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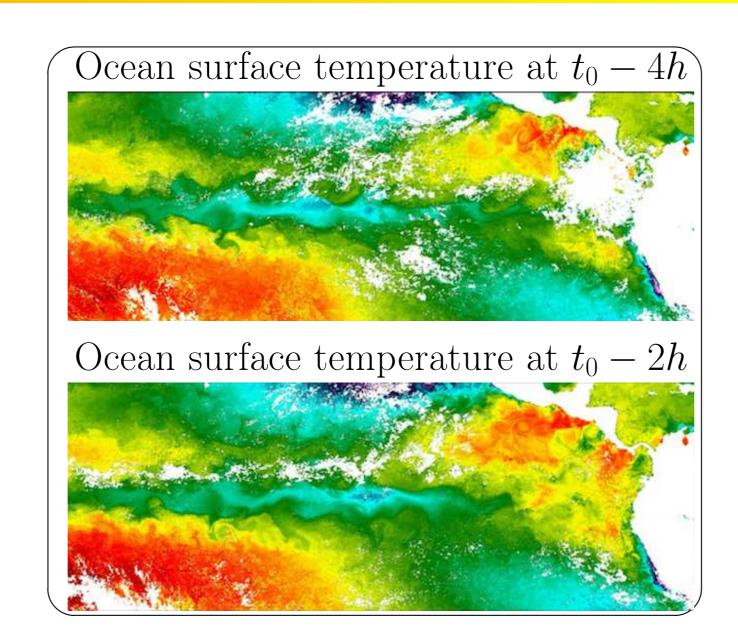
# Application of optimal transportation to data assimilation Nelson FEYEUX, Maëlle NODET, Arthur VIDARD



#### EXAMPLE OF OCEANOGRAPHIC DATA ASSIMILATION

Given a state  $\mathbf{x}$  whose evolution is simulated through a model  $\mathcal{M}$ , **data assimilation** is an inverse problem aiming at reconstructing an initial condition  $\mathbf{x}_0$  of  $\mathbf{x}$  given only several partial observations  $(\mathbf{y}_i^o)_i$  of  $\mathbf{x}$ . For example, the state  $\mathbf{x}(t, x, y, z)$  of the ocean gathers:

- Temperature T(t, x, y, z)
- Velocities  $\mathbf{u}(t, x, y, z)$
- Salinity S(t, x, y, z)
- Water height h(t, x, y)



The question is then: What is the complete state of the ocean at t = 0,  $\mathbf{x}_0(x, y, z)$ ?

In variational data assimilation this is done by minimizing a cost function  $\mathcal{J}$  defined as

$$\mathcal{J}(\mathbf{x}_0) = \underbrace{\frac{1}{2} \sum_{i} d\left(\mathcal{H}_i(\mathbf{x}_0), \mathbf{y}_i^o\right)^2}_{\approx d\left(\mathbf{x}_0, \mathbf{x}_0^t\right)^2} + \underbrace{\frac{\gamma}{2} d\left(\mathbf{x}_0, \mathbf{x}_0^b\right)^2}_{(1)},$$

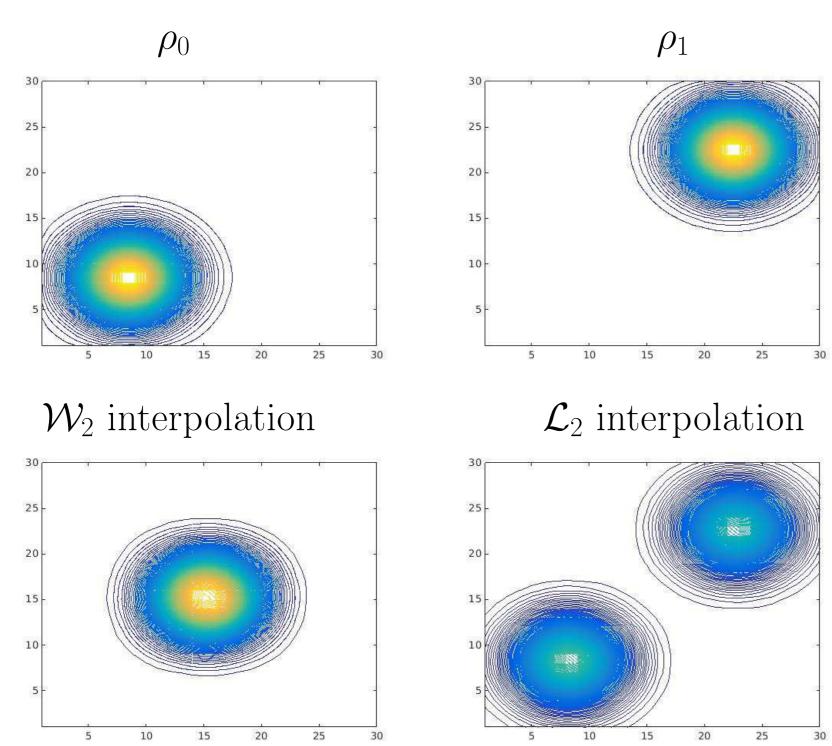
It is common for the distance d to be a weighted  $\mathcal{L}_2$  distance. Yet, with this  $\mathcal{L}_2$  distance, variational data assimilation cannot well cope with displacement errors. By choosing distances d that take into account the datas more in its entirety, we hope get a more realistic variational data assimilation. That is why we are interested in the Wasserstein distance  $d = \mathcal{W}_2$ .

