the procedure outlined in section II.

## IV. Practical Examples

### A. Simulated Triplexer

The first example is rather academical, as measurements are obtained from simulations. The considered triplexer is built from a 4-port manifold junction on which we connected 3 filters of degree 4, each implementing 2 transmission zeros thanks to a "quartet" coupling geometry. The triplexer is evaluated on a frequency band from 11.4 GHz to 11.6 GHz by means of full-wave simulations. All the scattering parameters of the filters too are therefore available for this test and they can be processed by a classical de-embedding technique [3] to provide a reference coupling matrix that we will confront with results obtained only from external simulated scattering parameters.

A rational approximation of the scattering parameters  $[G_{2,1}, G_{2,2}, G_{2,3}, G_{2,4}]$  is performed to de-embed filter 1: a degree 5 is needed in order to obtain a relative error less than 1%. Note here that, for all considered scattering parameters, the rest of the system (junction+remaining filters) is filtered out, away from the filter's passband: this explains the very moderate order needed for the rational approximation here. Then two finite transmission zeros are identified from the common zeros of rational representations of  $G_{2,1}, G_{2,3}, G_{2,4}$ . These and two additional zeros at infinity are retained as the filter's transmission zeros. The procedure in Sect. II can then be applied using the rational approximation of  $G_{2,2}$  in order to obtain possible characteristic polynomials of the filter and hence its scattering matrix. The latter is then realized as a circuit with the following coupling matrix in arrow form [10]:

0	+0.979	0	0	0	0
+0.979	+0.086	+0.716	-0.001	-0.167	0
0	+0.716	+0.100	+0.620	0	0
0	-0.001	+0.620	-0.027	+0.619	0
0	-0.167	0	+0.619	+0.240	+0.835
0	0	0	0	+0.835	0

that is to be compared to the one extracted directly from the scattering parameters of the filter:

0	+0.974	0	0	0	0
+0.974	+0.086	+0.718	-0.001	-0.174	0
0	+0.718	+0.094	+0.622	0	0
0	-0.001	+0.622	-0.002	+0.624	0
0	-0.174	0	+0.624	+0.070	+1.071
0	0	0	0	+1.071	0

As expected, our procedure recovers all filter's electrical parameters up to the frequency offset of the last cavity and the output coupling. Similar results were obtained for the two other filters.

### B. Measured Diplexer

We now consider a real-world example, made of a manifold diplexer manufactured in one piece shown in Figure 3. The ports are WR62 rectangular waveguides. The lower-band filter was initially tuned to provide  $|S_{1,1}| < -20$  dB from 14.478 GHz to 14.640 GHz, whereas the upper-band filter for  $|S_{1,1}| < -17$  dB from 14.896 GHz to 15.058 GHz (see Fig. 4). The minimum loss on  $S_{2,1}$  and  $S_{3,1}$  is about 0.8 dB.

Both filters are of order 5 and implement an in-line coupling topology: all transmission zeros are therefore at infinity. To test our approach we ran several measurements of all the device's scattering parameters, while screwing in and out some tuning screws. Tuning screws entering laterally into the coupling windows between two cavities act like small capacitors: when screwed in, the corresponding coupling tends to augment. One can also expect the resonating frequencies of the adjacent cavities to be slightly affected by the coupling screws. In order to get a reference, we therefore ran a first identification of filter 1 and found the following coupling matrix:

0	+1.015	0	0	0	0	0
+1.015	-0.006	+0.839	0	0	0	0
0	+0.839	+0.012	+0.631	0	0	0
0	0	+0.631	+0.032	+0.617	0	0
0	0	0	+0.617	+0.145	+0.860	0
0	0	0	0	+0.860	-0.136	+1.091
0	0	0	0	0	+1.091	0

We then screwed in the coupling screw between the 3<sup>rd</sup> and the 4<sup>th</sup> cavity, measured the diplexer (see Fig. 5) and identified the coupling matrix of the first filter as:

0	+1.014	0	0	0	0	0
+1.014	-0.006	+0.839	0	0	0	0
0	+0.839	+0.012	+0.628	0	0	0
0	0	+0.628	<span style="border: 1px solid black;">-0.023</span>	<span style="border: 1px solid black;">+0.692</span>	0	0
0	0	0	<span style="border: 1px solid black;">+0.692</span>	<span style="border: 1px solid black;">+0.080</span>	+0.864	0
0	0	0	0	+0.864	-0.151	+0.960
0	0	0	0	0	+0.960	0

The predicted increase of the coupling  $M_{3,4}$  is confirmed by the identification as well as a variation of the resonating frequency of the 3<sup>rd</sup> and 4<sup>th</sup> cavities. As explained in section II, the identified elements  $M_{5,5}$  and  $M_{N,L}$  should not be given attention to. We ran other similar experiments, also altering screws located at cavities centre, and always found the identified matrices to be coherent with the performed action.

