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# The Stochastic Porous Media Equations in $\mathbb{R}^d$

Viorel Barbu<sup>\*</sup> Michael Röckner<sup>†</sup> Fran

Francesco Russo<sup>‡</sup>

#### Abstract

Existence and uniqueness of solutions to the stochastic porous media equation  $dX - \Delta \psi(X)dt = XdW$  in  $\mathbb{R}^d$  are studied. Here, Wis a Wiener process,  $\psi$  is a maximal monotone graph in  $\mathbb{R} \times \mathbb{R}$  such that  $\psi(r) \leq C|r|^m$ ,  $\forall r \in \mathbb{R}$ , W is a coloured Wiener process. In this general case the dimension is restricted to  $d \geq 3$ , the main reason being the absence of a convenient multiplier result in the space  $\mathcal{H} = \{\varphi \in \mathcal{S}'(\mathbb{R}^d); |\xi|(\mathcal{F}\varphi)(\xi) \in L^2(\mathbb{R}^d)\}$ , for  $d \leq 2$ . When  $\psi$  is Lipschitz, the well-posedness, however, holds for all dimensions on the classical Sobolev space  $H^{-1}(\mathbb{R}^d)$ . If  $\psi(r)r \geq \rho|r|^{m+1}$  and  $m = \frac{d-2}{d+2}$ , we prove the finite time extinction with strictly positive probability. **Keywords:** stochastic, porous media, Wiener process, maximal mono-

tone graph, distributions.

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#### 1 Introduction

Consider the stochastic porous media equation

$$dX - \Delta \psi(X)dt = XdW \text{ in } (0,T) \times \mathbb{R}^d,$$
  

$$X(0) = x \text{ on } \mathbb{R}^d,$$
(1.1)

where  $\psi$  is a monotonically nondecreasing function on  $\mathbb{R}$  (eventually multivalued) and W(t) is a Wiener process of the form

<sup>\*</sup>Octav Mayer Institute of Mathematics (Romanian Academy), Blvd. Carol I no. 8, Iaşi, 700506, Romania

<sup>&</sup>lt;sup>†</sup>Facultät für Mathematik, Universität Bielefeld, D-33501 Bielefeld, Germany

 $<sup>^{\</sup>ddagger}\mathrm{ENSTA}$  Mathématiques Appliquées, 828, Bd. des Maréchaux, F-91120 Palaiseau, France

$$W(t) = \sum_{k=1}^{\infty} \mu_k e_k \beta_k(t), \ t \ge 0.$$
 (1.2)

Here  $\{\beta_k\}_{k=1}^{\infty}$  are independent Brownian motions on a stochastic basis  $\{\Omega, \mathcal{F}, \mathcal{F}_t, \mathbb{P}\}, \mu_k \in \mathbb{R}$  and  $\{e_k\}_{k=1}^{\infty}$  is an orthonormal basis in  $H^{-1}(\mathbb{R}^d)$  or  $\mathcal{H}^{-1}$  (see (2.2) below) to be made precise later on.

On bounded domains  $\mathcal{O} \subset \mathbb{R}^d$  with Dirichlet homogeneous boundary conditions, equation (1.1) was studied in [3], [4], [6], under general assumptions on  $\psi : \mathbb{R} \to \overline{\mathbb{R}}$  (namely, maximal monotone graph with polynomial growth, or even more general growth conditions in [4]). It should be said, however, that there is a principal difference between bounded and unbounded domains, mainly due to the multiplier problem in Sobolev spaces on  $\mathbb{R}^d$ . If  $d \geq 3$  and  $\mathcal{O} = \mathbb{R}^d$ , existence and uniqueness of solutions to (1.1) was proved in [21] (see, also, [23]) in a general setting which covers the case  $\mathcal{O} = \mathbb{R}^d$  (see Theorems 3.9, Proposition 3.1 and Example 3.4 in [21]). However, it should be said that in [21]  $\psi$  is assumed continuous, such that  $r\psi(r) \to \infty$  as  $r \to \infty$ , which we do not need in this paper.

We study the existence and uniqueness of (1.1) under two different sets of conditions requiring a different functional approach. The first one, which will be presented in Section 3, assumes that  $\psi$  is monotonically nondecreasing and Lipschitz. The state space for (1.1) is, in this case,  $H^{-1}(\mathbb{R}^d)$ , that is, the dual of the classical Sobolev space  $H^1(\mathbb{R}^d)$ . In spite of the apparent lack of generality ( $\psi$  Lipschitz), it should be mentioned that there are physical models described by such an equation as, for instance, the two phase Stefan transition problem perturbed by a stochastic Gaussian noise [2]; moreover, in this case there is no restriction on the dimension d.

The second case, which will be studied here in Section 4, is the one where  $\psi$  is a maximal monotone multivalued function with polynomial growth. An important physical problem covered by this case is the self-organized criticality model

$$dX - \Delta H(X - X_c)dt = (X - X_c)dW, \qquad (1.3)$$

where H is the Heaviside function and  $X_c$  is the critical state (see [6], [7], [10]). More generally, this case with discontinuous  $\psi$  covers the stochastic nonlinear diffusion equation with singular diffusivity  $D(u) = \psi'(u)$ .

It should be mentioned that, in this second case, the solution X(t) to (1.1) is defined in a certain distribution space  $\mathcal{H}^{-1}$  (see (2.2) below) on  $\mathbb{R}^d$  and the existence is obtained for  $d \geq 3$  only, as in the case of continuous  $\psi$ 

in [21]. The case  $1 \le d \le 2$  remains open due to the absence of a multiplier rule in the norm  $\|\cdot\|_{\mathcal{H}^{-1}}$  (see Lemma 4.1 below).

Finally, in Section 5 we prove the finite time extinction of the solution X to (1.1) with strictly positive probability under the assumption that  $\psi(r)r \ge \rho |r|^{m+1}$  and  $m = \frac{d-2}{d+2}$ .

### 2 Preliminaries

To begin with, let us briefly recall a few definitions pertaining distribution spaces on  $\mathbb{R}^d$ , whose classical Euclidean norm will be denoted by  $|\cdot|$ .

Denote by  $\mathcal{S}'(\mathbb{R}^d)$  the space of all temperate distributions on  $\mathbb{R}^d$  (see, e.g., [18]) and by  $\mathcal{H}$  the space

$$\mathcal{H} = \{ \varphi \in \mathcal{S}'(\mathbb{R}^d); \ \xi \mapsto |\xi| \mathcal{F}(\varphi)(\xi) \in L^2(\mathbb{R}^d) \},$$
(2.1)

where  $\mathcal{F}(\varphi)$  is the Fourier transform of  $\varphi$ . We denote by  $L^2(\mathbb{R}^d)$  the space of square integrable functions on  $\mathbb{R}^d$  with norm  $|\cdot|_2$  and scalar product  $\langle \cdot, \cdot \rangle_2$ . In general  $|\cdot|_p$  will denote the norm of  $L^p(\mathbb{R}^d)$  or  $L^p(\mathbb{R}^d; \mathbb{R}^d)$ ,  $1 \leq p \leq \infty$ . The dual space  $\mathcal{H}^{-1}$  of  $\mathcal{H}$  is given by

$$\mathcal{H}^{-1} = \{ \eta \in \mathcal{S}'(\mathbb{R}^d); \ \xi \mapsto \mathcal{F}(\eta)(\xi) |\xi|^{-1} \in L^2(\mathbb{R}^d) \}.$$
(2.2)

The duality between  $\mathcal{H}$  and  $\mathcal{H}^{-1}$  is denoted by  $\langle \cdot, \cdot \rangle$  and is given by

$$\langle \varphi, \eta \rangle = \int_{\mathbb{R}^d} \mathcal{F}(\varphi)(\xi) \overline{\mathcal{F}(\eta)}(\xi) d\xi$$
 (2.3)

and the norm of  $\mathcal{H}$  denoted by  $\|\cdot\|_1$  is given by

$$\|\varphi\|_{1} = \left(\int_{\mathbb{R}^{d}} |\mathcal{F}(\varphi)(\xi)|^{2} |\xi|^{2} d\xi\right)^{\frac{1}{2}} = \left(\int_{\mathbb{R}^{d}} |\nabla\varphi|^{2} d\xi\right)^{\frac{1}{2}}.$$
 (2.4)

The norm of  $\mathcal{H}^{-1}$ , denoted by  $\|\cdot\|_{-1}$  is given by

$$\|\eta\|_{-1} = \left(\int_{\mathbb{R}^d} |\xi|^{-2} |\mathcal{F}(\eta)(\xi)|^2 d\xi\right)^{\frac{1}{2}} = \left(\left\langle (-\Delta)^{-1} \eta, \eta \right\rangle\right)^{\frac{1}{2}}.$$
 (2.5)

(We note that the operator  $-\Delta$  is an isomorphism from  $\mathcal{H}$  onto  $\mathcal{H}^{-1}$ .) The scalar product of  $\mathcal{H}^{-1}$  is given by

$$\langle \eta_1, \eta_2 \rangle_{-1} = \left\langle (-\Delta)^{-1} \eta_1, \eta_2 \right\rangle.$$
 (2.6)

As regards the relationship of  $\mathcal{H}$  with the space  $L^p(\mathbb{R}^d)$  of *p*-summable functions on  $\mathbb{R}^d$ , we have the following. **Lemma 2.1** Let  $d \geq 3$ . Then we have

$$\mathcal{H} \subset L^{\frac{2d}{d-2}}(\mathbb{R}^d) \tag{2.7}$$

algebraically and topologically.

Indeed, by the Sobolev embedding theorem (see, e.g., [15], p. 278), we have

$$|\varphi|_{\frac{2d}{d-2}} \le C |\nabla \varphi|_2, \ \forall \varphi \in C_0^\infty(\mathbb{R}^d),$$

and, by density, this implies (2.7), as claimed.

It should be mentioned that (2.7) is no longer true for  $1 \leq d \leq 2$ . However, by duality, we have

$$L^{\frac{2d}{d+2}}(\mathbb{R}^d) \subset \mathcal{H}^{-1}, \ \forall d \ge 3.$$
 (2.8)

Denote by  $H^1(\mathbb{R}^d)$  the Sobolev space

$$H^{1}(\mathbb{R}^{d}) = \{ u \in L^{2}(\mathbb{R}^{d}); \ \nabla u \in L^{2}(\mathbb{R}^{d}) \}$$
$$= \{ u \in L^{2}(\mathbb{R}^{d}); \ \xi \mapsto \mathcal{F}(u)(\xi)(1+|\xi|^{2})^{\frac{1}{2}} \in L^{2}(\mathbb{R}^{d}) \}$$

with norm

$$|u|_{H^{1}(\mathbb{R}^{d})} = \left(\int_{\mathbb{R}^{d}} (u^{2} + |\nabla u|^{2}) d\xi\right)^{\frac{1}{2}} = \left(\int_{\mathbb{R}^{d}} |\mathcal{F}u(\xi)|^{2} (1 + |\xi|^{2}) d\xi\right)^{\frac{1}{2}}$$

and by  $H^{-1}(\mathbb{R}^d)$  its dual, that is,

$$H^{-1}(\mathbb{R}^d) = \{ u \in \mathcal{S}'(\mathbb{R}^d); \ \mathcal{F}(u)(\xi)(1+|\xi|^2)^{-\frac{1}{2}} \in L^2(\mathbb{R}^d) \}.$$

The norm of  $H^{-1}(\mathbb{R}^d)$  is denoted by  $|\cdot|_{-1}$  and its scalar product by  $\langle \cdot, \cdot \rangle_{-1}$ . We have the continuous and dense embeddings

$$H^1(\mathbb{R}^d) \subset \mathcal{H}, \ \mathcal{H}^{-1} \subset H^{-1}(\mathbb{R}^d).$$

It should be emphasized, however, that  $\mathcal{H}$  is not a subspace of  $L^2(\mathbb{R}^d)$  and so  $L^2(\mathbb{R}^d)$  is not the pivot space in the duality  $\langle \cdot, \cdot \rangle$  given by (2.3).

Given a Banach space Y, we denote by  $L^p(0,T;Y)$  the space of all Y-valued Bochner measurable *p*-integrable functions on (0,T) and by C([0,T];Y) the space of continuous Y-valued functions on [0,T]. For two Hilbert spaces  $H_1, H_2$  let  $L(H_1, H_2)$  and  $L_2(H_1, H_2)$  denote the set of all bounded linear and Hilbert-Schmidt operators, respectively. We refer to [17], [20] for definitions and basic results pertaining infinite dimensional stochastic processes.

### **3** Equation (1.1) with the Lipschitzian $\psi$

Consider here equation (1.1) under the following conditions.

- (i)  $\psi : \mathbb{R} \to \mathbb{R}$  is monotonically nondecreasing, Lipschitz such that  $\psi(0) = 0$ .
- (ii) W is a Wiener process as in (1.2), where  $e_k \in H^1(\mathbb{R}^d)$ , such that

$$C_{\infty}^{2} := 36 \sum_{k=1}^{\infty} \mu_{k}^{2} (|\nabla e_{k}|_{\infty}^{2} + |e_{k}|_{\infty}^{2} + 1) < \infty,$$
(3.1)

and  $\{e_k\}$  is an orthonormal basis in  $H^{-1}(\mathbb{R}^d)$ .

We insert the factor 36 for convenience here to avoid additional large numerical constants in subsequent estimates.

**Remark 3.1** By Lemma 4.1 below,  $|\nabla e_k|_{\infty}$  in (3.1) can be replaced by  $|\nabla e_k|_d$ , and all the results in this section remain true.