### Advantages of optimal transportation in case of displacement error

Knowing two densities  $\rho_0(x)$  and  $\rho_1(x)$ , the **Wasserstein distance**  $\mathcal{W}_2(\rho_0, \rho_1)$  is defined as

$$\mathcal{W}_{2}^{2}(\rho_{0}, \rho_{1}) := \inf_{\substack{(\rho(t, x), v(t, x)) \\ \partial_{t}\rho + \operatorname{div}(\rho v) = 0 \\ \rho(0, x) = \rho_{0}(x), \rho(1, x) = \rho_{1}(x)}} \int \int_{[0, 1] \times \Omega} \rho(t, x) |v(t, x)|^{2} dt dx.$$

Example of interpolation using Wasserstein and  $\mathcal{L}_2$  distances. (Interpolation in the sense of minimizing  $d(\rho, \rho_0)^2 + d(\rho, \rho_1)^2$ ).



For Wasserstein distance to be defined, one needs  $\rho_0 \geq 0$ ,  $\rho_1 \geq 0$  and  $\int_{\Omega} \rho_0 = \int_{\Omega} \rho_1$ .

# GRADIENT DESCENT WITH THE WASSERSTEIN DISTANCE

Assuming  $\mathbf{x}_0$ ,  $\mathcal{H}_i(\mathbf{x}_0)$  and  $\mathbf{y}_i^o$  are probability measures, the cost function  $\mathcal{J}_W$  is

$$\mathcal{J}_W(\mathbf{x}_0) = \frac{1}{2} \sum_i \mathcal{W}_2^2 \Big( \mathcal{H}_i(\mathbf{x}_0), \mathbf{y}_i^o \Big)$$

The gradient descent algorithm for minimizing  $\mathcal{J}_W$  consists in using the following iterative algorithm

$$\mathbf{x}_0^{n+1} = \mathbf{x}_0^n - \alpha \operatorname{grad} \mathcal{J}_W(\mathbf{x}_0^n)$$

so that  $\mathcal{J}_W(\mathbf{x}_0^{n+1}) < \mathcal{J}_W(\mathbf{x}_0^n)$ . To get the gradient, the **differential** must be computed and also an **inner product**. Indeed,  $\operatorname{grad} \mathcal{J}(\mathbf{x}_0)$  is such that for all  $\eta$ ,  $(\eta, \operatorname{grad} \mathcal{J}(\mathbf{x}_0)) = D\mathcal{J}[\mathbf{x}_0].\eta$ .

# The differential:

Let's differentiate  $\mathcal{J}_W$  by computing  $\mathcal{J}_W(\mathbf{x}_0 + \epsilon \eta)$ :

- $\mathcal{H}_i(\mathbf{x}_0 + \epsilon \eta) = \mathcal{H}_i(\mathbf{x}_0) + \epsilon L[t_i, \mathbf{x}_0].\eta$  with L the tangent model.
- $\frac{1}{2}W_2^2(\mathbf{y} + \epsilon \mu, \mathbf{y}') = \epsilon \langle \mu, \psi \rangle_2$  with  $\psi$  the **Kantorovich potential** of optimal transportation between  $\mathbf{y}$  and  $\mathbf{y}'$ .

