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# The Stochastic Porous Media Equations in $\mathbb{R}^{d}$ 

Viorel Barbu* Michael Röckner ${ }^{\dagger}$ Francesco Russo ${ }^{\ddagger}$


#### Abstract

Existence and uniqueness of solutions to the stochastic porous media equation $d X-\Delta \psi(X) d t=X d W$ in $\mathbb{R}^{d}$ are studied. Here, $W$ is a Wiener process, $\psi$ is a maximal monotone graph in $\mathbb{R} \times \mathbb{R}$ such that $\psi(r) \leq C|r|^{m}, \forall r \in \mathbb{R}, W$ is a coloured Wiener process. In this general case the dimension is restricted to $d \geq 3$, the main reason being the absence of a convenient multiplier result in the space $\mathcal{H}=\left\{\varphi \in \mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right) ;|\xi|(\mathcal{F} \varphi)(\xi) \in L^{2}\left(\mathbb{R}^{d}\right)\right\}$, for $d \leq 2$. When $\psi$ is Lipschitz, the well-posedness, however, holds for all dimensions on the classical Sobolev space $H^{-1}\left(\mathbb{R}^{d}\right)$. If $\psi(r) r \geq \rho|r|^{m+1}$ and $m=\frac{d-2}{d+2}$, we prove the finite time extinction with strictly positive probability.


Keywords: stochastic, porous media, Wiener process, maximal monotone graph, distributions.
2000 Mathematics Subject Classifications: 76505, 60H15, 47D007.

## 1 Introduction

Consider the stochastic porous media equation

$$
\begin{align*}
& d X-\Delta \psi(X) d t=X d W \text { in }(0, T) \times \mathbb{R}^{d}, \\
& X(0)=x \text { on } \mathbb{R}^{d} \tag{1.1}
\end{align*}
$$

where $\psi$ is a monotonically nondecreasing function on $\mathbb{R}$ (eventually multivalued) and $W(t)$ is a Wiener process of the form

[^0]\[

$$
\begin{equation*}
W(t)=\sum_{k=1}^{\infty} \mu_{k} e_{k} \beta_{k}(t), t \geq 0 \tag{1.2}
\end{equation*}
$$

\]

Here $\left\{\beta_{k}\right\}_{k=1}^{\infty}$ are independent Brownian motions on a stochastic basis $\left\{\Omega, \mathcal{F}, \mathcal{F}_{t}, \mathbb{P}\right\}, \mu_{k} \in \mathbb{R}$ and $\left\{e_{k}\right\}_{k=1}^{\infty}$ is an orthonormal basis in $H^{-1}\left(\mathbb{R}^{d}\right)$ or $\mathcal{H}^{-1}$ (see (2.2) below) to be made precise later on.

On bounded domains $\mathcal{O} \subset \mathbb{R}^{d}$ with Dirichlet homogeneous boundary conditions, equation (1.1) was studied in [3], [4], [6], under general assumptions on $\psi: \mathbb{R} \rightarrow \overline{\mathbb{R}}$ (namely, maximal monotone graph with polynomial growth, or even more general growth conditions in [4]). It should be said, however, that there is a principal difference between bounded and unbounded domains, mainly due to the multiplier problem in Sobolev spaces on $\mathbb{R}^{d}$. If $d \geq 3$ and $\mathcal{O}=\mathbb{R}^{d}$, existence and uniqueness of solutions to (1.1) was proved in [21] (see, also, [23]) in a general setting which covers the case $\mathcal{O}=\mathbb{R}^{d}$ (see Theorems 3.9, Proposition 3.1 and Example 3.4 in [21]). However, it should be said that in [21] $\psi$ is assumed continuous, such that $r \psi(r) \rightarrow \infty$ as $r \rightarrow \infty$, which we do not need in this paper.

We study the existence and uniqueness of (1.1) under two different sets of conditions requiring a different functional approach. The first one, which will be presented in Section 3, assumes that $\psi$ is monotonically nondecreasing and Lipschitz. The state space for $(1.1)$ is, in this case, $H^{-1}\left(\mathbb{R}^{d}\right)$, that is, the dual of the classical Sobolev space $H^{1}\left(\mathbb{R}^{d}\right)$. In spite of the apparent lack of generality ( $\psi$ Lipschitz), it should be mentioned that there are physical models described by such an equation as, for instance, the two phase Stefan transition problem perturbed by a stochastic Gaussian noise [2]; moreover, in this case there is no restriction on the dimension $d$.

The second case, which will be studied here in Section 4, is the one where $\psi$ is a maximal monotone multivalued function with polynomial growth. An important physical problem covered by this case is the self-organized criticality model

$$
\begin{equation*}
d X-\Delta H\left(X-X_{c}\right) d t=\left(X-X_{c}\right) d W \tag{1.3}
\end{equation*}
$$

where $H$ is the Heaviside function and $X_{c}$ is the critical state (see [6], [7], [10]). More generally, this case with discontinuous $\psi$ covers the stochastic nonlinear diffusion equation with singular diffusivity $D(u)=\psi^{\prime}(u)$.

It should be mentioned that, in this second case, the solution $X(t)$ to (1.1) is defined in a certain distribution space $\mathcal{H}^{-1}$ (see (2.2) below) on $\mathbb{R}^{d}$ and the existence is obtained for $d \geq 3$ only, as in the case of continuous $\psi$
in [21]. The case $1 \leq d \leq 2$ remains open due to the absence of a multiplier rule in the norm $\|\cdot\|_{\mathcal{H}^{-1}}$ (see Lemma 4.1 below).

Finally, in Section 5 we prove the finite time extinction of the solution $X$ to (1.1) with strictly positive probability under the assumption that $\psi(r) r \geq$ $\rho|r|^{m+1}$ and $m=\frac{d-2}{d+2}$.

## 2 Preliminaries

To begin with, let us briefly recall a few definitions pertaining distribution spaces on $\mathbb{R}^{d}$, whose classical Euclidean norm will be denoted by $|\cdot|$.

Denote by $\mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right)$ the space of all temperate distributions on $\mathbb{R}^{d}$ (see, e.g., [18]) and by $\mathcal{H}$ the space

$$
\begin{equation*}
\mathcal{H}=\left\{\varphi \in \mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right) ; \xi \mapsto|\xi| \mathcal{F}(\varphi)(\xi) \in L^{2}\left(\mathbb{R}^{d}\right)\right\} \tag{2.1}
\end{equation*}
$$

where $\mathcal{F}(\varphi)$ is the Fourier transform of $\varphi$. We denote by $L^{2}\left(\mathbb{R}^{d}\right)$ the space of square integrable functions on $\mathbb{R}^{d}$ with norm $|\cdot|_{2}$ and scalar product $\langle\cdot, \cdot\rangle_{2}$. In general $|\cdot|_{p}$ will denote the norm of $L^{p}\left(\mathbb{R}^{d}\right)$ or $L^{p}\left(\mathbb{R}^{d} ; \mathbb{R}^{d}\right), 1 \leq p \leq \infty$. The dual space $\mathcal{H}^{-1}$ of $\mathcal{H}$ is given by

$$
\begin{equation*}
\mathcal{H}^{-1}=\left\{\eta \in \mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right) ; \xi \mapsto \mathcal{F}(\eta)(\xi)|\xi|^{-1} \in L^{2}\left(\mathbb{R}^{d}\right)\right\} . \tag{2.2}
\end{equation*}
$$

The duality between $\mathcal{H}$ and $\mathcal{H}^{-1}$ is denoted by $\langle\cdot, \cdot\rangle$ and is given by

$$
\begin{equation*}
\langle\varphi, \eta\rangle=\int_{\mathbb{R}^{d}} \mathcal{F}(\varphi)(\xi) \overline{\mathcal{F}(\eta)}(\xi) d \xi \tag{2.3}
\end{equation*}
$$

and the norm of $\mathcal{H}$ denoted by $\|\cdot\|_{1}$ is given by

$$
\begin{equation*}
\|\varphi\|_{1}=\left(\int_{\mathbb{R}^{d}}|\mathcal{F}(\varphi)(\xi)|^{2}|\xi|^{2} d \xi\right)^{\frac{1}{2}}=\left(\int_{\mathbb{R}^{d}}|\nabla \varphi|^{2} d \xi\right)^{\frac{1}{2}} \tag{2.4}
\end{equation*}
$$

The norm of $\mathcal{H}^{-1}$, denoted by $\|\cdot\|_{-1}$ is given by

$$
\begin{equation*}
\|\eta\|_{-1}=\left(\int_{\mathbb{R}^{d}}|\xi|^{-2}|\mathcal{F}(\eta)(\xi)|^{2} d \xi\right)^{\frac{1}{2}}=\left(\left\langle(-\Delta)^{-1} \eta, \eta\right\rangle\right)^{\frac{1}{2}} \tag{2.5}
\end{equation*}
$$

(We note that the operator $-\Delta$ is an isomorphism from $\mathcal{H}$ onto $\mathcal{H}^{-1}$.) The scalar product of $\mathcal{H}^{-1}$ is given by

$$
\begin{equation*}
\left\langle\eta_{1}, \eta_{2}\right\rangle_{-1}=\left\langle(-\Delta)^{-1} \eta_{1}, \eta_{2}\right\rangle \tag{2.6}
\end{equation*}
$$

As regards the relationship of $\mathcal{H}$ with the space $L^{p}\left(\mathbb{R}^{d}\right)$ of $p$-summable functions on $\mathbb{R}^{d}$, we have the following.

Lemma 2.1 Let $d \geq 3$. Then we have

$$
\begin{equation*}
\mathcal{H} \subset L^{\frac{2 d}{d-2}}\left(\mathbb{R}^{d}\right) \tag{2.7}
\end{equation*}
$$

algebraically and topologically.
Indeed, by the Sobolev embedding theorem (see, e.g., [15], p. 278), we have

$$
|\varphi|_{\frac{2 d}{d-2}} \leq C|\nabla \varphi|_{2}, \forall \varphi \in C_{0}^{\infty}\left(\mathbb{R}^{d}\right)
$$

and, by density, this implies (2.7), as claimed.
It should be mentioned that (2.7) is no longer true for $1 \leq d \leq 2$. However, by duality, we have

$$
\begin{equation*}
L^{\frac{2 d}{d+2}}\left(\mathbb{R}^{d}\right) \subset \mathcal{H}^{-1}, \forall d \geq 3 \tag{2.8}
\end{equation*}
$$

Denote by $H^{1}\left(\mathbb{R}^{d}\right)$ the Sobolev space

$$
\begin{aligned}
H^{1}\left(\mathbb{R}^{d}\right) & =\left\{u \in L^{2}\left(\mathbb{R}^{d}\right) ; \nabla u \in L^{2}\left(\mathbb{R}^{d}\right)\right\} \\
& =\left\{u \in L^{2}\left(\mathbb{R}^{d}\right) ; \xi \mapsto \mathcal{F}(u)(\xi)\left(1+|\xi|^{2}\right)^{\frac{1}{2}} \in L^{2}\left(\mathbb{R}^{d}\right)\right\}
\end{aligned}
$$

with norm

$$
|u|_{H^{1}\left(\mathbb{R}^{d}\right)}=\left(\int_{\mathbb{R}^{d}}\left(u^{2}+|\nabla u|^{2}\right) d \xi\right)^{\frac{1}{2}}=\left(\int_{\mathbb{R}^{d}}|\mathcal{F} u(\xi)|^{2}\left(1+|\xi|^{2}\right) d \xi\right)^{\frac{1}{2}}
$$

and by $H^{-1}\left(\mathbb{R}^{d}\right)$ its dual, that is,

$$
H^{-1}\left(\mathbb{R}^{d}\right)=\left\{u \in \mathcal{S}^{\prime}\left(\mathbb{R}^{d}\right) ; \mathcal{F}(u)(\xi)\left(1+|\xi|^{2}\right)^{-\frac{1}{2}} \in L^{2}\left(\mathbb{R}^{d}\right)\right\}
$$

The norm of $H^{-1}\left(\mathbb{R}^{d}\right)$ is denoted by $|\cdot|_{-1}$ and its scalar product by $\langle\cdot, \cdot\rangle_{-1}$. We have the continuous and dense embeddings

$$
H^{1}\left(\mathbb{R}^{d}\right) \subset \mathcal{H}, \mathcal{H}^{-1} \subset H^{-1}\left(\mathbb{R}^{d}\right)
$$

It should be emphasized, however, that $\mathcal{H}$ is not a subspace of $L^{2}\left(\mathbb{R}^{d}\right)$ and so $L^{2}\left(\mathbb{R}^{d}\right)$ is not the pivot space in the duality $\langle\cdot, \cdot\rangle$ given by (2.3).

Given a Banach space $Y$, we denote by $L^{p}(0, T ; Y)$ the space of all $Y$ valued Bochner measurable $p$-integrable functions on $(0, T)$ and by $C([0, T] ; Y)$ the space of continuous $Y$-valued functions on $[0, T]$. For two Hilbert spaces $H_{1}, H_{2}$ let $L\left(H_{1}, H_{2}\right)$ and $L_{2}\left(H_{1}, H_{2}\right)$ denote the set of all bounded linear and Hilbert-Schmidt operators, respectively. We refer to [17], [20] for definitions and basic results pertaining infinite dimensional stochastic processes.

## 3 Equation (1.1) with the Lipschitzian $\psi$

Consider here equation (1.1) under the following conditions.
(i) $\psi: \mathbb{R} \rightarrow \mathbb{R}$ is monotonically nondecreasing, Lipschitz such that $\psi(0)=0$.
(ii) $W$ is a Wiener process as in (1.2), where $e_{k} \in H^{1}\left(\mathbb{R}^{d}\right)$, such that

$$
\begin{equation*}
C_{\infty}^{2}:=36 \sum_{k=1}^{\infty} \mu_{k}^{2}\left(\left|\nabla e_{k}\right|_{\infty}^{2}+\left|e_{k}\right|_{\infty}^{2}+1\right)<\infty, \tag{3.1}
\end{equation*}
$$

and $\left\{e_{k}\right\}$ is an orthonormal basis in $H^{-1}\left(\mathbb{R}^{d}\right)$.
We insert the factor 36 for convenience here to avoid additional large numerical constants in subsequent estimates.

Remark 3.1 By Lemma 4.1 below, $\left|\nabla e_{k}\right|_{\infty}$ in (3.1) can be replaced by $\left|\nabla e_{k}\right|_{d}$, and all the results in this section remain true.

