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Abstract: In this paper we consider an aerial vehicle equipped with a monocular camera and inertial sensors. Additionally, a laser pointer is mounted on the vehicle and it produces a laser spot. The laser spot is observed by the monocular camera and it is the unique point feature used in the proposed approach. We focus our attention to the case when the vehicle moves in proximity of a planar surface and in particular when the laser spot belongs to this surface. The paper provides two main contributions. The former is the analytical derivation of all the observable modes, i.e. all the physical quantities that can be determined by only using the inertial data and the camera observations of the laser spot during a short time-interval. Specifically, it is shown that the observable modes are: the distance of the vehicle from the planar surface; the component of the vehicle speed, which is orthogonal to the planar surface; the relative orientation of the vehicle with respect to the planar surface; the orientation of the planar surface with respect to the gravity. The second contribution is the introduction of a simple recursive method to perform the estimation of all the aforementioned observable modes. This method is based on a local decomposition of the original system, which separates the observable modes from the rest of the system. The method is validated by using synthetic data. Additionally, preliminary tests with real data are provided and more complete experiments are in progress. The presented approach can be integrated in the framework of autonomous take-off and landing, safe touch-down and low altitude manoeuvres even in dark or featureless environment.

Key-words: Sensor Fusion, Inertial Sensors, Vision, Non linear Observability, Aerial Robotics

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Résumé : Dans ce document nous considérons un véhicule aérien équipé avec une caméra monoculaire et capteurs inertiels. En plus, un pointeur laser est monté sur le véhicule et il produit un spot laser. Le spot laser est observé par la caméra monoculaire et c'est l'unique point feature utilisé dans l'approche proposée. Nous nous concentrons sur le cas lorsque le véhicule se déplace à proximité d'une surface plane et en particulier lorsque le laser spot appartient à cette surface. Ce document contient deux principales contributions. La première est la dérivation analytique de tous les modes observables, c'est à dire toutes les quantités physiques qui peuvent être déterminés en utilisant uniquement les données inertielles et les observations du laser spot par la caméra pendant un court intervalle de temps. Plus précisément, les modes observables sont: la distance du véhicule depuis la surface plane; le composant de la vitesse du véhicule, qui est orthogonale à la surface plane; l'orientation relative du véhicule par rapport à la surface plane; l'orientation de la surface plane par rapport à la gravité. La deuxième contribution est l'introduction d'une méthode récursive simple pour effectuer l'estimation de tous les modes observables susmentionnées. Cette méthode est basée sur une décomposition locale du système d'origine, qui sépare les modes observables du reste du système. La méthode est validée en utilisant des données synthétiques. En plus, des essais préliminaires avec des données réelles sont fournis et des expériences sont en cours. L'approche présentée peut être intégré dans le cadre de départ et atterrissage autonome, manuvres a basse altitude même dans un environnement sombre ou sans features.

Mots-clés : Fusion Sensoriel, Capteurs inertiels, Vision, Observabilité, Robotique Aérienne

1 Introduction

In recent years, vision and inertial sensing have received great attention by the mobile robotics community. These sensors require no external infrastructure and this is a key advantage for robots operating in unknown environments where GPS signals are shadowed. Additionally, these sensors have very interesting complementarities and together provide rich information to build a system capable of vision-aided inertial navigation and mapping.

A special issue of the International Journal of Robotics Research has recently been devoted to the problem of fusing vision and inertial data [6]. In [5], a tutorial introduction to the vision and inertial sensing is presented. This work provides a biological point of view and it illustrates how vision and inertial sensors have useful complementarities allowing them to cover the respective limitations and deficiencies. In [1], these sensors are used to perform egomotion estimation. The sensor fusion is obtained by an Extended Kalman Filter (EKF) and by an Unscented Kalman Filter (UKF). The approach proposed in [7] extends the previous one by also estimating the structure of the environment where the motion occurs. In particular, new landmarks are inserted on line into the estimated map. This approach has been validated by conducting experiments in a known environment where a ground truth was available. Also, in [21] an EKF has been adopted. In this case, the proposed algorithm estimates a state containing the robot speed, position and attitude, together with the inertial sensor biases and the location of the features of interest. In the framework of airbone SLAM, an EKF has been adopted in [12] to perform 3D-SLAM by fusing inertial and vision measurements. It was observed that any inconsistent attitude update severely affects any SLAM solution. The authors proposed to separate attitude update from position and velocity update. Alternatively, they proposed to use additional velocity observations, such as air velocity observation. More recently, a vision based navigation approach in unknown and unstructured environments has been suggested [3]. Natraj et al. [23] proposed a vision based approach, close to structured light, for roll, pitch and altitude estimation of UAV. They use a fisheve camera and a laser circle projector, assuming that the projected circle belongs to a planar surface. The latter must be orthogonal to the gravity vector in order to allow the estimation of the aforementioned quantities. The attitude estimation of the planar surface becomes crucial in order to extend the operational environment of UAVs. Shipboard operations, search and rescue cooperation between ground and aerial robots, low altitude manoeuvres, require to attenuate the position error and to track the platform attitude.

Recent works investigate the observability properties of the vision-aided inertial navigation system [9], [10], [11], [18], [19] and [20]. In particular, in [19], the observable modes are expressed in closed-form in terms of the sensor measurements acquired during a short time-interval.