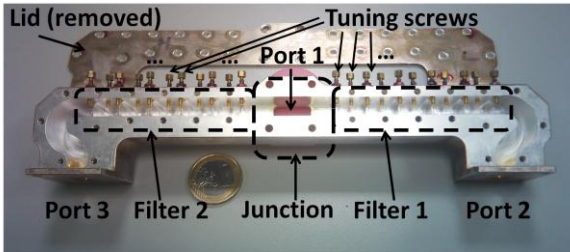


Fig. 3: Manufactured diplexer. Tuning screws are used to adjust coupling and resonating frequency of the cavities

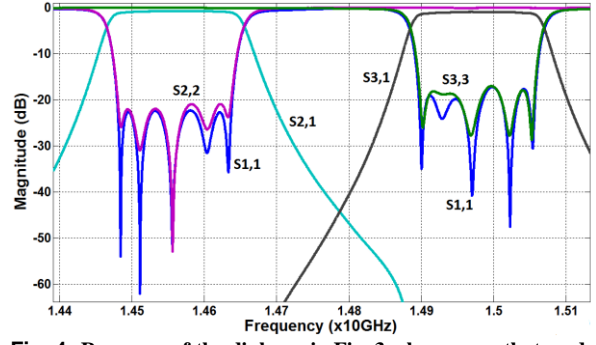


Fig. 4: Response of the diplexer in Fig. 3 when correctly tuned, used to extract the reference model

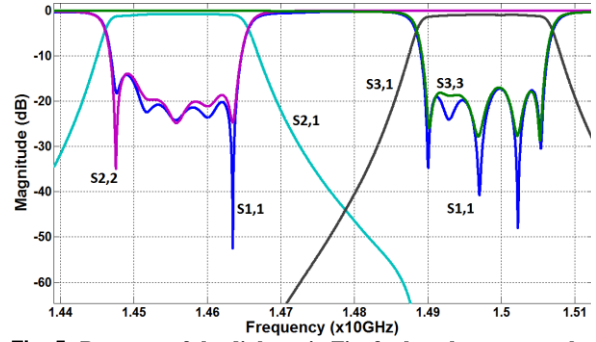


Fig. 5: Response of the diplexer in Fig. 3 when the screw on the coupling between 3<sup>rd</sup> and 4<sup>th</sup> resonator of the lower filter was inserted by half a turn from the original (tuned) position. The lower edge of the lower band is clearly modified from Fig. 4

## V. Conclusions

In this paper we presented a method to extract the electrical parameters of a multiplexer's filters through the sole knowledge of external scattering measurements. The method can be seen as an analogue, in the context of system identification, of Darlington's extraction procedure [11] used in the context of filter synthesis. The method was validated on practical examples and is suited for applications in the context of computer-assisted tuning and fault diagnosis or in the design and early prototyping. Although it has not been tested yet, application of the methodology to other complex devices combining filters accessible by only one of their ports seems promising.

## VI. Appendix

### A. Padé interpolation

The purpose here is to detail how the interpolation system (3) translates into a linear system of equations, and can therefore be tackled efficiently in practice. We consider one of the  $l$  interpolation series of the equation (3), w.l.o.g. the one occurring at  $\omega_1$ :

$$\forall i \in \{0 \dots 2m_1 - 1\} \left( \frac{p}{q} \right)^{(i)}(\omega_1) = S_{1,1}^{(i)}(\omega_1) \quad (7)$$

The latter specifies the value of the rational function  $p/q$  and of  $2m_1 - 1$  of its derivatives at the point  $\omega_1$ . Using the Taylor expansion of  $S_{1,1}$  at  $\omega_1$ , this can be rephrased as:

$$\frac{p}{q}(\omega) = \sum_{i=0}^{2m_1-1} \frac{S_{1,1}^{(i)}(\omega_1)}{i!} (\omega - \omega_1)^i + O((\omega - \omega_1)^{2m_1-1}) \quad (8)$$

which, after multiplication by the denominator  $q$  implies,

$$p(\omega) - \sum_{i=0}^{2m_1-1} \frac{S_{1,1}^{(i)}(\omega_1)}{i!} (\omega - \omega_1)^i q(\omega) = O((\omega - \omega_1)^{2m_1-1}). \quad (9)$$