Definition 3.2 Let $x \in H^{-1}\left(\mathbb{R}^{d}\right)$. A continuous, $\left(\mathcal{F}_{t}\right)_{t \geq 0^{-}}$adapted process $X:[0, T] \rightarrow H^{-1}\left(\mathbb{R}^{d}\right)$ is called strong solution to (1.1) if the following conditions hold:

$$
\begin{align*}
& X \in L^{2}\left(\Omega ; C\left([0, T] ; H^{-1}\left(\mathbb{R}^{d}\right)\right)\right) \cap L^{2}\left([0, T] \times \Omega ; L^{2}\left(\mathbb{R}^{d}\right)\right)  \tag{3.2}\\
& \int_{0}^{\bullet} \psi(X(s)) d s \in C\left([0, T] ; H^{1}\left(\mathbb{R}^{d}\right)\right), \mathbb{P} \text {-a.s. }  \tag{3.3}\\
& X(t)-\Delta \int_{0}^{t} \psi(X(s)) d s=x+\int_{0}^{t} X(s) d W(s)  \tag{3.4}\\
& \forall t \in[0, T], \mathbb{P} \text {-a.s. }
\end{align*}
$$

Remark 3.3 The stochastic (Itô-) integral in (3.4) is the standard one from [17] or [20]. In fact, in the terminology of these references, $W$ is a $Q$-Wiener process $W^{Q}$ on $H^{-1}$, where $Q: H^{-1} \rightarrow H^{-1}$ is the symmetric trace class operator defined by

$$
Q h:=\sum_{k=1}^{\infty} \mu_{k}^{2}\left\langle e_{k}, h\right\rangle_{-1} e_{k}, h \in H^{-1} .
$$

For $x \in H^{-1}$, define $\sigma(x): Q^{1 / 2} H^{-1} \rightarrow H^{-1}$ by

$$
\begin{equation*}
\sigma(x)\left(Q^{1 / 2} h\right)=\sum_{k=1}^{\infty}\left(\mu_{k}\left\langle e_{k}, h\right\rangle_{-1} e_{k} \cdot x\right), h \in H \tag{3.5}
\end{equation*}
$$

By (3.1), each $e_{k}$ is an $H^{-1}$-multiplier such that

$$
\begin{equation*}
\left|e_{k} \cdot x\right|_{-1} \leq 2\left(\left|e_{k}\right|_{\infty}+\left|\nabla e_{k}\right|_{\infty}\right)|x|_{-1}, x \in H^{-1} \tag{3.6}
\end{equation*}
$$

Hence, for all $x \in H^{-1}, h \in H^{-1}$,

$$
\begin{aligned}
\sum_{k=1}^{\infty}\left|\mu_{k}\left\langle e_{k}, h\right\rangle_{-1} e_{k} x\right|_{-1} & \leq\left(\sum_{k=1}^{\infty} \mu_{k}^{2}\left|e_{k} x\right|_{-1}^{2}\right)^{1 / 2}|h|_{-1} \\
& \leq 2 C_{\infty}|x|_{-1}|h|_{-1} \\
& =2 C_{\infty}|x|_{-1}\left|Q^{1 / 2} h\right|_{Q^{1 / 2} H^{-1}}
\end{aligned}
$$

and thus $\sigma(x)$ is well-defined and an element in $L\left(Q^{1 / 2} H^{-1}, H^{-1}\right)$. Moreover, for $x \in H^{-1}$, by (3.5), (3.6),

$$
\begin{align*}
\|\sigma(x)\|_{L_{2}\left(Q^{1 / 2} H^{-1}, H^{-1}\right)}^{2} & =\sum_{k=1}^{\infty}\left|\sigma(x)\left(Q^{1 / 2} e_{k}\right)\right|_{-1}^{2}=\sum_{k=1}^{\infty}\left|\mu_{k} e_{k} x\right|_{-1}^{2}  \tag{3.7}\\
& =\sum_{k=1}^{\infty} \mu_{k}^{2}\left|e_{k} x\right|_{-1}^{2} \leq C_{\infty}^{2}|x|_{-1}^{2} .
\end{align*}
$$

Since $\left\{Q^{1 / 2} e_{k} \mid k \in \mathbb{N}\right\}$ is an orthonormal basis of $Q^{1 / 2} H^{-1}$, it follows that $\sigma(x) \in L_{2}\left(Q^{1 / 2} H^{-1}, H^{-1}\right)$ and the map $x \mapsto \sigma(x)$ is linear and continuous (hence Lipschitz) from $H^{-1}$ to $L_{2}\left(Q^{1 / 2} H^{-1}, H^{-1}\right)$. Hence (e.g., according to [20, Section 2.3])

$$
\int_{0}^{t} X(s) d W(s):=\int_{0}^{t} \sigma(X(s)) d W^{Q}(s), t \in[0, T]
$$

is well-defined as a continuous $H^{-1}$-valued martingale and by Itô's isometry and (3.7)

$$
\begin{align*}
\mathbb{E}\left|\int_{0}^{t} X(s) d W(s)\right|_{-1}^{2} & =\sum_{k=1}^{\infty} \mu_{k}^{2} \mathbb{E} \int_{0}^{t}\left|X(s) e_{k}\right|_{-1}^{2} d s  \tag{3.8}\\
& \leq C_{\infty}^{2} \mathbb{E} \int_{0}^{t}|X(s)|_{-1}^{2} d s, t \in[0, T]
\end{align*}
$$

Furthermore, it follows that

$$
\begin{align*}
\int_{0}^{t} X(s) d W(s) & =\sum_{k=1}^{\infty} \int_{0}^{t} \sigma(X(s))\left(Q^{1 / 2} e_{k}\right) d \beta_{k}(s) \\
& =\sum_{k=1}^{\infty} \int_{0}^{t} \mu_{k} e_{k} X(s) d \beta_{k}(s), t \in[0, T] \tag{3.9}
\end{align*}
$$

where the series converges in $L^{2}\left(\Omega ; C\left([0, T] ; H^{-1}\right)\right)$.
In fact, $\int_{0}^{\bullet} X(s) d W(s)$ is a continuous $L^{2}$-valued martingale, because $X \in L^{2}\left([0, T] \times \Omega ; L^{2}\left(\mathbb{R}^{d}\right)\right)$ and, analogously to (3.7), we get

$$
\|\sigma(x)\|_{L_{2}\left(Q^{1 / 2} H^{-1}, L^{2}\right)}^{2} \leq C_{\infty}^{2}|x|_{2}^{2}, \quad x \in L^{2}\left(\mathbb{R}^{d}\right)
$$

In particular, by Itô's isometry,

$$
\mathbb{E}\left|\int_{0}^{t} X(s) d W(s)\right|_{2}^{2} \leq C_{\infty}^{2} \mathbb{E} \int_{0}^{t}|X(s)|_{2}^{2} d s, t \in[0, T]
$$

Furthermore, the series in (3.8) even converges in $L^{2}\left(\Omega ; C\left([0, T] ; L^{2}\left(\mathbb{R}^{d}\right)\right)\right)$.
We shall use the facts presented in this remark throughout this paper without further notice.

Theorem 3.4 Let $d \geq 1$ and $x \in L^{2}\left(\mathbb{R}^{d}\right)$. Then, under assumptions (i), (ii), there is a unique strong solution to equation (1.1). This solution satisfies

$$
\mathbb{E}\left[\sup _{t \in[0, T]}|X(t)|_{2}^{2}\right] \leq 2|x|_{2}^{2} e^{3 C_{\infty}^{2} t} .
$$

In particular, $X \in L^{2}\left(\Omega ; L^{\infty}\left([0, T] ; L^{2}\left(\mathbb{R}^{d}\right)\right)\right)$. Assume further that

$$
\begin{equation*}
\psi(r) r \geq \alpha r^{2}, \forall r \in \mathbb{R} \tag{3.10}
\end{equation*}
$$

where $\alpha>0$. Then, there is a unique strong solution $X$ to (1.1) for all $x \in H^{-1}\left(\mathbb{R}^{d}\right)$.

Proof of Theorem 3.4. We approximate (1.1) by

$$
\begin{align*}
& d X+(\nu-\Delta) \psi(X) d t=X d W(t), t \in(0, T), \\
& X(0)=x \text { on } \mathbb{R}^{d}, \tag{3.11}
\end{align*}
$$

where $\nu \in(0,1)$. We have the following.

Lemma 3.5 Assume that $\psi$ is as in assumption (i). Let $x \in L^{2}\left(\mathbb{R}^{d}\right)$. Then, there is a unique $\left(\mathcal{F}_{t}\right)_{t \geq 0}$-adapted solution $X=X^{\nu}$ to (3.11) in the following strong sense:

$$
\begin{equation*}
X^{\nu} \in L^{2}\left(\Omega, C\left([0, T] ; H^{-1}\left(\mathbb{R}^{d}\right)\right)\right) \cap L^{2}\left([0, T] \times \Omega ; L^{2}\left(\mathbb{R}^{d}\right)\right) \tag{3.12}
\end{equation*}
$$

and $\mathbb{P}$-a.s.

$$
\begin{equation*}
X^{\nu}(t)=x+(\Delta-\nu) \int_{0}^{t} \psi\left(X^{\nu}(s)\right) d s+\int_{0}^{t} X^{\nu}(s) d W(s), t \in[0, T] \tag{3.13}
\end{equation*}
$$

In addition, for all $\nu \in(0,1)$,

$$
\begin{equation*}
\mathbb{E}\left[\sup _{t \in[0, T]}\left|X^{\nu}(t)\right|_{2}^{2}\right] \leq 2|x|_{2}^{2} e^{3 C_{\infty}^{2} T} \tag{3.14}
\end{equation*}
$$

If, moreover, $\psi$ satisfies (3.10), then for each $x \in H^{-1}\left(\mathbb{R}^{d}\right)$ there is a unique solution $X^{\nu}$ satisfying (3.12) and (3.13).

Proof of Lemma 3.5. Let us start with the second part of the assertion, i.e., we assume that $\psi$ satisfies (3.10) and that $x \in H^{-1}\left(\mathbb{R}^{d}\right)$. Then the standard theory (see, e.g., [20, Sections 4.1 and 4.2]) applies to ensure that there exists a unique solution $X^{\nu}$ taking value in $H^{-1}\left(\mathbb{R}^{d}\right)$ satisfying (3.12) and (3.13) above. Indeed, it is easy to check that (H1)-(H4) from [20, Section 4.1] are satisfied with $V:=L^{2}\left(\mathbb{R}^{d}\right), H:=H^{-1}\left(\mathbb{R}^{d}\right), A u:=(\Delta-\nu)(\psi(u))$, $u \in V$, and $H^{-1}\left(\mathbb{R}^{d}\right)$ is equipped with the equivalent norm

$$
|\eta|_{-1, \nu}:=\left\langle\eta,(\nu-\Delta)^{-1} \eta\right\rangle^{1 / 2}, \eta \in H^{-1}\left(\mathbb{R}^{d}\right)
$$

(in which case, we also write $H_{\nu}^{-1}$ ). Here, as before, we use $\langle\cdot, \cdot\rangle$ also to denote the dualization between $H^{1}\left(\mathbb{R}^{d}\right)$ and $H^{-1}\left(\mathbb{R}^{d}\right)$. For details, we refer to the calculations in [20, Example 4.1.11], which since $p=2$, go through when the bounded domain $\Omega$ there is replaced by $\mathbb{R}^{d}$. Hence [20, Theorem 4.2.4] applies to give the above solution $X^{\nu}$.

In the case when $\psi$ does not satisfy (3.10), the above conditions (H1), (H2), (H4) from [20] still hold, but (H3) not in general. Therefore, we replace $\psi$ by $\psi+\lambda I, \lambda \in(0,1)$, and thus consider $A_{\lambda}(u):=(\Delta-\nu)(\psi(u)+\lambda u)$, $u \in V:=L^{2}\left(\mathbb{R}^{d}\right)$ and, as above, by [20, Theorem 4.2.4], obtain a solution $X_{\lambda}^{\nu}$, satisfying (3.12), (3.13), to

$$
\begin{align*}
& d X_{\lambda}^{\nu}(t)+(\nu-\Delta)\left(\psi\left(X_{\lambda}^{\nu}(t)\right)+\lambda X_{\lambda}^{\nu}(t)\right) d t=X_{\lambda}^{\nu}(t) d W(t), t \in[0, T],  \tag{3.15}\\
& X_{\lambda}^{\nu}(0)=x \in H^{-1}\left(\mathbb{R}^{d}\right) .
\end{align*}
$$

In particular, by (3.12),

$$
\begin{equation*}
\mathbb{E}\left[\sup _{t \in[0, T]}\left|X_{\lambda}^{\nu}(t)\right|_{-1}^{2}\right]<\infty \tag{3.16}
\end{equation*}
$$

We want to let $\lambda \rightarrow 0$ to obtain a solution to (3.11). To this end, in this case (i.e., without assuming (3.10)), we assume from now on that $x \in L^{2}\left(\mathbb{R}^{d}\right)$. The reason is that we need the following:

Claim 1. We have $X_{\lambda}^{\nu} \in L^{2}\left([0, T] \times \Omega ; H^{1}\left(\mathbb{R}^{d}\right)\right)$ and

$$
\begin{array}{r}
\mathbb{E}\left[\sup _{t \in[0, T]}\left|X_{\lambda}^{\nu}(t)\right|_{2}^{2}\right]+4 \lambda \mathbb{E} \int_{0}^{T}\left|\nabla X_{\lambda}^{\nu}(s)\right|_{2}^{2} d s \leq 2|x|_{2}^{2} e^{3 C_{\infty}^{2} T} \\
\text { for all } \nu, \lambda \in(0,1) .
\end{array}
$$

Furthermore, $X_{\lambda}^{\nu}$ has continuous sample paths in $L^{2}\left(\mathbb{R}^{d}\right), \mathbb{P}$-a.s.
Proof of Claim 1. We know that

$$
\begin{align*}
X_{\lambda}^{\nu}(t)=x+(\Delta-\nu) & \int_{0}^{t}\left(\psi\left(X_{\lambda}^{\nu}(s)\right)+\lambda X_{\lambda}^{\nu}(s)\right) d s \\
& +\int_{0}^{t} X_{\lambda}^{\nu}(s) d W(s), t \in[0, T] \tag{3.17}
\end{align*}
$$

Let $\alpha \in(\nu, \infty)$. Recalling that $(\alpha-\Delta)^{-\frac{1}{2}}: H^{-1}\left(\mathbb{R}^{d}\right) \rightarrow L^{2}\left(\mathbb{R}^{d}\right)$ and applying this operator to the above equation, we find

$$
\begin{align*}
& (\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(t) \\
& =(\alpha-\Delta)^{-\frac{1}{2}} x+\int_{0}^{t}(\Delta-\nu)(\alpha-\Delta)^{-\frac{1}{2}}\left(\psi\left(X_{\lambda}^{\nu}(s)\right)+\lambda X_{\lambda}^{\nu}(s)\right) d s  \tag{3.18}\\
& +\int_{0}^{t}(\alpha-\Delta)^{-\frac{1}{2}} \sigma\left(X_{\lambda}^{\nu}(s)\right) Q^{1 / 2} d W(s), t \in[0, T] .
\end{align*}
$$

Applying Itô's formula (see, e.g., [20, Theorem 4.2.5] with $H=L^{2}\left(\mathbb{R}^{d}\right)$ ) to $\left|(\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(t)\right|_{2}^{2}$, we obtain, for $t \in[0, T]$,

$$
\begin{align*}
& \left|(\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(t)\right|_{2}^{2}=\left|(\alpha-\Delta)^{-\frac{1}{2}} x\right|_{2}^{2} \\
& +2 \int_{0}^{t}\left\langle(\Delta-\nu)(\alpha-\Delta)^{-\frac{1}{2}} \psi\left(X_{\lambda}^{\nu}(s)\right),(\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s)\right\rangle d s \\
& -2 \lambda \int_{0}^{t}\left(\left|\nabla\left((\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s)\right)\right|_{2}^{2}+\nu\left|(\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s)\right|_{2}^{2}\right) d s  \tag{3.19}\\
& +\int_{0}^{t}\left\|(\alpha-\Delta)^{-\frac{1}{2}} \sigma\left(X_{\lambda}^{\nu}(s)\right) Q^{1 / 2}\right\|_{L_{2}\left(H^{-1}, L^{2}\right)}^{2} d s \\
& +2 \int_{0}^{t}\left\langle(\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s),(\alpha-\Delta)^{-\frac{1}{2}} \sigma\left(X_{\lambda}^{\nu}(s)\right) Q^{1 / 2} d W(s)\right\rangle_{2} .
\end{align*}
$$