**Definition 3.2** Let  $x \in H^{-1}(\mathbb{R}^d)$ . A continuous,  $(\mathcal{F}_t)_{t\geq 0}$ -adapted process  $X : [0,T] \to H^{-1}(\mathbb{R}^d)$  is called strong solution to (1.1) if the following conditions hold:

$$X \in L^{2}(\Omega; C([0, T]; H^{-1}(\mathbb{R}^{d}))) \cap L^{2}([0, T] \times \Omega; L^{2}(\mathbb{R}^{d}))$$
(3.2)

$$\int_0 \psi(X(s))ds \in C([0,T]; H^1(\mathbb{R}^d)), \ \mathbb{P}\text{-a.s.}$$
(3.3)

$$X(t) - \Delta \int_0^t \psi(X(s)) ds = x + \int_0^t X(s) dW(s), \qquad (3.4)$$
$$\forall t \in [0, T], \ \mathbb{P}\text{-a.s.}$$

**Remark 3.3** The stochastic (Itô-) integral in (3.4) is the standard one from [17] or [20]. In fact, in the terminology of these references, W is a Q-Wiener process  $W^Q$  on  $H^{-1}$ , where  $Q : H^{-1} \to H^{-1}$  is the symmetric trace class operator defined by

$$Qh := \sum_{k=1}^{\infty} \mu_k^2 \langle e_k, h \rangle_{-1} e_k, \ h \in H^{-1}.$$

For  $x \in H^{-1}$ , define  $\sigma(x): Q^{1/2}H^{-1} \to H^{-1}$  by

$$\sigma(x)(Q^{1/2}h) = \sum_{k=1}^{\infty} (\mu_k \langle e_k, h \rangle_{-1} e_k \cdot x), \ h \in H.$$
(3.5)

By (3.1), each  $e_k$  is an  $H^{-1}$ -multiplier such that

$$|e_k \cdot x|_{-1} \le 2\left(|e_k|_{\infty} + |\nabla e_k|_{\infty}\right)|x|_{-1}, \ x \in H^{-1}.$$
(3.6)

Hence, for all  $x \in H^{-1}$ ,  $h \in H^{-1}$ ,

$$\sum_{k=1}^{\infty} \left| \mu_k \left\langle e_k, h \right\rangle_{-1} e_k x \right|_{-1} \leq \left( \sum_{k=1}^{\infty} \mu_k^2 |e_k x|_{-1}^2 \right)^{1/2} |h|_{-1}$$
$$\leq 2C_{\infty} |x|_{-1} |h|_{-1}$$
$$= 2C_{\infty} |x|_{-1} |Q^{1/2} h|_{Q^{1/2} H^{-1}},$$

and thus  $\sigma(x)$  is well-defined and an element in  $L(Q^{1/2}H^{-1}, H^{-1})$ . Moreover, for  $x \in H^{-1}$ , by (3.5), (3.6),

$$\begin{aligned} \|\sigma(x)\|_{L_2(Q^{1/2}H^{-1},H^{-1})}^2 &= \sum_{k=1}^{\infty} |\sigma(x)(Q^{1/2}e_k)|_{-1}^2 = \sum_{k=1}^{\infty} |\mu_k e_k x|_{-1}^2 \\ &= \sum_{k=1}^{\infty} \mu_k^2 |e_k x|_{-1}^2 \le C_{\infty}^2 |x|_{-1}^2. \end{aligned}$$
(3.7)

Since  $\{Q^{1/2}e_k \mid k \in \mathbb{N}\}$  is an orthonormal basis of  $Q^{1/2}H^{-1}$ , it follows that  $\sigma(x) \in L_2(Q^{1/2}H^{-1}, H^{-1})$  and the map  $x \mapsto \sigma(x)$  is linear and continuous (hence Lipschitz) from  $H^{-1}$  to  $L_2(Q^{1/2}H^{-1}, H^{-1})$ . Hence (e.g., according to [20, Section 2.3])

$$\int_{0}^{t} X(s) dW(s) := \int_{0}^{t} \sigma(X(s)) dW^{Q}(s), \ t \in [0, T],$$

is well-defined as a continuous  $H^{-1}$ -valued martingale and by Itô's isometry and (3.7)

$$\mathbb{E} \left| \int_{0}^{t} X(s) dW(s) \right|_{-1}^{2} = \sum_{k=1}^{\infty} \mu_{k}^{2} \mathbb{E} \int_{0}^{t} |X(s)e_{k}|_{-1}^{2} ds \\ \leq C_{\infty}^{2} \mathbb{E} \int_{0}^{t} |X(s)|_{-1}^{2} ds, \ t \in [0, T].$$
(3.8)

Furthermore, it follows that

$$\int_{0}^{t} X(s) dW(s) = \sum_{k=1}^{\infty} \int_{0}^{t} \sigma(X(s)) (Q^{1/2} e_{k}) d\beta_{k}(s)$$

$$= \sum_{k=1}^{\infty} \int_{0}^{t} \mu_{k} e_{k} X(s) d\beta_{k}(s), \ t \in [0, T],$$
(3.9)

where the series converges in  $L^2(\Omega; C([0, T]; H^{-1}))$ . In fact,  $\int_0^{\bullet} X(s) dW(s)$  is a continuous  $L^2$ -valued martingale, because  $X \in L^2([0, T] \times \Omega; L^2(\mathbb{R}^d))$  and, analogously to (3.7), we get

$$\|\sigma(x)\|_{L_2(Q^{1/2}H^{-1},L^2)}^2 \le C_\infty^2 |x|_2^2, \ x \in L^2(\mathbb{R}^d).$$

In particular, by Itô's isometry,

$$\mathbb{E}\left|\int_0^t X(s)dW(s)\right|_2^2 \le C_\infty^2 \mathbb{E}\int_0^t |X(s)|_2^2 ds, \ t \in [0,T].$$

Furthermore, the series in (3.8) even converges in  $L^2(\Omega; C([0, T]; L^2(\mathbb{R}^d)))$ .

We shall use the facts presented in this remark throughout this paper without further notice.

**Theorem 3.4** Let  $d \ge 1$  and  $x \in L^2(\mathbb{R}^d)$ . Then, under assumptions (i), (ii), there is a unique strong solution to equation (1.1). This solution satisfies

$$\mathbb{E}\left[\sup_{t\in[0,T]}|X(t)|_{2}^{2}\right] \leq 2|x|_{2}^{2}e^{3C_{\infty}^{2}t}.$$

In particular,  $X \in L^2(\Omega; L^{\infty}([0, T]; L^2(\mathbb{R}^d)))$ . Assume further that

$$\psi(r)r \ge \alpha r^2, \ \forall r \in \mathbb{R},$$
(3.10)

where  $\alpha > 0$ . Then, there is a unique strong solution X to (1.1) for all  $x \in H^{-1}(\mathbb{R}^d).$ 

**Proof of Theorem 3.4.** We approximate (1.1) by

$$dX + (\nu - \Delta)\psi(X)dt = XdW(t), \ t \in (0,T),$$
  

$$X(0) = x \text{ on } \mathbb{R}^d,$$
(3.11)

where  $\nu \in (0, 1)$ . We have the following.

**Lemma 3.5** Assume that  $\psi$  is as in assumption (i). Let  $x \in L^2(\mathbb{R}^d)$ . Then, there is a unique  $(\mathcal{F}_t)_{t\geq 0}$ -adapted solution  $X = X^{\nu}$  to (3.11) in the following strong sense:

$$X^{\nu} \in L^{2}(\Omega, C([0, T]; H^{-1}(\mathbb{R}^{d}))) \cap L^{2}([0, T] \times \Omega; L^{2}(\mathbb{R}^{d})),$$
(3.12)

and  $\mathbb{P}$ -a.s.

$$X^{\nu}(t) = x + (\Delta - \nu) \int_0^t \psi(X^{\nu}(s)) ds + \int_0^t X^{\nu}(s) dW(s), \ t \in [0, T].$$
(3.13)

In addition, for all  $\nu \in (0, 1)$ ,

$$\mathbb{E}\left[\sup_{t\in[0,T]}|X^{\nu}(t)|_{2}^{2}\right] \leq 2|x|_{2}^{2}e^{3C_{\infty}^{2}T}.$$
(3.14)

If, moreover,  $\psi$  satisfies (3.10), then for each  $x \in H^{-1}(\mathbb{R}^d)$  there is a unique solution  $X^{\nu}$  satisfying (3.12) and (3.13).

**Proof of Lemma 3.5.** Let us start with the second part of the assertion, i.e., we assume that  $\psi$  satisfies (3.10) and that  $x \in H^{-1}(\mathbb{R}^d)$ . Then the standard theory (see, e.g., [20, Sections 4.1 and 4.2]) applies to ensure that there exists a unique solution  $X^{\nu}$  taking value in  $H^{-1}(\mathbb{R}^d)$  satisfying (3.12) and (3.13) above. Indeed, it is easy to check that (H1)–(H4) from [20, Section 4.1] are satisfied with  $V := L^2(\mathbb{R}^d)$ ,  $H := H^{-1}(\mathbb{R}^d)$ ,  $Au := (\Delta - \nu)(\psi(u))$ ,  $u \in V$ , and  $H^{-1}(\mathbb{R}^d)$  is equipped with the equivalent norm

$$|\eta|_{-1,\nu} := \left\langle \eta, (\nu - \Delta)^{-1} \eta \right\rangle^{1/2}, \ \eta \in H^{-1}(\mathbb{R}^d)$$

(in which case, we also write  $H_{\nu}^{-1}$ ). Here, as before, we use  $\langle \cdot, \cdot \rangle$  also to denote the dualization between  $H^1(\mathbb{R}^d)$  and  $H^{-1}(\mathbb{R}^d)$ . For details, we refer to the calculations in [20, Example 4.1.11], which since p = 2, go through when the bounded domain  $\Omega$  there is replaced by  $\mathbb{R}^d$ . Hence [20, Theorem 4.2.4] applies to give the above solution  $X^{\nu}$ .

In the case when  $\psi$  does not satisfy (3.10), the above conditions (H1), (H2), (H4) from [20] still hold, but (H3) not in general. Therefore, we replace  $\psi$  by  $\psi + \lambda I$ ,  $\lambda \in (0, 1)$ , and thus consider  $A_{\lambda}(u) := (\Delta - \nu)(\psi(u) + \lambda u)$ ,  $u \in V := L^2(\mathbb{R}^d)$  and, as above, by [20, Theorem 4.2.4], obtain a solution  $X_{\lambda}^{\nu}$ , satisfying (3.12), (3.13), to

$$dX^{\nu}_{\lambda}(t) + (\nu - \Delta)(\psi(X^{\nu}_{\lambda}(t)) + \lambda X^{\nu}_{\lambda}(t))dt = X^{\nu}_{\lambda}(t)dW(t), \ t \in [0, T],$$
  
$$X^{\nu}_{\lambda}(0) = x \in H^{-1}(\mathbb{R}^d).$$
(3.15)

In particular, by (3.12),

$$\mathbb{E}\left[\sup_{t\in[0,T]}|X_{\lambda}^{\nu}(t)|_{-1}^{2}\right]<\infty.$$
(3.16)

We want to let  $\lambda \to 0$  to obtain a solution to (3.11). To this end, in this case (i.e., without assuming (3.10)), we assume from now on that  $x \in L^2(\mathbb{R}^d)$ . The reason is that we need the following:

Claim 1. We have  $X_{\lambda}^{\nu} \in L^2([0,T] \times \Omega; H^1(\mathbb{R}^d))$  and

$$\mathbb{E}\left[\sup_{t\in[0,T]}|X_{\lambda}^{\nu}(t)|_{2}^{2}\right] + 4\lambda\mathbb{E}\int_{0}^{T}|\nabla X_{\lambda}^{\nu}(s)|_{2}^{2}ds \leq 2|x|_{2}^{2}e^{3C_{\infty}^{2}T}$$
  
for all  $\nu, \lambda \in (0,1)$ .

Furthermore,  $X^{\nu}_{\lambda}$  has continuous sample paths in  $L^{2}(\mathbb{R}^{d})$ ,  $\mathbb{P}$ -a.s.