In this paper we consider a vehicle which accomplishes a 3D-trajectory in the surrounding of a planar surface. The vehicle is equipped with a monocular camera and inertial sensors. This is the typical navigation problem in an indoor environment or in a city-like environment. All the approaches previously mentioned, require to extract natural features from the images provided by the camera and in particular to detect the same features in different images. The feature matching task becomes critical in outdoor environment because of possible illumination changes. In order to significantly reduce the computational burden required to perform these tasks and to make the feature matching more robust, we introduce a *virtual* feature by equipping our vehicle with a laser pointer. The laser beam produces a laser spot on the planar surface. This laser spot is observed by the monocular camera and it is the unique point feature used by the proposed approach.

To the best of our knowledge, this problem has never been considered so far. Compared to classical vision and IMU data fusion problems, the feature is moving in the environment but we exploit the hypothesis that it moves on a planar surface. The first question which arises is to understand which are the observable modes, i.e. the physical quantities that can be determined by only using the inertial data and the camera observation of the laser spot during a short time-interval. The results provided in section 3 address precisely this issue (more details about the analytical computation can be found in appendix B). Then, the second step we consider is to analytically determine the link between the observable modes and the sensor data. This is obtained by performing a local decomposition of the original system (section 4). This decomposition separates the observable modes from the rest of the system and will allow us to introduce a simple recursive method to perform the estimation of all the observable modes (second part of section 4). The method is validated by using synthetic data (section 5). Preliminary tests with real data are also provided and more complete experiments are in progress.

2 The Considered System

Let us consider an aerial vehicle equipped with a monocular camera and IMUsensors. The vehicle is also equipped with a laser pointer. The configuration of the laser pointer in the camera reference frame is known. The vehicle moves in the surrounding of a planar surface and we assume that the laser spot produced by the laser beam belongs to this planar surface (see fig. 1). The position and the orientation of this planar surface are unknown. The camera observations consist in the position of the laser spot in the camera frame up to a scale factor. The IMU consists of three orthogonal accelerometers and three orthogonal gyroscopes. We assume that the monocular camera is intrinsically calibrated and that the transformations among the camera frame and the IMU frames are known (we can assume that the vehicle frame coincides with the camera frame). The *IMU* provides the vehicle angular speed and acceleration. Actually, regarding the acceleration, the one perceived by the accelerometer (\mathbf{A}) is not simply the vehicle acceleration (A_v) . It also contains the gravity acceleration (G). In particular, we have $A = A_v - G$ since, when the camera does not accelerate (i.e. A_v is zero) the accelerometer perceives an acceleration which is the same of an object accelerated upward in the absence of gravity.

We will use uppercase letters when the vectors are expressed in the local frame and lowercase letters when they are expressed in the global frame. Hence, regarding the gravity we have: $\boldsymbol{g} = [0, 0, -g]^T$, being $g \simeq 9.8 \ ms^{-2}$.

Finally, we will adopt a quaternion to represent the vehicle orientation. Indeed, even if this representation is redundant, it is very powerful since the dynamics can be expressed in a very easy and compact notation [13].

Our system is characterized by the state $[\mathbf{r}, \mathbf{v}, q]^T$ where $\mathbf{r} = [r_x, r_y, r_z]^T$ is the 3D vehicle position, \mathbf{v} is its time derivative, i.e. the vehicle speed in



Figure 1: Quadrotor equipped with a monocular camera, IMU and a laser pointer. The laser spot is on a planar surface and its position in the camera frame is obtained by the camera up to a scale factor.

the global frame $(\boldsymbol{v} \equiv \frac{d\boldsymbol{r}}{dt})$, $q = q_t + iq_x + jq_y + kq_z$ is a unitary quaternion (i.e. satisfying $q_t^2 + q_x^2 + q_y^2 + q_z^2 = 1$) and characterizes the vehicle orientation. The analytical expression of the dynamics and the camera observations can be easily provided by expressing all the 3D vectors as imaginary quaternions. In practice, given a 3D vector $\boldsymbol{w} = [w_x, w_y, w_z]^T$ we associate with it the imaginary quaternion $w_q \equiv 0 + iw_x + jw_y + kw_z$. The dynamics of the state $[r_q, v_q, q]^T$ are:

$$\begin{cases} \dot{r}_q = v_q \\ \dot{v}_q = qA_{vq}q^* = qA_qq^* + g_q \\ \dot{q} = \frac{1}{2}q\Omega_q \end{cases}$$
(1)

being q^* the conjugate of q, $q^* = q_t - iq_x - jq_y - kq_z$ and Ω the angular velocity.



Figure 2: The original camera frame XYZ, the chosen camera frame X'Y'Z'and the laser module at the position $[L_x, L_y, 0]$ and the direction (θ, ϕ) (in the original frame) and position [L, 0, 0] and the direction (0, 0) (in the chosen camera frame).

We derive the expression of the camera observation consisting in the position of the laser spot in the camera frame up to a scale factor. The laser spot is on a planar surface whose configuration is unknown. Without loss of generality, we choose the camera frame with the z-axis parallel to the laser pointer (see figure 2). In addition, the camera frame is such that the laser beam intercept the xy-plane in [L, 0, 0]. In appendix A we introduce a simple and efficient method in order to determine the parameter L together with the rotation to transform vectors from the original camera frame into the chosen camera frame.

Finally, we characterize the planar surface in the global frame with the equation z = ky, where k is an unknown parameter.