But, for $f \in L^{2}\left(\mathbb{R}^{d}\right)$, we have

$$
(\alpha-\Delta)^{-\frac{1}{2}}(\Delta-\nu)(\alpha-\Delta)^{-\frac{1}{2}} f=(P-I) f
$$

where

$$
P:=(\alpha-\nu)(\alpha-\Delta)^{-1} .
$$

For the Green function $g_{\alpha}$ of $(\alpha-\Delta)$, we then have, for $f \in L^{2}\left(\mathbb{R}^{d}\right)$,

$$
P f=(\alpha-\nu) \int_{\mathbb{R}^{d}} f(\xi) g_{\alpha}(\cdot, \xi) d \xi
$$

Hence, by [23, Lemma 5.1], the integrand of the second term on the right hand side of (3.19) with $f:=X_{\lambda}^{\nu}(s)\left(\in L^{2}\left(\mathbb{R}^{d}\right)\right.$ for $d s$-a.e. $\left.s \in[0, T]\right)$ can be rewritten as

$$
\begin{aligned}
\langle\psi(f),(P-I) f\rangle_{2}= & -\frac{1}{2} \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}}[\psi(f(\widetilde{\xi}))-\psi(f(\xi))][f(\widetilde{\xi})-f(\xi)] g_{\alpha}(\xi, \widetilde{\xi}) d \widetilde{\xi} d \xi \\
& -\int_{\mathbb{R}^{d}}(1-P 1(\xi)) \cdot \psi(f(\xi)) f(\xi) d \xi
\end{aligned}
$$

Since $\psi$ is monotone, $\psi(0)=0$ and $P 1 \leq 1$, we deduce that

$$
\langle\psi(f),(P-I) f\rangle \leq 0 .
$$

Hence, after a multiplication by $\alpha$, (3.19) implies that, for all $t \in[0, T]$ (see Remark 3.3),

$$
\begin{aligned}
& \alpha\left|(\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(t)\right|_{2}^{2}+2 \lambda \int_{0}^{t}\left|\nabla\left(\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s)\right)\right|_{2}^{2} d s \\
& \leq \alpha\left|(\alpha-\Delta)^{-\frac{1}{2}} x\right|_{2}^{2}+\int_{0}^{t} \sum_{k=1}^{\infty} \mu_{k}^{2}\left\langle\alpha(\alpha-\Delta)^{-1}\left(e_{k} X_{\lambda}^{\nu}(s)\right), e_{k} X_{\lambda}^{\nu}(s)\right\rangle_{2} d s \\
& +2 \int_{0}^{t}\left\langle\alpha(\alpha-\Delta)^{-1} X_{\lambda}^{\nu}(s), \sigma\left(X_{\lambda}^{\nu}(s)\right) Q^{1 / 2} d W(s)\right\rangle_{2} .
\end{aligned}
$$

Hence, by the Burkholder-Davis-Dundy (BDG) inequality (with $p=1$ ) and since $\alpha(\alpha-\Delta)^{-1}$ is a contraction on $L^{2}\left(\mathbb{R}^{d}\right)$,

$$
\begin{align*}
& \mathbb{E}\left[\sup _{s \in[0, t]}\left|\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s)\right|_{2}^{2}\right] \\
& \quad+2 \lambda \mathbb{E} \int_{0}^{t}\left|\nabla\left(\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s)\right)\right|_{2}^{2} d s \\
& \quad \leq\left|\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}} x\right|_{2}^{2}+C_{\infty}^{2} \mathbb{E} \int_{0}^{t}\left|X_{\lambda}^{\nu}(s)\right|_{2}^{2} d s  \tag{3.20}\\
& \quad+6 \mathbb{E}\left(\int_{0}^{t} \sum_{k=1}^{\infty} \mu_{k}^{2}\left\langle\alpha(\alpha-\Delta)^{-1} X_{\lambda}^{\nu}(s), e_{k} X_{\lambda}^{\nu}(s)\right\rangle_{2}^{2} d s\right)^{1 / 2} .
\end{align*}
$$

The latter term can be estimated by

$$
\begin{align*}
& C_{\infty} \mathbb{E}\left[\sup _{s \in[0, t]}\left|\alpha(\alpha-\Delta)^{-1} X_{\lambda}^{\nu}(s)\right|_{2}\left(\int_{0}^{t}\left|X_{\lambda}^{\nu}(s)\right|_{2}^{2} d s\right)^{1 / 2}\right]  \tag{3.21}\\
& \leq \frac{1}{2} \mathbb{E}\left\{\sup _{s \in[0, t]}\left|\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s)\right|_{2}^{2}\right]+\frac{1}{2} C_{\infty}^{2} \mathbb{E} \int_{0}^{t}\left|X_{\lambda}^{\nu}(s)\right|_{2}^{2} d s,
\end{align*}
$$

where we used that $\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}}$ is a contraction on $L^{2}\left(\mathbb{R}^{d}\right)$. Note that the first summand on the right hand side is finite by (3.16), since the norm $\left|\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}} \cdot\right|_{2}$ is equivalent to $|\cdot|_{-1}$. Hence, we can subtract this term after substituting (3.21) into (3.20) to obtain

$$
\begin{align*}
\mathbb{E}\left[\sup _{s \in[0, t]} \mid\right. & \left.\left.\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s)\right|_{2} ^{2}\right] \\
& +4 \lambda \mathbb{E} \int_{0}^{t} \left\lvert\, \nabla\left(\left.\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s)\right|_{2} ^{2} d s\right.\right.  \tag{3.22}\\
& \leq 2\left|\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}} x\right|_{2}^{2}+3 C_{\infty}^{2} \mathbb{E} \int_{0}^{t}\left|X_{\lambda}^{\nu}(s)\right|_{2}^{2} d s, t \in[0, T]
\end{align*}
$$

Obviously, the quantity under the $\sup _{s \in[0, t]}$ on the left hand side of (3.22) is increasing in $\alpha$. So, by the monotone convergence theorem, we may let $\alpha \rightarrow \infty$ in (3.22) and then, except for its last part, Claim 1 immediately follows by Gronwall's lemma, since $\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}}$ is a contraction in $L^{2}\left(\mathbb{R}^{d}\right)$ and $x \in L^{2}\left(\mathbb{R}^{d}\right)$. The last part of Claim 1 then immediately follows from [19, Theorem 2.1].

Applying Itô's formula to $\left|X_{\lambda}^{\nu}(t)-X_{\lambda^{\prime}}^{\nu}(t)\right|_{-1, \nu}^{2}$ (see [20, Theorem 4.2.5]), it follows from (3.17) that, for $\lambda, \lambda^{\prime} \in(0,1)$ and $t \in[0, T]$,

$$
\begin{align*}
& \left|X_{\lambda}^{\nu}(t)-X_{\lambda^{\prime}}^{\nu}(t)\right|_{-1, \nu}^{2} \\
& \quad+2 \int_{0}^{t}\left\langle\psi\left(X_{\lambda}^{\nu}\right)-\psi\left(X_{\lambda^{\prime}}^{\nu}\right)+\left(\lambda X_{\lambda}^{\nu}-\lambda^{\prime} X_{\lambda^{\prime}}^{\nu}\right), X_{\lambda}^{\nu}-X_{\lambda^{\prime}}^{\nu}\right\rangle_{2} d s \\
& \quad=\int_{0}^{t} \| \sigma\left(X_{\lambda}^{\nu}(s)-X_{\lambda^{\prime}}^{\nu}(s) \|_{L_{2}\left(Q^{1 / 2} H^{-1}, H_{\nu^{-1}}^{-1}\right)}^{2} d s\right.  \tag{3.23}\\
& \quad+2 \int_{0}^{t}\left\langle X_{\lambda}^{\nu}(s)-X_{\lambda^{\prime}}^{\nu}(s), \sigma\left(X_{\lambda}^{\nu}(s)-X_{\lambda^{\prime}}^{\nu}(s)\right) d W^{Q}(s)\right\rangle_{-1, \nu} .
\end{align*}
$$

Our assumption (i) on $\psi$ implies that

$$
\left(\psi(r)-\psi\left(r^{\prime}\right)\right)\left(r-r^{\prime}\right) \geq(\operatorname{Lip} \psi+1)^{-1}\left|\psi(r)-\psi\left(r^{\prime}\right)\right|^{2}, \text { for } r, r^{\prime} \in \mathbb{R}
$$

where $\operatorname{Lip} \psi$ is the Lipschitz constant of $\psi$. Hence (3.23), (3.7) and the BDG inequality (for $p=1$ imply that, for all $t \in[0, T]$,

$$
\begin{aligned}
& \mathbb{E}\left[\sup _{s \in[0, t]}\left|X_{\lambda}^{\nu}(s)-X_{\lambda^{\prime}}^{\nu}(s)\right|_{-1, \nu}^{2}\right] \\
& +2(\operatorname{Lip} \psi+1)^{-1} \mathbb{E} \int_{0}^{t}\left|\psi\left(X_{\lambda}^{\nu}(s)\right)-\psi\left(X_{\lambda^{\prime}}^{\nu}(s)\right)\right|_{2}^{2} d s \\
& \leq 2\left(\lambda+\lambda^{\prime}\right) \mathbb{E} \int_{0}^{t}\left(\left|X_{\lambda}^{\nu}(s)\right|_{2}^{2}+\left|X_{\lambda^{\prime}}^{\nu}(s)\right|_{2}^{2}\right) d s+C_{\infty}^{2} \int_{0}^{t}\left|X_{\lambda}^{\nu}(s)-X_{\lambda^{\prime}}^{\nu}(s)\right|_{-1, \nu}^{2} d s \\
& +2 \mathbb{E}\left(\int_{0}^{t} \sum_{k=1}^{\infty} \mu_{k}^{2}\left\langle X_{\lambda}^{\nu}(s)-X_{\lambda^{\prime}}^{\nu}(s),\left(X_{\lambda}^{\nu}(s)-X_{\lambda^{\prime}}^{\nu}(s)\right) e_{k}\right\rangle_{-1, \nu}^{2} d s\right)^{1 / 2} .
\end{aligned}
$$

By (3.7) and Young's inequality, the latter term is dominated by

$$
\frac{1}{2} \mathbb{E}\left[\sup _{s \in[0, t]}\left|X_{\lambda}^{\nu}(s)-X_{\lambda^{\prime}}^{\nu}(s)\right|_{-1, \nu}^{2}\right]+\frac{1}{2} C_{\infty}^{2} \mathbb{E} \int_{0}^{t}\left|X_{\lambda}^{\nu}(s)-X_{\lambda^{\prime}}^{\nu}(s)\right|_{-1, \nu}^{2} d s
$$

Hence, because of $x \in L^{2}\left(\mathbb{R}^{d}\right)$ and Claim 1, we may now apply Gronwall's lemma to obtain that, for some constant $C$ independent of $\lambda^{\prime}, \lambda$ (and $\nu$ ),

$$
\begin{align*}
& \mathbb{E}\left[\sup _{t \in[0, T]}\left|X_{\lambda}^{\nu}(t)-X_{\lambda^{\prime}}^{\nu}(t)\right|_{-1, \nu}^{2}\right]  \tag{3.24}\\
& \quad+\mathbb{E} \int_{0}^{T}\left|\psi\left(X_{\lambda}^{\nu}(s)\right)-\psi\left(X_{\lambda^{\prime}}^{\nu}(s)\right)\right|_{2}^{2} d s \leq C\left(\lambda+\lambda^{\prime}\right)
\end{align*}
$$

Hence there exists an $\left(\mathcal{F}_{t}\right)$-adapted continuous $H^{-1}$-valued process $X^{\nu}=$ $\left(X^{\nu}(t)\right)_{t \in[0, T]}$ such that $X^{\nu} \in L^{2}\left(\Omega ; C\left([0, T] ; H^{-1}\right)\right)$. Now, by Claim 1, it follows that

$$
X^{\nu} \in L^{2}\left([0, T] \times \Omega ; L^{2}\left(\mathbb{R}^{d}\right)\right)
$$

Claim 2. $X^{\nu}$ satisfies equation (3.13) (i.e., we can pass to the limit in (3.17) as $\lambda \rightarrow 0$ ).

Proof of Claim 2. We already know that

$$
X_{\lambda}^{\nu} \longrightarrow X^{\nu} \text { and } \int_{0}^{\bullet} X_{\lambda}^{\nu}(s) d W(s) \longrightarrow \int_{0}^{\bullet} X^{\nu}(s) d W(s)
$$

in $L^{2}\left(\Omega ; C\left([0, T] ; H^{-1}\right)\right)$ as $\lambda \rightarrow 0$ (for the second convergence see the above argument using (3.7) and the BDG inequality). So, by (3.17) it follows that

$$
\left.\int_{0}^{\bullet}\left(\psi\left(X_{\lambda}^{\nu}(s)\right)+\lambda X_{\lambda}^{\nu}(s)\right)\right) d s, \lambda>0
$$

converges as $\lambda \rightarrow 0$ to an element in $L^{2}\left(\Omega ; C\left([0, T] ; H^{1}\right)\right)$. But, by (3.24) and Claim 1, it follows that

$$
\begin{equation*}
\int_{0}^{\bullet}\left(\psi\left(X_{\lambda}^{\nu}(s)\right)+\lambda X_{\lambda}^{\nu}(s)\right) d s \longrightarrow \int_{0}^{\bullet} \psi\left(X^{\nu}(s)\right) d s \tag{3.25}
\end{equation*}
$$

as $\lambda \rightarrow 0$ in $L^{2}\left(\Omega ; L^{2}\left([0, T] ; L^{2}\left(\mathbb{R}^{d}\right)\right)\right)$. Hence Claim 2 is proved.
Now, (3.14) follows from Claim 1 by lower semicontinuity. This completes the proof of Lemma 3.5.

Proof of Theorem 3.4 (continued). We are going to use Lemma 3.5 and let $\nu \rightarrow 0$. The arguments are similar to those in the proof of Lemma 3.5. So, we shall not repeat all the details.