**Proof of Claim 1.** We know that

$$X_{\lambda}^{\nu}(t) = x + (\Delta - \nu) \int_{0}^{t} (\psi(X_{\lambda}^{\nu}(s)) + \lambda X_{\lambda}^{\nu}(s)) ds + \int_{0}^{t} X_{\lambda}^{\nu}(s) dW(s), \ t \in [0, T].$$

$$(3.17)$$

Let  $\alpha \in (\nu, \infty)$ . Recalling that  $(\alpha - \Delta)^{-\frac{1}{2}} : H^{-1}(\mathbb{R}^d) \to L^2(\mathbb{R}^d)$  and applying this operator to the above equation, we find

$$(\alpha - \Delta)^{-\frac{1}{2}} X^{\nu}_{\lambda}(t)$$

$$= (\alpha - \Delta)^{-\frac{1}{2}} x + \int_{0}^{t} (\Delta - \nu)(\alpha - \Delta)^{-\frac{1}{2}} (\psi(X^{\nu}_{\lambda}(s)) + \lambda X^{\nu}_{\lambda}(s)) ds \qquad (3.18)$$

$$+ \int_{0}^{t} (\alpha - \Delta)^{-\frac{1}{2}} \sigma(X^{\nu}_{\lambda}(s)) Q^{1/2} dW(s), \ t \in [0, T].$$

Applying Itô's formula (see, e.g., [20, Theorem 4.2.5] with  $H = L^2(\mathbb{R}^d)$ ) to  $|(\alpha - \Delta)^{-\frac{1}{2}} X^{\nu}_{\lambda}(t)|_2^2$ , we obtain, for  $t \in [0, T]$ ,

$$\begin{split} &|(\alpha - \Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(t)|_{2}^{2} = |(\alpha - \Delta)^{-\frac{1}{2}} x|_{2}^{2} \\ &+ 2 \int_{0}^{t} \left\langle (\Delta - \nu)(\alpha - \Delta)^{-\frac{1}{2}} \psi(X_{\lambda}^{\nu}(s)), (\alpha - \Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s) \right\rangle ds \\ &- 2\lambda \int_{0}^{t} (|\nabla((\alpha - \Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s))|_{2}^{2} + \nu |(\alpha - \Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s)|_{2}^{2}) ds \\ &+ \int_{0}^{t} ||(\alpha - \Delta)^{-\frac{1}{2}} \sigma(X_{\lambda}^{\nu}(s)) Q^{1/2} ||_{L_{2}(H^{-1}, L^{2})}^{2} ds \\ &+ 2 \int_{0}^{t} \left\langle (\alpha - \Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s), (\alpha - \Delta)^{-\frac{1}{2}} \sigma(X_{\lambda}^{\nu}(s)) Q^{1/2} dW(s) \right\rangle_{2}. \end{split}$$
(3.19)

But, for  $f \in L^2(\mathbb{R}^d)$ , we have

$$(\alpha - \Delta)^{-\frac{1}{2}} (\Delta - \nu) (\alpha - \Delta)^{-\frac{1}{2}} f = (P - I) f,$$

where

$$P := (\alpha - \nu)(\alpha - \Delta)^{-1}.$$

For the Green function  $g_{\alpha}$  of  $(\alpha - \Delta)$ , we then have, for  $f \in L^2(\mathbb{R}^d)$ ,

$$Pf = (\alpha - \nu) \int_{\mathbb{R}^d} f(\xi) g_\alpha(\cdot, \xi) d\xi.$$

Hence, by [23, Lemma 5.1], the integrand of the second term on the right hand side of (3.19) with  $f := X_{\lambda}^{\nu}(s) \ (\in L^2(\mathbb{R}^d)$  for *ds*-a.e.  $s \in [0,T]$ ) can be rewritten as

$$\langle \psi(f), (P-I)f \rangle_2 = -\frac{1}{2} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} [\psi(f(\widetilde{\xi})) - \psi(f(\xi))] [f(\widetilde{\xi}) - f(\xi)] g_\alpha(\xi, \widetilde{\xi}) d\widetilde{\xi} d\xi - \int_{\mathbb{R}^d} (1 - P1(\xi)) \cdot \psi(f(\xi)) f(\xi) d\xi.$$

Since  $\psi$  is monotone,  $\psi(0) = 0$  and  $P1 \leq 1$ , we deduce that

$$\langle \psi(f), (P-I)f \rangle \le 0.$$

Hence, after a multiplication by  $\alpha$ , (3.19) implies that, for all  $t \in [0, T]$  (see Remark 3.3),

$$\begin{aligned} &\alpha |(\alpha - \Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(t)|_{2}^{2} + 2\lambda \int_{0}^{t} |\nabla(\sqrt{\alpha}(\alpha - \Delta)^{-\frac{1}{2}} X_{\lambda}^{\nu}(s))|_{2}^{2} ds \\ &\leq \alpha |(\alpha - \Delta)^{-\frac{1}{2}} x|_{2}^{2} + \int_{0}^{t} \sum_{k=1}^{\infty} \mu_{k}^{2} \left\langle \alpha(\alpha - \Delta)^{-1}(e_{k} X_{\lambda}^{\nu}(s)), e_{k} X_{\lambda}^{\nu}(s) \right\rangle_{2} ds \\ &+ 2 \int_{0}^{t} \left\langle \alpha(\alpha - \Delta)^{-1} X_{\lambda}^{\nu}(s), \sigma(X_{\lambda}^{\nu}(s)) Q^{1/2} dW(s) \right\rangle_{2}. \end{aligned}$$

Hence, by the Burkholder–Davis–Dundy (BDG) inequality (with p = 1) and since  $\alpha(\alpha - \Delta)^{-1}$  is a contraction on  $L^2(\mathbb{R}^d)$ ,

$$\mathbb{E}\left[\sup_{s\in[0,t]} |\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}}X_{\lambda}^{\nu}(s)|_{2}^{2}\right] + 2\lambda\mathbb{E}\int_{0}^{t} |\nabla(\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}}X_{\lambda}^{\nu}(s))|_{2}^{2}ds \\ \leq |\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}}x|_{2}^{2} + C_{\infty}^{2}\mathbb{E}\int_{0}^{t} |X_{\lambda}^{\nu}(s)|_{2}^{2}ds \\ + 6\mathbb{E}\left(\int_{0}^{t}\sum_{k=1}^{\infty}\mu_{k}^{2}\langle\alpha(\alpha-\Delta)^{-1}X_{\lambda}^{\nu}(s),e_{k}X_{\lambda}^{\nu}(s)\rangle_{2}^{2}ds\right)^{1/2}.$$
(3.20)

The latter term can be estimated by

$$C_{\infty}\mathbb{E}\left[\sup_{s\in[0,t]}|\alpha(\alpha-\Delta)^{-1}X_{\lambda}^{\nu}(s)|_{2}\left(\int_{0}^{t}|X_{\lambda}^{\nu}(s)|_{2}^{2}ds\right)^{1/2}\right]$$

$$\leq\frac{1}{2}\mathbb{E}\left\{\sup_{s\in[0,t]}|\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}}X_{\lambda}^{\nu}(s)|_{2}^{2}\right]+\frac{1}{2}C_{\infty}^{2}\mathbb{E}\int_{0}^{t}|X_{\lambda}^{\nu}(s)|_{2}^{2}ds,$$
(3.21)

where we used that  $\sqrt{\alpha}(\alpha - \Delta)^{-\frac{1}{2}}$  is a contraction on  $L^2(\mathbb{R}^d)$ . Note that the first summand on the right hand side is finite by (3.16), since the norm  $|\sqrt{\alpha}(\alpha - \Delta)^{-\frac{1}{2}} \cdot |_2$  is equivalent to  $|\cdot|_{-1}$ . Hence, we can subtract this term after substituting (3.21) into (3.20) to obtain

$$\mathbb{E}\left[\sup_{s\in[0,t]} |\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}}X_{\lambda}^{\nu}(s)|_{2}^{2}\right] + 4\lambda\mathbb{E}\int_{0}^{t} |\nabla(\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}}X_{\lambda}^{\nu}(s)|_{2}^{2}ds \qquad (3.22) \\ \leq 2|\sqrt{\alpha}(\alpha-\Delta)^{-\frac{1}{2}}x|_{2}^{2} + 3C_{\infty}^{2}\mathbb{E}\int_{0}^{t} |X_{\lambda}^{\nu}(s)|_{2}^{2}ds, \ t\in[0,T].$$

Obviously, the quantity under the sup on the left hand side of (3.22) is increasing in  $\alpha$ . So, by the monotone convergence theorem, we may let  $\alpha \to \infty$  in (3.22) and then, except for its last part, Claim 1 immediately follows by Gronwall's lemma, since  $\sqrt{\alpha}(\alpha - \Delta)^{-\frac{1}{2}}$  is a contraction in  $L^2(\mathbb{R}^d)$ and  $x \in L^2(\mathbb{R}^d)$ . The last part of Claim 1 then immediately follows from [19, Theorem 2.1].

Applying Itô's formula to  $|X_{\lambda}^{\nu}(t) - X_{\lambda'}^{\nu}(t)|_{-1,\nu}^2$  (see [20, Theorem 4.2.5]), it follows from (3.17) that, for  $\lambda, \lambda' \in (0, 1)$  and  $t \in [0, T]$ ,

$$X_{\lambda}^{\nu}(t) - X_{\lambda'}^{\nu}(t)|_{-1,\nu}^{2} + 2\int_{0}^{t} \langle \psi(X_{\lambda}^{\nu}) - \psi(X_{\lambda'}^{\nu}) + (\lambda X_{\lambda}^{\nu} - \lambda' X_{\lambda'}^{\nu}), X_{\lambda}^{\nu} - X_{\lambda'}^{\nu} \rangle_{2} ds$$

$$= \int_{0}^{t} \|\sigma(X_{\lambda}^{\nu}(s) - X_{\lambda'}^{\nu}(s))\|_{L_{2}(Q^{1/2}H^{-1}, H_{\nu}^{-1})}^{2} ds$$

$$+ 2\int_{0}^{t} \langle X_{\lambda}^{\nu}(s) - X_{\lambda'}^{\nu}(s), \sigma(X_{\lambda}^{\nu}(s) - X_{\lambda'}^{\nu}(s)) dW^{Q}(s) \rangle_{-1,\nu}.$$
(3.23)

Our assumption (i) on  $\psi$  implies that

$$(\psi(r) - \psi(r'))(r - r') \ge (\operatorname{Lip} \psi + 1)^{-1} |\psi(r) - \psi(r')|^2, \text{ for } r, r' \in \mathbb{R},$$

where  $\operatorname{Lip} \psi$  is the Lipschitz constant of  $\psi$ . Hence (3.23), (3.7) and the BDG inequality (for p = 1 imply that, for all  $t \in [0, T]$ ,

$$\begin{split} & \mathbb{E}\left[\sup_{s\in[0,t]}|X_{\lambda}^{\nu}(s)-X_{\lambda'}^{\nu}(s)|^{2}_{-1,\nu}\right] \\ &+2(\operatorname{Lip}\psi+1)^{-1}\mathbb{E}\int_{0}^{t}|\psi(X_{\lambda}^{\nu}(s))-\psi(X_{\lambda'}^{\nu}(s))|^{2}_{2}ds \\ &\leq 2(\lambda+\lambda')\mathbb{E}\int_{0}^{t}(|X_{\lambda}^{\nu}(s)|^{2}_{2}+|X_{\lambda'}^{\nu}(s)|^{2}_{2})ds+C_{\infty}^{2}\int_{0}^{t}|X_{\lambda}^{\nu}(s)-X_{\lambda'}^{\nu}(s)|^{2}_{-1,\nu}ds \\ &+2\mathbb{E}\left(\int_{0}^{t}\sum_{k=1}^{\infty}\mu_{k}^{2}\langle X_{\lambda}^{\nu}(s)-X_{\lambda'}^{\nu}(s),(X_{\lambda}^{\nu}(s)-X_{\lambda'}^{\nu}(s))e_{k}\rangle^{2}_{-1,\nu}ds\right)^{1/2}. \end{split}$$

By (3.7) and Young's inequality, the latter term is dominated by

$$\frac{1}{2}\mathbb{E}\left[\sup_{s\in[0,t]}|X_{\lambda}^{\nu}(s)-X_{\lambda'}^{\nu}(s)|^{2}_{-1,\nu}\right]+\frac{1}{2}C_{\infty}^{2}\mathbb{E}\int_{0}^{t}|X_{\lambda}^{\nu}(s)-X_{\lambda'}^{\nu}(s)|^{2}_{-1,\nu}ds.$$

Hence, because of  $x \in L^2(\mathbb{R}^d)$  and Claim 1, we may now apply Gronwall's lemma to obtain that, for some constant C independent of  $\lambda', \lambda$  (and  $\nu$ ),

$$\mathbb{E}\left[\sup_{t\in[0,T]} |X_{\lambda}^{\nu}(t) - X_{\lambda'}^{\nu}(t)|^{2}_{-1,\nu}\right] \\
+ \mathbb{E}\int_{0}^{T} |\psi(X_{\lambda}^{\nu}(s)) - \psi(X_{\lambda'}^{\nu}(s))|^{2}_{2}ds \leq C(\lambda + \lambda').$$
(3.24)

Hence there exists an  $(\mathcal{F}_t)$ -adapted continuous  $H^{-1}$ -valued process  $X^{\nu} = (X^{\nu}(t))_{t \in [0,T]}$  such that  $X^{\nu} \in L^2(\Omega; C([0,T]; H^{-1}))$ . Now, by Claim 1, it follows that

$$X^{\nu} \in L^2([0,T] \times \Omega; L^2(\mathbb{R}^d)).$$

Claim 2.  $X^{\nu}$  satisfies equation (3.13) (i.e., we can pass to the limit in (3.17) as  $\lambda \to 0$ ).

**Proof of Claim 2.** We already know that

$$X^{\nu}_{\lambda} \longrightarrow X^{\nu} \text{ and } \int_{0}^{\bullet} X^{\nu}_{\lambda}(s) dW(s) \longrightarrow \int_{0}^{\bullet} X^{\nu}(s) dW(s)$$

in  $L^2(\Omega; C([0, T]; H^{-1}))$  as  $\lambda \to 0$  (for the second convergence see the above argument using (3.7) and the BDG inequality). So, by (3.17) it follows that

$$\int_0^{\bullet} (\psi(X_{\lambda}^{\nu}(s)) + \lambda X_{\lambda}^{\nu}(s))) ds, \ \lambda > 0,$$

converges as  $\lambda \to 0$  to an element in  $L^2(\Omega; C([0, T]; H^1))$ . But, by (3.24) and Claim 1, it follows that

$$\int_0^{\bullet} (\psi(X_{\lambda}^{\nu}(s)) + \lambda X_{\lambda}^{\nu}(s)) ds \longrightarrow \int_0^{\bullet} \psi(X^{\nu}(s)) ds$$
(3.25)

as  $\lambda \to 0$  in  $L^2(\Omega; L^2([0, T]; L^2(\mathbb{R}^d)))$ . Hence Claim 2 is proved.