In these settings, by carrying out analytical computation (which uses the basic quaternion rules) we obtain the analytical expression of the position $[X_s, Y_s, X_s]$ of the laser spot in the camera reference frame. We have:

$$\begin{cases} X_s = L \\ Y_s = 0 \\ Z_s = \frac{r_z + 2q_z q_x L - 2q_y kq_x L - 2q_y q_t L - 2q_t L kq_z - kr_y}{2kq_z q_y - 2kq_t q_x - q_z^2 - q_t^2 + q_y^2 + q_x^2} \end{cases}$$
(2)

The camera provides the vector $[X_s, Y_s, Z_s]$ up to a scale factor. This is equivalent to the two ratios $\frac{X_s}{Z_s}$ and $\frac{Y_s}{Z_s}$. Hence, since the latter is identically zero, the camera observation is given by $h_{cam} = \frac{X_s}{Z_s}$, which is:

$$h_{cam} = \frac{L(2kq_zq_y - 2kq_tq_x - q_z^2 - q_t^2 + q_y^2 + q_x^2)}{r_z + 2q_zq_xL - 2q_ykq_xL - 2q_yq_tL - 2q_tLkq_z - kr_y}$$
(3)

3 Observability Properties

We investigate the observability properties of the system whose dynamics are given in (1) and whose observation function is given in (3). We have also to consider the constraint $q^*q = 1$. This can be dealt as a further observation (system output):

$$h_{const}(r_q, v_q, q) = q^*q \tag{4}$$

Finally, we want to investigate whether the parameter k is identifiable or not. This is done by performing an observability analysis on the extended state $\boldsymbol{S} = [r_q, v_q, q, k]^T$, whose dynamics are given in (1) and by the additional equation $\dot{k} = 0$.

We apply the method introduced in [17]. This will allow us to detect all the observable modes, i.e. all the physical quantities that can be determined by only using the information contained in the data provided by the IMU and the camera in a given time-interval. In appendix B we perform this analytical computation and we show that the system has six observable modes. Additionally, in appendix B, we provide the physical meaning of these observable modes. Specifically, we found convenient to express the observable modes in a new global reference frame, denoted with \tilde{x} , \tilde{y} , \tilde{z} . In this frame the \tilde{z} -axis coincides with the axis normal to the planar surface. From now on, we will adopt this global frame to characterize the vehicle configuration. The \tilde{x} , \tilde{y} , \tilde{z} -frame is obtained by rotating the x, y, z-frame about the x-axis of the angle $\arctan(k)$. Hence, it is characterized by the quaternion:

$$p = \cos\left(\frac{\arctan(k)}{2}\right) + i\sin\left(\frac{\arctan(k)}{2}\right) \tag{5}$$

The state in this frame is $\tilde{\boldsymbol{S}} = [\tilde{r}_q \ \tilde{v}_q \ \tilde{q}, \ k]^T$, where:

$$\tilde{q} = p^* q \quad \tilde{r}_q = p^* r_q p \quad \tilde{v}_q = p^* v_q p \tag{6}$$

The observable modes are (see the derivation in appendix B):

$$m_{1} = \tilde{r}_{z}$$

$$m_{2} = \tilde{v}_{z}$$

$$m_{3} = 2(\tilde{q}_{t}\tilde{q}_{x} + \tilde{q}_{y}\tilde{q}_{z})$$

$$m_{4} = 2(\tilde{q}_{t}\tilde{q}_{y} - \tilde{q}_{x}\tilde{q}_{z})$$

$$m_{5} = k$$

$$m_{6} = \tilde{q}^{*}\tilde{q}$$

$$(7)$$

 m_1 is the \tilde{z} -coordinate of the vehicle, m_2 the component of the vehicle speed along the \tilde{z} -axis, m_3 and m_4 are related to the roll and pitch angles of the vehicle in this frame. In particular, the roll is $\arctan\left(\frac{m_3}{\sqrt{1-m_3^2-m_4^2}}\right)$ and the pitch is $\arcsin(m_4)$. m_5 is related to the orientation of the $\tilde{x}\tilde{y}$ -plane with respect to the gravity. In particular, the \tilde{z} -axis makes an angle $\arctan(k) = \arctan(m_5)$ with the gravity. m_6 is trivially the magnitude of the quaternion \tilde{q} , which is 1 since it describes a rotation.

By summarizing the results of the observability analysis performed in this section (and in appendix B) we say that the information contained in the data provided by the IMU and the camera during a given time-interval, allows us to determine the six modes m_1, \ldots, m_6 . For this reason, in the rest of the paper, we will focus our attention only on these six quantities (actually, on the first five, since m_6 trivially expresses the constraint of having a unitary quaternion).

4 Local Decomposition and Recursive Estimation

The goal of this section is to provide a method able to estimate the observable modes in (7). To achieve this goal, the first step is to determine the link between the observable modes and the sensor data. By adopting the terminology introduced in [8], we have to perform a local decomposition of our system. We remind the reader that the local decomposition is the extension of the Kalman canonical decomposition [4] to the case of a non linear system. It consists in writing the equations characterizing the dynamics and the observation only in terms of the observable modes. We also remind the reader that in the non linear case it is often impossible to characterize the system with a unique decomposition. The decomposition only holds in a local region of the space of states. This is the reason why it is called *local* decomposition. To cover the entire space of states more than one decomposition is required (see [8]). In the following, we will show that for our system the number of decompositions is two. We first provide the dynamics of the state $\tilde{S} = [\tilde{r}_q \ \tilde{v}_q \ \tilde{q}, \ k]^T$. We obtain (appendix B):

$$\begin{cases} \tilde{r}_q = \tilde{v}_q \\ \tilde{v}_q = \tilde{q}A_q\tilde{q}^* + p^*g_qp \\ \tilde{q} = \frac{1}{2}\tilde{q}\Omega_q \\ \tilde{k} = 0 \end{cases}$$

$$\tag{8}$$

A local decomposition for the dynamics is:

$$\begin{cases} \dot{m}_1 = m_2 \\ \dot{m}_2 = -m_4 A_x + m_3 A_y + \xi(m_3, \ m_4) A_z + g_z \\ \dot{m}_3 = \Omega_x \ \xi(m_3, \ m_4) + \Omega_z m_4 \\ \dot{m}_4 = \Omega_y \ \xi(m_3, \ m_4) - \Omega_z m_3 \\ \dot{g}_z = 0 \end{cases}$$
(9)

where g_z is the component of the gravity along the \tilde{z} -axis, i.e. $g_z = -g \cos(\arctan(k)) = \frac{-g}{\sqrt{1+k^2}} = \frac{-g}{\sqrt{1+m_5^2}}$; the function $\xi(m_3, m_4)$ depends on the original state and in particular changes its sign depending on the sign of $\tilde{q}_x^2 + \tilde{q}_y^2 - \frac{1}{2}$:

$$\xi(m_3, m_4) \equiv \begin{cases} \sqrt{1 - m_3^2 - m_4^2} & if \quad \tilde{q}_x^2 + \tilde{q}_y^2 < \frac{1}{2} \\ -\sqrt{1 - m_3^2 - m_4^2} & if \quad \tilde{q}_x^2 + \tilde{q}_y^2 > \frac{1}{2} \end{cases}$$
(10)

Hence, as previously said, we have two local decompositions for our original system. The validity of (9) can be checked by using (7) and (8). Note that deriving (9) is troublesome. In contrast, checking its validity is very simple since only demands to perform differentiation.

To complete the local decomposition we need to express the camera observation function in (3) in terms of the observable modes. We obtain:

$$h_{cam} = \frac{L\xi(m_3, m_4)}{m_4 L - m_1} \tag{11}$$

The validity of (11) can be checked by using (3), (13), (14), (7) and (10).

The equations (9) and (11) represent a local decomposition for our system. They provide the analytical expression of the link between the observable modes and the sensor data. Specifically, equation (9) provides the link between the observable modes and the IMU data. Equation (11) provides the link between the observable modes and the data delivered by the monocular camera. Having these equations allows us to perform the estimation of the state $[m_1, m_2, m_3, m_4, g_z]$. An efficient and simple approach is obtained by using an Extended Kalman Filter, EKF. To implement this filter it suffices to compute the Jacobian of the dynamics in (9) and the Jacobian of the observation function in (11) (see [2]). We provide these Jacobians in appendix C.

5 Performance Evaluation

We evaluate the performance of the proposed strategy by using both synthetic and real data. The advantage of simulations is that the ground truth is perfectly known and this allows us a quantitative evaluation of the proposed strategy.

5.1 Simulations

We simulate many different trajectories in 3D and many different scenarios corresponding to different orientation of the planar surface. For all the simulations we use the proposed strategy to estimate the observable modes, i.e.:

- 1. the distance of the vehicle from the planar surface $(d = |m_1|)$;
- 2. the component of the vehicle speed orthogonal to the planar surface $(v_o = m_2);$
- 3. the roll (R) and the pitch (P) angles in the \tilde{x} , \tilde{y} , \tilde{z} -frame (i.e. the frame where the \tilde{x} , \tilde{y} -plane coincides with the planar surface);
- 4. the orientation of the plane with respect to the gravity (α) .

Specifically, in all the simulations the values of the estimated d, v_o , R, P and α will be compared with the ground truth values.

5.1.1 Simulated Trajectories

The trajectories are generated by randomly generating the linear and angular acceleration of the vehicle at 100 Hz. In particular, at each time step, the three components of the linear acceleration and the angular speed are generated as Gaussian independent variables whose mean values will be denoted respectively with μ_a and μ_{ω} and whose variances will be denoted respectively with σ_a^2 and σ_{ω}^2 . We set the parameters in order to be close to a real case: $\mu_a = 0 \ ms^{-2}$, $\sigma_a = 1 \ ms^{-2}$, $\mu_{\omega} = 0 \ deg \ s^{-1}$ and $\sigma_{\omega} = 10 \ deg \ s^{-1}$. The initial vehicle position is at $\tilde{x} = 0$, $\tilde{y} = 0$, $\tilde{z} = 1m$. The initial vehicle speed is $[1, 0, 0]ms^{-1}$ in the \tilde{x} , \tilde{y} , \tilde{z} -frame.

5.1.2 Simulated Sensors

Starting from the performed trajectory, the true angular speed and the linear acceleration are computed at each time step of 0.01s (respectively, at the time step *i*, we denote them with Ω_i^{true} and $A_{v\ i}^{true}$). Starting from them, the IMU sensors are simulated by randomly generating the angular speed and the linear acceleration at each step according to the following: $\Omega_i = N \left(\Omega_i^{true} - \Omega_{bias}, P_{\Omega_i} \right)$ and $A_i = N \left(A_{v\ i}^{true} - A_{gi} - A_{bias}, P_{A_i} \right)$ where:

- N indicates the Normal distribution whose first entry is the mean value and the second its covariance matrix;
- P_{Ω_i} and P_{A_i} are the covariance matrices characterizing the accuracy of the IMU;

- A_{gi} is the gravitational acceleration in the local frame and A_{bias} is the bias affecting the data from the accelerometer;
- Ω_{bias} is the bias affecting the data from the gyroscope.

In all the simulations we set both the matrices P_{Ω_i} and P_{A_i} diagonal and in particular: $P_{\Omega_i} = \sigma_{gyro}^2 I_3$ and $P_{A_i} = \sigma_{acc}^2 I_3$, where I_3 is the identity 3×3 matrix. We considered several values for σ_{gyro} and σ_{acc} , in particular: $\sigma_{gyro} = 1 \ deg \ s^{-1}$ and $\sigma_{acc} = 0.01 \ ms^{-2}$.