Then,

$$\mathrm{D}\mathcal{J}_W[\mathbf{x}_0].\eta = \sum_i \left\langle L[t_i,\mathbf{x}_0].\eta , \psi_i \right\rangle_2$$

with  $\psi_i$  the Kantorovich potential of the optimal transportation between  $\mathcal{H}_i(\mathbf{x}_0)$  and  $\mathbf{y}_i^o$ .

# The inner product:

For convergence reason, instead of using the  $\mathcal{L}_2$  inner product, we use the  $\mathcal{W}_2$  inner product defined in  $\mathbf{x}_0$  by

$$(\eta, \eta') = \int_{\Omega} \mathbf{x}_0 \nabla \Phi \cdot \nabla \Phi',$$
 with  $\Phi$ ,  $\Phi'$  s.t.  $\eta = -\text{div}(\mathbf{x}_0 \nabla \Phi), \qquad \eta' = -\text{div}(\mathbf{x}_0 \nabla \Phi')$ 

Finally, the gradient of  $\mathcal{J}_W$  w.r.t.  $\mathcal{W}_2$  inner product is,  $(L^*)$  is the **adjoint model**),

$$\operatorname{grad} \mathcal{J}_W(\mathbf{x}_0) = -\operatorname{div} \left( \mathbf{x}_0 \nabla \left( \sum_i L^*[t_i, \mathbf{x}_0]. \psi_i \right) \right).$$

#### Assimilation of non-probability measure variables

In the case where  $\mathcal{H}_i(\mathbf{x}_0)$  and  $\mathbf{y}_0$  are probability measures, but not  $\mathbf{x}_0$ , it is not possible to write the background term  $\mathcal{J}^b$  with  $\mathcal{W}_2$ . For example, for the Shallow-water system

$$\begin{cases} \partial_t h + \operatorname{div}(hu) = 0 \\ \partial_t u + (u \cdot \nabla)u + g\nabla h = \mu \Delta u \end{cases}$$
 (2)

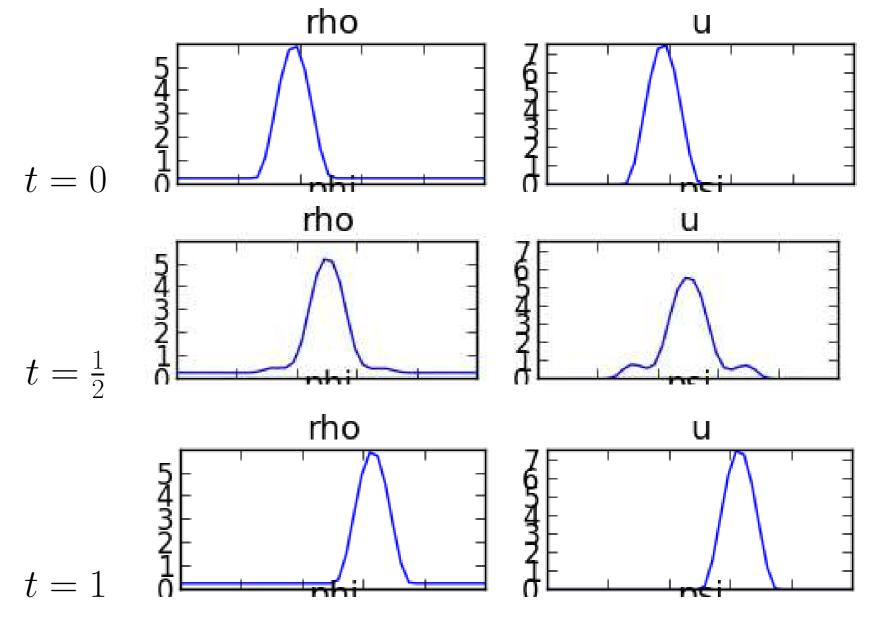
where  $(h_0, u_0)$  are to be assimilated using observations of h only. The cost function writes

$$\mathcal{J}(h_0, u_0) = \frac{1}{2} \sum_i \mathcal{W}_2^2 \Big( \mathcal{H}_i(h_0, \mathbf{u}_0), h_i^o \Big) + \gamma \mathcal{J}^b(h_0, u_0).$$

