Data assimilation in glaciology
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Data assimilation in glaciology

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1 Data assimilation in glaciology

1.1 Introduction

In this short paper, we will give one example of an inverse problem in glaciology. This problem is fairly simple to state: how to infer a climatic scenario (i.e. how to reconstruct past polar temperature) from ice volume records?

Ice-sheet dynamics is quite complex. Ice is a non-newtonian fluid, its flow follows a highly non-linear Stokes equation. One key factor of ice dynamics is surface mass balance, that is the amount of ice accumulating at the surface (from snow precipitation) minus the amount of melting ice. This mass budget is related to temperature at the surface of the ice-sheet, depending in particular on temperature. See textbooks (Greve and Blatter, 2009) and (Paterson, 1981) for detailed modelling of ice-sheets.

Ice volume observations are available from oceanic records, that give information about sea level, and therefore about the amount of water stored in ice-sheets, ice-caps and glaciers.

As far as we know, there are very few past works on this subject. (Bintanja, van de Wal and Oerlemans, 2004) and (Bintanja, van de Wal and Oerlemans, 2005) use some a very simple correction method, equivalent to the nudging method, with quite good results.

The idea of this work is to explore the ability of the adjoint method to solve the inverse problem of reconstructing past temperature given all available observations. We start here with a simplified ice-sheet model and perform twin experiments.

We present the ice-sheet model in section 2, the adjoint method in section 3, and give numerical illustrations in section 4.

1.2 Ice-sheet model

1.2.1 Ice dynamics processes

As the full Stokes model is very costful to run over paleoclimatology time scales (thousands of years), large scale ice-sheet models have been developed, using the so-called Shallow Ice Approximation (SIA). Such models are able to reproduce ice-sheet evolution over hundreds of thousands of years, for large ice-sheets such as Antarctica, Greenland, Laurentide and Fennoscandian ice-sheets. One such model, GRISLI, is developed by LGGE (Ritz, Rommelaere and Dumas, 2001). It is a three dimensional thermomechanically coupled model, whose main physical processes are summed up in figure 1.1.
1.2.2 Model equations

As GRISLI model is still costful to run, we will use a simplified, flowline, isothermal prototype of this model, on a flat bedrock topography, without floating shelves, which is faster to run, in order to study the paleoclimate inverse problem. Flowline means that the model is one-dimensional in space.

The first equation is the mass balance equation:

\[
\frac{\partial H}{\partial t} = \dot{b}_m - \frac{\partial (U H)}{\partial x}, \quad H_{|t=0} = H_0
\]

with

- \( t \) time, \( x \) latitude
- \( H(t, x) \) ice thickness, \( H_0(x) \) initial ice thickness
- \( U(t, x) \) Eulerian velocity averaged over the ice thickness
- \( \dot{b}_m(t, x) \) surface mass balance rate

Vertically averaged ice velocity \( \overline{U}(t) \) is a diagnostic variable, that is to say it depends only on geometry variables \( (S, H) \) at time \( t \), and it is given by the following equation:

\[
\overline{U} = u_B - a_1 \frac{\partial S H^2}{\partial x} - a_2 \left( \frac{\partial S}{\partial x} \right)^3 \frac{H^4}{3}
\]

\[
S = B + H \quad H \geq 0
\]

with
1.2.3 Climatic scenario and surface mass balance

The surface mass balance is given as a function of the ice surface temperature:

\[
\begin{align*}
\dot{b}_m &= \text{Acc} + \text{Abl} \\
\text{Acc} &= f(T_s) \\
\text{Abl} &= g(T_s)
\end{align*}
\]

with

- \(\text{Acc}(t, x)\) accumulation rate
- \(\text{Abl}(t, x)\) ablation rate
- \(T_s(t, x)\) surface temperature
- \(f, g\) given functions

See figure 1.2 to see how accumulation and ablation depend on surface temperature.

And finally, the surface temperature itself depends on latitude \(x\), on altitude \(S\), and on the (poorly known) temperature at the northern point at sea-level height \(T_{\text{clim}}(t)\):

\[
T_s(t, x) = T_{\text{clim}}(t) + b \cdot x + c \cdot S(t, x)
\]

with \(b, c\) given coefficients.