Regarding the camera, the provided readings are generated in the following way. By knowing the true trajectory, and the position and the orientation of the planar surface, the true bearing angles of the laser spot in the camera frame are computed¹. They are computed each 0.1s. The parameter L is set equal to 0.3m. Then, the camera readings are generated by adding to the true values zero-mean Gaussian errors whose variance is equal to $(1 \ deg)^2$ for all the readings.

5.1.3 Simulation Results

Figure 3 displays a typical 3D trajectory obtained in our simulations. The figure also displays the planar surface, consisting of a plane, which makes an angle of $\alpha = \frac{\pi}{8} rad = 22.5 deg$ with the gravity.



Figure 3: A typical vehicle trajectory in our simulations.

¹This is obtained also by knowing that the laser pointer has the same orientation as the camera and that it is located at the position [L, 0, 0]

Figures 4 *a* and *b* display the estimated α respectively in the case without and with bias. The values of the biases adopted in our simulations are: $\Omega_{bias} = [0.03 \ 0.03 \ 0.03]^T (deg \ s^{-1})$ and $A_{bias} = [0.03 \ 0.03 \ 0.03]^T (ms^{-2})$. As expected, the estimation in presence of bias becomes worse. However, the error on the estimated α in presence of bias is smaller than 1 deg.



Figure 4: Estimated α in absence (a) and in presence of bias (b) on the inertial data. Blue dots indicate ground true values while red discs indicate the estimated values.

Figures 5, 6, 7 and 8 a and b display respectively the estimated P, R, v_o and d. In each figure, both the cases of unbiased and biased inertial measurements are displayed. We initialized the filter by using a value of the initial observable state which differs from the ground truth by a relative error in the range [10, 20]%.



Figure 5: Estimated P in absence (a) and in presence of bias (b) on the inertial data. Blue dots indicate ground true values while red discs indicate the estimated values.

We also evaluated the robustness of the filter with respect to systematic errors on the imu-camera calibration and laser-camera calibration. Specifically, we performed simulations by introducing errors of one cm and one deg on the



Figure 6: Estimated R in absence (a) and in presence of bias (b) on the inertial data. Blue dots indicate ground true values while red discs indicate the estimated values.



Figure 7: Estimated v_0 in absence (a) and in presence of bias (b) on the inertial data. Blue dots indicate ground true values while red discs indicate the estimated values.



Figure 8: Estimated d in absence (a) and in presence of bias (b) on the inertial data. Blue dots indicate ground true values while red discs indicate the estimated values.

calibration parameters. These systematic errors affect the estimated α (the difference with respect to the ground truth is in the range [4, 6]deg) while for all the other observable modes the effect is negligible (less than one deg for R and P and less than $1 cm s^{-1}$ and 1 cm respectively for v_o and d).

5.2 Preliminary experiments

In this section we provide preliminary results obtained by using a data set provided by the autonomous system laboratory at ETHZ in Zurich. The data are provided together with a reliable ground-truth, which has been obtained by performing the experiments at the ETH Zurich Flying Machine Arena [16], which is equipped with a Vicon motion capture system. As previously said, the observations of the laser spot are simulated. This was possible thanks to the fact that a reliable ground truth was provided together with the inertial data. In particular, given the true trajectory, we simulated the same planar surface described in the previous section. By having the true vehicle configuration it was possible to create the observations performed by the camera on the laser spot produced by a laser pointer as in the simulations (see the last paragraph in 5.1.2).

Figure 9 displays the estimated α . Figures 10 *a* and *b* display the estimated *P* and *R* and figures 11 *a* and *b* display the estimated v_o and *d*. All the observable modes are estimated with very good accuracy. Additionally, we remark that the convergence of the filter occurs in less than half second for all the observable modes.

5.3 Experiments in progress

In this last section we describe our vehicle and sensors. The robot platform is a *Pelican* from *Ascending Technologies* equipped with an Intel Atom processor board (1.6 GHz, 1 GB RAM). Our sensor suite is composed by an Inertial Measurement Unit (3-Axis Gyro, 3-Axis Accelerometer), a monocular camera (Matrix Vision mvBlueFOX, FOV: 130 deg) and a Laser Module (SparkFun



Figure 9: Estimated α in the experiment. Blue dots indicate ground true values while red discs indicate the estimated values.



Figure 10: Estimated P(a) and R(b) in the experiment. Blue dots indicate ground true values while red discs indicate the estimated values.



Figure 11: Estimated $v_o(a)$ and d(b) in the experiment. Blue dots indicate ground true values while red discs indicate the estimated values.

TTL, wavelength: 650nm, poweroutput: 0.45-0.8mW). The Laser module and the monocular camera are mounted on a fixed baseline, and the latter is calibrated using the Camera Calibration Toolbox for Matlab [22]. The calibration between IMU and camera has been performed using the Inertial Measurement Unit and Camera Calibration Toolbox by Lobo [15], whereas the Camera-Laser Module calibration technique we used is described in appendix A. By using this latter technique we obtained the following values: $\theta = 47.1 \ deg, \phi = -3.1 \ deg$ and $L_x = -0.146 \ m, L_y = -0.005 \ m$. The resulting L is 0.100 m.

The IMU provides measurements update at a rate of 100Hz, while the camera framerate is 13Hz. We are currently equipping our lab with a motion capture system to perform experiments.