Now, we rewrite (3.11) as

$$
\begin{equation*}
d X^{\nu}+(I-\Delta) \psi\left(X^{\nu}\right) d t=(1-\nu) \psi\left(X^{\nu}\right) d t+X^{\nu} d W(t) \tag{3.26}
\end{equation*}
$$

and apply Itô's formula to $\varphi(x)=\frac{1}{2}|x|_{-1}^{2}$ (see, e.g., [20, Theorem 4.2.5]). We get, for $x \in H^{-1}$, by (3.8) and after taking expectation,

$$
\begin{aligned}
& \frac{1}{2} \mathbb{E}\left|X^{\nu}(t)\right|_{-1}^{2}+\mathbb{E} \int_{0}^{t} \int_{\mathbb{R}^{d}} \psi\left(X^{\nu}(s)\right) X^{\nu}(s) d \xi d s \\
& \quad=\frac{1}{2}|x|_{-1}^{2}+(1-\nu) \mathbb{E} \int_{0}^{t}\left\langle\psi\left(X^{\nu}(s)\right), X^{\nu}(s)\right\rangle_{-1} d s \\
& \quad+\frac{1}{2} \mathbb{E} \int_{0}^{t} \sum_{k=1}^{\infty} \mu_{k}^{2}\left|X^{\nu} e_{k}\right|_{-1}^{2} d s \\
& \quad \leq \frac{1}{2}|x|_{-1}^{2}+\mathbb{E} \int_{0}^{t}\left|\psi\left(X^{\nu}\right)\right|_{-1}\left|X^{\nu}\right|_{-1} d s \\
& \quad+\frac{1}{2} C_{\infty}^{2} \mathbb{E} \int_{0}^{t}\left|X^{\nu}(s)\right|_{-1}^{2} d s, \quad \forall t \in[0, T] .
\end{aligned}
$$

Recalling that $|\cdot|_{-1} \leq|\cdot|_{2}$, we get, via Young's and Gronwall's inequalities, for some $C \in(0, \infty)$ that

$$
\begin{equation*}
\mathbb{E}\left|X^{\nu}(t)\right|_{-1}^{2}+\frac{\alpha}{2} \mathbb{E} \int_{0}^{T}\left|X^{\nu}(s)\right|_{2}^{2} d s \leq C|x|_{-1}^{2}, t \in[0, T], \nu \in(0,1), \tag{3.27}
\end{equation*}
$$

because, by assumption (i), $\psi(r) r \geq \widetilde{\alpha}|\psi(r)|^{2}, \forall r \in \mathbb{R}$, with $\widetilde{\alpha}:=(\operatorname{Lip} \psi+1)^{-1}$. Here we set $\alpha=0$ if (3.10) does not hold.

Now, by a similar calculus, for $X^{\nu}-X^{\nu^{\prime}}$ we get

$$
\begin{aligned}
& \left|X^{\nu}(t)-X^{\nu^{\prime}}(t)\right|_{-1}^{2}+2 \int_{0}^{t} \int_{\mathbb{R}^{d}}\left(\psi\left(X^{\nu}\right)-\psi\left(X^{\nu^{\prime}}\right)\right)\left(X^{\nu}-X^{\nu^{\prime}}\right) d \xi d s \\
& \quad \leq C \int_{0}^{t}\left\langle\psi\left(X^{\nu}\right)-\psi\left(X^{\nu^{\prime}}\right), X^{\nu}-X^{\nu^{\prime}}\right\rangle_{-1} d s \\
& \quad+C \int_{0}^{t}\left(\nu\left|\psi\left(X^{\nu}\right)\right|_{2}+\nu^{\prime}\left|\psi\left(X^{\nu^{\prime}}\right)\right|_{2}\right)\left|X^{\nu}-X^{\nu^{\prime}}\right|_{-1} d s \\
& \quad+C \int_{0}^{t}\left|X^{\nu}-X^{\nu^{\prime}}\right|_{-1}^{2} d s+\sum_{k=1}^{\infty} \int_{0}^{t} \mu_{k}\left\langle\left(X^{\nu}-X^{\nu^{\prime}}\right), e_{k}\left(X^{\nu}-X^{\nu^{\prime}}\right) d \beta_{k}\right\rangle_{-1} \\
& \quad t \in[0, T]
\end{aligned}
$$

Taking into account that, by assumption (i),

$$
(\psi(x)-\psi(y))(x-y) \geq \widetilde{\alpha}|\psi(x)-\psi(y)|^{2}, \forall x, y \in \mathbb{R}^{d},
$$

we get, for all $\nu, \nu^{\prime}>0$,

$$
\begin{aligned}
& \left|X^{\nu}(t)-X^{\nu^{\prime}}(t)\right|_{-1}^{2}+\tilde{\alpha} \int_{0}^{t}\left|\psi\left(X^{\nu}(s)\right)-\psi\left(X^{\nu^{\prime}}(s)\right)\right|_{2}^{2} d s \\
& \quad \leq C_{1} \int_{0}^{t}\left|X^{\nu}(s)-X^{\nu^{\prime}}(s)\right|_{-1}^{2} d s+\frac{\tilde{\alpha}}{2} \int_{0}^{t}\left|\psi\left(X^{\nu}(s)\right)-\psi\left(X^{\nu^{\prime}}(s)\right)\right|_{2}^{2} d s \\
& \quad+C_{2}\left(\nu+\nu^{\prime}\right) \int_{0}^{t}\left(\left|\psi\left(X^{\nu}(s)\right)\right|_{2}^{2}+\left|\psi\left(X^{\nu^{\prime}}(s)\right)\right|_{2}^{2}\right) d s \\
& \quad+\sum_{k=1}^{\infty} \int_{0}^{t} \mu_{k}\left\langle\left(X^{\nu}(s)-X^{\nu^{\prime}}(s)\right), e_{k}\left(X^{\nu}(s)-X^{\nu^{\prime}}(s)\right)\right\rangle_{-1} d \beta_{k}(s), \\
& \quad t \in[0, T] .
\end{aligned}
$$

So, similarly to showing (3.24) in the proof of Lemma 3.5, by (3.14), if $x \in L^{2}\left(\mathbb{R}^{d}\right)$, and by (3.27), if $x \in H^{-1}\left(\mathbb{R}^{d}\right)$ and $\psi$ satisfies (3.10), by the Burkholder-Davis-Gundy inequality, for $p=1$, we get, for all $\nu, \nu^{\prime} \in(0,1)$,

$$
\mathbb{E} \sup _{t \in[0, T]}\left|X^{\nu}(t)-X^{\nu^{\prime}}(t)\right|_{-1}^{2}+\mathbb{E} \int_{0}^{T}\left|\psi\left(X^{\nu}(s)\right)-\psi\left(X^{\nu^{\prime}}(s)\right)\right|_{2}^{2} d s \leq C\left(\nu+\nu^{\prime}\right)
$$

The remaining part of the proof is now exactly the same as the last part of the proof of Lemma 3.5.

Remark 3.6 Theorem 3.4 is a basic tool for the probabilistic (double) representation of equation (1.1), which holds when $\psi$ is Lipschitz, as it is proved in [9]. If (1.1) is not perturbed by noise, and $\psi$ is possibly discontinuous, its probabilistic representation was performed in [14, 8] with extensions and numerical simulations located in [11, 12].

## 4 Equation (1.1) for maximal monotone functions $\psi$ with polynomial growth

In this section, we assume $d \geq 3$ and we shall study the existence for equation (1.1) under the following assumptions:
(j) $\psi: \mathbb{R} \rightarrow 2^{\mathbb{R}}$ is a maximal monotone graph such that $0 \in \psi(0)$ and

$$
\begin{equation*}
\sup \{|\eta| ; \eta \in \psi(r)\} \leq C\left(1+|r|^{m}\right), \forall r \in \mathbb{R} \tag{4.1}
\end{equation*}
$$

where $1 \leq m<\infty$.
(jj) $W(t)=\sum_{k=1}^{\infty} \mu_{k} e_{k} \beta_{k}(t), t \geq 0$, where $\left\{\beta_{k}\right\}_{k=1}^{\infty}$ are independent Brownian motions on a stochastic basis $\left\{\Omega, \mathcal{F}, \mathcal{F}_{t}, \mathbb{P}\right\}, \mu_{k} \in \mathbb{R}$, and $e_{k} \in C^{1}\left(\mathbb{R}^{d}\right) \cap$ $\mathcal{H}^{-1}$ are such that $\left\{e_{k}\right\}$ is an orthonormal basis in $\mathcal{H}^{-1}$ and

$$
\begin{equation*}
\sum_{k=1}^{\infty} \mu_{k}^{2}\left(\left|e_{k}\right|_{\infty}^{2}+\left|\nabla e_{k}\right|_{d}^{2}+1\right)<\infty \tag{4.2}
\end{equation*}
$$

The existence of $\left\{e_{k}\right\}$ as in ( jj ) is ensured by the following lemma.
Lemma 4.1 Let $d \geq 3$ and let $e \in L^{\infty}\left(\mathbb{R}^{d} ; \mathbb{R}^{d}\right)$ be such that $\nabla e \in L^{d}\left(\mathbb{R}^{d} ; \mathbb{R}^{d}\right)$. Then

$$
\begin{equation*}
\|x e\|_{-1} \leq\|x\|_{-1}\left(|e|_{\infty}+C|\nabla e|_{d}\right), \quad \forall x \in \mathcal{H}^{-1} \tag{4.3}
\end{equation*}
$$

where $C$ is independent of $x$ and $e$.
Proof. We have

$$
\begin{equation*}
\|x e\|_{-1}=\sup \left\{\langle x, e \varphi\rangle ;\|\varphi\|_{1} \leq 1\right\} \leq\|x\|_{-1} \sup \left\{\|e \varphi\|_{1} ;\|\varphi\|_{1} \leq 1\right\} . \tag{4.4}
\end{equation*}
$$

On the other hand, by Lemma 2.1 we have, for all $\varphi \in C_{0}^{\infty}\left(\mathbb{R}^{d}\right)$,

$$
\begin{aligned}
\|e \varphi\|_{1} & \leq|e \nabla \varphi+\varphi \nabla e|_{2} \leq|e \nabla \varphi|_{2}+|\varphi \nabla e|_{2} \\
& \leq|e|_{\infty}|\nabla \varphi|_{2}+|\varphi|_{p}|\nabla e|_{d} \leq|e|_{\infty}\|\varphi\|_{1}+C\|\varphi\|_{1}|\nabla e|_{d},
\end{aligned}
$$

where $p=\frac{2 d}{d-2}$. Then, by (4.4), (4.3) follows, as claimed.

## Remark 4.2

(i) It should be mentioned that, for $d=2$, Lemma 4.1 fails and this is the main reason why our treatment of equation (1.1) under assumptions (j), ( jj ) is constrained to $d \geq 3$.
(ii) We note that Remark 3.3 with the rôle of $H^{-1}\left(\mathbb{R}^{d}\right)$ replaced by $\mathcal{H}^{-1}$ remains true in all its parts under condition (jj) above. We shall use this below without further notice.

We denote by $j: \mathbb{R} \rightarrow \mathbb{R}$ the potential associated with $\psi$, that is, a continuous convex function on $\mathbb{R}$ such that $\partial j=\psi$, i.e.,

$$
j(r) \leq \zeta(r-\bar{r})+j(\bar{r}), \forall \zeta \in \psi(r), r, \bar{r} \in \mathbb{R}
$$

Definition 4.3 Let $x \in \mathcal{H}^{-1}$ and $p:=\max (2,2 m)$. An $\mathcal{H}^{-1}$-valued adapted process $X=X(t)$ is called strong solution to (1.1) if the following conditions hold:

$$
\begin{align*}
& X \text { is } \mathcal{H}^{-1} \text {-valued continuous on }[0, T], \mathbb{P} \text {-a.s., }  \tag{4.5}\\
& X \in L^{p}\left(\Omega \times(0, T) \times \mathbb{R}^{d}\right) \text {. } \tag{4.6}
\end{align*}
$$

There is $\eta \in L^{\frac{p}{m}}\left(\Omega \times(0, T) \times \mathbb{R}^{d}\right)$ such that

$$
\begin{equation*}
\eta \in \psi(X), d t \otimes \mathbb{P} \otimes d \xi-\text { a.e. on }(0, T) \times \Omega \times \mathbb{R}^{d} \tag{4.7}
\end{equation*}
$$

and $\mathbb{P}$-a.s.

$$
\begin{gather*}
X(t)=x+\Delta \int_{0}^{t} \eta(s) d s+\sum_{k=1}^{\infty} \mu_{k} \int_{0}^{t} X(s) e_{k} d \beta_{k}(s)  \tag{4.8}\\
\text { in } \mathcal{D}^{\prime}\left(\mathbb{R}^{d}\right), \quad t \in[0, T] .
\end{gather*}
$$

Here $\mathcal{D}^{\prime}\left(\mathbb{R}^{d}\right)$ is the standard space of distributions on $\mathbb{R}^{d}$.
Theorem 4.4 below is the main existence result for equation (1.1).
Theorem 4.4 Assume that $d \geq 3$ and that $x \in L^{p}\left(\mathbb{R}^{d}\right) \cap L^{2}\left(\mathbb{R}^{d}\right) \cap \mathcal{H}^{-1}$, $p:=\max (2,2 m)$. Then, under assumptions $(\mathrm{j}),(\mathrm{jj})$, there is a unique strong solution $X$ to (1.1) such that

$$
\begin{equation*}
X \in L^{2}\left(\Omega ; C\left([0, T] ; \mathcal{H}^{-1}\right)\right) \tag{4.9}
\end{equation*}
$$

Moreover, if $x \geq 0$, a.e. in $\mathbb{R}^{d}$, then $X \geq 0$, a.e. on $(0, T) \times \mathbb{R}^{d} \times \Omega$.

Theorem 4.4 is applicable to a large class of nonlinearities $\psi: \mathbb{R} \rightarrow 2^{\mathbb{R}}$ and, in particular, to

$$
\psi(r)=\rho H(r)+\alpha r, \forall r \in \mathbb{R}, \psi(r)=\rho H\left(r-r_{c}\right) r
$$

where $\rho>0, \alpha, r_{c} \geq 0$, which models the dynamics of self-organized criticality (see [6], [7], [10]). Here $H$ is the Heaviside function.

As mentioned earlier, Theorem 4.4 can be compared most closely to the main existence result of [21]. But there are, however, a few notable differences we explain below. The function $\psi$ arising in [21] is monotonically increasing, continuous and is assumed to satisfy a growth condition of the form $N(r) \leq r \psi(r) \leq C(N(r)+1) r, \forall r \in \mathbb{R}$, where $N$ is a smooth and $\Delta_{2}$-regular Young function defining the Orlicz class $L_{N}$. In contrast to this, here $\psi$ is any maximal monotone graph (multivalued) with arbitrary polynomial growth.