As it is impossible to write  $\mathcal{W}_2^2(u_0, u_0^b)$ , we rather use the following background term

$$J^{b}(h_{0}, u_{0}) = \inf_{\begin{subarray}{l}\partial_{t}\rho + \operatorname{div}(\rho v) = 0\\ \partial_{t}u + \operatorname{div}(\rho w) = 0\end{subarray}} \int \int_{[0,1]\times\Omega} (|v|^{2} + |w|^{2}) \rho \, \mathrm{d}t \, \mathrm{d}x, \tag{3}$$
$$\rho(0, x) = h_{0}, \rho(1, x) = h_{0}^{b}$$
$$u(0, x) = u_{0}, u(1, x) = u_{0}^{b}$$

Using this, the interpolation of two shifted  $(h_0, u_0)$  and  $(h_1, u_1)$  is in-between:



The inner product to be chosen for computing the gradient will be

$$\left(\begin{pmatrix} \eta \\ v \end{pmatrix}, \begin{pmatrix} \eta' \\ v' \end{pmatrix}\right) = \int_{\Omega} h_0 \nabla \Phi \cdot \nabla \Phi' + \int_{\Omega} h_0 \nabla \Psi \cdot \nabla \Psi'$$

with  $\Phi$ ,  $\Phi'$ ,  $\Psi$ ,  $\Psi'$  such that

$$\eta = -\operatorname{div}(h_0 \nabla \Phi), \qquad \eta' = -\operatorname{div}(h_0 \nabla \Phi')$$

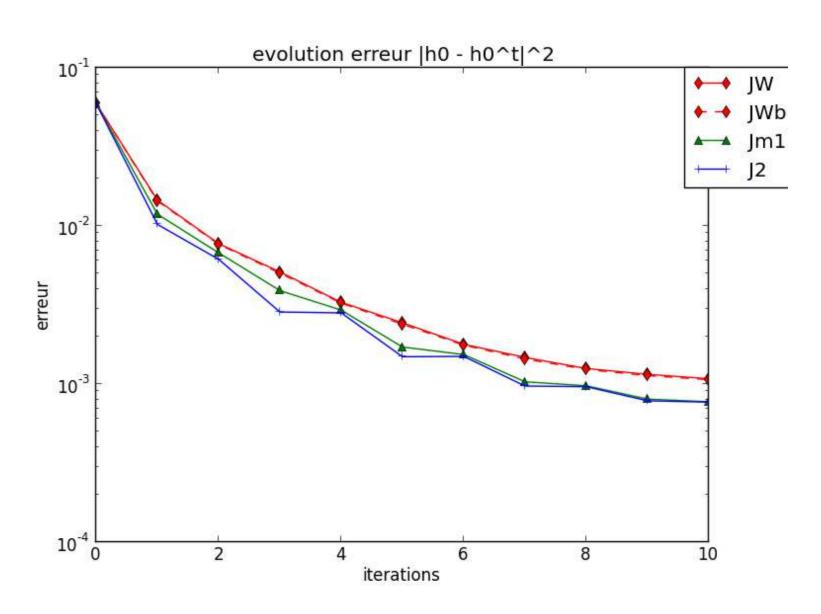
$$v = -\operatorname{div}(h_0 \nabla \Psi), \qquad v' = -\operatorname{div}(h_0 \nabla \Psi').$$

# Difficulties of using the Wasserstein distance

- The Wasserstein distance is only defined for probability measures.
- When  $\rho_0, \rho_1 \approx 1$ , the  $W_2$  interpolation looks like  $\mathcal{L}_2$  interpolation...
- When  $\mathcal{J}(h_0^n) \to \min_{h_0} \mathcal{J}(h_0)$ , one only has  $h_0^n \rightharpoonup \arg\min_{h_0} \mathcal{J}(h_0)$ .
- The computing time is larger for  $W_2$  than for  $\mathcal{L}_2$  [Peyré, Papadakis, Oudet, 2013].

# RESULTS AND PROPECTS

With some tests on the assimilation of (2), we compare the error  $||h_0 - h_0^t||_2^2$  by using  $d = \mathcal{L}_2$  and  $d = \mathcal{W}_2$  in the cost function. The behaviors seem correct.



The background term (3) is still to be implemented.