6 Conclusion

In this paper we considered an aerial vehicle equipped with a monocular camera and inertial sensors. Additionally, a laser pointer is mounted on the vehicle and it produces a laser spot. The laser spot is observed by the monocular camera and it is the unique point feature used in the proposed approach. We focused our attention to the case when the vehicle moves in proximity of a planar surface and in particular when the laser spot belongs to this surface. The paper provided two main contributions. The former is the analytical derivation of all the observable modes, i.e. all the physical quantities that can be determined by only using the inertial data and the camera observation of the laser spot during a short time-interval. Specifically, it is shown that the observable modes are: the distance of the vehicle from the planar surface; the component of the vehicle speed, which is orthogonal to the planar surface; the relative orientation of the vehicle with respect to the planar surface; the orientation of the planar surface with respect to the gravity. Once the observed modes have been derived, a local decomposition of the original system has also been provided. This decomposition separates the observable modes from the rest of the system and allowed us to introduce a simple recursive method to perform the estimation of all the observable modes (second contribution). The use of a virtual laser spot feature



Figure 12: The Pelican quadcopter equipped with a monocular camera, a laser module and passive markers which will be used to collect sensors data and ground truth.

overcomes the limits of feature tracking algorithms and makes our approach suitable to work even in dark or featureless environment. The method is validated by using synthetic data. The validation with real data is in progress. We presented a low-cost low-weight sensor suite and a low computational complexity approach in the framework of aerial navigation. It can be integrated in the framework of autonomous takeoff and landing, safe touch-down and low altitude manoeuvres. However, we want to emphasize that both the paper contributions are very general and can be applied in other frameworks. In particular, in all the environment where GPS is denied and where the most of objects have planar surfaces (e.g. in an indoor or city-like environment). For instance, these paper contributions could be used in the framework of humanoid robotics (where visual and inertial sensing are often adopted and the navigation usually occurs in an indoor environment).

A very important extension of the proposed strategy is to consider laser patterns in order to improve the precision. So far, the simplest case of a pointspot was considered. The performance of the considered system depends on the extrinsic parameters of the laser module with respect to the distance from the planar surface. Therefore we want to consider the option of a laser module with a servomotor in order to improve precision according to the trajectory. We finally want to implement this approach on our quadrotor platform to perform autonomous landing and safe touchdown tasks.

A Camera-laser module calibration

In figure 2, we display the position and the direction of the laser pointer in the original camera frame. The calibration consists in estimating the four parameters L_x , L_y , θ , ϕ . In other words, it consists in estimating the line made by the laser beam in the original camera frame. This line is determined starting from the position of the laser spot in the original camera frame for at least two spots. To have an accurate estimate, the two spots must be as far as possible one each other. The precision can be further improved by considering more than two spots and by finding the best line fit. In order to have the Cartesian coordinates of a single spot in the original camera frame, it suffices to project the spot on a checkerboard. By using the Camera Calibration Toolbox for Matlab [22], it is possible to get the equation of the plane containing the checkerboard in the original camera frame and, known the direction of the spot from the camera measurement, the 3D position is finally obtained.

The chosen camera frame is obtained by rotating the original frame such that in the new frame the z-axis has the same orientation of the laser beam. Additionally, we also require that the laser beam intersects the new x - axis. In other words, we require that the laser beam intersects the new xy-plane in the point $[L, 0, 0]^T$, for a given L. We want to obtain the quaternion q which characterizes this rotation. This will allow us to express the vectors provided by the camera in the chosen frame. Note that, since the two frames only differ by a rotation (i.e., they share the same origin), we are allowed to express the vectors provided by the camera in the new frame, even if these vectors are defined up to a scale. Finally, in this section we want to determine the value of L. As we will see, both the quaternion q and the parameter L only depend on the calibration parameters: L_x , L_y , θ , ϕ .

We start by rotating the original frame of ϕ about its z-axis. The quaternion characterizing this rotation is $q_{z-axis}(\phi) = \cos\left(\frac{\phi}{2}\right) + k\sin\left(\frac{\phi}{2}\right)$. Then, we rotate the frame obtained with this rotation of θ about its y-axis. The quaternion characterizing this rotation is $q_{y-axis}(\theta) = \cos\left(\frac{\theta}{2}\right) + j\sin\left(\frac{\theta}{2}\right)$. Hence, the previous two rotations are characterized by the quaternion $q_{zy} \equiv q_{z-axis}(\phi)q_{y-axis}(\theta)$. The obtained frame has the z-axis aligned with the laser beam. On the other hand, the laser beam does not intersect necessarily the x-axis. In order to obtain this, we have to rotate again the frame about its current z-axis. Let us compute the intersection of the laser beam with the xy-plane. We compute this point in the original frame. By a direct computation we obtain: $r^{inters} = [L_x - \tau_x^2 L_x - \tau_x \tau_y L_y, L_y - \tau_x \tau_y L_x - \tau_y^2 L_y, - \tau_x \tau_z L_x - \tau_y \tau_z L_y]^T$, where $\tau_x = \sin(\theta) \cos(\phi)$, $\tau_y = \sin(\theta) \sin(\phi)$, $\tau_z = \cos(\theta)$. We then compute this vector in the rotated frame by doing the quaternion product: $R_q^{inters} = q_{zy}^* r_q^{inters} q_{zy}$. Let us denote R^{inters} with $[L'_x, L'_y, 0]^T$. We finally have: $L = \sqrt{L'_x^2 + L'_y^2}$ and $q = q_{zy}q_{z-axis}(\alpha)$, where $q_{z-axis}(\alpha) = \cos\left(\frac{\alpha}{2}\right) + k\sin\left(\frac{\alpha}{2}\right)$ and $\alpha = \arctan\left(\frac{L'_y}{L'_x}\right)$.