Proof of Theorem 4.4. Consider the approximating equation

$$
\begin{align*}
& d X_{\lambda}-\Delta\left(\psi_{\lambda}\left(X_{\lambda}\right)+\lambda X_{\lambda}\right) d t=X_{\lambda} d W, t \in(0, T)  \tag{4.10}\\
& X_{\lambda}(0)=x
\end{align*}
$$

where $\psi_{\lambda}=\frac{1}{\lambda}\left(1-(1+\lambda \psi)^{-1}\right), \lambda>0$. We note that $\psi_{\lambda}=\partial j_{\lambda}$, where (see, e.g., [1])

$$
j_{\lambda}(r)=\inf \left\{\frac{|r-\bar{r}|^{2}}{2 \lambda}+j(\bar{r}) ; \bar{r} \in \mathbb{R}\right\}, \forall r \in \mathbb{R}
$$

We have the following.
Lemma 4.5 Let $x \in \mathcal{H}^{-1} \cap L^{p}\left(\mathbb{R}^{d}\right) \cap L^{2}\left(\mathbb{R}^{d}\right), p:=2 m, d \geq 3$. Then (4.10) has a unique solution

$$
\begin{equation*}
X_{\lambda} \in L^{2}\left(\Omega ; C\left([0, T] ; \mathcal{H}^{-1}\right)\right) \cap L^{\infty}\left([0, T] ; L^{p}\left(\Omega \times \mathbb{R}^{d}\right)\right) \tag{4.11}
\end{equation*}
$$

Moreover, for all $\lambda, \mu>0$, we have

$$
\begin{align*}
& \mathbb{E} \sup _{0 \leq t \leq T}\left\|\left(X_{\lambda}(t)-X_{\mu}(t)\right)\right\|_{-1}^{2} \leq C(\lambda+\mu),  \tag{4.12}\\
& \quad \mathbb{E}\left|X_{\lambda}(t)\right|_{p}^{p} \leq C|x|_{p}^{p}, \forall t \in[0, T],  \tag{4.13}\\
& \mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}}\left|\psi_{\lambda}\left(X_{\lambda}\right)\right|^{\frac{p}{m}} d t d \xi \leq C|x|_{p}^{p}, \forall \lambda>0,  \tag{4.14}\\
& \mathbb{E}\left[\sup _{0 \leq t \leq T}\left\|X_{\lambda}(t)\right\|_{-1}^{2}\right] \leq C\|x\|_{-1}^{2}, \forall \lambda>0, \tag{4.15}
\end{align*}
$$

where $C$ is independent of $\lambda, \mu$.

Proof. We consider for each fixed $\lambda$ the equation (see (3.11))

$$
\begin{align*}
& d X_{\lambda}^{\nu}+(\nu-\Delta)\left(\psi_{\lambda}\left(X_{\lambda}^{\nu}\right)+\lambda X_{\lambda}^{\nu}\right) d t=X_{\lambda}^{\nu} d W  \tag{4.16}\\
& X_{\lambda}^{\nu}(0)=x
\end{align*}
$$

where $\nu>0$. Let $x \in L^{2}\left(\mathbb{R}^{d}\right) \cap L^{p}\left(\mathbb{R}^{d}\right) \cap \mathcal{H}^{-1}$. By Claim 1 in the proof of Lemma 3.5, (4.16) has a unique solution $X_{\lambda}^{\nu} \in L^{2}\left(\Omega ; L^{\infty}\left([0, T] ; L^{2}\left(\mathbb{R}^{d}\right)\right)\right) \cap$ $L^{2}\left(\Omega \times[0, T] ; H^{1}\left(\mathbb{R}^{d}\right)\right)$ with continuous sample paths in $L^{2}\left(\mathbb{R}^{d}\right)$.

As seen in the proof of Theorem 3.4, we have, for $\nu \rightarrow 0$,

$$
\begin{aligned}
X_{\lambda}^{\nu} \rightarrow X_{\lambda} & \text { strongly in } L^{2}\left(\Omega ; C\left([0, T] ; H^{-1}\left(\mathbb{R}^{d}\right)\right)\right) \\
& \text { and, by }(3.14), \text { along a subsequence also, } \\
& \text { weakly* in } L^{2}\left(\Omega ; L^{\infty}\left([0, T] ; L^{2}\left(\mathbb{R}^{d}\right)\right)\right),
\end{aligned}
$$

where $X_{\lambda}$ is the solution to (4.10). It remains to be shown that $X_{\lambda}$ satisfies (4.11)-(4.15). In order to explain the ideas, we apply first (formally) Itô's formula to (4.16) for the function $\varphi(x)=\frac{1}{p}|x|_{p}^{p}$. We obtain

$$
\begin{gather*}
\frac{1}{p} \mathbb{E}\left|X_{\lambda}^{\nu}(t)\right|_{p}^{p}+\mathbb{E} \int_{0}^{t} \int_{\mathbb{R}^{d}}(\nu-\Delta)\left(\psi_{\lambda}\left(X_{\lambda}^{\nu}\right)+\lambda X_{\lambda}^{\nu}\right)\left|X_{\lambda}^{\nu}\right|^{p-2} X_{\lambda}^{\nu} d s d \xi \\
=\frac{1}{p}|x|_{p}^{p}+\frac{p-1}{2} \mathbb{E} \int_{0}^{t} \int_{\mathbb{R}^{d}} \sum_{k=1}^{\infty} \mu_{k}^{2}\left|X_{\lambda}^{\nu} e_{k}\right|^{2}\left|X_{\lambda}^{\nu}\right|^{p-2} d t d \xi \tag{4.17}
\end{gather*}
$$

Taking into account that $X_{\lambda}^{\nu}, \psi_{\lambda}\left(X_{\lambda}^{\nu}\right) \in L^{2}\left(0, T ; H^{1}\left(\mathbb{R}^{d}\right)\right)$, $\mathbb{P}$-a.s., by Claim 1 in the proof of Lemma 3.5, we have

$$
\begin{gathered}
\int_{0}^{t} \int_{\mathbb{R}^{d}}(\nu-\Delta)\left(\psi_{\lambda}\left(X_{\lambda}^{\nu}\right)+\lambda X_{\lambda}^{\nu}\right)\left|X_{\lambda}^{\nu}\right|^{p-2} X_{\lambda}^{\nu} d s d \xi \\
\geq \lambda(p-1) \int_{0}^{t} \int_{\mathbb{R}^{d}}\left|\nabla X_{\lambda}^{\nu}\right|^{2}\left|X_{\lambda}^{\nu}\right|^{p-2} d \xi d s
\end{gathered}
$$

and by (4.2) we have

$$
\mathbb{E} \int_{0}^{t} \int_{\mathbb{R}^{d}} \sum_{k=1}^{\infty} \mu_{k}^{2}\left|X_{\lambda}^{\nu} e_{k}\right|^{2}\left|X_{\lambda}^{\nu}\right|^{p-2} d s d \xi \leq C_{\infty} \mathbb{E} \int_{0}^{t} \int_{\mathbb{R}^{d}}\left|X_{\lambda}^{\nu}\right|^{p} d \xi d s<\infty
$$

Then, we obtain by (4.17) via Gronwall's lemma

$$
\begin{equation*}
\mathbb{E}\left|X_{\lambda}^{\nu}(t)\right|_{p}^{p} \leq C|x|_{p}^{p}, t \in(0, T) \tag{4.18}
\end{equation*}
$$

and, by (4.1),

$$
\begin{equation*}
\mathbb{E} \int_{0}^{t} \int_{\mathbb{R}^{d}}\left|\psi_{\lambda}\left(X_{\lambda}^{\nu}\right)\right|^{\frac{p}{m}} d t d \xi \leq C|x|_{p}^{p}, t \in[0, T] . \tag{4.19}
\end{equation*}
$$

It should be said, however, that the above argument is formal, because the function $\varphi$ is not of class $C^{2}$ on $L^{2}\left(\mathbb{R}^{d}\right)$ and we do not know a priori if the integral in the left side of (4.17) makes sense, that is, whether $\left|X_{\lambda}^{\nu}\right|^{p-2} X_{\lambda}^{\nu} \in L^{2}\left(0, T ; L^{2}\left(\Omega ; H^{1}\left(\mathbb{R}^{d}\right)\right)\right)$. To make it rigorous, we approximate $X_{\lambda}^{\nu}$ by a sequence $\left\{X_{\lambda}^{\nu, \varepsilon}\right\}$ of solutions to the equation

$$
\begin{align*}
& d X_{\lambda}^{\nu, \varepsilon}+A_{\lambda}^{\nu, \varepsilon}\left(X_{\lambda}^{\nu, \varepsilon}\right) d t=X_{\lambda}^{\nu, \varepsilon} d W,  \tag{4.20}\\
& X_{\lambda}^{\nu, \varepsilon}(0)=x .
\end{align*}
$$

Here, $A_{\lambda}^{\nu, \varepsilon}=\frac{1}{\varepsilon}\left(I-\left(I+\varepsilon A_{\lambda}^{\nu}\right)^{-1}\right), \varepsilon \in(0,1)$, is the Yosida approximation of the operator $A_{\lambda}^{\nu} x=(\nu-\Delta)\left(\psi_{\lambda}(x)+\lambda x\right), \forall x \in D\left(A_{\lambda}^{\nu}\right)=H^{1}\left(\mathbb{R}^{d}\right)$. We set $J_{\varepsilon}=\left(I+\varepsilon A_{\lambda}^{\nu}\right)^{-1}$ and note that $J_{\varepsilon}$ is Lipschitz in $H=H^{-1}\left(\mathbb{R}^{d}\right)$ as well as in all $L^{q}\left(\mathbb{R}^{d}\right)$ for $1<q<\infty$. Moreover, we have

$$
\begin{equation*}
\left|J_{\varepsilon}(x)\right|_{q} \leq|x|_{q}, \forall x \in L^{q}\left(\mathbb{R}^{d}\right) \tag{4.21}
\end{equation*}
$$

(See [3], Lemma 3.1.) Since $A_{\lambda}^{\nu, \varepsilon}$ is Lipschitz in $H$, equation (3.1) has a unique adapted solution $X_{\lambda}^{\nu, \varepsilon} \in L^{2}(\Omega ; C([0, T] ; H)$ and by Itô's formula we have

$$
\frac{1}{2} \mathbb{E}\left|X_{\lambda}^{\nu, \varepsilon}(t)\right|_{-1}^{2} \leq \frac{1}{2}|x|_{-1}^{2}+C_{1} \sum_{k=1}^{\infty} \mu_{k}^{2} \mathbb{E} \int_{0}^{t}\left|X_{\lambda}^{\nu, \varepsilon}(s) e_{k}\right|_{-1}^{2} d s,
$$

which, by virtue of ( jj ), yields

$$
\begin{equation*}
\mathbb{E}\left|X_{\lambda}^{\nu, \varepsilon}(t)\right|_{-1}^{2} \leq C_{2}|x|_{-1}^{2}, \forall \varepsilon>0, x \in H \tag{4.22}
\end{equation*}
$$

Similarly, since $A_{\lambda}^{\nu, \varepsilon}$ is Lipschitz in $L^{2}\left(\mathbb{R}^{d}\right)$ (see Lemma 4.6 below), we have also that $X_{\lambda}^{\nu, \varepsilon} \in L^{2}\left(\Omega ; C\left([0, T] ; L^{2}\left(\mathbb{R}^{2}\right)\right)\right)$ and, again by Itô's formula applied to the function $\left|X_{\lambda}^{\nu, \varepsilon}(t)\right|_{2}^{2}$, we obtain that

$$
\mathbb{E}\left|X_{\lambda}^{\nu, \varepsilon}(t)\right|_{2}^{2} \leq \frac{1}{2}|x|_{2}^{2}+C_{3} \sum_{k=1}^{\infty} \mu_{k}^{2} \mathbb{E} \int_{0}^{\infty}\left|X_{\lambda}^{\nu, \varepsilon}(s) e_{k}\right|_{2}^{2} d s
$$

which yields, by virtue of ( jj ),

$$
\begin{equation*}
\mathbb{E}\left|X_{\lambda}^{\nu, \varepsilon}(t)\right|_{2}^{2} \leq C_{4}|x|_{2}^{2}, \quad \forall t \in[0, T] . \tag{4.23}
\end{equation*}
$$

Claim 1. For $p \in[2, \infty)$ and $x \in L^{p}\left(\mathbb{R}^{d}\right)$, we have that $X_{\lambda}^{\nu, \varepsilon} \in L_{W}^{\infty}([0, T]$; $L^{p}\left(\Omega ; L^{p}\left(\mathbb{R}^{d}\right)\right) \cap L^{2}\left(\Omega ; L^{2}\left(\mathbb{R}^{d}\right)\right)$ ), where here and below the subscript $W$ refers to $\left(\mathcal{F}_{t}\right)$-adapted processes.

Proof. For $R>0$, consider the set

$$
\begin{aligned}
\mathcal{K}_{R}= & \left\{X \in L_{W}^{\infty}\left([0, T] ; L^{p}\left(\Omega ; L^{p}\left(\mathbb{R}^{d}\right)\right) \cap L^{2}\left(\Omega ; L^{2}\left(\mathbb{R}^{d}\right)\right)\right),\right. \\
& \left.e^{-p \alpha t} \mathbb{E}|X(t)|_{p}^{p} \leq R^{p}, e^{-2 \alpha t} \mathbb{E}|X(t)|_{2}^{2} \leq R^{2}, t \in[0, T]\right\} .
\end{aligned}
$$

Since, by (4.20), $X_{\lambda}^{\nu, \varepsilon}$ is a fixed point of the map

$$
F: X \mapsto e^{-\frac{t}{\varepsilon}} X+\frac{1}{\varepsilon} \int_{0}^{t} e^{-\frac{t-s}{\varepsilon}} J_{\varepsilon}(X(s)) d s+\int_{0}^{t} e^{-\frac{(t-s)}{\varepsilon}} X(s) d W(s)
$$

obtained by iteration in $C_{W}\left([0, T] ; L^{2}\left(\Omega ; L^{2}\left(\mathbb{R}^{d}\right)\right)\right)$, it suffices to show that $F$ leaves the set $\mathcal{K}_{R}$ invariant for $R>0$ large enough. By (4.21), we have

$$
\begin{align*}
& \left(e^{-p \alpha t} \mathbb{E}\left|e^{-\frac{t}{\varepsilon}} x+\frac{1}{\varepsilon} \int_{0}^{t} e^{-\frac{t-s}{\varepsilon}} J_{\varepsilon}(X(s)) d s\right|_{p}^{p}\right)^{\frac{1}{p}} \\
& \quad \leq e^{-\left(\frac{1}{\varepsilon}+\alpha\right) t}|x|_{p}+e^{-\alpha t} \int_{0}^{t} \frac{1}{\varepsilon} e^{-\frac{(t-s)}{\varepsilon}}\left(\mathbb{E}|X(s)|_{p}^{p}\right)^{\frac{1}{p}} d s  \tag{4.24}\\
& \quad \leq e^{-\left(\frac{1}{\varepsilon}+\alpha\right) t}|x|_{p}+\frac{R}{1+\alpha \varepsilon},
\end{align*}
$$

and, similarly, that

$$
\begin{align*}
& \left(e^{-2 \alpha t} \mathbb{E}\left|e^{-\frac{t}{\varepsilon}} x+\frac{1}{\varepsilon} \int_{0}^{t} e^{-\frac{(t-s)}{\varepsilon}} J_{\varepsilon}(X(s)) d s\right|_{2}^{2}\right)^{\frac{1}{2}}  \tag{4.25}\\
& \leq e^{-\left(\frac{1}{\varepsilon}+\alpha\right) t}|x|_{2}+\frac{R}{1+\alpha \varepsilon} .
\end{align*}
$$

Now, we set

$$
Y(t)=\int_{0}^{t} e^{-\frac{(t-s)}{\varepsilon}} X(s) d W(s), t \geq 0
$$

We have

$$
\begin{aligned}
& d Y+\frac{1}{\varepsilon} Y d t=X d W, \quad t \geq 0 \\
& Y(0)=0
\end{aligned}
$$

