



**HAL**  
open science

## Numerical Analysis of a reduced formulation of an elasto-acoustic scattering problem

Hélène Barucq, Juliette Chabassier, Julien Diaz, Elodie Estecahandy

► **To cite this version:**

Hélène Barucq, Juliette Chabassier, Julien Diaz, Elodie Estecahandy. Numerical Analysis of a reduced formulation of an elasto-acoustic scattering problem. WAVES 13: 11th International Conference on Mathematical and Numerical Aspects of Waves, Jun 2013, Gammarth, Tunisia. hal-00873633

**HAL Id: hal-00873633**

**<https://inria.hal.science/hal-00873633v1>**

Submitted on 16 Oct 2013

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Numerical Analysis of a reduced formulation of an elasto-acoustic scattering problem

H. Barucq<sup>1,\*</sup>, J. Chabassier<sup>1</sup>, J. Diaz<sup>1</sup>, E. Estecahandy<sup>1</sup>

<sup>1</sup> INRIA Bordeaux Sud-Ouest Research Center, Team Project Magique-3D, LMAP, Université de Pau et des Pays de l'Adour, PAU, France.

\*Emails: firstname.lastname@inria.fr

## Abstract

We present a reduced formulation for the elasto-acoustic scattering problem. The modeling is based on the On Surface Radiation Condition method. It leads to a discrete system which solution is assessed when using Discontinuous Finite Elements.

## Introduction

The numerical reconstruction of the shape of a solid immersed into a fluid is an interesting issue both from a mathematical and a practical point of view. Indeed, if engineers have solved this problem from a while in relatively simple configurations, it continues to deserve attention because its solution requires to invert a sparse linear system composed of two discretized Helmholtz equations which are known to be very sensitive to the values of the frequency [1]. The computational costs are then very high and become quickly prohibitive in particular in 3D. Solution methodologies that are able to decrease the computational burden are thus welcome. Obviously, they can be minimized by choosing suitable finite elements [2], [?]. In this work, we propose to investigate a different approach by using a reduced problem that can be solved with lower computational costs. For the construction of the reduced problem, we propose to play with the boundary condition that is used to limit the fluid domain. This is an absorbing boundary condition that is set on a surface which can be more or less far from the solid. Now in the simplest case of a sound-soft or hard obstacle, Kriegsmann *et al.* [4], have shown that it is possible to compute an approximate solution by setting the ABC directly on the surface of the scatterer. The corresponding reduced problem is then given by an equation called On Surface Radiation Condition (OSRC) that is set on the boundary of the scatterer and that requires less computations than the initial problem. The interest of OSRC methodology has been demonstrated later for solving the scattering problem of a penetrable object immersed into a fluid [5] and to the best of our knowledge, it has not been investigated in the case of an elasto-acoustic problem. The accuracy of the OSRC

method depends on the geometry of the scatterer and on the frequency regime and it can not provide an accurate solution in all the situations. Nevertheless, it can certainly quickly deliver a solution that could be used as an initial guess for the solution of the inverse problem.

## 1 Problem setting

Let  $\Omega_s$  be a bounded domain representing the solid and let  $\Gamma$  be the boundary of  $\Omega_s$ . We denote by  $n$  the normal vector defined on  $\Gamma$  and outwardly directed to  $\Omega_s$ . The solid is immersed into a fluid  $\Omega_f$  which is limited by a surface  $\Sigma$  that has been introduced for numerical reasons. We then consider the mixed boundary value problem:

$$\nabla \cdot \sigma(u) + \omega^2 \rho_s u = 0 \quad \text{in } \Omega_s \quad (1)$$

$$\Delta p + (\omega^2/c_f^2) p = 0 \quad \text{in } \Omega_f \quad (2)$$

$$\omega^2 \rho_f u \cdot n = \partial_n p + \partial_n p^{inc} \quad \text{on } \Gamma \quad (3)$$