Now, (3.14) follows from Claim 1 by lower semicontinuity. This completes the proof of Lemma 3.5.  $\blacksquare$ 

**Proof of Theorem 3.4 (continued).** We are going to use Lemma 3.5 and let  $\nu \to 0$ . The arguments are similar to those in the proof of Lemma 3.5. So, we shall not repeat all the details.

Now, we rewrite (3.11) as

$$dX^{\nu} + (I - \Delta)\psi(X^{\nu})dt = (1 - \nu)\psi(X^{\nu})dt + X^{\nu}dW(t)$$
(3.26)

and apply Itô's formula to  $\varphi(x) = \frac{1}{2} |x|_{-1}^2$  (see, e.g., [20, Theorem 4.2.5]). We get, for  $x \in H^{-1}$ , by (3.8) and after taking expectation,

$$\begin{split} \frac{1}{2} \, \mathbb{E} |X^{\nu}(t)|_{-1}^2 + \mathbb{E} \int_0^t \int_{\mathbb{R}^d} \psi(X^{\nu}(s)) X^{\nu}(s) d\xi \, ds \\ &= \frac{1}{2} \, |x|_{-1}^2 + (1-\nu) \mathbb{E} \int_0^t \langle \psi(X^{\nu}(s)), X^{\nu}(s) \rangle_{-1} \, ds \\ &+ \frac{1}{2} \, \mathbb{E} \int_0^t \sum_{k=1}^\infty \mu_k^2 |X^{\nu} e_k|_{-1}^2 ds \\ &\leq \frac{1}{2} \, |x|_{-1}^2 + \mathbb{E} \int_0^t |\psi(X^{\nu})|_{-1} |X^{\nu}|_{-1} ds \\ &+ \frac{1}{2} C_{\infty}^2 \mathbb{E} \int_0^t |X^{\nu}(s)|_{-1}^2 ds, \ \forall t \in [0,T]. \end{split}$$

Recalling that  $|\cdot|_{-1} \leq |\cdot|_2$ , we get, via Young's and Gronwall's inequalities, for some  $C \in (0, \infty)$  that

$$\mathbb{E}|X^{\nu}(t)|_{-1}^{2} + \frac{\alpha}{2} \mathbb{E} \int_{0}^{T} |X^{\nu}(s)|_{2}^{2} ds \le C|x|_{-1}^{2}, \ t \in [0,T], \ \nu \in (0,1), \quad (3.27)$$

because, by assumption (i),  $\psi(r)r \geq \widetilde{\alpha}|\psi(r)|^2, \forall r \in \mathbb{R}$ , with  $\widetilde{\alpha} := (\operatorname{Lip} \psi + 1)^{-1}$ . Here we set  $\alpha = 0$  if (3.10) does not hold. Now, by a similar calculus, for  $X^{\nu} - X^{\nu'}$  we get

$$\begin{split} |X^{\nu}(t) - X^{\nu'}(t)|^{2}_{-1} &+ 2\int_{0}^{t}\int_{\mathbb{R}^{d}}(\psi(X^{\nu}) - \psi(X^{\nu'}))(X^{\nu} - X^{\nu'})d\xi \, ds \\ &\leq C\int_{0}^{t}\left\langle\psi(X^{\nu}) - \psi(X^{\nu'}), X^{\nu} - X^{\nu'}\right\rangle_{-1} ds \\ &+ C\int_{0}^{t}(\nu|\psi(X^{\nu})|_{2} + \nu'|\psi(X^{\nu'})|_{2})|X^{\nu} - X^{\nu'}|_{-1} ds \\ &+ C\int_{0}^{t}|X^{\nu} - X^{\nu'}|^{2}_{-1} ds + \sum_{k=1}^{\infty}\int_{0}^{t}\mu_{k}\left\langle(X^{\nu} - X^{\nu'}), e_{k}(X^{\nu} - X^{\nu'})d\beta_{k}\right\rangle_{-1}, \\ &\quad t \in [0, T]. \end{split}$$

Taking into account that, by assumption (i),

$$(\psi(x) - \psi(y))(x - y) \ge \widetilde{\alpha} |\psi(x) - \psi(y)|^2, \ \forall x, y \in \mathbb{R}^d,$$

we get, for all  $\nu, \nu' > 0$ ,

$$\begin{aligned} |X^{\nu}(t) - X^{\nu'}(t)|^{2}_{-1} + \tilde{\alpha} \int_{0}^{t} |\psi(X^{\nu}(s)) - \psi(X^{\nu'}(s))|^{2}_{2} ds \\ &\leq C_{1} \int_{0}^{t} |X^{\nu}(s) - X^{\nu'}(s)|^{2}_{-1} ds + \frac{\tilde{\alpha}}{2} \int_{0}^{t} |\psi(X^{\nu}(s)) - \psi(X^{\nu'}(s))|^{2}_{2} ds \\ &+ C_{2}(\nu + \nu') \int_{0}^{t} (|\psi(X^{\nu}(s))|^{2}_{2} + |\psi(X^{\nu'}(s))|^{2}_{2}) ds \\ &+ \sum_{k=1}^{\infty} \int_{0}^{t} \mu_{k} \langle (X^{\nu}(s) - X^{\nu'}(s)), e_{k}(X^{\nu}(s) - X^{\nu'}(s)) \rangle_{-1} d\beta_{k}(s), \\ &\quad t \in [0, T]. \end{aligned}$$

So, similarly to showing (3.24) in the proof of Lemma 3.5, by (3.14), if  $x \in L^2(\mathbb{R}^d)$ , and by (3.27), if  $x \in H^{-1}(\mathbb{R}^d)$  and  $\psi$  satisfies (3.10), by the Burkholder-Davis-Gundy inequality, for p = 1, we get, for all  $\nu, \nu' \in (0, 1)$ ,

$$\mathbb{E}\sup_{t\in[0,T]}|X^{\nu}(t)-X^{\nu'}(t)|^{2}_{-1}+\mathbb{E}\int_{0}^{T}|\psi(X^{\nu}(s))-\psi(X^{\nu'}(s))|^{2}_{2}ds\leq C(\nu+\nu').$$

The remaining part of the proof is now exactly the same as the last part of the proof of Lemma 3.5.  $\blacksquare$ 

**Remark 3.6** Theorem 3.4 is a basic tool for the probabilistic (double) representation of equation (1.1), which holds when  $\psi$  is Lipschitz, as it is proved in [9]. If (1.1) is not perturbed by noise, and  $\psi$  is possibly discontinuous, its probabilistic representation was performed in [14, 8] with extensions and numerical simulations located in [11, 12].

# 4 Equation (1.1) for maximal monotone functions $\psi$ with polynomial growth

In this section, we assume  $d \ge 3$  and we shall study the existence for equation (1.1) under the following assumptions:

(j)  $\psi: \mathbb{R} \to 2^{\mathbb{R}}$  is a maximal monotone graph such that  $0 \in \psi(0)$  and

$$\sup\{|\eta|; \ \eta \in \psi(r)\} \le C(1+|r|^m), \ \forall r \in \mathbb{R},$$
(4.1)

where  $1 \leq m < \infty$ .

(jj) 
$$W(t) = \sum_{k=1}^{\infty} \mu_k e_k \beta_k(t), \ t \ge 0$$
, where  $\{\beta_k\}_{k=1}^{\infty}$  are independent Brownian

motions on a stochastic basis  $\{\Omega, \mathcal{F}, \mathcal{F}_t, \mathbb{P}\}, \mu_k \in \mathbb{R}, \text{ and } e_k \in C^1(\mathbb{R}^d) \cap \mathcal{H}^{-1}$  are such that  $\{e_k\}$  is an orthonormal basis in  $\mathcal{H}^{-1}$  and

$$\sum_{k=1}^{\infty} \mu_k^2 (|e_k|_{\infty}^2 + |\nabla e_k|_d^2 + 1) < \infty.$$
(4.2)

The existence of  $\{e_k\}$  as in (jj) is ensured by the following lemma.

**Lemma 4.1** Let  $d \geq 3$  and let  $e \in L^{\infty}(\mathbb{R}^d; \mathbb{R}^d)$  be such that  $\nabla e \in L^d(\mathbb{R}^d; \mathbb{R}^d)$ . Then

$$\|xe\|_{-1} \le \|x\|_{-1} (|e|_{\infty} + C|\nabla e|_{d}), \ \forall x \in \mathcal{H}^{-1},$$
(4.3)

where C is independent of x and e.

#### **Proof.** We have

$$||xe||_{-1} = \sup\{\langle x, e\varphi\rangle; ||\varphi||_1 \le 1\} \le ||x||_{-1} \sup\{||e\varphi||_1; ||\varphi||_1 \le 1\}.$$
(4.4)

On the other hand, by Lemma 2.1 we have, for all  $\varphi \in C_0^{\infty}(\mathbb{R}^d)$ ,

$$\begin{aligned} \|e\varphi\|_1 &\leq |e\nabla\varphi + \varphi\nabla e|_2 \leq |e\nabla\varphi|_2 + |\varphi\nabla e|_2 \\ &\leq |e|_{\infty}|\nabla\varphi|_2 + |\varphi|_p|\nabla e|_d \leq |e|_{\infty}\|\varphi\|_1 + C\|\varphi\|_1|\nabla e|_d, \end{aligned}$$

where  $p = \frac{2d}{d-2}$ . Then, by (4.4), (4.3) follows, as claimed.

#### Remark 4.2

(i) It should be mentioned that, for d = 2, Lemma 4.1 fails and this is the main reason why our treatment of equation (1.1) under assumptions (j), (jj) is constrained to  $d \ge 3$ .

(ii) We note that Remark 3.3 with the rôle of  $H^{-1}(\mathbb{R}^d)$  replaced by  $\mathcal{H}^{-1}$  remains true in all its parts under condition (jj) above. We shall use this below without further notice.

We denote by  $j : \mathbb{R} \to \mathbb{R}$  the potential associated with  $\psi$ , that is, a continuous convex function on  $\mathbb{R}$  such that  $\partial j = \psi$ , i.e.,

$$j(r) \leq \zeta(r-\overline{r}) + j(\overline{r}), \ \forall \zeta \in \psi(r), \ r, \overline{r} \in \mathbb{R}.$$

**Definition 4.3** Let  $x \in \mathcal{H}^{-1}$  and  $p := \max(2, 2m)$ . An  $\mathcal{H}^{-1}$ -valued adapted process X = X(t) is called strong solution to (1.1) if the following conditions hold:

X is 
$$\mathcal{H}^{-1}$$
-valued continuous on  $[0, T]$ ,  $\mathbb{P}$ -a.s., (4.5)

$$X \in L^p(\Omega \times (0,T) \times \mathbb{R}^d).$$
(4.6)

There is  $\eta \in L^{\frac{p}{m}}(\Omega \times (0,T) \times \mathbb{R}^d)$  such that

$$\eta \in \psi(X), \ dt \otimes \mathbb{P} \otimes d\xi \ - \text{a.e. on } (0,T) \times \Omega \times \mathbb{R}^d$$

$$(4.7)$$

and  $\mathbb{P}$ -a.s.

$$X(t) = x + \Delta \int_0^t \eta(s) ds + \sum_{k=1}^\infty \mu_k \int_0^t X(s) e_k d\beta_k(s)$$
(4.8)  
in  $\mathcal{D}'(\mathbb{R}^d), \quad t \in [0, T].$ 

Here  $\mathcal{D}'(\mathbb{R}^d)$  is the standard space of distributions on  $\mathbb{R}^d$ .

Theorem 4.4 below is the main existence result for equation (1.1).

**Theorem 4.4** Assume that  $d \geq 3$  and that  $x \in L^p(\mathbb{R}^d) \cap L^2(\mathbb{R}^d) \cap \mathcal{H}^{-1}$ ,  $p := \max(2, 2m)$ . Then, under assumptions (j), (jj), there is a unique strong solution X to (1.1) such that

$$X \in L^{2}(\Omega; C([0, T]; \mathcal{H}^{-1})).$$
(4.9)

Moreover, if  $x \ge 0$ , a.e. in  $\mathbb{R}^d$ , then  $X \ge 0$ , a.e. on  $(0,T) \times \mathbb{R}^d \times \Omega$ .

Theorem 4.4 is applicable to a large class of nonlinearities  $\psi : \mathbb{R} \to 2^{\mathbb{R}}$ and, in particular, to

$$\psi(r) = \rho H(r) + \alpha r, \ \forall r \in \mathbb{R}, \psi(r) = \rho H(r - r_c)r_c$$

where  $\rho > 0$ ,  $\alpha, r_c \ge 0$ , which models the dynamics of self-organized criticality (see [6], [7], [10]). Here *H* is the Heaviside function.

As mentioned earlier, Theorem 4.4 can be compared most closely to the main existence result of [21]. But there are, however, a few notable differences we explain below. The function  $\psi$  arising in [21] is monotonically increasing, continuous and is assumed to satisfy a growth condition of the form  $N(r) \leq r\psi(r) \leq C(N(r)+1)r, \forall r \in \mathbb{R}$ , where N is a smooth and  $\Delta_2$ -regular Young function defining the Orlicz class  $L_N$ . In contrast to this, here  $\psi$  is any maximal monotone graph (multivalued) with arbitrary polynomial growth.

Proof of Theorem 4.4. Consider the approximating equation

$$dX_{\lambda} - \Delta(\psi_{\lambda}(X_{\lambda}) + \lambda X_{\lambda})dt = X_{\lambda}dW, \ t \in (0, T),$$
  
$$X_{\lambda}(0) = x,$$
  
(4.10)

where  $\psi_{\lambda} = \frac{1}{\lambda} (1 - (1 + \lambda \psi)^{-1}), \ \lambda > 0$ . We note that  $\psi_{\lambda} = \partial j_{\lambda}$ , where (see, e.g., [1])

$$j_{\lambda}(r) = \inf\left\{\frac{|r-\bar{r}|^2}{2\lambda} + j(\bar{r}); \ \bar{r} \in \mathbb{R}\right\}, \ \forall r \in \mathbb{R}.$$

We have the following.