Equivalently,

$$
d\left(e^{\frac{t}{\varepsilon}} Y(t)\right)=e^{\frac{t}{\varepsilon}} X(t) d W(t), t>0 ; Y(0)=0 .
$$

By Lemma 5.1 in [19], it follows that $e^{\frac{t}{\varepsilon}} Y$ is an $L^{p}\left(\mathbb{R}^{d}\right)$-valued $\left(\mathcal{F}_{t}\right)$-adapted continuous process on $[0, \infty)$ and

$$
\mathbb{E}\left|e^{\frac{t}{\varepsilon}} Y(t)\right|_{p}^{p}=\frac{1}{2} p(p-1) \sum_{k=1}^{\infty} \mu_{k}^{2} \mathbb{E} \int_{0}^{t} \int_{\mathbb{R}^{d}}\left|e^{\frac{s}{\varepsilon}} Y(s)\right|^{p-2}\left|e^{\frac{s}{\varepsilon}} X(s) e_{k}\right|^{2} d s
$$

This yields via Hypothesis (jj)

$$
\mathbb{E}\left|e^{\frac{t}{\varepsilon}} Y(t)\right|_{p}^{p} \leq \frac{1}{2}(p-1) \mathbb{E} \int_{0}^{t}\left|e^{\frac{s}{\varepsilon}} Y(s)\right|_{p}^{p} d s+\left.\left.C \mathbb{E} \int_{0}^{t}\right|^{\frac{s}{\varepsilon}} X(s)\right|_{p} ^{p} d s, \quad \forall t \in[0, T],
$$

and, therefore,

$$
\mathbb{E}|Y(t)|_{p}^{p} \leq C_{1} e^{-\left(\alpha+\frac{1}{\varepsilon}\right) p t} \mathbb{E} \int_{0}^{t}\left|e^{\frac{s}{\varepsilon}} X(s)\right|_{p}^{p} d s \leq \frac{R^{p} e^{-p \alpha t} \varepsilon C_{1}}{p(1+\varepsilon \alpha)}, \quad \forall t \in[0, T] .
$$

Similarly, we get

$$
e^{-2 \alpha t} \mathbb{E}|Y(t)|_{2}^{2} \leq \frac{R^{2} \varepsilon C_{1}}{2(1+\varepsilon \alpha)}, \quad \forall t \in[0, T]
$$

Then, by formulae (4.24), (4.25), we infer that, for $\alpha$ large enough and $R>2\left(|x|_{p}+|x|_{2}\right), F$ leaves $\mathcal{K}_{R}$ invariant, which proves Claim 1.

Claim 2. We have, for all $p \in[2, \infty)$ and $x \in L^{p}\left(\mathbb{R}^{d}\right)$, that there exists $C_{p} \in(0, \infty)$ such that

$$
\begin{equation*}
\underset{t \in[0, T]}{\operatorname{ess} \sup } \mathbb{E}\left|X_{\lambda}^{\nu, \varepsilon}(t)\right|_{p}^{p} \leq C_{p} \text { for all } \varepsilon, \lambda, \nu \in(0,1) \tag{4.26}
\end{equation*}
$$

Proof. Again invoking Lemma 5.1 in [19], we have by (4.20) that $X_{\lambda}^{\nu, \varepsilon}$ satisfies

$$
\begin{align*}
& \mathbb{E}\left|X_{\lambda}^{\nu, \varepsilon}(t)\right|_{p}^{p}=|x|_{p}^{p}-p \mathbb{E} \int_{0}^{t} \int_{\mathbb{R}^{d}} A_{\lambda}^{\nu, \varepsilon}\left(X_{\lambda}^{\nu, \varepsilon}\right) X_{\lambda}^{\nu, \varepsilon}\left|X_{\lambda}^{\nu, \varepsilon}\right|^{p-2} d \xi d s \\
& +p(p-1) \sum_{k=1}^{\infty} \mu_{k}^{2} \mathbb{E} \int_{0}^{t} \int_{\mathbb{R}^{d}}\left|X_{\lambda}^{\nu, \varepsilon}\right|^{p-2}\left|X_{\lambda}^{\nu, \varepsilon} e_{k}\right|^{2} d \xi d s . \tag{4.27}
\end{align*}
$$

On the other hand, $A_{\lambda}^{\nu, \varepsilon}\left(X_{\lambda}^{\nu, \varepsilon}\right)=\frac{1}{\varepsilon}\left(X_{\lambda}^{\nu, \varepsilon}-J_{\varepsilon}\left(X_{\lambda}^{\nu, \varepsilon}\right)\right)$ and so we have

$$
\begin{aligned}
\int_{\mathbb{R}^{d}} & A_{\lambda}^{\nu, \varepsilon}\left(X_{\lambda}^{\nu, \varepsilon}\right) X_{\lambda}^{\nu, \varepsilon}\left|X_{\lambda}^{\nu, \varepsilon}\right|^{p-2} d \xi \\
& =\frac{1}{\varepsilon} \int_{\mathbb{R}^{d}}\left|X_{\lambda}^{\nu, \varepsilon}\right|^{p} d \xi-\frac{1}{\varepsilon} \int_{\mathbb{R}^{r}} J_{\varepsilon}\left(X_{\lambda}^{\nu, \varepsilon}\right)\left|X_{\lambda}^{\nu, \varepsilon}\right|^{p-2} X_{\lambda}^{\nu, \varepsilon} d \xi
\end{aligned}
$$

Recalling (4.21), we get, via the Hölder inequality,

$$
\int_{\mathbb{R}^{d}} A_{\lambda}^{\nu, \varepsilon}\left(X_{\lambda}^{\nu, \varepsilon}\right) X_{\lambda}^{\nu, \varepsilon}\left|X_{\lambda}^{\nu, \varepsilon}\right|^{p-2} d \xi \geq 0
$$

and so, by (4.27) and Hypothesis (jj), we obtain, via Gronwall's lemma, estimate (4.26), as claimed.
Claim 3. We have, for $\varepsilon \rightarrow 0$,

$$
X_{\lambda}^{\nu, \varepsilon} \longrightarrow X_{\lambda}^{\nu} \text { strongly in } L_{W}^{\infty}\left([0, T] ; L^{2}(\Omega ; H)\right)
$$

and weakly* in $L^{\infty}\left([0, T] ; L^{p}\left(\Omega ; L^{p}\left(\mathbb{R}^{d}\right)\right) \cap L^{2}\left(\Omega ; L^{2}\left(\mathbb{R}^{d}\right)\right)\right)$.
Proof. For simplicity, we write $X_{\varepsilon}$ instead of $X_{\lambda}^{\nu, \varepsilon}$ and $X$ instead of $X_{\lambda}^{\nu}$. Also, we set $\gamma(r) \equiv \psi_{\lambda}(r)+\lambda r$.

Subtracting equations (4.20) and (4.16), we get via Itô's formula and because $A_{\lambda}^{\nu, \varepsilon}$ is monotone on $H$

$$
\begin{aligned}
\frac{1}{2} \mathbb{E}\left|X_{\varepsilon}(t)-X(t)\right|_{-1, \nu}^{2}+\mathbb{E} \int_{0}^{t} \int_{\mathbb{R}^{d}}( & \left.\left(J_{\varepsilon}(X)\right)-\gamma(X)\right)\left(X_{\varepsilon}-X\right) d \xi d s \\
& \leq C \mathbb{E} \int_{0}^{t}\left|X_{\varepsilon}(s)-X(s)\right|_{-1, \nu}^{2} d s
\end{aligned}
$$

and hence, by Gronwall's lemma, we obtain

$$
\begin{equation*}
\mathbb{E}\left|X_{\varepsilon}(t)-X(t)\right|_{-1, \nu}^{2} \leq C \mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}}\left|\gamma\left(J_{\varepsilon}(X)\right)-\gamma(X)\right|\left|X_{\varepsilon}-X\right| d \xi d s \tag{4.28}
\end{equation*}
$$

On the other hand, it follows by (4.21) that

$$
\int_{\Omega \times[0, T] \times \mathbb{R}^{d}}\left|J_{\varepsilon}(X)\right|^{2} \mathbb{P}(d \omega) d t d \xi \leq \int_{\Omega \times[0, T] \times \mathbb{R}^{d}}|X|^{2} \mathbb{P}(d \omega) d t d \xi,
$$

while, for $\varepsilon \rightarrow 0$,

$$
J_{\varepsilon}(y) \longrightarrow y \text { in } H^{-1}, \quad \forall y \in H^{-1}
$$

(because $A_{\lambda}^{\nu, \varepsilon}$ is maximal monotone in $H^{-1}\left(\mathbb{R}^{d}\right)$ ) and so, $J_{\varepsilon}(X(t, \omega)) \longrightarrow$ $X(t, \omega)$ in $H^{-1}\left(\mathbb{R}^{d}\right)$ for all $(t, \omega) \in(0, T) \times \Omega$. Hence, as $\varepsilon \rightarrow 0$,

$$
\begin{equation*}
J_{\varepsilon}(X) \longrightarrow X \text { weakly in } L^{2}\left(\Omega \times[0, T] \times \mathbb{R}^{d}\right), \tag{4.29}
\end{equation*}
$$

and, according to the inequality above, this implies that, for $\varepsilon \rightarrow 0$,

$$
\left|J_{\varepsilon}(X)\right|_{L^{2}\left((0, T) \times \Omega \times \mathbb{R}^{d}\right)} \longrightarrow|X|_{L^{2}\left((0, T) \times \Omega \times \mathbb{R}^{d}\right)} .
$$

Hence, $J_{\varepsilon}(X) \longrightarrow X$ strongly in $L^{2}\left(\Omega \times[0, T] \times \mathbb{R}^{d}\right)$ as $\varepsilon \rightarrow 0$. Now, taking into account that $\gamma$ is Lipschitz, we conclude by (4.28), (4.29) and by estimates (4.23), (4.26) that Claim 3 is true.

Now, we can complete the proof of Lemma 4.5. Namely, letting first $\varepsilon \rightarrow 0$ and then $\nu \rightarrow \infty$ in (4.26), we get (4.13) and hence (4.14) as desired.

Now, let us prove (4.12) and (4.15). Arguing as in the proof of Theorem 3.4, we obtain

$$
\begin{align*}
& \frac{1}{2}\left|X_{\lambda}^{\nu}(t)\right|_{-1, \nu}^{2}+\int_{0}^{t} \int_{\mathbb{R}^{d}}\left(\psi_{\lambda}\left(X_{\lambda}^{\nu}\right)+\lambda X_{\lambda}^{\nu}\right) X_{\lambda}^{\nu} d \xi d s \\
& \quad=\frac{1}{2}|x|_{-1, \nu}^{2}+\frac{1}{2} \int_{0}^{t} \int_{\mathbb{R}^{d}} \sum_{k=1}^{\infty} \mu_{k}^{2}\left|X_{\lambda}^{\nu} e_{k}\right|_{-1, \nu}^{2} d \xi d s  \tag{4.30}\\
& \quad+\int_{0}^{t}\left\langle X_{\lambda}^{\nu}, X_{\lambda}^{\nu} d W\right\rangle_{-1, \nu} d s
\end{align*}
$$

Keeping in mind that, by (4.3), $\left|X_{\lambda}^{\nu} e_{k}\right|_{-1, \nu} \leq C\left|X_{\lambda}^{\nu}\right|_{-1, \nu}\left(\left|e_{k}\right|_{\infty}+\left|\nabla e_{k}\right|_{d}\right)$, where $C$ is independent of $\nu$, we obtain by the Burkholder-Davis-Gundy inequality for $p=1$ (cf. the proof of Theorem 3.4)

$$
\mathbb{E} \sup _{t \in[0, T]}\left|X_{\lambda}^{\nu}(t)\right|_{-1, \nu}^{2}+\lambda \mathbb{E} \int_{0}^{T}\left|X_{\lambda}^{\nu}\right|_{2}^{2} d s \leq C|x|_{-1, \nu}^{2}
$$

Taking into account that

$$
\lim _{\nu \rightarrow 0}|y|_{-1, \nu}=\|y\|_{-1}, \forall y \in \mathcal{H}^{-1}
$$

we obtain, as in Theorem 3.4 (see the part following (3.26)), that

$$
\begin{equation*}
\mathbb{E}\left[\sup _{t \in[0, T]}\left\|X_{\lambda}(t)\right\|_{-1}^{2}\right]+\lambda \mathbb{E} \int_{0}^{T}\left|X_{\lambda}(t)\right|_{2}^{2} d t \leq C\|x\|_{-1}^{2}, \forall \lambda>0, \tag{4.31}
\end{equation*}
$$

where $C$ is independent of $\lambda$. In particular, (4.15) holds.
Completely similarly, one proves (4.12). Namely, we have
$d\left(X_{\lambda}^{\nu}-X_{\mu}^{\nu}\right)+(\nu-\Delta)\left(\psi_{\lambda}\left(X_{\lambda}^{\nu}\right)+\lambda X_{\lambda}^{\nu}-\psi_{\mu}\left(X_{\mu}^{\nu}\right)-\mu X_{\mu}^{\nu}\right) d t=\left(X_{\lambda}^{\nu}-X_{\mu}^{\nu}\right) d W$
and again proceeding as in the proof of Theorem 3.4, we obtain as above that

$$
\begin{aligned}
& \frac{1}{2}\left|X_{\lambda}^{\nu}(t)-X_{\mu}^{\nu}(t)\right|_{-1, \nu}^{2} \\
& \quad+\int_{0}^{t} \int_{\mathbb{R}^{d}}\left(\psi_{\lambda}\left(X_{\lambda}^{\nu}\right)+\lambda X_{\lambda}^{\nu}-\psi_{\mu}\left(X_{\mu}^{\nu}\right)-\mu X_{\mu}^{\nu}\right)\left(X_{\lambda}^{\nu}-X_{\mu}^{\nu}\right) d \xi d s \\
& \quad=\frac{1}{2} \int_{0}^{t} \int_{\mathbb{R}^{d}} \sum_{k=1}^{\infty} \mu_{k}^{2}\left|\left(X_{\lambda}^{\nu}-X_{\mu}^{\nu}\right) e_{k}\right|_{-1, \nu}^{2} d s \\
& \quad+\int_{0}^{t}\left\langle X_{\lambda}^{\nu}-X_{\mu}^{\nu},\left(X_{\lambda}^{\nu}-X_{\mu}^{\nu}\right) d W\right\rangle_{-1, \nu}, t \in[0, T] .
\end{aligned}
$$

Then, applying once again the Burkholder-Davis-Gundy inequality for $p=1$, and the fact that, by Hypothesis $(\mathrm{j}),\left|\psi_{\lambda}(r)\right| \leq C|r|^{m}, \forall \in \mathbb{R}$ with $C$ independent of $\lambda$, we get, proceeding as in the proof of Theorem 3.4, that

$$
\mathbb{E}\left[\sup _{t \in[0, T]}\left|X_{\lambda}^{\nu}(t)-X_{\mu}^{\nu}(t)\right|_{-1}^{2}\right] \leq C(\lambda+\mu),
$$

where $C$ is independent of $\nu, \lambda, \mu$. (For details, we refer to the proof of (3.10), (3.14) in [6].) Letting $\nu \rightarrow 0$ as in the previous case, we obtain (4.12), as claimed. This completes the proof of Lemma 4.5.