The aim of the problem is to infer \(T_{\text{clim}}(t)\) from volume measurements.

1.3 Adjoint method and adjoint model

We refer to main and specialized lectures regarding variational data assimilation and adjoint code construction.
1.3.1 Observations and control variables

As we said before, we will use ice volume observations, available from sea level records, see figure 1.3 from (Paillard and Parrenin, 2004) as an example. As we do twin experiments, we will generate ourselves our data, from a reference “true state”.

In the framework of four-dimensional variational assimilation, we write our inverse problem as a minimization problem: we want to find the “best” temperature $T_{clim}(t)$, so that the misfit between modeled and observed volume is smallest. This means that our control variable is the climatic temperature $T_{clim}(t)$, and our cost function writes as follows:

$$ J(T_{clim}) = J^b + \int_{t_0}^{t_f} \|Vol(t) - Vol^{obs}(t)\|^2 dt, \quad Vol(t) = \int_x H(t, x) dx $$

where the background term $J^b$ is a classical regularising term, involving a background temperature and background error covariance matrix:

$$ J^b(T_{clim}) = \|T_{clim} - T_{clim}^b\|_k^2 $$

The norm $\|.\|_k$ is classically given by:

$$ \|X\|_k^2 = \int_{t=t_0}^{t_f} \int_{s=t_0}^{t_f} k(s, t) X(s) X(t) ds dt $$

with $k$ a regularizing kernel. In discrete form, this would writes:

$$ \|X\|_{B^{-1}}^2 = X^T B^{-1} X $$

where $B$ is the background error covariance matrix, and $X$ contains $T_{clim}(t_i) - T_{clim}^b(t_i)$ for all discrete control times $t_i$.

1.3.2 Adjoint code validation

We minimize our cost function using quasi-Newton descent algorithm. This requires the gradient of $J$, obtained thanks to the adjoint model. As our flowline model is simple and in Matlab, we implemented by hand the adjoint code using classical recipes.
for adjoint code construction, see e.g. \cite{Giering1998}. To validate our gradient, we perform first- and second-order tests, that is we check that the 1st and 2nd order Taylor approximation formula

\[ J(x + \alpha z) = J(x) + \alpha \nabla J(x) z + \frac{\alpha^2}{2} z^T \nabla^2 J(x) z + o(\alpha^2) \]

is valid. To do so we form

\[ \delta = \nabla J(x) z, \quad \tau = \frac{J(x + \alpha z) - J(x)}{\alpha}, \quad \varepsilon = \frac{\tau(\alpha, z) - \delta(z)}{|\delta(z)|}, \quad r = \frac{\tau(\alpha, z) - \delta(z)}{\alpha} \]

and we check that \( \varepsilon \) tends to zero as \( \alpha \) tends to zero, and \( r \) tends to a constant depending only on \( z \), see figure 1.4 for an illustration with our model.

1.4 Numerical results on twin experiments

In order to validate our method and adjoint model, we present a first illustration on twin experiments (see figures 1.5 and 1.6). The principle of such experiment is as follows: first we choose a reference climatic scenario \( T_{\text{clim}} \), called the “true” temperature forcing. Second we perform a direct run of the model with the “true” forcing and we generate synthetic observations (figure 1.5). And finally, we forget that we know the true temperature, we start with a constant first guess for \( T_{\text{clim}} \), and we perform data assimilation thanks to the adjoint method, in order to get an assimilated \( T_{\text{clim}} \). We can then compare our result with the true \( T_{\text{clim}} \) (figure 1.6).
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![Fig. 1.5](image1.jpg) Twin assimilation results: True climatic scenario $T_{\text{clim}}$ (top), generated volume data (bottom).

![Fig. 1.6](image2.jpg) Twin assimilation results: $T_{\text{clim}}$ recovered thanks to data assimilation.
References


Greve, Ralph and Blatter, Heinz (2009, August). *Dynamics of ice-sheets and glaciers*. Springer.