$$\sigma(u)n = -pn - p^{inc}n \quad \text{on } \Gamma \quad (4)$$

$$\partial_n p + \alpha p - \beta \Delta_\Sigma p = 0 \quad \text{on } \Sigma \quad (5)$$

to model the behavior of the pair  $(u, p)$  representing the displacement  $u$  into the solid and the pressure  $p$  in the fluid.  $\rho_f$  and  $\rho_s$  are the density moduli of the fluid and the solid,  $c_f$  is the propagation velocity in the fluid and the positive constant  $\omega$  is the pulsation. In the following, we set  $k_f = \omega/c_f$ . The stress tensor  $\sigma(u)$  is defined by  $\sigma(u) = C\varepsilon(u) = \frac{1}{2}C(\nabla u + \nabla^t u)$ , where  $C$  is the elasticity tensor. The solid  $S$  is illuminated by the incident wave  $p^{inc}$  propagating into the fluid and impinging the surface  $\Gamma$  according to the transmission conditions. Now let  $x$  be a generic point in  $\mathbb{R}^2$  and  $\gamma_\delta$  be the level set defined by  $\gamma_\delta = \{x := \tau + \delta n\}$  where  $\tau := \tau(s)$  denotes the orthogonal projection of  $x$  onto  $\gamma_\delta$  and  $s$  is the curvilinear abscissa. We then have  $\gamma_0 = \Gamma$  and we choose  $\Sigma = \gamma_R$  for a given  $R > 0$ . By this way,  $\Gamma$  and  $\Sigma$  are parallel and the parameter  $\delta$  measures the distance between the two surfaces. The boundary condition on  $\Sigma$  is an ABC which involves the Laplace-Beltrami operator  $\Delta_\Sigma$  and the coefficient  $\alpha$  incorporates the geometry of  $\Gamma$  and depends on  $\delta$ . For instance, fol-

lowing the same ideas than in [6], we get :

$$\alpha = (\xi - \kappa_\delta/2 + ik_f)^{-1}(\xi(\kappa_\delta + ik_f) - k_f^2) \quad (6)$$

where  $\kappa_\delta$  denotes the curvature of  $\Sigma$  defined by:

$$\kappa_\delta = \frac{\kappa}{1 + \delta\kappa}$$

with  $\kappa$  the curvature of  $\Gamma$ . Regarding  $\beta$ , we have:

$$\beta = (\xi - \kappa_\delta + ik_f)^{-1}. \quad (7)$$

In the above definitions,  $\xi$  is a positive parameter which is determined empirically.

## 2 Reduced problem

The construction of a reduced formulation begins with assuming that  $\gamma$  and  $\Sigma$  are close enough to be merged. It is well-known that if  $p$  and  $\partial_n p$  can be computed on  $\Gamma$ , it is possible to reconstruct the pressure field  $p$  in the fluid by using an integral representation. Based on this remark, we propose to replace  $\partial_n p$  on  $\Gamma$  in the transmission condition (3) by using the ABC (5). By this way, we get the reduced problem composed of Eq.(1) combined with a modified version of (3) given by

$$\omega^2 \rho_f u \cdot n = -\alpha p + \beta \Delta_\Gamma p + \partial_n p^{inc}, \quad (8)$$

and Eq.(4). The solution to the reduced problem is then given by the pair  $(u, p|_\Gamma)$  and the computation are performed in  $\Omega_s$  only.

To compute the solution to the elasto-acoustic scattering problem, we thus propose to solve a discrete system related to Eq. (1) set in the bounded domain  $\Omega_s$  combined with the boundary conditions (8) and (4) at first. The pressure field can next be reconstructed by computing  $\partial_n p|_\Gamma$  thanks to (5) set on  $\Gamma$ , which enables to get the expression of  $p$  in  $\Omega_f$  thanks to the Kirchhoff integral formulation. While (4) is included in the variational formulation of (1), Condition (8) is weakly taken into account by a surfacic variational formulation on  $\Gamma$ .