Above we have used the following lemma.
Lemma 4.6 $A_{\lambda}^{\nu, \varepsilon}$ is Lipschitz in $L^{2}\left(\mathbb{R}^{d}\right)$.
Proof. It suffices to check that $J_{\varepsilon}$ is Lipschitz in $L^{2}\left(\mathbb{R}^{d}\right)$. We set $\gamma(r)=$ $\psi_{\lambda}(r)+\lambda r$. We have, for $x, \bar{x} \in L^{2}\left(\mathbb{R}^{d}\right)$,

$$
J_{\varepsilon}(x)-J_{\varepsilon}(\bar{x})-\varepsilon \Delta\left(\gamma\left(J_{\varepsilon}(x)\right)-\gamma\left(J_{\varepsilon}(\bar{x})\right)\right)=x-\bar{x} .
$$

Multiplying by $\gamma\left(J_{\varepsilon}(x)\right)-\gamma\left(J_{\varepsilon}(\bar{x})\right)$ in $L^{2}\left(\mathbb{R}^{d}\right)$, we get

$$
\left\langle J_{\varepsilon}(x)-J_{\varepsilon}(\bar{x}), \gamma\left(J_{\varepsilon}(x)\right)-\gamma\left(J_{\varepsilon}(\bar{x})\right)\right\rangle_{2} \leq\left|\gamma\left(J_{\varepsilon}(x)\right)-\gamma\left(J_{\varepsilon}(\bar{x})\right)\right|_{2}|x-\bar{x}|_{2} .
$$

Taking into account that $(\gamma(r)-\gamma(\bar{r}))(r-\bar{r}) \geq L|r-\bar{r}|, \forall r, \bar{r} \in \mathbb{R}$, and that $\gamma$ is Lipschitz, we get

$$
\left|J_{\varepsilon}(x)-J_{\varepsilon}(\bar{x})\right|_{2} \leq C|x-\bar{x}|_{2},
$$

as claimed.
Proof of Theorem 4.4 (continued). By (4.12)-(4.15), it follows that there is a process $X \in L^{\infty}\left([0, T] ; L^{p}\left(\Omega \times \mathbb{R}^{d}\right)\right)$ such that, for $\lambda \rightarrow 0$,

$$
\begin{array}{rll}
X_{\lambda} & \rightarrow X & \text { weak-star in } L^{\infty}\left([0, T] ; L^{p}\left(\Omega \times \mathbb{R}^{d}\right)\right) \\
\lambda X_{\lambda} & \rightarrow 0 & \text { strongly in } L^{2}\left([0, T] ; L^{2}\left(\Omega \times \mathbb{R}^{d}\right)\right) \\
\psi_{\lambda}\left(X_{\lambda}\right) & \rightarrow \eta & \text { weakly in } L^{\frac{p}{m}}\left([0, T] \times \Omega \times \mathbb{R}^{d}\right)  \tag{4.32}\\
X_{\lambda} & \rightarrow X & \text { strongly in } L^{2}\left(\Omega ; C\left([0, T] ; \mathcal{H}^{-1}\right)\right) .
\end{array}
$$

It remains to be shown that $X$ is a strong solution to (1.1) in the sense of Definition 4.3.

By (4.10) and (4.32), we see that

$$
\begin{align*}
& d X-\Delta \eta d t=X d W, t \in(0, T)  \tag{4.33}\\
& X(0)=x
\end{align*}
$$

To prove that $\eta \in \psi(X)$, a.e. in $\Omega \times(0, T) \times \mathbb{R}^{d}$, it suffices to show that, for each $\varphi \in C_{0}^{\infty}\left(\mathbb{R}^{d}\right)$, we have

$$
\begin{equation*}
\limsup _{\lambda \rightarrow 0} \mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}} \varphi^{2} \psi_{\lambda}\left(X_{\lambda}\right) X_{\lambda} d t d \xi \leq \mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}} \varphi^{2} \eta X d \xi d t \tag{4.34}
\end{equation*}
$$

Indeed, we have by convexity of $j_{\lambda}$

$$
\begin{array}{r}
\mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}} \varphi^{2} \psi_{\lambda}\left(X_{\lambda}\right)\left(X_{\lambda}-Z\right) d \xi d t \geq \mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}} \varphi^{2}\left(j_{\lambda}\left(X_{\lambda}\right)-j_{\lambda}(Z)\right) d \xi d t, \\
\forall Z \in L^{p}\left((0, T) \times \Omega \times \mathbb{R}^{d}\right),
\end{array}
$$

and so, by (4.32) and (4.34), we see that

$$
\begin{array}{r}
\mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}} \varphi^{2}(\eta(X-Z)) d t d \xi \geq \mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}} \varphi^{2}(j(X)-j(Z)) d \xi d t, \\
\forall Z \in L^{p}\left((0, T) \times \Omega \times \mathbb{R}^{d}\right),
\end{array}
$$

because, for $\lambda \rightarrow 0, j_{\lambda}(Z) \rightarrow j(Z)$, and $j_{\lambda}\left(X_{\lambda}\right) \rightarrow j(X)$, a.e. and thus, by Fatou's lemma

$$
\liminf _{\lambda \rightarrow 0} \mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}} \varphi^{2} j_{\lambda}\left(X_{\lambda}\right) d \xi d t \geq \mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}} \varphi^{2} j(X) d \xi d t
$$

Now, we take $\varphi \in C_{0}^{\infty}\left(\mathbb{R}^{d}\right)$ to be non-negative, such that $\varphi=1$ on $B_{N}$ and $\varphi=0$, outside $B_{N+1}$ where for a given $N \in \mathbb{N}, B_{N}$ is the closed ball of $\mathbb{R}^{d}$ with radius $N$. We get

$$
\begin{array}{r}
\mathbb{E} \int_{0}^{T} \int_{B_{N+1}} \varphi^{2}(\eta(X-Z)) d \xi d t \geq \mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}} \varphi^{2}(j(X)-j(Z)) d \xi d t,  \tag{4.35}\\
\forall Z \in L^{p}\left((0, T) \times \Omega \times \mathbb{R}^{d}\right) .
\end{array}
$$

This yields

$$
\begin{equation*}
\mathbb{E} \int_{0}^{T} \int_{B_{N+1}} \varphi^{2} \eta(X-Z) d \xi d t \geq \mathbb{E} \int_{0}^{T} \int_{B_{N+1}} \varphi^{2} \zeta(X-Z) d \xi d t \tag{4.36}
\end{equation*}
$$

for all $Z \in L^{p}\left((0, T) \times \Omega \times B_{N+1}\right)$ and $\zeta \in L^{p^{\prime}}\left((0, T) \times \Omega \times B_{N+1}\right)$ such that $\zeta \in \psi(Z)$, a.e. in $(0, T) \times \Omega \times B_{N+1}$.

We denote by $\widetilde{\psi}: L^{p}\left((0, T) \times \Omega \times B_{N+1}\right) \rightarrow L^{p^{\prime}}\left((0, T) \times \Omega \times B_{N+1}\right)$ the realization of the mapping $\psi$ in $L^{p}\left((0, T) \times \Omega \times B_{N+1}\right)$, that is,

$$
\widetilde{\psi}(Z)=\left\{\zeta \in L^{p^{\prime}}\left((0, T) \times \Omega \times B_{N+1}\right), \zeta \in \psi(Z), \text { a.e. }\right\} .
$$

Since $\frac{m}{p} \leq p^{\prime}$ with $\frac{1}{p^{\prime}}=1-\frac{1}{p}$, by virtue of assumption ( j ), $\widetilde{\psi}$ is maximal monotone in $L^{p}\left((0, T) \times \Omega \times B_{N+1}\right) \times L^{p^{\prime}}\left((0, T) \times \Omega \times B_{N+1}\right)$, and so, the equation

$$
\begin{equation*}
J(Z)+\widetilde{\psi}(Z) \ni J(X)+\eta \tag{4.37}
\end{equation*}
$$

where $J(Z)=|Z|^{p-2} Z$, has a unique solution $(Z, \eta)$ (see, e.g., [1], p. 31).
If, in (4.36), we take $Z$ the solution to (4.37), we obtain that

$$
\mathbb{E} \int_{0}^{T} \int_{B_{N+1}} \varphi^{2}(J(X)-J(Z))(X-Z) d t d \xi \leq 0
$$

Then, choosing $\alpha=\frac{2}{p}$, yields

$$
\mathbb{E} \int_{0}^{T} \int_{B_{N+1}}\left(\left|\varphi^{\alpha} X\right|^{p-2} \varphi^{\alpha} X-\left|\varphi^{\alpha} Z\right|^{p-2} \varphi^{\alpha} Z\right)\left(\varphi^{\alpha} X-\varphi^{\alpha} Z\right) d t d \xi \leq 0
$$

Consequently, this gives

$$
\begin{equation*}
\mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}}\left(J\left(\varphi^{\alpha} X\right)-J\left(\varphi^{\alpha} Z\right)\right)\left(\varphi^{\alpha} X-\varphi^{\alpha} Z\right) d t d \xi \leq 0 \tag{4.38}
\end{equation*}
$$

On the other hand, we have

$$
J\left(\varphi^{\alpha} X\right)-J\left(\varphi^{\alpha} Z\right)=(p-1)\left|\lambda \varphi^{\alpha} X+(1-\lambda) \varphi^{\alpha} Z\right|^{p-2}(X-Z)
$$

for some $\lambda=\lambda(X, Z) \in[0,1]$. Substituting into (4.38) yields

$$
\left|\varphi^{\alpha}(X-Z)\right|^{2}=0 \text { a.e. in }(0, T) \times \Omega \times B_{N+1},
$$

Hence, $X=Z$ on $(0, T) \times \Omega \times B_{N}$.
Coming back to (4.37), this gives $\eta \in \psi(X), d t d P d \xi$, a.e., because $N$ is arbitrary.

To prove (4.34), we use the Itô formula in (4.16) to $x \mapsto \frac{1}{2}\|\varphi x\|_{-1}^{2}$ to get, as in (4.30),

$$
\begin{aligned}
& \frac{1}{2} \mathbb{E}\left\|\varphi X_{\lambda}^{\nu}(t)\right\|_{-1}^{2}+\mathbb{E} \int_{0}^{t}\left\langle(-\Delta)^{-1}(\nu-\Delta)\left(\psi_{\lambda}\left(X_{\lambda}^{\nu}\right)+\lambda X_{\lambda}^{\nu}, \varphi^{2} X_{\lambda}^{\nu}\right)\right\rangle d s \\
& \leq \frac{1}{2}\|\varphi x\|_{-1}^{2}+\frac{1}{2} \mathbb{E} \int_{0}^{t} \sum_{k=1}^{\infty} \mu_{k}^{2}\left\|\varphi X_{\lambda}^{\nu} e_{k}\right\|_{-1}^{2} d s
\end{aligned}
$$

Then, letting $\nu \rightarrow 0$, we obtain

$$
\begin{gather*}
\frac{1}{2} \mathbb{E}\left\|\varphi X_{\lambda}(t)\right\|_{-1}^{2}+\mathbb{E} \int_{0}^{t}\left\langle\psi_{\lambda}\left(X_{\lambda}\right)+\lambda X_{\lambda}, \varphi^{2} X_{\lambda}\right\rangle_{2} d s \\
\leq \frac{1}{2}\|\varphi x\|_{-1}^{2}+\frac{1}{2} \mathbb{E} \int_{0}^{t} \sum_{k=1}^{\infty} \mu_{k}^{2}\left\|\varphi X_{\lambda} e_{k}\right\|_{-1}^{2} d s \tag{4.39}
\end{gather*}
$$

On the other hand, by (4.33) we get similarly

$$
\begin{aligned}
& \frac{1}{2} \mathbb{E}\|\varphi X(t)\|_{-1}^{2}+\mathbb{E} \int_{0}^{t}\left\langle\eta(s), \varphi^{2} X\right\rangle_{2} d s \\
& \quad=\frac{1}{2}\|\varphi x\|_{-1}^{2}+\frac{1}{2} \mathbb{E} \int_{0}^{1} \sum_{k=1}^{\infty} \mu_{k}^{2}\left\|\varphi X e_{k}\right\|_{-1}^{2}, t \in[0, T]
\end{aligned}
$$

Comparing with (4.39), we obtain (4.34), as claimed.
If $x \geq 0$, a.e. in $\mathbb{R}^{d}$, it follows that $X \geq 0$, a.e. in in $\Omega \times(0, T) \times \mathbb{R}^{d}$. To prove this, one applies Itô's formula in (4.10) to the function $x \mapsto\left|x^{-}\right|_{2}^{2}$ and get $\left(X_{\lambda}^{\nu}\right)^{-}=0$, a.e. in $\Omega \times(0, T) \times \mathbb{R}^{d}$. Then, for $\nu \rightarrow 0$, we obtain the desired result. This completes the existence proof for $x \in L^{2}\left(\mathbb{R}^{d}\right) \cap L^{p}\left(\mathbb{R}^{d}\right) \cap \mathcal{H}^{-1}$.

Uniqueness. If $X_{1}, X_{2}$ are two solutions, we have

$$
\begin{aligned}
& d\left(X_{1}-X_{2}\right)-\Delta\left(\eta_{1}-\eta_{2}\right) d t=\left(X_{1}-X_{2}\right) d W, t \in(0, T), \\
& \left(X_{1}-X_{2}\right)(0)=0
\end{aligned}
$$

where $\eta_{i} \in \psi\left(X_{i}\right), i=1,2$, a.e. in $\Omega \times(0, T) \times \mathbb{R}^{d}$.
Applying again, as above (that is, via the approximating device) Itô's formula in $\mathcal{H}^{-1}$ to $\frac{1}{2}\left\|\varphi\left(X_{1}-X_{2}\right)\right\|_{-1}^{2}$, where $\varphi \in C_{0}^{\infty}\left(\mathbb{R}^{d}\right)$, we get that

$$
\begin{aligned}
& \frac{1}{2} d\left\|\varphi\left(X_{1}-X_{2}\right)\right\|_{-1}^{2}-\left\langle\Delta\left(\eta_{1}-\eta_{2}\right), \varphi\left(X_{1}-X_{2}\right)\right\rangle_{-1} \\
& =\frac{1}{2} \sum_{k=1}^{\infty} \mu_{k}^{2}\left\|\varphi\left(X_{1}-X_{2}\right) e_{k}\right\|_{-1}^{2} d t+\left\langle\left(X_{1}-X_{2}\right), \varphi\left(X_{1}-X_{2}\right) d W\right\rangle_{-1}=0 .
\end{aligned}
$$

Note that, since $\eta_{1}-\eta_{2} \in L^{\frac{p}{m}}\left(\Omega \times(0, T) \times \mathbb{R}^{d}\right)$, we have

$$
\begin{aligned}
& -\mathbb{E} \int_{0}^{T}\left\langle\Delta\left(\eta_{1}-\eta_{2}\right), \varphi\left(X_{1}-X_{2}\right)\right\rangle_{-1} d t \\
& \quad=\mathbb{E} \int_{0}^{T} \int_{\mathbb{R}^{d}}\left(\eta_{1}-\eta_{2}\right), \varphi\left(X_{1}-X_{2}\right) d t d \xi \geq 0
\end{aligned}
$$

and, therefore,

$$
\mathbb{E}\left\|\varphi\left(X_{1}(t)-X_{2}(t)\right)\right\|_{-1}^{2} \leq C \int_{0}^{t} \mathbb{E}\left\|\varphi\left(X_{1}-X_{2}\right)\right\|_{-1}^{2} d s, \forall t \in[0, T]
$$

and, since $\varphi$ was arbitrary in $C_{0}^{\infty}\left(\mathbb{R}^{d}\right)$, we get $X_{1} \equiv X_{2}$, as claimed.
Remark 4.7 The self-organized criticality model (1.3), that is, $\psi(r) \equiv H(r)=$ Heaviside function, which is not covered by Theorem 4.4 for $1 \leq d \leq 2$, can, however, be treated in the special case

$$
W(t)=\sum_{j=1}^{N} \mu_{j} \beta_{j}(t), \mu_{j} \in \mathbb{R},
$$

(i.e., spatially independent noise) via the rescaling transformation $X=e^{W} Y$, which reduces it to the random parabolic equation

$$
\frac{\partial}{\partial t} Y-e^{-W} \Delta \psi(Y)+\frac{1}{2} \sum_{j=1}^{N} \mu_{j}^{2} Y=0
$$

By approximating $W$ by a smooth $W_{\varepsilon} \in C^{1}([0, T] ; \mathbb{R})$ and letting $\varepsilon \rightarrow 0$, after some calculation one concludes that the latter equation has a unique strong solution $Y$. We omit the details, but refer to [5] for a related treatment.