**Lemma 4.5** Let  $x \in \mathcal{H}^{-1} \cap L^p(\mathbb{R}^d) \cap L^2(\mathbb{R}^d)$ ,  $p := 2m, d \ge 3$ . Then (4.10) has a unique solution

$$X_{\lambda} \in L^{2}(\Omega; C([0,T]; \mathcal{H}^{-1})) \cap L^{\infty}([0,T]; L^{p}(\Omega \times \mathbb{R}^{d})).$$

$$(4.11)$$

Moreover, for all  $\lambda, \mu > 0$ , we have

$$\mathbb{E} \sup_{0 \le t \le T} \| (X_{\lambda}(t) - X_{\mu}(t)) \|_{-1}^{2} \le C(\lambda + \mu),$$
(4.12)

$$\mathbb{E}|X_{\lambda}(t)|_{p}^{p} \leq C|x|_{p}^{p}, \ \forall t \in [0,T],$$

$$(4.13)$$

$$\mathbb{E}\int_{0}^{T}\int_{\mathbb{R}^{d}}|\psi_{\lambda}(X_{\lambda})|^{\frac{p}{m}}dt\,d\xi \leq C|x|_{p}^{p},\;\forall\lambda>0,$$
(4.14)

$$\mathbb{E}\left[\sup_{0 \le t \le T} \|X_{\lambda}(t)\|_{-1}^{2}\right] \le C \|x\|_{-1}^{2}, \ \forall \lambda > 0,$$
(4.15)

where C is independent of  $\lambda, \mu$ .

**Proof.** We consider for each fixed  $\lambda$  the equation (see (3.11))

$$dX^{\nu}_{\lambda} + (\nu - \Delta)(\psi_{\lambda}(X^{\nu}_{\lambda}) + \lambda X^{\nu}_{\lambda})dt = X^{\nu}_{\lambda}dW$$
  
$$X^{\nu}_{\lambda}(0) = x,$$
(4.16)

where  $\nu > 0$ . Let  $x \in L^2(\mathbb{R}^d) \cap L^p(\mathbb{R}^d) \cap \mathcal{H}^{-1}$ . By Claim 1 in the proof of Lemma 3.5, (4.16) has a unique solution  $X^{\nu}_{\lambda} \in L^2(\Omega; L^{\infty}([0,T]; L^2(\mathbb{R}^d))) \cap L^2(\Omega \times [0,T]; H^1(\mathbb{R}^d))$  with continuous sample paths in  $L^2(\mathbb{R}^d)$ .

As seen in the proof of Theorem 3.4, we have, for  $\nu \to 0,$ 

$$\begin{split} X_{\lambda}^{\nu} \to X_{\lambda} & \text{strongly in } L^2(\Omega; C([0, T]; H^{-1}(\mathbb{R}^d))) \\ & \text{and, by (3.14), along a subsequence also,} \\ & \text{weakly}^* \text{ in } L^2(\Omega; L^{\infty}([0, T]; L^2(\mathbb{R}^d))), \end{split}$$

where  $X_{\lambda}$  is the solution to (4.10). It remains to be shown that  $X_{\lambda}$  satisfies (4.11)–(4.15). In order to explain the ideas, we apply first (formally) Itô's formula to (4.16) for the function  $\varphi(x) = \frac{1}{p} |x|_p^p$ . We obtain

$$\frac{1}{p} \mathbb{E} |X_{\lambda}^{\nu}(t)|_{p}^{p} + \mathbb{E} \int_{0}^{t} \int_{\mathbb{R}^{d}} (\nu - \Delta) (\psi_{\lambda}(X_{\lambda}^{\nu}) + \lambda X_{\lambda}^{\nu}) |X_{\lambda}^{\nu}|^{p-2} X_{\lambda}^{\nu} ds \, d\xi$$

$$= \frac{1}{p} |x|_{p}^{p} + \frac{p-1}{2} \mathbb{E} \int_{0}^{t} \int_{\mathbb{R}^{d}} \sum_{k=1}^{\infty} \mu_{k}^{2} |X_{\lambda}^{\nu}e_{k}|^{2} |X_{\lambda}^{\nu}|^{p-2} dt \, d\xi.$$
(4.17)

Taking into account that  $X^{\nu}_{\lambda}, \psi_{\lambda}(X^{\nu}_{\lambda}) \in L^2(0,T; H^1(\mathbb{R}^d))$ , P-a.s., by Claim 1 in the proof of Lemma 3.5, we have

$$\int_0^t \int_{\mathbb{R}^d} (\nu - \Delta) (\psi_\lambda(X_\lambda^\nu) + \lambda X_\lambda^\nu) |X_\lambda^\nu|^{p-2} X_\lambda^\nu ds \, d\xi$$
$$\geq \lambda(p-1) \int_0^t \int_{\mathbb{R}^d} |\nabla X_\lambda^\nu|^2 |X_\lambda^\nu|^{p-2} d\xi \, ds,$$

and by (4.2) we have

$$\mathbb{E}\int_0^t \int_{\mathbb{R}^d} \sum_{k=1}^\infty \mu_k^2 |X_\lambda^\nu e_k|^2 |X_\lambda^\nu|^{p-2} ds \, d\xi \le C_\infty \mathbb{E}\int_0^t \int_{\mathbb{R}^d} |X_\lambda^\nu|^p d\xi \, ds < \infty.$$

Then, we obtain by (4.17) via Gronwall's lemma

$$\mathbb{E}|X_{\lambda}^{\nu}(t)|_{p}^{p} \le C|x|_{p}^{p}, \ t \in (0,T),$$
(4.18)

and, by (4.1),

$$\mathbb{E}\int_0^t \int_{\mathbb{R}^d} |\psi_\lambda(X_\lambda^\nu)|^{\frac{p}{m}} dt \, d\xi \le C |x|_p^p, \ t \in [0,T].$$

$$(4.19)$$

It should be said, however, that the above argument is formal, because the function  $\varphi$  is not of class  $C^2$  on  $L^2(\mathbb{R}^d)$  and we do not know a priori if the integral in the left side of (4.17) makes sense, that is, whether  $|X_{\lambda}^{\nu}|^{p-2}X_{\lambda}^{\nu} \in L^2(0,T; L^2(\Omega; H^1(\mathbb{R}^d)))$ . To make it rigorous, we approximate  $X_{\lambda}^{\nu}$  by a sequence  $\{X_{\lambda}^{\nu,\varepsilon}\}$  of solutions to the equation

$$dX_{\lambda}^{\nu,\varepsilon} + A_{\lambda}^{\nu,\varepsilon}(X_{\lambda}^{\nu,\varepsilon})dt = X_{\lambda}^{\nu,\varepsilon}dW,$$
  

$$X_{\lambda}^{\nu,\varepsilon}(0) = x.$$
(4.20)

Here,  $A_{\lambda}^{\nu,\varepsilon} = \frac{1}{\varepsilon} \left( I - (I + \varepsilon A_{\lambda}^{\nu})^{-1} \right)$ ,  $\varepsilon \in (0, 1)$ , is the Yosida approximation of the operator  $A_{\lambda}^{\nu}x = (\nu - \Delta)(\psi_{\lambda}(x) + \lambda x)$ ,  $\forall x \in D(A_{\lambda}^{\nu}) = H^{1}(\mathbb{R}^{d})$ . We set  $J_{\varepsilon} = (I + \varepsilon A_{\lambda}^{\nu})^{-1}$  and note that  $J_{\varepsilon}$  is Lipschitz in  $H = H^{-1}(\mathbb{R}^{d})$  as well as in all  $L^{q}(\mathbb{R}^{d})$  for  $1 < q < \infty$ . Moreover, we have

$$|J_{\varepsilon}(x)|_q \le |x|_q, \ \forall x \in L^q(\mathbb{R}^d).$$
(4.21)

(See [3], Lemma 3.1.) Since  $A_{\lambda}^{\nu,\varepsilon}$  is Lipschitz in H, equation (3.1) has a unique adapted solution  $X_{\lambda}^{\nu,\varepsilon} \in L^2(\Omega; C([0,T]; H) \text{ and by Itô's formula we have}$ 

$$\frac{1}{2}\mathbb{E}|X_{\lambda}^{\nu,\varepsilon}(t)|_{-1}^{2} \leq \frac{1}{2}|x|_{-1}^{2} + C_{1}\sum_{k=1}^{\infty}\mu_{k}^{2}\mathbb{E}\int_{0}^{t}|X_{\lambda}^{\nu,\varepsilon}(s)e_{k}|_{-1}^{2}ds,$$

which, by virtue of (jj), yields

$$\mathbb{E}|X_{\lambda}^{\nu,\varepsilon}(t)|_{-1}^{2} \leq C_{2}|x|_{-1}^{2}, \ \forall \varepsilon > 0, \ x \in H.$$

$$(4.22)$$

Similarly, since  $A_{\lambda}^{\nu,\varepsilon}$  is Lipschitz in  $L^2(\mathbb{R}^d)$  (see Lemma 4.6 below), we have also that  $X_{\lambda}^{\nu,\varepsilon} \in L^2(\Omega; C([0,T]; L^2(\mathbb{R}^2)))$  and, again by Itô's formula applied to the function  $|X_{\lambda}^{\nu,\varepsilon}(t)|_2^2$ , we obtain that

$$\mathbb{E}|X_{\lambda}^{\nu,\varepsilon}(t)|_2^2 \leq \frac{1}{2} |x|_2^2 + C_3 \sum_{k=1}^{\infty} \mu_k^2 \mathbb{E} \int_0^\infty |X_{\lambda}^{\nu,\varepsilon}(s)e_k|_2^2 ds,$$

which yields, by virtue of (jj),

$$\mathbb{E}|X_{\lambda}^{\nu,\varepsilon}(t)|_2^2 \le C_4|x|_2^2, \ \forall t \in [0,T].$$

$$(4.23)$$

**Claim 1.** For  $p \in [2, \infty)$  and  $x \in L^p(\mathbb{R}^d)$ , we have that  $X^{\nu,\varepsilon}_{\lambda} \in L^{\infty}_W([0,T]; L^p(\Omega; L^p(\mathbb{R}^d)) \cap L^2(\Omega; L^2(\mathbb{R}^d)))$ , where here and below the subscript W refers to  $(\mathcal{F}_t)$ -adapted processes.

**Proof.** For R > 0, consider the set

$$\mathcal{K}_{R} = \{ X \in L^{\infty}_{W}([0,T]; L^{p}(\Omega; L^{p}(\mathbb{R}^{d})) \cap L^{2}(\Omega; L^{2}(\mathbb{R}^{d}))), \\ e^{-p\alpha t} \mathbb{E}|X(t)|_{p}^{p} \leq R^{p}, \ e^{-2\alpha t} \mathbb{E}|X(t)|_{2}^{2} \leq R^{2}, t \in [0,T] \}.$$

Since, by (4.20),  $X_{\lambda}^{\nu,\varepsilon}$  is a fixed point of the map

$$F: X \mapsto e^{-\frac{t}{\varepsilon}}X + \frac{1}{\varepsilon}\int_0^t e^{-\frac{t-s}{\varepsilon}}J_{\varepsilon}(X(s))ds + \int_0^t e^{-\frac{(t-s)}{\varepsilon}}X(s)dW(s),$$

obtained by iteration in  $C_W([0,T]; L^2(\Omega; L^2(\mathbb{R}^d)))$ , it suffices to show that F leaves the set  $\mathcal{K}_R$  invariant for R > 0 large enough. By (4.21), we have

$$\left(e^{-p\alpha t}\mathbb{E}\left|e^{-\frac{t}{\varepsilon}}x+\frac{1}{\varepsilon}\int_{0}^{t}e^{-\frac{t-s}{\varepsilon}}J_{\varepsilon}(X(s))ds\right|_{p}^{p}\right)^{\frac{1}{p}} \leq e^{-\left(\frac{1}{\varepsilon}+\alpha\right)t}|x|_{p}+e^{-\alpha t}\int_{0}^{t}\frac{1}{\varepsilon}e^{-\frac{(t-s)}{\varepsilon}}(\mathbb{E}|X(s)|_{p}^{p})^{\frac{1}{p}}ds \qquad (4.24)$$

$$\leq e^{-\left(\frac{1}{\varepsilon}+\alpha\right)t}|x|_{p}+\frac{R}{1+\alpha\varepsilon},$$

and, similarly, that

$$\left(e^{-2\alpha t}\mathbb{E}\left|e^{-\frac{t}{\varepsilon}}x+\frac{1}{\varepsilon}\int_{0}^{t}e^{-\frac{(t-s)}{\varepsilon}}J_{\varepsilon}(X(s))ds\right|_{2}^{2}\right)^{\frac{1}{2}} \leq e^{-\left(\frac{1}{\varepsilon}+\alpha\right)t}|x|_{2}+\frac{R}{1+\alpha\varepsilon}.$$
(4.25)

Now, we set

$$Y(t) = \int_0^t e^{-\frac{(t-s)}{\varepsilon}} X(s) dW(s), \ t \ge 0.$$

We have

$$dY + \frac{1}{\varepsilon} Y dt = X dW, \quad t \ge 0,$$
  
$$Y(0) = 0.$$

Equivalently,

$$d(e^{\frac{t}{\varepsilon}}Y(t)) = e^{\frac{t}{\varepsilon}}X(t)dW(t), \ t > 0; \ Y(0) = 0.$$