To get an approximate solution of this problem, we propose to use an Interior Penalty Discontinuous Galerkin (IPDG) method. We denote by  $\Omega_h$  (resp.  $\Gamma_h$ ) a triangulation of  $\Omega$  (resp.  $\Gamma$ ) and we consider the spaces

$$V_h = \{u \in L^2(\Omega) \mid u|_T \in P^r(T), \forall T \in \Omega_h\} \text{ and}$$

$$W_h = \{u \in L^2(\Gamma) \mid p|_\Sigma \in P^r(\Sigma), \forall \Sigma \in \Gamma_h\},$$

where  $P^r(T)$  denotes the space of polynomials of degree  $r$  on  $T$ . We denote by  $\Phi = (\phi_i)_{i=1..N_s}$  (resp.  $\Psi = (\psi_i)_{i=1..N_f}$ ) a basis of  $V_h$  (resp.  $W_h$ ). It is worth noting that  $N_f$  is very small compared to  $N_s$  since we only compute an approximation of  $p$  on the boundary  $\Gamma$ . The linear system to be solved reads as

$$\begin{aligned} (M_s + K_s)U + B_{sf}P &= F_1, \\ B_{fs}U + (M_f + K_f)P &= F_2, \end{aligned} \quad (9)$$

where  $U$  and  $P$  are two vectors of size  $N_s$  and  $N_f$  containing the components of the approximation of  $u$  and  $p$  in the basis  $\Phi$  and  $\Psi$ .  $M_s$  and  $M_f$  are block-diagonal mass matrices of size  $N_s \times N_s$  and  $N_f \times N_f$  defined by

$$(M_s)_{i,j} = \int_\Omega \phi_i \cdot \phi_j d\Omega \text{ and } (M_f)_{i,j} = \int_\Gamma \alpha \psi_i \psi_j d\Gamma.$$

$K_s$  and  $K_f$  are stiffness matrices of size  $N_s \times N_s$  and  $N_f \times N_f$  defined by

$$(K_s)_{i,j} = a(\phi_i, \phi_j) \text{ and } (K_f)_{i,j} = b(\psi_i, \psi_j),$$

where  $a(.,.)$  and  $b(.,.)$  are the bilinear forms obtained by the IPDG discretization of the operators  $\nabla \cdot \sigma(u)$  and  $-\beta \Delta_\Gamma$ . Finally,  $B_{sf}$  and  $B_{fs}$  are two coupling matrices of size  $N_s \times N_f$  and  $N_f \times N_s$  defined by

$$(B_{sf})_{i,j} = -\int_\Gamma \phi_i \cdot \psi_j n d\Gamma \text{ and } (B_{fs})_{i,j} = \omega^2 \rho_f (B_{sf})_{j,i}.$$

We present numerical results that illustrate the performance of the reduced model. We also assess the impact of parameter  $\delta$  on the accuracy of the solution.

## References

- [1] I. M. Babuska, and S. A. Sauter, *SIAM Review*, **42(3)** (2000), pp. 451–484.
- [2] P. Monk and D.Q. Wang, *Comput Methods Appl Mech Eng*, **175** (1999), pp. 121–@136.
- [3] M. Amara, H. Calandra, R. Djellouli, and M. Grigoroscuta, *Computers and Structures* (2012), **106-107**, pp. 258–272.
- [4] G. A. Kriegsmann, A. Taflove and K. R. Umashankar, *IEEE Trans. Antennas Propag.*, **35** (1987), pp. 153–161
- [5] X. Antoine, H. Barucq, and L. Vernhet, *Asymptotic Analysis*; **26**, 3-4 (2001), pp. 257–283.
- [6] H. Barucq, J. Diaz and V. Duprat, *Communications in Computational Physics* **11**, 2 (2012) pp. 674-690.