## 5 The finite time extinction

Assume here that $\psi$ satisfies condition ( j ) of the beginning of Section 4 and that $W$ is of the form ( jj ). Moreover, one assumes that

$$
\begin{equation*}
\psi(r) r \geq \rho|r|^{m+1}, \forall r \in \mathbb{R}, \tag{5.1}
\end{equation*}
$$

where $m$ is as in Hypothesis ( j ).
Theorem 5.1 Let $d \geq 3$ and $m=\frac{d-2}{d+2}$. Let $x \in L^{m+1}\left(\mathbb{R}^{d}\right) \cap L^{2}\left(\mathbb{R}^{d}\right) \cap \mathcal{H}^{-1}$ and let $X=X(t, \cdot)$ be the solution to (1.1) given by Theorem 4.4. We set

$$
\begin{equation*}
\tau=\inf \left\{t \geq 0 ;\|X(t, \cdot)\|_{-1}=0\right\} . \tag{5.2}
\end{equation*}
$$

Then, for every $t>0$,

$$
\begin{equation*}
X(t, \cdot)=0, \forall t \geq \tau \tag{5.3}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathbb{P}[\tau \leq t] \geq 1-\|x\|_{-1}^{1-m} \frac{C^{*}}{\rho \gamma^{m+1}\left(1-e^{-C^{*}(1-m) t}\right)} \tag{5.4}
\end{equation*}
$$

where $\gamma^{-1}=\sup \left\{\|u\|_{-1}|u|_{m+1}^{-1}\right\}$ and $C^{*}>0$ is independent of the initial condition $x$.

Proof. We follow the arguments of [7]. The basic inequality is

$$
\begin{align*}
& \|X(t)\|_{-1}^{1-m}+\rho(1-m) \gamma^{m+1} \int_{r}^{t} \mathbb{1}_{[\|X(s)\|-1>0]} d s \\
& \leq\|X(r)\|_{-1}^{1-m}+C^{*}(1-m) \int_{r}^{t}\|X(s)\|_{-1}^{1-m} d s  \tag{5.5}\\
& +(1-m) \int_{r}^{t}\left\langle\|X(s)\|_{-1}^{(m+1)} X(s), X(s) d W(s)\right\rangle_{-1}, \\
& \quad \mathbb{P} \text {-a.s., } 0<r<t<\infty,
\end{align*}
$$

where $C^{*}$ is a suitable constant. (We note that, by virtue of $(2.8), \gamma^{-1}<\infty$.) To get (5.5), we apply the Itô formula in (4.10) to the semimartingale $\left\|X_{\lambda}(t)\right\|_{-1}^{2}$ and to the function $\varphi_{\varepsilon}(r)=\left(r+\varepsilon^{2}\right)^{\frac{1-m}{2}}, r>-\varepsilon^{2}$, where $X_{\lambda}$ is the solution to (4.10).

We have

$$
\begin{aligned}
& d \varphi_{\varepsilon}\left(\left\|X_{\lambda}(t)\right\|_{-1}^{2}\right) \\
& +(1-m)\left(\left\|X_{\lambda}(t)\right\|_{-1}^{2}+\varepsilon^{2}\right)^{-\frac{m+1}{2}}\left\langle X_{\lambda}(t), \psi_{\lambda}\left(X_{\lambda}(t)\right)+\lambda X_{\lambda}(t)\right\rangle_{2} d t \\
& =\frac{1}{2} \sum_{k=1}^{\infty} \mu_{k}^{2}\left[\frac{(1-m)\left\|X_{\lambda}(t) e_{k}\right\|_{-1}^{2}}{\left(\left\|X_{\lambda}(t)\right\|_{-1}^{2}+\varepsilon^{2}\right)^{\frac{m+1}{2}}}-\left(1-m^{2}\right) \frac{\left\|X_{\lambda}(t) e_{k}\right\|_{-1}^{2}\left\|X_{\lambda}(t)\right\|_{-1}^{2}}{\left(\left\|X_{\lambda}(t)\right\|_{-1}^{2}+\varepsilon^{2}\right)^{\frac{m+1}{2}}}\right] d t \\
& +2\left\langle\varphi_{\varepsilon}^{\prime}\left(\left\|X_{\lambda}(t)\right\|_{-1}^{2}\right) X_{\lambda}(t), X_{\lambda}(t) d W(t)\right\rangle .
\end{aligned}
$$

This yields

$$
\begin{aligned}
& \varphi_{\varepsilon}\left(\left\|X_{\lambda}(t)\right\|_{-1}^{2}\right)+\rho(1-m) \int_{r}^{t}\left(\left\|X_{\lambda}(s)\right\|_{-1}^{2}+\varepsilon^{2}\right)^{-\frac{m+1}{2}} \int_{\mathbb{R}^{d}}\left|X_{\lambda}\right|^{m+1} d s d \xi \\
& \quad \leq \varphi_{\varepsilon}\left(\left\|X_{\lambda}(r)\right\|_{-1}^{2}\right)+C^{*} \int_{r}^{t}\left\|\varphi^{\frac{1}{2}} X_{\lambda}(s)\right\|_{-1}^{2}\left(\left\|X_{\lambda}(s)\right\|_{-1}^{2}+\varepsilon^{2}\right)^{-\frac{1+m}{2}} d s \\
& \quad+2 \int_{r}^{t}\left\langle\varphi_{\varepsilon}^{\prime}\left(\left\|X_{\lambda}(s)\right\|_{-1}^{2}\right) X_{\lambda}(s), X_{\lambda}(s) d W(s)\right\rangle_{-1} .
\end{aligned}
$$

Now, letting $\lambda \rightarrow 0$, we obtain that $X$ satisfies the estimate

$$
\begin{align*}
& \left.\varphi_{\varepsilon}\left(\|X(t)\|_{-1}^{2}\right)+\rho(1-m) \int_{r}^{t}\left(\|X(s)\|_{-1}^{2}+\varepsilon^{2}\right)^{-\frac{m+1}{2}} \int_{\mathbb{R}^{d}}|X(s, \xi)|^{m+1} d \xi\right) d s \\
& \leq \varphi_{\varepsilon}\left(\|\widetilde{X}(t)\|_{-1}^{2}\right)+C^{*} \int_{r}^{t}\|X(s)\|_{-1}^{2}\left(\|\widetilde{X}(s)\|_{-1}^{2}+\varepsilon^{2}\right)^{\frac{m+1}{2}} d s  \tag{5.6}\\
& +2 \int_{r}^{t}\left\langle\varphi_{\varepsilon}^{\prime}\left(\|X(s)\|_{-1}^{2}\right) X(s), X(s) d W(s)\right\rangle_{-1} .
\end{align*}
$$

Here, we have used the fact that, by Lemma 4.5, for $\lambda \rightarrow 0$,

$$
X_{\lambda} \rightarrow X \text { in } \mathcal{H}^{-1}
$$

and, by (4.32) it follows, via Fatou's lemma,

$$
\liminf _{\lambda \rightarrow 0} \int_{\mathbb{R}^{d}}\left|X_{\lambda}\right|^{m+1} d \xi \geq \int_{\mathbb{R}^{d}}|X|^{m+1} d \xi
$$

and

$$
\begin{aligned}
& \left(\|X(t)\|_{-1}^{2}+\varepsilon^{2}\right)^{\frac{1-m}{2}}+\rho(1-m) \gamma^{m+1} \int_{r}^{t}\left(\|X(s)\|_{-1}^{2}+\varepsilon^{2}\right)^{-\frac{m+1}{2}}\|X(s)\|_{-1}^{m+1} d s \\
& \quad \leq\left(\|\left. X(r)\right|_{-1} ^{2}+\varepsilon^{2}\right)^{\frac{1-m}{2}}+C^{*} \int_{r}^{t}\|X(s)\|_{-1}^{2}\left(\|X(s)\|_{-1}^{2}+\varepsilon^{2}\right)^{-\frac{m+1}{2}} d s \\
& \quad+2 \int_{r}^{t}\left\langle\varphi_{\varepsilon}^{\prime}\left(\|X(s)\|_{-1}^{2} X(s), X(s) d W(s)\right\rangle_{-1}, 0 \leq r \leq t<\infty\right.
\end{aligned}
$$

because, by (2.7), $\|x\|_{-1} \leq \gamma^{-1}|x|_{m+1}, \forall x \in L^{m+1}\left(\mathbb{R}^{d}\right)$. Letting $\varepsilon \rightarrow 0$, we get (5.5), as claimed.

Now, we conclude the proof as in [7]. Namely, by (5.5), it follows that

$$
\begin{aligned}
& e^{-C^{*}(1-m) t}\|X(t)\|_{-1}^{1-m}+\rho(1-m) \gamma^{m+1} \int_{r}^{t} e^{-C^{*}(1-m) s} \mathbb{1}_{\left.\| \| X s \|_{-1}>0\right]} d s \\
& \quad \leq e^{-C^{*}(1-m) r}\|X(r)\|_{-1}^{1-m} \\
& \quad+(1-m) \int_{r}^{t} e^{C^{*}(1-m) s}\left\langle\|X(s)\|_{-1}^{-(m+1)} X(s), X(s) d W(s)\right\rangle_{-1}
\end{aligned}
$$

and, therefore, $t \rightarrow e^{-C^{*}(1-m) t}\|X(t)\|_{-1}^{1-m}$ is an $\left\{\mathcal{F}_{t}\right\}$ supermartingale. Hence, $\|X(t)\|_{-1}=0$ for $t \geq \tau$, because of Proposition 3.4, Chap. 2 of [22]. Moreover, taking expectation for $r=0$, we get

$$
e^{-C^{*}(1-m) t} \mathbb{E}\|X(t)\|_{-1}^{1-m}+\rho(1-m) \gamma^{m+1} \int_{0}^{t} e^{-C^{*}(1-m) s} \mathbb{P}(\tau>s) d s \leq\|x\|_{-1}^{1-m}
$$

This implies that

$$
\mathbb{P}(\tau>t) \frac{1-e^{-C^{*}(1-m) t}}{C^{*}(1-m)} \leq \int_{0}^{t} e^{-C^{*}(1-m) s} \mathbb{P}(\tau>s) d s \leq \frac{\|x\|_{-1}^{1-m}}{\rho(1-m) \gamma^{m+1}}
$$

and so (5.4) follows. This completes the proof.
Corollary 5.2 Let $x \in \mathcal{H}^{-1} \cap L^{m+1}\left(\mathbb{R}^{d}\right) \cap L^{2}\left(\mathbb{R}^{d}\right)$ be such that $\|x\|_{-1}<\frac{\rho \gamma^{m+1}}{C^{*}}$. Let $\tau$ be the stopping time defined in (5.2). Then there is $t_{0}>0$ such that $\mathbb{P}\left(\tau>t_{0}\right)>0$. In other words, there is extinction in finite time with positive probability.

Remark 5.3 In the case of bounded domain, Theorem 5.1 remains true for $m \in\left[\frac{d-2}{d+2}, 1\right)$ (see [7]). One might suspect that also in this case the extinction property (5.4) holds for a larger class of exponents $m$. However, the analysis carried out in [24] for deterministic fast diffusion equations in $\mathbb{R}^{d}$ shows that the extinction property is dependent not only on the exponent $m$, but also on the space $L^{p}\left(\mathbb{R}^{d}\right)$, where the solution exists (the so called extinction space).

Remark 5.4 - The analysis in this section holds in particular if all the coefficients $\mu_{k}$ do vanish, i.e. in the deterministic framework. In that case, Theorem 5.1 implies the existence of a deterministic time $\tau>0$ so that

$$
t \geq \tau \Rightarrow\|X(t)\|_{-1}=0
$$

and so $X(t)=0$ for all $t \geq \tau$.

- Let us set for instance $\psi(u)=u^{m}, d \geq 3, m=\frac{d-2}{d+2}$. Observe that $\left(L^{1} \cap L^{\infty}\right)\left(\mathbb{R}^{d}\right) \subset\left(L^{m+1} \cap L^{2}\right)\left(\mathbb{R}^{d}\right) \cap \mathcal{H}^{-1}$. Consider for instance as initial condition $x \in\left(L^{1} \cap L^{\infty}\right)\left(\mathbb{R}^{d}\right)$.
- By Benilan-Crandall approach, see e.g. Theorem 1 of [13], there is a solution $u:[0, T] \times \mathbb{R}^{d} \rightarrow \mathbb{R}$, of

$$
\begin{align*}
& d X-\Delta \psi(X) d t=0 \text { in }(0, T) \times \mathbb{R}^{d} \\
& X(0)=x \text { on } \mathbb{R}^{d} \tag{5.7}
\end{align*}
$$

in the sense of distributions. $u$ belongs to $\left(L^{1} \cap L^{\infty}\right)\left((0, T) \times \mathbb{R}^{d}\right)$ and also the $\eta_{u}=\psi(u)$. In this case, $u$ fulfills mass conservation.

- By use of Theorem 4.4, there is another solution $v:[0, T] \times \mathbb{R}^{d} \rightarrow \mathbb{R}^{d}$ in the sense of distributions, such that $v \in L^{p}\left((0, T) \times \mathbb{R}^{d}\right)$, with $p=$ $\max (2,2 m)$. Also $\eta_{v}=\psi(v) \in L^{\frac{p}{m}}\left((0, T) \times \mathbb{R}^{d}\right)$. By Theorem 5.1, if $x$ is small enough, there will be extinction and so $v$ does not fulfill any mass conservation.
- In particular there is no uniqueness for (5.7) in the sense of distributions. Remark that according to [16] uniqueness is guaranteed in the class $\left(L^{1} \cap L^{\infty}\right)\left((0, T) \times \mathbb{R}^{d}\right)$.

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