By Lemma 5.1 in [19], it follows that  $e^{\frac{t}{\varepsilon}}Y$  is an  $L^p(\mathbb{R}^d)$ -valued  $(\mathcal{F}_t)$ -adapted continuous process on  $[0, \infty)$  and

$$\mathbb{E}|e^{\frac{t}{\varepsilon}}Y(t)|_p^p = \frac{1}{2}p(p-1)\sum_{k=1}^{\infty}\mu_k^2\mathbb{E}\int_0^t\int_{\mathbb{R}^d}|e^{\frac{s}{\varepsilon}}Y(s)|^{p-2}|e^{\frac{s}{\varepsilon}}X(s)e_k|^2ds.$$

This yields via Hypothesis (jj)

$$\mathbb{E}|e^{\frac{t}{\varepsilon}}Y(t)|_p^p \leq \frac{1}{2}(p-1)\mathbb{E}\int_0^t |e^{\frac{s}{\varepsilon}}Y(s)|_p^p ds + C\mathbb{E}\int_0^t |\frac{s}{\varepsilon}X(s)|_p^p ds, \ \forall t \in [0,T],$$

and, therefore,

$$\mathbb{E}|Y(t)|_p^p \le C_1 e^{-\left(\alpha + \frac{1}{\varepsilon}\right)pt} \mathbb{E}\int_0^t |e^{\frac{s}{\varepsilon}} X(s)|_p^p ds \le \frac{R^p e^{-p\alpha t} \varepsilon C_1}{p(1 + \varepsilon \alpha)}, \ \forall t \in [0, T].$$

Similarly, we get

$$e^{-2\alpha t}\mathbb{E}|Y(t)|_2^2 \le \frac{R^2 \varepsilon C_1}{2(1+\varepsilon\alpha)}, \ \forall t \in [0,T].$$

Then, by formulae (4.24), (4.25), we infer that, for  $\alpha$  large enough and  $R > 2(|x|_p + |x|_2)$ , F leaves  $\mathcal{K}_R$  invariant, which proves Claim 1.

**Claim 2.** We have, for all  $p \in [2, \infty)$  and  $x \in L^p(\mathbb{R}^d)$ , that there exists  $C_p \in (0, \infty)$  such that

$$\operatorname{ess\,sup}_{t\in[0,T]} \mathbb{E}|X_{\lambda}^{\nu,\varepsilon}(t)|_{p}^{p} \leq C_{p} \quad for \ all \ \varepsilon, \lambda, \nu \in (0,1).$$

$$(4.26)$$

**Proof.** Again invoking Lemma 5.1 in [19], we have by (4.20) that  $X_{\lambda}^{\nu,\varepsilon}$  satisfies

$$\mathbb{E}|X_{\lambda}^{\nu,\varepsilon}(t)|_{p}^{p} = |x|_{p}^{p} - p \mathbb{E}\int_{0}^{t} \int_{\mathbb{R}^{d}} A_{\lambda}^{\nu,\varepsilon}(X_{\lambda}^{\nu,\varepsilon}) X_{\lambda}^{\nu,\varepsilon} |X_{\lambda}^{\nu,\varepsilon}|^{p-2} d\xi \, ds + p(p-1) \sum_{k=1}^{\infty} \mu_{k}^{2} \mathbb{E}\int_{0}^{t} \int_{\mathbb{R}^{d}} |X_{\lambda}^{\nu,\varepsilon}|^{p-2} |X_{\lambda}^{\nu,\varepsilon}e_{k}|^{2} d\xi \, ds.$$

$$(4.27)$$

On the other hand,  $A_{\lambda}^{\nu,\varepsilon}(X_{\lambda}^{\nu,\varepsilon}) = \frac{1}{\varepsilon} \left( X_{\lambda}^{\nu,\varepsilon} - J_{\varepsilon}(X_{\lambda}^{\nu,\varepsilon}) \right)$  and so we have

$$\int_{\mathbb{R}^d} A_{\lambda}^{\nu,\varepsilon}(X_{\lambda}^{\nu,\varepsilon}) X_{\lambda}^{\nu,\varepsilon} |X_{\lambda}^{\nu,\varepsilon}|^{p-2} d\xi$$
$$= \frac{1}{\varepsilon} \int_{\mathbb{R}^d} |X_{\lambda}^{\nu,\varepsilon}|^p d\xi - \frac{1}{\varepsilon} \int_{\mathbb{R}^r} J_{\varepsilon}(X_{\lambda}^{\nu,\varepsilon}) |X_{\lambda}^{\nu,\varepsilon}|^{p-2} X_{\lambda}^{\nu,\varepsilon} d\xi.$$

Recalling (4.21), we get, via the Hölder inequality,

$$\int_{\mathbb{R}^d} A_{\lambda}^{\nu,\varepsilon}(X_{\lambda}^{\nu,\varepsilon}) X_{\lambda}^{\nu,\varepsilon} |X_{\lambda}^{\nu,\varepsilon}|^{p-2} d\xi \ge 0,$$

and so, by (4.27) and Hypothesis (jj), we obtain, via Gronwall's lemma, estimate (4.26), as claimed.  $\blacksquare$ 

Claim 3. We have, for  $\varepsilon \to 0$ ,

$$X^{\nu,\varepsilon}_{\lambda} \longrightarrow X^{\nu}_{\lambda} \text{ strongly in } L^{\infty}_{W}([0,T]; L^{2}(\Omega; H))$$

and weakly<sup>\*</sup> in  $L^{\infty}([0,T]; L^p(\Omega; L^p(\mathbb{R}^d)) \cap L^2(\Omega; L^2(\mathbb{R}^d))).$ 

**Proof.** For simplicity, we write  $X_{\varepsilon}$  instead of  $X_{\lambda}^{\nu,\varepsilon}$  and X instead of  $X_{\lambda}^{\nu}$ . Also, we set  $\gamma(r) \equiv \psi_{\lambda}(r) + \lambda r$ .

Subtracting equations (4.20) and (4.16), we get via Itô's formula and because  $A^{\nu,\varepsilon}_\lambda$  is monotone on H

$$\frac{1}{2} \mathbb{E} |X_{\varepsilon}(t) - X(t)|^2_{-1,\nu} + \mathbb{E} \int_0^t \int_{\mathbb{R}^d} (\gamma(J_{\varepsilon}(X)) - \gamma(X)) (X_{\varepsilon} - X) d\xi \, ds$$
$$\leq C \mathbb{E} \int_0^t |X_{\varepsilon}(s) - X(s)|^2_{-1,\nu} ds,$$

and hence, by Gronwall's lemma, we obtain

$$\mathbb{E}|X_{\varepsilon}(t) - X(t)|^{2}_{-1,\nu} \le C\mathbb{E}\int_{0}^{T}\int_{\mathbb{R}^{d}}|\gamma(J_{\varepsilon}(X)) - \gamma(X)||X_{\varepsilon} - X|d\xi\,ds.$$
(4.28)

On the other hand, it follows by (4.21) that

$$\int_{\Omega \times [0,T] \times \mathbb{R}^d} |J_{\varepsilon}(X)|^2 \mathbb{P}(d\omega) dt \, d\xi \le \int_{\Omega \times [0,T] \times \mathbb{R}^d} |X|^2 \mathbb{P}(d\omega) dt \, d\xi,$$

while, for  $\varepsilon \to 0$ ,

$$J_{\varepsilon}(y) \longrightarrow y \text{ in } H^{-1}, \quad \forall y \in H^{-1},$$

(because  $A_{\lambda}^{\nu,\varepsilon}$  is maximal monotone in  $H^{-1}(\mathbb{R}^d)$ ) and so,  $J_{\varepsilon}(X(t,\omega)) \longrightarrow X(t,\omega)$  in  $H^{-1}(\mathbb{R}^d)$  for all  $(t,\omega) \in (0,T) \times \Omega$ . Hence, as  $\varepsilon \to 0$ ,

$$J_{\varepsilon}(X) \longrightarrow X$$
 weakly in  $L^2(\Omega \times [0,T] \times \mathbb{R}^d)$ , (4.29)

and, according to the inequality above, this implies that, for  $\varepsilon \to 0$ ,

$$|J_{\varepsilon}(X)|_{L^{2}((0,T)\times\Omega\times\mathbb{R}^{d})}\longrightarrow|X|_{L^{2}((0,T)\times\Omega\times\mathbb{R}^{d})}.$$

Hence,  $J_{\varepsilon}(X) \longrightarrow X$  strongly in  $L^2(\Omega \times [0,T] \times \mathbb{R}^d)$  as  $\varepsilon \to 0$ . Now, taking into account that  $\gamma$  is Lipschitz, we conclude by (4.28), (4.29) and by estimates (4.23), (4.26) that Claim 3 is true.

Now, we can complete the proof of Lemma 4.5. Namely, letting first  $\varepsilon \to 0$  and then  $\nu \to \infty$  in (4.26), we get (4.13) and hence (4.14) as desired.

Now, let us prove (4.12) and (4.15). Arguing as in the proof of Theorem 3.4, we obtain

$$\frac{1}{2} |X_{\lambda}^{\nu}(t)|^{2}_{-1,\nu} + \int_{0}^{t} \int_{\mathbb{R}^{d}} (\psi_{\lambda}(X_{\lambda}^{\nu}) + \lambda X_{\lambda}^{\nu}) X_{\lambda}^{\nu} d\xi \, ds \\
= \frac{1}{2} |x|^{2}_{-1,\nu} + \frac{1}{2} \int_{0}^{t} \int_{\mathbb{R}^{d}} \sum_{k=1}^{\infty} \mu_{k}^{2} |X_{\lambda}^{\nu} e_{k}|^{2}_{-1,\nu} d\xi \, ds \\
+ \int_{0}^{t} \langle X_{\lambda}^{\nu}, X_{\lambda}^{\nu} dW \rangle_{-1,\nu} \, ds.$$
(4.30)

Keeping in mind that, by (4.3),  $|X_{\lambda}^{\nu}e_k|_{-1,\nu} \leq C|X_{\lambda}^{\nu}|_{-1,\nu}(|e_k|_{\infty}+|\nabla e_k|_d)$ , where C is independent of  $\nu$ , we obtain by the Burkholder-Davis-Gundy inequality for p = 1 (cf. the proof of Theorem 3.4)

$$\mathbb{E} \sup_{t \in [0,T]} |X_{\lambda}^{\nu}(t)|^{2}_{-1,\nu} + \lambda \mathbb{E} \int_{0}^{T} |X_{\lambda}^{\nu}|^{2}_{2} ds \leq C |x|^{2}_{-1,\nu}.$$

Taking into account that

$$\lim_{\nu \to 0} |y|_{-1,\nu} = ||y||_{-1}, \ \forall y \in \mathcal{H}^{-1},$$

we obtain, as in Theorem 3.4 (see the part following (3.26)), that

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$$\mathbb{E}\left[\sup_{t\in[0,T]} \|X_{\lambda}(t)\|_{-1}^{2}\right] + \lambda \mathbb{E}\int_{0}^{T} |X_{\lambda}(t)|_{2}^{2} dt \leq C \|x\|_{-1}^{2}, \ \forall \lambda > 0, \qquad (4.31)$$

where C is independent of  $\lambda$ . In particular, (4.15) holds.

Completely similarly, one proves (4.12). Namely, we have

$$d(X^{\nu}_{\lambda} - X^{\nu}_{\mu}) + (\nu - \Delta)(\psi_{\lambda}(X^{\nu}_{\lambda}) + \lambda X^{\nu}_{\lambda} - \psi_{\mu}(X^{\nu}_{\mu}) - \mu X^{\nu}_{\mu})dt = (X^{\nu}_{\lambda} - X^{\nu}_{\mu})dW$$

and again proceeding as in the proof of Theorem 3.4, we obtain as above that

$$\begin{split} \frac{1}{2} \, |X_{\lambda}^{\nu}(t) - X_{\mu}^{\nu}(t)|^{2}_{-1,\nu} \\ &+ \int_{0}^{t} \int_{\mathbb{R}^{d}} (\psi_{\lambda}(X_{\lambda}^{\nu}) + \lambda X_{\lambda}^{\nu} - \psi_{\mu}(X_{\mu}^{\nu}) - \mu X_{\mu}^{\nu})(X_{\lambda}^{\nu} - X_{\mu}^{\nu})d\xi \, ds \\ &= \frac{1}{2} \int_{0}^{t} \int_{\mathbb{R}^{d}} \sum_{k=1}^{\infty} \mu_{k}^{2} |(X_{\lambda}^{\nu} - X_{\mu}^{\nu})e_{k}|^{2}_{-1,\nu} ds \\ &+ \int_{0}^{t} \left\langle X_{\lambda}^{\nu} - X_{\mu}^{\nu}, (X_{\lambda}^{\nu} - X_{\mu}^{\nu})dW \right\rangle_{-1,\nu}, \ t \in [0,T]. \end{split}$$

Then, applying once again the Burkholder-Davis-Gundy inequality for p = 1, and the fact that, by Hypothesis (j),  $|\psi_{\lambda}(r)| \leq C|r|^m$ ,  $\forall \in \mathbb{R}$  with C independent of  $\lambda$ , we get, proceeding as in the proof of Theorem 3.4, that

$$\mathbb{E}\left[\sup_{t\in[0,T]}|X_{\lambda}^{\nu}(t)-X_{\mu}^{\nu}(t)|^{2}_{-1}\right]\leq C(\lambda+\mu),$$

where C is independent of  $\nu, \lambda, \mu$ . (For details, we refer to the proof of (3.10), (3.14) in [6].) Letting  $\nu \to 0$  as in the previous case, we obtain (4.12), as claimed. This completes the proof of Lemma 4.5.