By eq. (7), the left-hand term of (9) therefore vanishes at  $\omega_1$  as well as its first  $2m_1-1$  derivatives. Differentiating the latter  $j$  times ( $j < 2m_1$ ) and evaluating it at  $\omega = \omega_1$  leads to,  $\forall j \in \{0 \dots 2m_1 - 1\}$

$$(p(\omega_1))^{(j)} - \sum_{i=0}^j \binom{j}{i} S_{1,1}^{(i)}(\omega_1) (q(\omega_1))^{(j-i)} = 0 \quad (10)$$

where  $\binom{j}{i} = \frac{j!}{i!(j-i)!}$  is the binomial coefficient. If we set

$p = \sum_{k=0}^N p_k \omega^k$  and  $q = \sum_{k=0}^N q_k \omega^k$ , equations (10) express as  $2m_1$  homogeneous linear equations in the unknowns  $(p_k), (q_k)$ . The interpolation problem (3) amounts therefore to find the kernel of an homogeneous system of  $2N$  linear equations in the  $2N+2$  unknowns  $(p_k, q_k)$ .

In practice this can be tackled, for example, by a numerical procedure computing the singular value decomposition of a matrix (ex. `svd()` from Matlab). One can in fact define a column vector composed by the coefficients of  $p$  and  $q$  (both in decreasing order of power) as unknown quantity. Then, each eq. (10) can be written as a row of numerical coefficients multiplying the unknown vector to give 0. Thus the first transmission zero  $\omega_1$  will result in  $2m_1$  rows and, similarly,  $2m_k$  rows will correspond to all the other zeros  $\omega_k$ . The resulting matrix  $(M, 2N \text{ rows by } 2(N+1) \text{ columns})$  can be expressed by its SVD decomposition as  $USV^T = M$ . The matrix  $S$  will include two null singular values in its diagonal, whose associated right vectors (corresponding columns of  $V$ ) represent one set of coefficients of  $p$  and  $q$  each. Any linear combination of these solutions is also a valid solution of the interpolation problem.

We have presented here the elementary notions for Padé interpolation, showing how rational interpolation amounts to solve a linear problem and how it develops. For a survey about this well studied topic, see for example [12].

#### B. Sketched proof of Proposition 2.1

The first part of proposition 2.1 is a consequence of the preceding section, that transforms rational interpolation into solving a linear homogeneous system. The dimension of the solution space (here 2) is derived from the presence of  $2N+2$  unknowns when compared to  $2N$  linear equations (which can be shown to be generically of full rank  $2N$ ). Now if  $p_1/q_1$  and  $p_2/q_2$  are two solutions to problem (3) we have:

$$\forall k \in \{1 \dots L\} \quad \frac{p_1}{q_1} - \frac{p_2}{q_2} = O((\omega - \omega_k)^{2m_k-1}) \quad (11)$$

which leads to

$$\forall k \in \{1 \dots L\} \quad p_1 q_2 - p_2 q_1 = O((\omega - \omega_k)^{2m_k-1}) \quad (12)$$

This indicates that the polynomial  $p_1 q_2 - p_2 q_1$  has zeros of orders  $2m_k$  in each of the  $\omega_k$ 's. This polynomial is however of order at most  $2N$  (and  $\sum m_k = N$ ) and it is therefore proportional to the polynomial  $r^2$  of equation (5).

#### C. Sketched proof of Proposition 2.2

The assertion about the MacMillan degree of rational matrix  $F_{a,b}$  is a consequence of the divisibility of the determinant of its numerator matrix by the common denominator  $q$  [13]. This holds by construction of the matrix  $F_{a,b}$  as defined by eq. (5). Eventually, if  $T_{a,b}$  and  $T_{a',b'}$  are the chain matrices associated to  $F_{a,b}$  and  $F_{a',b'}$ , straightforward algebraic computations show that  $T_{a,b} \cdot T_{a',b'}$  is a constant chain matrix, proving the last assertion of the proposition.

#### D. Sketched proof of Proposition 2.3

The retained filter model is here the classical low-pass model made of coupled resonators shown in Figure 6. The electrical equations governing this model are given in state-space form as follows:

$$\begin{cases} \dot{x} = Ax + B \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} \\ \begin{pmatrix} b_1 \\ b_2 \end{pmatrix} = Cx + D \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} \end{cases} \quad (13)$$

where  $a_1, a_2, b_1, b_2$  are the I/O power waves and  $x$  is a vector containing resonators voltages. The matrices  $A, B, C, D$  are defined in terms of the circuitual parameters as follows:

$$\begin{aligned} C &= \begin{pmatrix} jM_{s,1} & 0 & 0 & 0 \\ 0 & 0 & 0 & jM_{N,L} \end{pmatrix} & B &= C' \\ A &= jM + C \cdot B & D &= I \end{aligned} \quad (14)$$