B Continuous symmetries and observable modes for the system defined in section 2

The system is characterized by the state: $\mathbf{S} = [r_q \ v_q \ q, \ k]^T$, whose dimension is 11 (r_q and v_q are imaginary quaternions, i.e. they are characterized by 3 real numbers; q contains 4 real numbers; k is a real number). The dynamics are given in (1) together with the equation $\dot{k} = 0$ and the observations are given in (3) and (4). In order to compute the Lie derivatives, we need to express the dynamics as follows:

$$\dot{\boldsymbol{S}} = \boldsymbol{f}(\boldsymbol{S}, \boldsymbol{u}) = \boldsymbol{f}_{\boldsymbol{0}}(\boldsymbol{S}) + \sum_{i=1}^{L} \boldsymbol{f}_{i}(\boldsymbol{S})u_{i}$$
(12)

We have L = 6 and the six inputs are the three components of the acceleration, A, and the three components of the angular speed, Ω . Namely: $u_1 = A_x$, $u_2 = A_y$, $u_3 = A_z$, $u_4 = \Omega_x$, $u_5 = \Omega_y$, $u_6 = \Omega_z$. The seven vector functions f_0 , f_1 , ..., f_6 are:

$$\begin{aligned} \mathbf{f_0} &= [v_x, v_y, v_z, 0, 0, -g, \mathbf{0_5}]^T \\ \mathbf{f_1} &= [\mathbf{0_3}, q_t^2 + q_x^2 - q_y^2 - q_z^2, 2q_tq_z + 2q_yq_x, -2q_tq_y + 2q_zq_x, \mathbf{0_5}]^T \\ \mathbf{f_2} &= [\mathbf{0_3}, -2q_tq_z + 2q_yq_x, q_t^2 + q_y^2 - q_z^2 - q_x^2, 2q_tq_x + 2q_zq_y, \mathbf{0_5}]^T \\ \mathbf{f_3} &= [\mathbf{0_3}, 2q_tq_y + 2q_zq_x, -2q_tq_x + 2q_zq_y, q_t^2 + q_z^2 - q_x^2 - q_y^2, \mathbf{0_5}]^T \\ \mathbf{f_4} &= [\mathbf{0_6}, -1/2q_x, 1/2q_t, 1/2q_z, -1/2q_y, 0]^T \\ \mathbf{f_5} &= [\mathbf{0_6}, -1/2q_y, -1/2q_z, 1/2q_t, 1/2q_x, 0]^T \\ \mathbf{f_6} &= [\mathbf{0_6}, -1/2q_z, 1/2q_y, -1/2q_x, 1/2q_t, 0]^T \end{aligned}$$

where we denoted with $\mathbf{0}_n$ the vector line whose dimension is n and whose entries are all zeros.

We must compute the Lie derivatives of the two observation functions given in (3) and (4) with respect to all the vector fields. By a direct computation, performed by using the symbolic Matlab computational tool, we were able to find not more than 6 independent Lie derivatives².

In particular, according to the notation introduced in [17], the system has 5 continuous symmetries which are:

$$\boldsymbol{w_s^1} = \begin{bmatrix} 0, \ 1, \ k, \ \mathbf{0_8} \end{bmatrix}^T$$
$$\boldsymbol{w_s^2} = \begin{bmatrix} 1, \ \mathbf{0_{10}} \end{bmatrix}^T$$
$$\boldsymbol{w_s^3} = \begin{bmatrix} \mathbf{0_4}, \ 1, \ k, \ \mathbf{0_5} \end{bmatrix}^T$$
$$\boldsymbol{w_s^4} = \begin{bmatrix} \mathbf{0_3}, \ 1, \ \mathbf{0_7} \end{bmatrix}^T$$
$$\boldsymbol{w_s^5} = \begin{bmatrix} \mathbf{0_6}, \ q_z - kq_y, \ q_y + kq_z, \ -q_x + kq_t, \ -q_t - kq_x, \ 0 \end{bmatrix}^T$$

²A possible choice of 6 independent Lie derivatives is: $L^0 h_{const}$, $L^0 h_{cam}$, $L^1_{f_0} h_{cam}$, $L^1_{f_4} h_{cam}$, $L^1_{f_5} h_{cam}$, $L^2_{f_0} f_0 h_{cam}$.

The observable modes are all the solutions of the system of partial differential equations associated with the five symmetries. For instance, the equation associated with w_s^1 is $\frac{\partial}{\partial r_y} + k \frac{\partial}{\partial r_z} = 0$ (see [17] for more details). Since this system of partial differential equations consists of 5 equations on a manifold whose dimension is 11, the number of independent solutions is 6 = 11 - 5 [14]. A possible choice of these solutions is:

$$\begin{split} r_z - kr_y \\ v_z - kv_y \\ 2[c_k s_k (-q_t^2 + q_x^2 - q_y^2 + q_z^2) + (2c_k^2 - 1)(q_t q_x + q_y q_z)] \\ 4c_k s_k (q_t q_z + q_x q_y) + 2(2c_k^2 - 1)(q_t q_y - q_x q_z) \\ k \\ q^* q \end{split}$$

where $c_k \equiv \cos\left(\frac{\arctan(k)}{2}\right)$ and $s_k \equiv \sin\left(\frac{\tan(k)}{2}\right)$. By knowing the value of the first solution, $r_z - kr_y$ and the value of the fifth solution, k, we can determine the quantity $\frac{|r_z - kr_y|}{\sqrt{1+k^2}}$, which is the distance of the vehicle from the planar surface³. Hence, to better visualize the physical meaning, it is convenient to select the following 6 observable modes:

$$m_{1} = \frac{r_{z} - kr_{y}}{\sqrt{1 + k^{2}}}$$

$$m_{2} = \frac{v_{z} - kv_{y}}{\sqrt{1 + k^{2}}}$$

$$m_{3} = 2[c_{k}s_{k}(-q_{t}^{2} + q_{x}^{2} - q_{y}^{2} + q_{z}^{2}) + (2c_{k}^{2} - 1)(q_{t}q_{x} + q_{y}q_{z})]$$

$$m_{4} = 4c_{k}s_{k}(q_{t}q_{z} + q_{x}q_{y}) + 2(2c_{k}^{2} - 1)(q_{t}q_{y} - q_{x}q_{z})$$

$$m_{5} = k$$

$$m_{6} = q^{*}q$$

Finally, the physical meaning of also m_3 and m_4 becomes clear by referring to a new global frame \tilde{x} , \tilde{y} , \tilde{z} . This frame has the $\tilde{x}\tilde{y}$ -plane coincident with the planar surface. In other words, this new global frame has the vertical axis coincident with the axis orthogonal to the planar surface. In this new frame, m_1 is the \tilde{z} -coordinate of the vehicle, m_2 is the component of the vehicle speed along the \tilde{z} -axis. m_3 and m_4 are related to the roll and pitch angles of the vehicle in the new frame. In particular, the roll angle is $\arctan\left(\frac{m_3}{\sqrt{1-m_3^2-m_4^2}}\right)$ and the pitch is $\arcsin(m_4)$. m_5 is related to the orientation of the $\tilde{x}\tilde{y}$ -plane with respect to the gravity. In particular, the \tilde{z} -axis makes an angle $\arctan(k) = \arctan(m_5)$ with the gravity. m_6 is trivially the magnitude of the quaternion, which is 1 since it describes a rotation.

From now on, we will adopt the new frame to characterize the vehicle configuration. The state in this frame is $\tilde{\boldsymbol{S}} = [\tilde{r}_q \ \tilde{v}_q \ \tilde{q}, \ k]^T$. The $\tilde{x}, \ \tilde{y}, \ \tilde{z}$ -frame is

³In other words, also $\frac{r_z - kr_y}{\sqrt{1 + k^2}}$ is a solution of the system of partial differential equations.

obtained by rotating the x, y, z-frame about the x-axis of the angle $\arctan(k)$. Hence, it is characterized by the quaternion:

$$p = \cos\left(\frac{\arctan(k)}{2}\right) + i\sin\left(\frac{\arctan(k)}{2}\right) \tag{13}$$

Therefore, $q = p\tilde{q}$ or:

$$\tilde{q} = p^* q \tag{14}$$

By using the quaternion p it is also possible to obtain:

$$\tilde{r}_q = p^* r_q p \qquad \tilde{v}_q = p^* v_q p \tag{15}$$

By using (14) and (15), we obtain the expressions of the observable modes in the new coordinates. We have:

$$m_1 = \tilde{r}_z$$

$$m_2 = \tilde{v}_z$$

$$m_3 = 2(\tilde{q}_t \tilde{q}_x + \tilde{q}_y \tilde{q}_z)$$

$$m_4 = 2(\tilde{q}_t \tilde{q}_y - \tilde{q}_x \tilde{q}_z)$$

$$m_5 = k$$

$$m_6 = \tilde{q}^* \tilde{q}$$

C EKF-equations

Let us consider the state $\boldsymbol{m} = [m_1, m_2, m_3, m_4, g_z]$. The basic ingredients to implement an *EKF*, which estimates \boldsymbol{m} , are the analytical expression of the Jacobians of the dynamics, and the observation [2].

The Jacobian of the observation is obtained by differentiating the expression of h_{cam} given in (11) with respect to m, i.e.

$$H \equiv \frac{\partial h_{cam}}{\partial m} = \frac{L\xi(m_3, m_4)}{(m_4 L - m_1)^2} \times \times \left[1 \quad 0 \quad \frac{-m_3(m_4 L - m_1)}{1 - m_3^2 - m_4^2} \quad \frac{Lm_3^2 - L + m_1 m_4}{1 - m_3^2 - m_4^2} \quad 0 \right]$$

where the function $\xi(m_3, m_4)$ is defined in (10). Regarding the Jacobian of the dynamics, we need first of all to discretize the equations in (9). Let us denote with δt the discretization time step. The Jacobian of the dynamics with respect to the state m is:

$$F_m = \begin{bmatrix} 1 & \delta t & 0 & 0 & 0 \\ 0 & 0 & \delta t \left(A_y - A_z \overline{m}_3\right) & -\delta t \left(A_x + A_z \overline{m}_4\right) & 0 \\ 0 & 0 & -\delta t \Omega_x \overline{m}_3 & \delta t (\Omega_z - \Omega_x \overline{m}_4) & 0 \\ 0 & 0 & -\delta t (\Omega_z + \Omega_y \overline{m}_3) & -\delta t \Omega_y \overline{m}_4 & 0 \\ 0 & 0 & 0 & 0 & 1 \end{bmatrix}$$

where $\overline{m}_3 = \frac{m_3}{\xi(m_3, m_4)}$ and $\overline{m}_4 = \frac{m_4}{\xi(m_3, m_4)}$. The Jacobian of the dynamics with respect to the input $\boldsymbol{u} = [A_x, A_y, A_z, \Omega_x, \Omega_y, \Omega_z]^T$ is:

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$$F_u = \left[\begin{array}{ccccccc} 0 & 0 & 0 & 0 & 0 & 0 \\ -m_4 & m_3 & \xi & 0 & 0 & 0 \\ 0 & 0 & 0 & \xi & 0 & m_4 \\ 0 & 0 & 0 & 0 & \xi & -m_3 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{array} \right]$$

where $\xi \equiv \xi(m_3, m_4)$.

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