Above we have used the following lemma.

**Lemma 4.6**  $A_{\lambda}^{\nu,\varepsilon}$  is Lipschitz in  $L^2(\mathbb{R}^d)$ .

**Proof.** It suffices to check that  $J_{\varepsilon}$  is Lipschitz in  $L^2(\mathbb{R}^d)$ . We set  $\gamma(r) = \psi_{\lambda}(r) + \lambda r$ . We have, for  $x, \bar{x} \in L^2(\mathbb{R}^d)$ ,

$$J_{\varepsilon}(x) - J_{\varepsilon}(\bar{x}) - \varepsilon \Delta(\gamma(J_{\varepsilon}(x)) - \gamma(J_{\varepsilon}(\bar{x}))) = x - \bar{x}.$$

Multiplying by  $\gamma(J_{\varepsilon}(x)) - \gamma(J_{\varepsilon}(\bar{x}))$  in  $L^{2}(\mathbb{R}^{d})$ , we get

$$\langle J_{\varepsilon}(x) - J_{\varepsilon}(\bar{x}), \gamma(J_{\varepsilon}(x)) - \gamma(J_{\varepsilon}(\bar{x})) \rangle_{2} \leq |\gamma(J_{\varepsilon}(x)) - \gamma(J_{\varepsilon}(\bar{x}))|_{2} |x - \bar{x}|_{2}.$$

Taking into account that  $(\gamma(r) - \gamma(\bar{r}))(r - \bar{r}) \ge L|r - \bar{r}|, \forall r, \bar{r} \in \mathbb{R}$ , and that  $\gamma$  is Lipschitz, we get

$$|J_{\varepsilon}(x) - J_{\varepsilon}(\bar{x})|_2 \le C|x - \bar{x}|_2,$$

as claimed.  $\blacksquare$ 

**Proof of Theorem 4.4 (continued).** By (4.12)-(4.15), it follows that there is a process  $X \in L^{\infty}([0,T]; L^p(\Omega \times \mathbb{R}^d))$  such that, for  $\lambda \to 0$ ,

$$X_{\lambda} \to X \quad \text{weak-star in } L^{\infty}([0,T]; L^{p}(\Omega \times \mathbb{R}^{d}))$$
  

$$\lambda X_{\lambda} \to 0 \quad \text{strongly in } L^{2}([0,T]; L^{2}(\Omega \times \mathbb{R}^{d}))$$
  

$$\psi_{\lambda}(X_{\lambda}) \to \eta \quad \text{weakly in } L^{\frac{p}{m}}([0,T] \times \Omega \times \mathbb{R}^{d})$$
  

$$X_{\lambda} \to X \quad \text{strongly in } L^{2}(\Omega; C([0,T]; \mathcal{H}^{-1})).$$

$$(4.32)$$

It remains to be shown that X is a strong solution to (1.1) in the sense of Definition 4.3.

By (4.10) and (4.32), we see that

$$dX - \Delta \eta dt = X dW, \ t \in (0, T)$$
  
X(0) = x. (4.33)

To prove that  $\eta \in \psi(X)$ , a.e. in  $\Omega \times (0,T) \times \mathbb{R}^d$ , it suffices to show that, for each  $\varphi \in C_0^{\infty}(\mathbb{R}^d)$ , we have

$$\limsup_{\lambda \to 0} \mathbb{E} \int_0^T \int_{\mathbb{R}^d} \varphi^2 \psi_\lambda(X_\lambda) X_\lambda dt \, d\xi \le \mathbb{E} \int_0^T \int_{\mathbb{R}^d} \varphi^2 \eta X \, d\xi \, dt.$$
(4.34)

Indeed, we have by convexity of  $j_{\lambda}$ 

and so, by (4.32) and (4.34), we see that

$$\mathbb{E}\int_0^T \int_{\mathbb{R}^d} \varphi^2(\eta(X-Z)) dt d\xi \ge \mathbb{E}\int_0^T \int_{\mathbb{R}^d} \varphi^2(j(X)-j(Z)) d\xi dt, \\ \forall Z \in L^p((0,T) \times \Omega \times \mathbb{R}^d),$$

because, for  $\lambda \to 0$ ,  $j_{\lambda}(Z) \to j(Z)$ , and  $j_{\lambda}(X_{\lambda}) \to j(X)$ , a.e. and thus, by Fatou's lemma

$$\liminf_{\lambda \to 0} \mathbb{E} \int_0^T \int_{\mathbb{R}^d} \varphi^2 j_\lambda(X_\lambda) d\xi \, dt \ge \mathbb{E} \int_0^T \int_{\mathbb{R}^d} \varphi^2 j(X) d\xi \, dt.$$

Now, we take  $\varphi \in C_0^{\infty}(\mathbb{R}^d)$  to be non-negative, such that  $\varphi = 1$  on  $B_N$  and  $\varphi = 0$ , outside  $B_{N+1}$  where for a given  $N \in \mathbb{N}$ ,  $B_N$  is the closed ball of  $\mathbb{R}^d$  with radius N. We get

$$\mathbb{E} \int_0^T \int_{B_{N+1}} \varphi^2(\eta(X-Z)) d\xi \, dt \ge \mathbb{E} \int_0^T \int_{\mathbb{R}^d} \varphi^2(j(X) - j(Z)) d\xi \, dt, \qquad (4.35)$$
$$\forall Z \in L^p((0,T) \times \Omega \times \mathbb{R}^d).$$

This yields

$$\mathbb{E}\int_0^T \int_{B_{N+1}} \varphi^2 \eta(X-Z) d\xi \, dt \ge \mathbb{E}\int_0^T \int_{B_{N+1}} \varphi^2 \zeta(X-Z) d\xi \, dt, \qquad (4.36)$$

for all  $Z \in L^p((0,T) \times \Omega \times B_{N+1})$  and  $\zeta \in L^{p'}((0,T) \times \Omega \times B_{N+1})$  such that  $\zeta \in \psi(Z)$ , a.e. in  $(0,T) \times \Omega \times B_{N+1}$ .

We denote by  $\widetilde{\psi} : L^p((0,T) \times \Omega \times B_{N+1}) \to L^{p'}((0,T) \times \Omega \times B_{N+1})$  the realization of the mapping  $\psi$  in  $L^p((0,T) \times \Omega \times B_{N+1})$ , that is,

$$\widetilde{\psi}(Z) = \left\{ \zeta \in L^{p'}((0,T) \times \Omega \times B_{N+1}), \ \zeta \in \psi(Z), \ \text{a.e.} \right\}.$$

Since  $\frac{m}{p} \leq p'$  with  $\frac{1}{p'} = 1 - \frac{1}{p}$ , by virtue of assumption (j),  $\tilde{\psi}$  is maximal monotone in  $L^p((0,T) \times \Omega \times B_{N+1}) \times L^{p'}((0,T) \times \Omega \times B_{N+1})$ , and so, the equation

$$J(Z) + \tilde{\psi}(Z) \ni J(X) + \eta, \qquad (4.37)$$

where  $J(Z) = |Z|^{p-2}Z$ , has a unique solution  $(Z, \eta)$  (see, e.g., [1], p. 31).

If, in (4.36), we take Z the solution to (4.37), we obtain that

$$\mathbb{E}\int_0^T \int_{B_{N+1}} \varphi^2 (J(X) - J(Z))(X - Z) dt d\xi \le 0.$$

Then, choosing  $\alpha = \frac{2}{p}$ , yields

$$\mathbb{E}\int_0^T \int_{B_{N+1}} \left( |\varphi^{\alpha}X|^{p-2} \varphi^{\alpha}X - |\varphi^{\alpha}Z|^{p-2} \varphi^{\alpha}Z \right) (\varphi^{\alpha}X - \varphi^{\alpha}Z) dt d\xi \le 0.$$

Consequently, this gives

$$\mathbb{E}\int_0^T \int_{\mathbb{R}^d} (J(\varphi^{\alpha}X) - J(\varphi^{\alpha}Z))(\varphi^{\alpha}X - \varphi^{\alpha}Z)dtd\xi \le 0.$$
(4.38)

On the other hand, we have

$$J(\varphi^{\alpha}X) - J(\varphi^{\alpha}Z) = (p-1)|\lambda\varphi^{\alpha}X + (1-\lambda)\varphi^{\alpha}Z|^{p-2}(X-Z),$$

for some  $\lambda = \lambda(X, Z) \in [0, 1]$ . Substituting into (4.38) yields

$$|\varphi^{\alpha}(X-Z)|^2 = 0$$
 a.e. in  $(0,T) \times \Omega \times B_{N+1}$ ,

Hence, X = Z on  $(0, T) \times \Omega \times B_N$ .

Coming back to (4.37), this gives  $\eta \in \psi(X)$ ,  $dtdPd\xi$ , a.e., because N is arbitrary.

To prove (4.34), we use the Itô formula in (4.16) to  $x \mapsto \frac{1}{2} \|\varphi x\|_{-1}^2$  to get, as in (4.30),

$$\begin{aligned} \frac{1}{2} \mathbb{E} \|\varphi X_{\lambda}^{\nu}(t)\|_{-1}^{2} + \mathbb{E} \int_{0}^{t} \left\langle (-\Delta)^{-1} (\nu - \Delta) (\psi_{\lambda}(X_{\lambda}^{\nu}) + \lambda X_{\lambda}^{\nu}, \varphi^{2} X_{\lambda}^{\nu}) \right\rangle ds \\ & \leq \frac{1}{2} \|\varphi x\|_{-1}^{2} + \frac{1}{2} \mathbb{E} \int_{0}^{t} \sum_{k=1}^{\infty} \mu_{k}^{2} \|\varphi X_{\lambda}^{\nu} e_{k}\|_{-1}^{2} ds. \end{aligned}$$

Then, letting  $\nu \to 0$ , we obtain

$$\frac{1}{2} \mathbb{E} \|\varphi X_{\lambda}(t)\|_{-1}^{2} + \mathbb{E} \int_{0}^{t} \left\langle \psi_{\lambda}(X_{\lambda}) + \lambda X_{\lambda}, \varphi^{2} X_{\lambda} \right\rangle_{2} ds 
\leq \frac{1}{2} \|\varphi x\|_{-1}^{2} + \frac{1}{2} \mathbb{E} \int_{0}^{t} \sum_{k=1}^{\infty} \mu_{k}^{2} \|\varphi X_{\lambda} e_{k}\|_{-1}^{2} ds.$$
(4.39)

On the other hand, by (4.33) we get similarly

$$\frac{1}{2} \mathbb{E} \|\varphi X(t)\|_{-1}^2 + \mathbb{E} \int_0^t \left\langle \eta(s), \varphi^2 X \right\rangle_2 ds \\ = \frac{1}{2} \|\varphi x\|_{-1}^2 + \frac{1}{2} \mathbb{E} \int_0^1 \sum_{k=1}^\infty \mu_k^2 \|\varphi X e_k\|_{-1}^2, \ t \in [0, T].$$

Comparing with (4.39), we obtain (4.34), as claimed.

If  $x \ge 0$ , a.e. in  $\mathbb{R}^d$ , it follows that  $X \ge 0$ , a.e. in in  $\Omega \times (0,T) \times \mathbb{R}^d$ . To prove this, one applies Itô's formula in (4.10) to the function  $x \mapsto |x^-|_2^2$  and get  $(X^{\nu}_{\lambda})^- = 0$ , a.e. in  $\Omega \times (0,T) \times \mathbb{R}^d$ . Then, for  $\nu \to 0$ , we obtain the desired result. This completes the existence proof for  $x \in L^2(\mathbb{R}^d) \cap L^p(\mathbb{R}^d) \cap \mathcal{H}^{-1}$ .

**Uniqueness.** If  $X_1, X_2$  are two solutions, we have

$$d(X_1 - X_2) - \Delta(\eta_1 - \eta_2)dt = (X_1 - X_2)dW, \ t \in (0, T),$$
  
$$(X_1 - X_2)(0) = 0,$$

where  $\eta_i \in \psi(X_i)$ , i = 1, 2, a.e. in  $\Omega \times (0, T) \times \mathbb{R}^d$ .

Applying again, as above (that is, via the approximating device) Itô's formula in  $\mathcal{H}^{-1}$  to  $\frac{1}{2} \|\varphi(X_1 - X_2)\|_{-1}^2$ , where  $\varphi \in C_0^{\infty}(\mathbb{R}^d)$ , we get that

$$\frac{1}{2} d \|\varphi(X_1 - X_2)\|_{-1}^2 - \langle \Delta(\eta_1 - \eta_2), \varphi(X_1 - X_2) \rangle_{-1} \\ = \frac{1}{2} \sum_{k=1}^{\infty} \mu_k^2 \|\varphi(X_1 - X_2)e_k\|_{-1}^2 dt + \langle (X_1 - X_2), \varphi(X_1 - X_2)dW \rangle_{-1} = 0.$$

Note that, since  $\eta_1 - \eta_2 \in L^{\frac{p}{m}}(\Omega \times (0,T) \times \mathbb{R}^d)$ , we have

$$-\mathbb{E}\int_{0}^{T} \left\langle \Delta(\eta_{1}-\eta_{2}), \varphi(X_{1}-X_{2}) \right\rangle_{-1} dt$$
$$= \mathbb{E}\int_{0}^{T} \int_{\mathbb{R}^{d}} (\eta_{1}-\eta_{2}), \varphi(X_{1}-X_{2}) dt \, d\xi \ge 0$$

and, therefore,

$$\mathbb{E}\|\varphi(X_1(t) - X_2(t))\|_{-1}^2 \le C \int_0^t \mathbb{E}\|\varphi(X_1 - X_2)\|_{-1}^2 ds, \ \forall t \in [0, T],$$

and, since  $\varphi$  was arbitrary in  $C_0^{\infty}(\mathbb{R}^d)$ , we get  $X_1 \equiv X_2$ , as claimed.