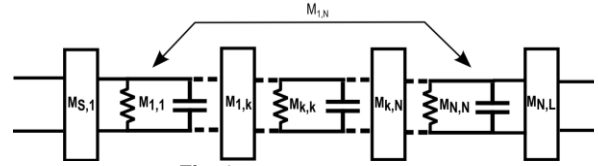


Fig. 6: Low-pass prototype



Fig. 7: Filter with a constant chain matrix on port 2

where  $M$  is the coupling matrix of the low-pass network and  $I$  is the  $2 \times 2$  identity matrix. The scattering matrix of the network is given as:

$$S = C(s \cdot I - A)^{-1} B + D \quad (15)$$

From what precedes, we know that we are able to recover the filter's scattering matrix  $S$  up to a constant chain matrix  $T$  connected at its second port (see Figure 7). We therefore derive a circuitual representation of the system composed by the filter followed by the chain matrix  $T$ . To do this, we seek a state-space representation of the dynamical system having  $(a_1, a_3)$  as inputs and  $(b_1, b_3)$  as outputs. We first express  $(a_2, b_2)$  in terms of  $(a_3, b_3)$  by means of the chain matrix  $T$ :

$$\begin{aligned} a_2 &= T_{1,1} a_3 + T_{1,2} b_3 \\ b_2 &= T_{2,1} a_3 + T_{2,2} b_3 \end{aligned} \quad (16)$$

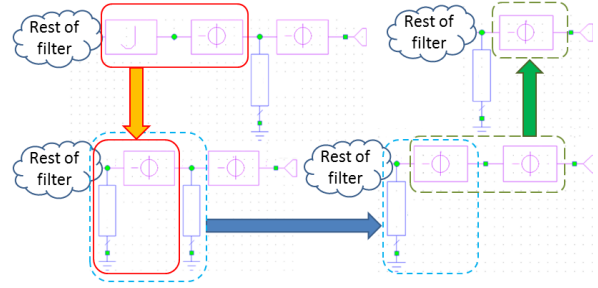


and replace the latter in the state space form (13). Simple algebraic manipulation yields a state-space realization ( $A'$ ,  $B'$ ,  $C'$ ,  $D'$ ) of our overall system that expresses as:

$$\left\{ \begin{array}{l} A' = A + B(:,2) \cdot C(2,:) \frac{T_{1,2}}{T_{2,2} - T_{1,2}} \\ C' = \begin{pmatrix} 1 & 0 \\ 0 & \frac{1}{T_{2,2} - T_{1,2}} \end{pmatrix} \cdot C \\ B' = C'^T \\ D' = \begin{pmatrix} 1 & 0 \\ 0 & \frac{T_{1,1} - T_{2,1}}{T_{2,2} - T_{1,2}} \end{pmatrix} \cdot D \end{array} \right. \quad (17)$$

The sparse form of matrices  $B$  and  $C$  implies that the perturbation  $B(:,2) \cdot C(2,:)$  only impacts the term  $A_{N,N}$  of the dynamic matrix: it represents the frequency offset of the last resonator. For  $B'$  and  $C'$  the perturbed terms are also those related to the output coupling. The modification in  $D$  only affects the value at infinity of the reflection at port 2, and can be seen as a change of reference plane at this access. We have therefore derived a circuital realization of the perturbed filter (by means of  $T$ ) which coincides with the original one up to the coupling  $M_{N,L}$  and the offset  $M_{N,N}$ . All other electrical parameters of the filter can therefore be read out of a circuital realization of any matrix of Prop. 2.2.

An alternative circuital interpretation can be used to reach the same conclusions. In fact, any constant 2-port network can be represented as a  $\Pi$  network of admittances; the central series admittance is in turn equivalent to a shunt one between two equal and opposite admittance inverters. These two and their respective outer admittances can be converted into phase shifters just by adding residual shunt admittances in the center. The resulting model of a general frequency-independent network is therefore made by two phase shifters with an equivalent single shunt admittance in between. When such a general component is connected as in Fig. 7 to port 2 of a filter modeled as in Fig. 6, simple circuital transformations (Fig. 8) involving scaling and equivalencies allow to collapse the  $M_{N,L}$  coupling with the closest phase shifter; the effect of the shunt admittance can then be included into the obtained phase shifter to give an equivalent phase shifter cascaded with the outer one. The resulting phase shifter can be finally decomposed as an external phase shifter and an inverter with a shunt admittance in parallel with the last resonator. The external phase shifter represents the reference plane change, whereas the modified inverter accounts for the varied external coupling coefficient and the shunt admittance for the last resonator's shifted frequency. The rest of the filter, after this complete transformation, remains unaffected.



**Fig. 8: Circuital transformation to embed a constant chain matrix into a filter's output coupling;  $M_{N,L}$  is implemented by the  $J$ -inverter whereas phase shifters by  $-\phi$  blocks.**

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