**Remark 4.7** The self-organized criticality model (1.3), that is,  $\psi(r) \equiv H(r) =$  Heaviside function, which is not covered by Theorem 4.4 for  $1 \leq d \leq 2$ , can, however, be treated in the special case

$$W(t) = \sum_{j=1}^{N} \mu_j \beta_j(t), \ \mu_j \in \mathbb{R},$$

(i.e., spatially independent noise) via the rescaling transformation  $X = e^W Y$ , which reduces it to the random parabolic equation

$$\frac{\partial}{\partial t}Y - e^{-W}\Delta\psi(Y) + \frac{1}{2}\sum_{j=1}^{N}\mu_j^2 Y = 0.$$

By approximating W by a smooth  $W_{\varepsilon} \in C^1([0,T];\mathbb{R})$  and letting  $\varepsilon \to 0$ , after some calculation one concludes that the latter equation has a unique strong solution Y. We omit the details, but refer to [5] for a related treatment.

### 5 The finite time extinction

Assume here that  $\psi$  satisfies condition (j) of the beginning of Section 4 and that W is of the form (jj). Moreover, one assumes that

$$\psi(r)r \ge \rho |r|^{m+1}, \ \forall r \in \mathbb{R},$$
(5.1)

where m is as in Hypothesis (j).

**Theorem 5.1** Let  $d \geq 3$  and  $m = \frac{d-2}{d+2}$ . Let  $x \in L^{m+1}(\mathbb{R}^d) \cap L^2(\mathbb{R}^d) \cap \mathcal{H}^{-1}$ and let  $X = X(t, \cdot)$  be the solution to (1.1) given by Theorem 4.4. We set

 $\tau = \inf\{t \ge 0; \ \|X(t, \cdot)\|_{-1} = 0\}.$ (5.2)

Then, for every t > 0,

$$X(t, \cdot) = 0, \ \forall t \ge \tau, \tag{5.3}$$

and

$$\mathbb{P}[\tau \le t] \ge 1 - \|x\|_{-1}^{1-m} \frac{C^*}{\rho \gamma^{m+1} (1 - e^{-C^*(1-m)t})}.$$
(5.4)

where  $\gamma^{-1} = \sup\{||u||_{-1}|u|_{m+1}^{-1}\}$  and  $C^* > 0$  is independent of the initial condition x.

**Proof.** We follow the arguments of [7]. The basic inequality is

$$\begin{aligned} \|X(t)\|_{-1}^{1-m} + \rho(1-m)\gamma^{m+1} \int_{r}^{t} \mathbf{1}_{[\|X(s)\|_{-1}>0]} ds \\ &\leq \|X(r)\|_{-1}^{1-m} + C^{*}(1-m) \int_{r}^{t} \|X(s)\|_{-1}^{1-m} ds \\ &+ (1-m) \int_{r}^{t} \left\langle \|X(s)\|_{-1}^{(m+1)} X(s), X(s) dW(s) \right\rangle_{-1}, \\ &\qquad \mathbb{P}\text{-a.s.}, \ 0 < r < t < \infty, \end{aligned}$$
(5.5)

where  $C^*$  is a suitable constant. (We note that, by virtue of (2.8),  $\gamma^{-1} < \infty$ .) To get (5.5), we apply the Itô formula in (4.10) to the semimartingale  $||X_{\lambda}(t)||_{-1}^2$  and to the function  $\varphi_{\varepsilon}(r) = (r + \varepsilon^2)^{\frac{1-m}{2}}$ ,  $r > -\varepsilon^2$ , where  $X_{\lambda}$  is the solution to (4.10).

We have

$$\begin{aligned} d\varphi_{\varepsilon}(\|X_{\lambda}(t)\|_{-1}^{2}) \\ + (1-m)(\|X_{\lambda}(t)\|_{-1}^{2} + \varepsilon^{2})^{-\frac{m+1}{2}} \langle X_{\lambda}(t), \psi_{\lambda}(X_{\lambda}(t)) + \lambda X_{\lambda}(t) \rangle_{2} dt \\ &= \frac{1}{2} \sum_{k=1}^{\infty} \mu_{k}^{2} \left[ \frac{(1-m)\|X_{\lambda}(t)e_{k}\|_{-1}^{2}}{(\|X_{\lambda}(t)\|_{-1}^{2} + \varepsilon^{2})^{\frac{m+1}{2}}} - (1-m^{2}) \frac{\|X_{\lambda}(t)e_{k}\|_{-1}^{2} \|X_{\lambda}(t)\|_{-1}^{2}}{(\|X_{\lambda}(t)\|_{-1}^{2} + \varepsilon^{2})^{\frac{m+1}{2}}} \right] dt \\ &+ 2 \left\langle \varphi_{\varepsilon}'(\|X_{\lambda}(t)\|_{-1}^{2}) X_{\lambda}(t), X_{\lambda}(t) dW(t) \right\rangle. \end{aligned}$$

This yields

$$\varphi_{\varepsilon}(\|X_{\lambda}(t)\|_{-1}^{2}) + \rho(1-m) \int_{r}^{t} (\|X_{\lambda}(s)\|_{-1}^{2} + \varepsilon^{2})^{-\frac{m+1}{2}} \int_{\mathbb{R}^{d}} |X_{\lambda}|^{m+1} ds \, d\xi$$
  

$$\leq \varphi_{\varepsilon}(\|X_{\lambda}(r)\|_{-1}^{2}) + C^{*} \int_{r}^{t} \|\varphi^{\frac{1}{2}}X_{\lambda}(s)\|_{-1}^{2} (\|X_{\lambda}(s)\|_{-1}^{2} + \varepsilon^{2})^{-\frac{1+m}{2}} ds$$
  

$$+2 \int_{r}^{t} \left\langle \varphi_{\varepsilon}'(\|X_{\lambda}(s)\|_{-1}^{2}) X_{\lambda}(s), X_{\lambda}(s) dW(s) \right\rangle_{-1}.$$

Now, letting  $\lambda \to 0$ , we obtain that X satisfies the estimate

$$\varphi_{\varepsilon}(\|X(t)\|_{-1}^{2}) + \rho(1-m) \int_{r}^{t} \left(\|X(s)\|_{-1}^{2} + \varepsilon^{2})^{-\frac{m+1}{2}} \int_{\mathbb{R}^{d}} |X(s,\xi)|^{m+1} d\xi \right) ds 
\leq \varphi_{\varepsilon}(\|\widetilde{X}(t)\|_{-1}^{2}) + C^{*} \int_{r}^{t} \|X(s)\|_{-1}^{2} (\|\widetilde{X}(s)\|_{-1}^{2} + \varepsilon^{2})^{\frac{m+1}{2}} ds$$

$$+ 2 \int_{r}^{t} \left\langle \varphi_{\varepsilon}'(\|X(s)\|_{-1}^{2}) X(s), X(s) dW(s) \right\rangle_{-1}.$$
(5.6)

Here, we have used the fact that, by Lemma 4.5, for  $\lambda \to 0,$ 

$$X_{\lambda} \to X \text{ in } \mathcal{H}^{-1},$$

and, by (4.32) it follows, via Fatou's lemma,

$$\liminf_{\lambda \to 0} \int_{\mathbb{R}^d} |X_\lambda|^{m+1} d\xi \ge \int_{\mathbb{R}^d} |X|^{m+1} d\xi,$$

and

$$\begin{aligned} (\|X(t)\|_{-1}^2 + \varepsilon^2)^{\frac{1-m}{2}} + \rho(1-m)\gamma^{m+1} \int_r^t (\|X(s)\|_{-1}^2 + \varepsilon^2)^{-\frac{m+1}{2}} \|X(s)\|_{-1}^{m+1} ds \\ &\leq (\|X(r)\|_{-1}^2 + \varepsilon^2)^{\frac{1-m}{2}} + C^* \int_r^t \|X(s)\|_{-1}^2 (\|X(s)\|_{-1}^2 + \varepsilon^2)^{-\frac{m+1}{2}} ds \\ &+ 2 \int_r^t \left\langle \varphi_{\varepsilon}'(\|X(s)\|_{-1}^2 X(s), X(s) dW(s) \right\rangle_{-1}, \ 0 \le r \le t < \infty, \end{aligned}$$

because, by (2.7),  $||x||_{-1} \leq \gamma^{-1} |x|_{m+1}$ ,  $\forall x \in L^{m+1}(\mathbb{R}^d)$ . Letting  $\varepsilon \to 0$ , we get (5.5), as claimed.

Now, we conclude the proof as in [7]. Namely, by (5.5), it follows that

$$e^{-C^*(1-m)t} \|X(t)\|_{-1}^{1-m} + \rho(1-m)\gamma^{m+1} \int_r^t e^{-C^*(1-m)s} \mathbb{1}_{[\|X_s\|_{-1}>0]} ds$$
  

$$\leq e^{-C^*(1-m)r} \|X(r)\|_{-1}^{1-m}$$
  

$$+ (1-m) \int_r^t e^{C^*(1-m)s} \left\langle \|X(s)\|_{-1}^{-(m+1)} X(s), X(s) dW(s) \right\rangle_{-1}$$

and, therefore,  $t \to e^{-C^*(1-m)t} \|X(t)\|_{-1}^{1-m}$  is an  $\{\mathcal{F}_t\}$  supermartingale. Hence,  $\|X(t)\|_{-1} = 0$  for  $t \ge \tau$ , because of Proposition 3.4, Chap. 2 of [22]. Moreover, taking expectation for r = 0, we get

$$e^{-C^*(1-m)t}\mathbb{E}||X(t)||_{-1}^{1-m} + \rho(1-m)\gamma^{m+1}\int_0^t e^{-C^*(1-m)s}\mathbb{P}(\tau > s)ds \le ||x||_{-1}^{1-m}.$$

This implies that

$$\mathbb{P}(\tau > t) \frac{1 - e^{-C^*(1-m)t}}{C^*(1-m)} \le \int_0^t e^{-C^*(1-m)s} \mathbb{P}(\tau > s) ds \le \frac{\|x\|_{-1}^{1-m}}{\rho(1-m)\gamma^{m+1}},$$

and so (5.4) follows. This completes the proof.  $\blacksquare$ 

**Corollary 5.2** Let  $x \in \mathcal{H}^{-1} \cap L^{m+1}(\mathbb{R}^d) \cap L^2(\mathbb{R}^d)$  be such that  $||x||_{-1} < \frac{\rho\gamma^{m+1}}{C^*}$ . Let  $\tau$  be the stopping time defined in (5.2). Then there is  $t_0 > 0$  such that  $\mathbb{P}(\tau > t_0) > 0$ . In other words, there is extinction in finite time with positive probability. **Remark 5.3** In the case of bounded domain, Theorem 5.1 remains true for  $m \in \left[\frac{d-2}{d+2}, 1\right)$  (see [7]). One might suspect that also in this case the extinction property (5.4) holds for a larger class of exponents m. However, the analysis carried out in [24] for deterministic fast diffusion equations in  $\mathbb{R}^d$  shows that the extinction property is dependent not only on the exponent m, but also on the space  $L^p(\mathbb{R}^d)$ , where the solution exists (the so called extinction space).

**Remark 5.4** • The analysis in this section holds in particular if all the coefficients  $\mu_k$  do vanish, i.e. in the deterministic framework. In that case, Theorem 5.1 implies the existence of a deterministic time  $\tau > 0$  so that

$$t \ge \tau \Rightarrow \|X(t)\|_{-1} = 0,$$

and so X(t) = 0 for all  $t \ge \tau$ .

- Let us set for instance  $\psi(u) = u^m, d \geq 3, m = \frac{d-2}{d+2}$ . Observe that  $(L^1 \cap L^\infty)(\mathbb{R}^d) \subset (L^{m+1} \cap L^2)(\mathbb{R}^d) \cap \mathcal{H}^{-1}$ . Consider for instance as initial condition  $x \in (L^1 \cap L^\infty)(\mathbb{R}^d)$ .
- By Benilan-Crandall approach, see e.g. Theorem 1 of [13], there is a solution  $u: [0,T] \times \mathbb{R}^d \to \mathbb{R}$ , of

$$dX - \Delta \psi(X)dt = 0 \text{ in } (0,T) \times \mathbb{R}^d, X(0) = x \text{ on } \mathbb{R}^d,$$
(5.7)

in the sense of distributions. u belongs to  $(L^1 \cap L^\infty)((0, T) \times \mathbb{R}^d)$  and also the  $\eta_u = \psi(u)$ . In this case, u fulfills mass conservation.

- By use of Theorem 4.4, there is another solution  $v : [0, T] \times \mathbb{R}^d \to \mathbb{R}^d$ in the sense of distributions, such that  $v \in L^p((0, T) \times \mathbb{R}^d)$ , with  $p = \max(2, 2m)$ . Also  $\eta_v = \psi(v) \in L^{\frac{p}{m}}((0, T) \times \mathbb{R}^d)$ . By Theorem 5.1, if x is small enough, there will be extinction and so v does not fulfill any mass conservation.
- In particular there is no uniqueness for (5.7) in the sense of distributions. Remark that according to [16] uniqueness is guaranteed in the class  $(L^1 \cap L^{\infty})((0,T) \times \mathbb{R}^d)$ .

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